

Dynamics of diatoms in a turbulent flow

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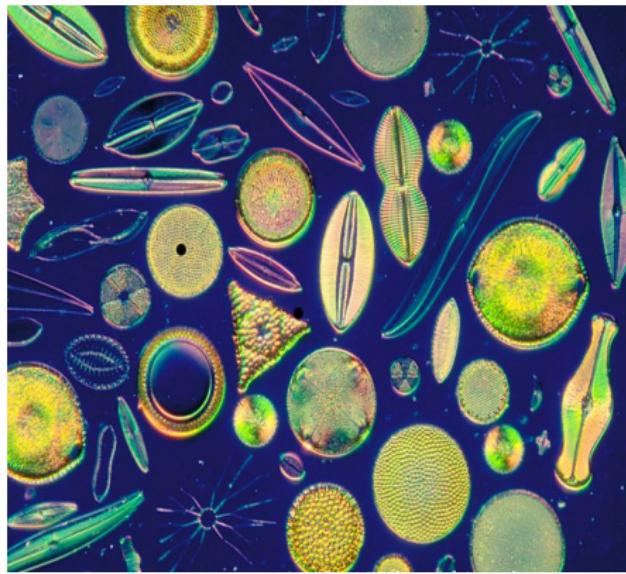
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Introduction

Diatoms

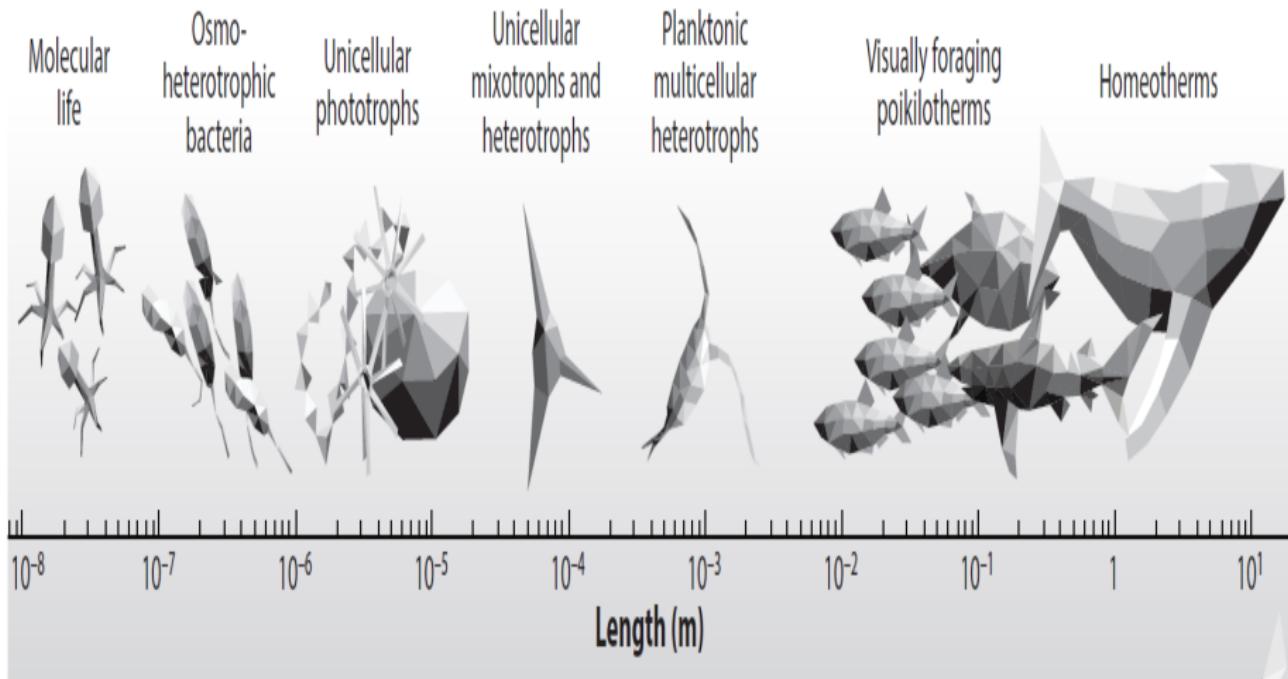


A family of plankton called **Phytoplankton**.

Why Phytoplanktons ?
As they are **unicellular**,
free floating and **non swimming**
microrganism.

How big are diatoms ?

■ Roughly $2 \mu\text{m}$ to $500 \mu\text{m}$

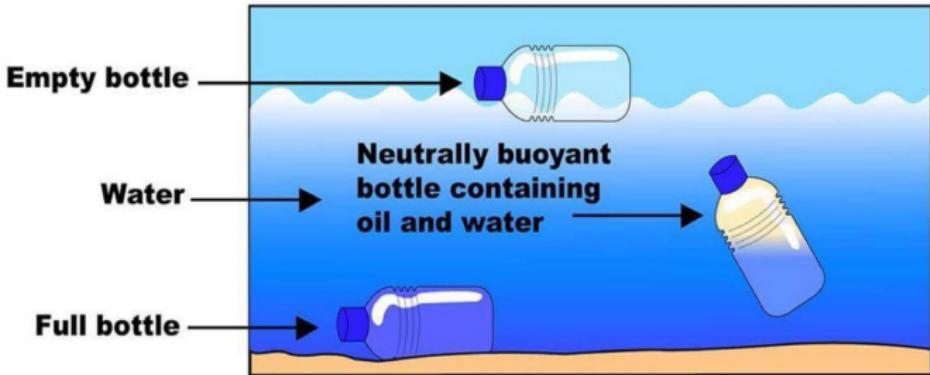


Non motile and neutrally buoyant

- Move passively by water turbulence.
- Neither sinks nor floats.

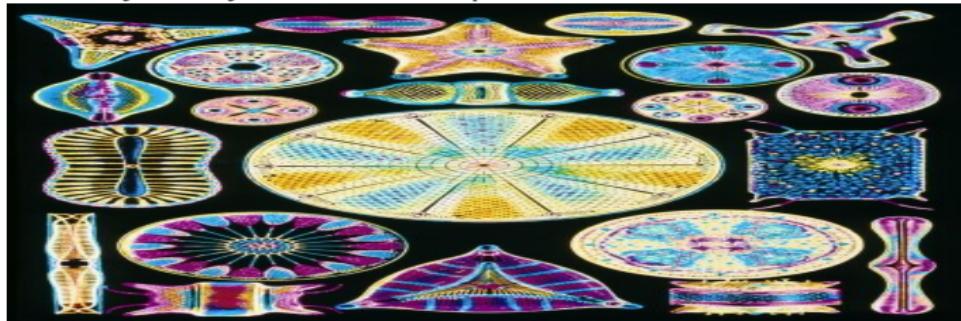
Neutrally buoyant

weight of object = Upward buoyant force exerted by water



Shape and classes of diatoms ?

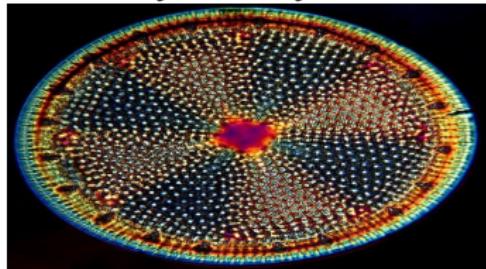
- Usually axisymmetric shape.



Two general category of diatoms.

Centrales

Radial symmetry



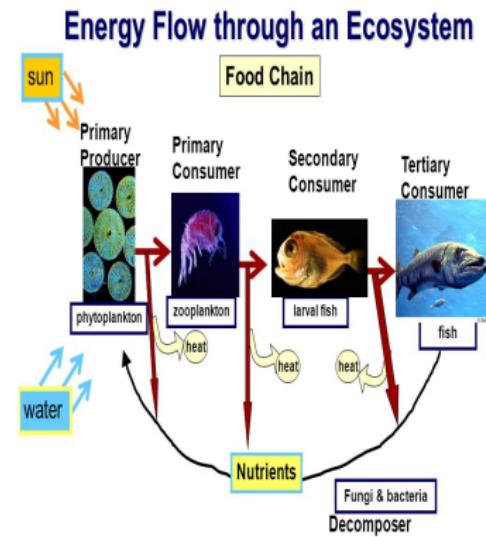
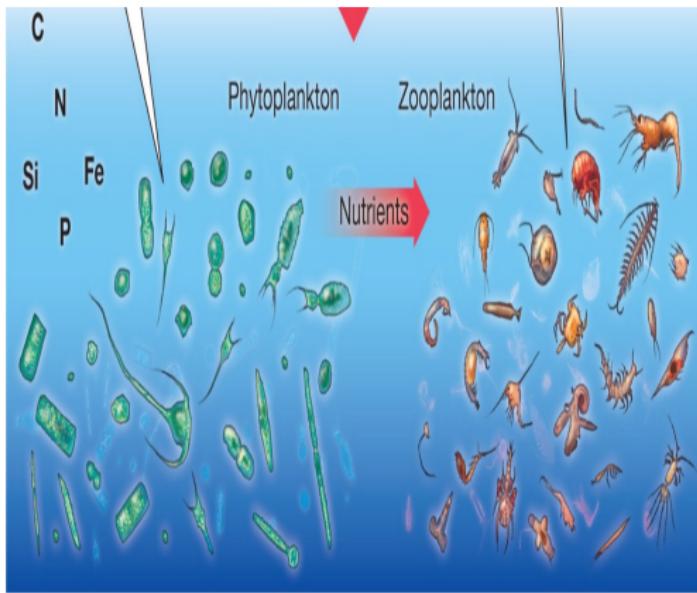
Pennales

Bilateral symmetry



Feeding habits

- Phytoplankton eats nutrients, nutrients are more smaller than phytoplankton.
- Can be modelled as passive tracers or scalar field which are subjected to diffusion.



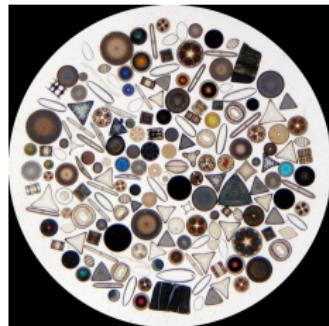
Biophysical habitat of plankton particles

- Homogeneous turbulent flow
- $Re_\lambda \approx 100$ and $\eta \approx 200\mu m$
- Size of the nutrients
- Passive scalar



Why this shape ?

- For collection nutrients
- Adaptation to turbulent environment.
- Evolution



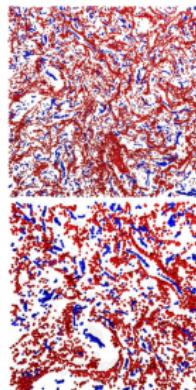
Shapes of planktons



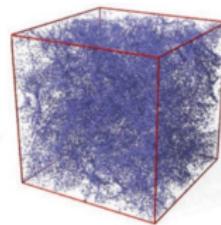
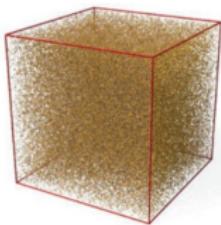
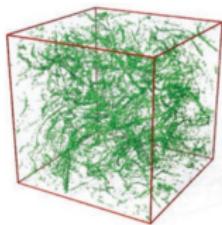
Turbulent environment

Dispersion of inertial particles in turbulent flow

- Preferential concentration based on inertia
- Alignment of the diatoms



Distribution of the particles w.r.t density



Lower to higher density of particles

Modeling of environment

- Navier stokes for fluid simulation + forcing

$$\frac{\partial u_i}{\partial t} + u_j \frac{\partial u_i}{\partial x_j} = -\frac{1}{\rho} \frac{\partial p}{\partial x_i} + \nu \frac{\partial^2 u_i}{\partial x_j^2} + F \quad (1)$$

$$\frac{\partial u_i}{\partial x_i} = 0 \quad (2)$$

- Advection diffusion equation

$$\frac{\partial}{\partial t} C + u_j \cdot \frac{\partial}{\partial x_j} C = \alpha \frac{\partial^2}{\partial x_j^2} C + S \quad (3)$$

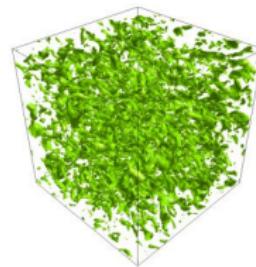
where F is the forcing term and S is the source term for the scalar

- Refinement due to variation of force and source in the terms with depth

Assumptions and Consideration for simulation

■ Fluid assumption

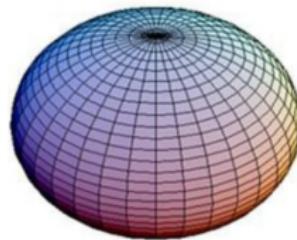
- Incompressible
- Homogeneous
- Isotropic



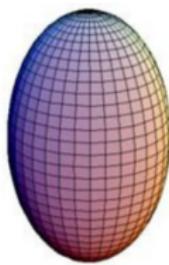
Simulation of turbulence
(Visualization of vorticity)

■ Particle assumption

- Shape of particle
(Oblate and Prolate)
- Symmetry is shape
- One way coupling



oblate spheroid



prolate spheroid

Modeling of diatoms w.r.t motion of center of mass

- Newton equation

$$m \cdot \frac{dv_p}{dt} = F \quad (4)$$

- Drag force³

$$F = \mu\pi a \cdot R K^b R^T (u_f - v_p) \quad (5)$$

where K is the resistance tensor.

$$K_{xx}^b = K_{yy}^b = \frac{16(\beta^2 - 1)}{(2\beta^2 - 3) \ln \frac{\beta + \sqrt{\beta^2 - 1}}{\sqrt{\beta^2 - 1}} + \beta}$$

$$K_{zz}^b = \frac{8(\beta^2 - 1)}{(2\beta^2 - 1) \ln \frac{\beta + \sqrt{\beta^2 - 1}}{\sqrt{\beta^2 - 1}} - \beta}$$

β = aspect ratio

3. Gallily, I. and Cohen, A. (1979). On the orderly nature of the motion of nonspherical aerosol particles. II. Inertial collision between a spherical large droplet and an axially symmetrical elongated particle. Journal of Colloid and Interface Science, 68(2), 338(356).

Modeling of diatoms w.r.t orientation

- Euler equation

$$\frac{d}{dt}(I \cdot \omega) + \omega \times (I \cdot \omega) = T \quad (6)$$

- Torque⁴

$$T_i = \frac{16 \times \pi \times \mu \times a^3 \times \beta}{3 \times (\beta_0 + \beta^2 \cdot \gamma_0)} \times ((1 - \beta^2) \times S_{kj} + (1 + \beta^2) \times (\omega_{jk} - \omega_i)) \quad (7)$$

where,

β = aspect ratio

ω_{jk} = strain rate tensor $(\frac{1}{2}(\frac{du_k}{dx_j} + \frac{du_j}{dx_k}))$

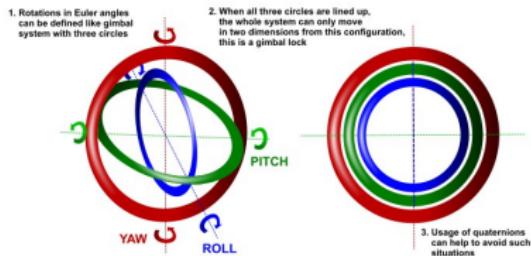
S_{kj} = rotation tensor $(\frac{1}{2}(\frac{du_k}{dx_j} - \frac{du_j}{dx_k}))$

- Homogeneous in density

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4. Jeffery, G. B. (1922). The motion of ellipsoidal particles immersed in a viscous fluid. Royal Society of London. Series A., 102(715), 161(179).

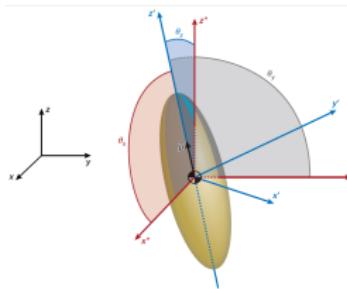
Difficulty in the numerical simulation

- 12 unknown variables
 $(x_i, v_i, \omega_i, \theta, \psi, \phi)$
- Lab frame to body frame
- Singularity problem (Rotation Matrix)
- Gimbal lock (Euler angles)



Gimbal lock configuration

Lab frame(x, y, z) and body frame(x'', y'', z'')



Coordinate frames

Solution and optimization

■ Quaternions

$$q = [q_0, q_1, q_2, q_3] \quad (8)$$

Property of quaternions

$$q_0^2 + q_1^2 + q_2^2 + q_3^2 = 1 \quad (9)$$

Rotation by quaternions⁵

$$K' = RKR^T \quad (10)$$

$$R = \begin{bmatrix} 1 - 2(q_2^2 - q_3^2) & 2 \times (q_1 \times q_2 + q_0 q_3) & 2(q_1 q_3 - q_0 q_2) \\ 2(q_1 q_2 - q_0 q_3) & 1 - 2(q_1^2 + q_3^2) & 2(q_2 q_3 + q_0 q_1) \\ 2(q_1 q_3 + q_0 q_2) & 2(q_2 q_3 - q_0 q_1) & 1 - 2(q_1^2 + q_2^2) \end{bmatrix} \quad (11)$$

■ Integration method

$$\dot{q} = \frac{1}{2}q \times \omega \quad (12)$$

-
5. Betsch, P. and Siebert, R. (2009). Rigid body dynamics in terms of quaternions : Hamiltonian formulation and conserving numerical integration. International Journal for Numerical Methods in Engineering, 79(4), 444(473)

Simpler approach

What is the simplified way to approach this problem ?

- Introducing the problem by neglecting inertia.

$$\frac{dx}{dt} = u \quad (13)$$

u – is the velocity of fluid.

Jeffery equation (1922)

$$\dot{p}_i = \Omega_{ij} p_j + \frac{\alpha^2 - 1}{\alpha^2 + 1} (S_{ij} p_j - p_i p_k S_{kl} p_l) \quad (14)$$

where,

$\alpha = I/d$ is the aspect ratio

Ω_{ij} is the rate of rotation tensor

S_{ij} is the rate of strain tensor

\dot{p}_i is the component of orientation vector which is along the axis of symmetry of particle

Simpler approach

- These will produce particle with complex dynamics.

6. [Parsa, Calzavarini et al PRL 2012,
Rotation rate of rods in turbulent fluid flow]

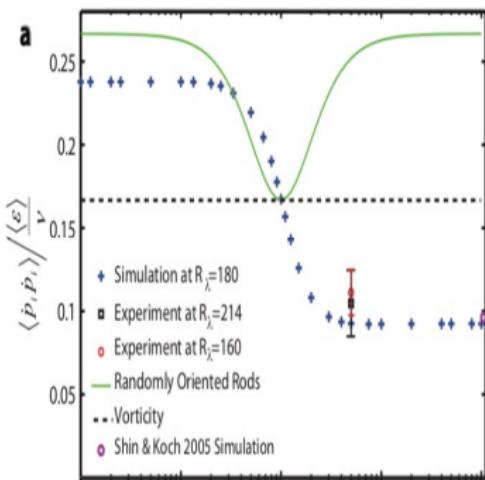
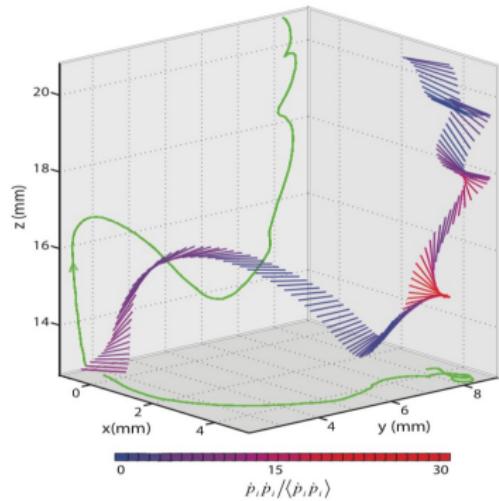
Simpler approach

- These will produce particle with complex dynamics.
- Particles distribute **homogeneously**, so is need to use **inertia**, also the orientation is not revealed.
- Particles gets aligned to the vorticity vector.

6. [Parsa, Calzavarini et al PRL 2012, Rotation rate of rods in turbulent fluid flow]

Simpler approach

- These will produce particle with complex dynamics.
- Particles distribute **homogeneously**, so is need to use **inertia**, also the orientation is not revealed.
- Particles gets aligned to the vorticity vector.
- Disks like particle rotate much faster then elongated particles.



6. [Parsa, Calzavarini et al PRL 2012, Rotation rate of rods in turbulent fluid flow]

Simpler approach,(with gravity)

- Another approach could be account for inertia only in center of mass.
- Based on observation that rotational relaxation time, are shorter than transnational relaxation time. It is not known in general, but few papers reveals this key information.

In recent works on diatoms obeys :⁷

$$\frac{dx}{dt} = u|_x + v_s(p) \quad (15)$$

$$v_s(p) = v_s^{\min} \hat{e}_g + \left(v_s^{\max} - v_s^{\min} \right) (\hat{e}_g \cdot p) p \quad (16)$$

where,

x is the spheroid position

$u|x$ is the fluid velocity at spheriod position.

v_s is the Stokes settling velocity.

p is the spheroid orientation position.

7. [M.Niazi Ardekani et al J.Fluid Mech. 2017,

Sedimentation of inertialess prolate,
spheroid in HIT with application to non motile phytoplankton]

Simpler approach,(with drag)

The fact that $\tau_R < \tau_T$ ⁸ allows us to propose this model.

- Newton equation

$$m \cdot \frac{d\boldsymbol{v}_p}{dt} = \mathbf{F} \quad (17)$$

- Drag force

$$\mathbf{F} = \mu \pi a \cdot \mathbf{R} K^b \mathbf{R}^T (\boldsymbol{u}_f - \boldsymbol{v}_p) \quad (18)$$

where, K is the resistance tensor.

For the rotation we use Jeffery equation.⁹

$$\dot{p}_i = \Omega_{ij} p_j + \frac{\alpha^2 - 1}{\alpha^2 + 1} (S_{ij} p_j - p_i p_k S_{kl} p_l) \quad (19)$$

8. Zhao et al, PRL 2015, Rotation of non spherical particles in turbulent channel flow

9. Jeffery,1922, The motion of ellipsoidal particle immersed in viscous fluid

Summary and Perspectives

■ Summary

- Proposed a model in turbulent environment with nutrients
- The model is Eulerian and Lagrangian
- We proposed a simplified model based on the idea that $\tau_R < \tau_T$

■ Perspective

- Parametric study of aspect ratio and Stokes number
- Study the difference between two models (select better model)
- Draw ecological conclusion

Questions

Thank you for your attention.

Do you have any questions ?