A Performance and Energy Study of the Hyperbolic PDE Solver Engine ExaHyPE

Master's Thesis in Computational Science and Engineering

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September 2016

Supervisor: Univ.-Prof. Dr. Michael Bader
Dr. Tobias Weinzierl
Advisor: Dr. Vasco Varduhn

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Abstract

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Chapter 1

Introduction

- Challenges of exascale
- The ExaHyPE project (numerics, resilience, profiling) as an answer
- On the importance of profiling and performance measuring

Chapter 2

Theory

2.1 A D-dimensional ADER-DG scheme with MUSCL-Hancock a-posteriori subcell limiting for non-linear hyperbolic conservation laws

Arbitrary High Order Derivatives Discontinuous Galerkin (ADER-DG)

2.1.1 Notation

We use vector notation whenever possible. Advantage: Complete derivation, direct conversion to code.

2.1.2 PDE

Task: Solve the PDE

$$\frac{\partial}{\partial t}\left[\boldsymbol{u}\right]_v + \frac{\partial}{\partial x_d}\left[\boldsymbol{F}(\boldsymbol{u})\right]_{vd} = \left[\boldsymbol{s}(\boldsymbol{u})\right]_v \text{ on } \boldsymbol{\Omega} \times (0, T) \tag{2.1}$$

with initial conditions

$$\left[\boldsymbol{u}(\boldsymbol{x}, 0)\right]_v = \left[\boldsymbol{u}_0(\boldsymbol{x})\right]_v \ \forall \boldsymbol{x} \in \boldsymbol{\Omega}, \tag{2.2}$$

and boundary conditions

$$\left[\boldsymbol{u}(\boldsymbol{x}, t)\right]_v = \left[\boldsymbol{u}_B(\boldsymbol{x}, t)\right]_v \ \forall \boldsymbol{x} \in \partial\boldsymbol{\Omega}, t \in (0, T), \tag{2.3}$$

for all $v \in \{1, 2, \ldots, V\}$, where V is the number of quantities involved in the system, $\boldsymbol{\Omega} \subset \mathbb{R}^D$ is the spatial domain, D the number of space dimensions, and $(0, T)$ a time interval. The function $\boldsymbol{F} : \mathbb{R}^V \to \mathbb{R}^{V \times D}, \boldsymbol{u} \mapsto \boldsymbol{F}(\boldsymbol{u}) = \left[\boldsymbol{f}_1(\boldsymbol{u}), \boldsymbol{f}_2(\boldsymbol{u}), \ldots, \boldsymbol{f}_D(\boldsymbol{u})\right]$ is called the flux function.

For the problem to be hyperbolic we require that all Jacobian matrices $\boldsymbol{A}_d(\boldsymbol{x}, t)$, $d \in \{1, 2, \ldots, D\}$, defined as

$$[\boldsymbol{A}_d]_{ij} = \frac{\partial \left[\boldsymbol{f}_d\right]_i}{\partial x_j}, \tag{2.4}$$

have D real eigenvalues in each admissible state $(\boldsymbol{x}, t) \in \boldsymbol{\Omega} \times (0, T)$.

2.1.3 Mesh

Let $\mathcal{T}_h$ be a quadrilateral partition of $\boldsymbol{\Omega}$, i.e.

$$K \cap J = \emptyset \, \forall K, J \in \mathcal{T}_h, K \neq J \tag{2.5}$$

$$\bigcup_{K \in \mathcal{T}_h} K = \boldsymbol{\Omega}. \tag{2.6}$$

Let $\{t_i\}_{i=0,1,\ldots I}$ be a partition of the time interval $(0, T)$ such that

$$0 = t_0 < t_1 < \ldots < t_I = T, \tag{2.7}$$

where I is the number of sub intervals. We furthermore define

$$\Delta t_i = t_{i+1} - t_i, \; i \; in \, \{0, 1, \ldots, I-1\}, \tag{2.8}$$

so that the interval (t_i, t_{i+1}) can be written as $(t_i, t_i + \Delta t_i)$.

2.1.4 Weak formulation

Let $L^2(\boldsymbol{\Omega})^V$ be the space of vector-valued, square-integrable functions on $\boldsymbol{\Omega}$, i.e.

$$L^2(\boldsymbol{\Omega})^V = \left\{ \boldsymbol{w} : \Omega \to \mathbb{R}^V \mid \int_{\boldsymbol{\Omega}} \|\boldsymbol{w}\| \, d\boldsymbol{x} < \infty \right\}. \tag{2.9}$$

Let $\boldsymbol{w} \in L^2(\boldsymbol{\Omega})^V$ be a spatial test function. Multiplication of the original PDE (2.1) and integration over a space-time cell $K \times (t_i, t_i + \Delta t_i)$ yields a weak, element local formulation of the problem

$$\begin{aligned} &\int_{t_i}^{t_i+\Delta t_i} \int_K \frac{\partial}{\partial t} \left[\boldsymbol{u}\right]_v \left[\boldsymbol{w}\right]_v \, d\boldsymbol{x} \, dt + \int_{t_i}^{t_i+\Delta t_i} \int_K \frac{\partial}{\partial x_d} \left[\boldsymbol{F}(\boldsymbol{u})\right]_{vd} \left[\boldsymbol{w}\right]_v \, d\boldsymbol{x} \, dt = \\ &\int_{t_i}^{t_i+\Delta t_i} \int_K \left[\boldsymbol{s}(\boldsymbol{u})\right]_v \left[\boldsymbol{w}\right]_v \, d\boldsymbol{x} \, dt, \end{aligned} \tag{2.10}$$

which we require to hold for $v \in \{1, 2, \ldots, V\}$, $\boldsymbol{w} \in L^2(\Omega)^V$, $K \in \mathcal{T}_h$ and $i \in \{0, 1, \ldots, I-1\}$.

Integration by parts of the spatial integral in the second term yields

$$\begin{aligned}&\int_K \frac{\partial}{\partial x_d} \left[\boldsymbol{F}(\boldsymbol{u})\right]_{vd} [\boldsymbol{w}]_v \, d\boldsymbol{x} = \\ &\int_K \frac{\partial}{\partial x_d} \left(\left[\boldsymbol{F}(\boldsymbol{u})\right]_{vd} [\boldsymbol{w}]_v \right) d\boldsymbol{x} - \int_K \left[\boldsymbol{F}(\boldsymbol{u})\right]_{vd} \frac{\partial}{\partial x_d} [\boldsymbol{w}]_v \, d\boldsymbol{x}.\end{aligned} \tag{2.11}$$

Application of the divergence theorem to the first term on the right-hand side of (2.11) yields

$$\int_K \frac{\partial}{\partial x_d} \left(\left[\boldsymbol{F}(\boldsymbol{u})\right]_{vd} [\boldsymbol{w}]_v \right) d\boldsymbol{x} = \int_{\partial K} \left[\boldsymbol{F}(\boldsymbol{u})\right]_{vd} [\boldsymbol{w}]_v [\boldsymbol{n}]_d \, ds(\boldsymbol{x}), \tag{2.12}$$

where $\boldsymbol{n} \in \mathbb{R}^D$ is the unit-length, outward-pointing normal vector at a point $\boldsymbol{x}$ on the surface of K, which we denote by ∂K.

Inserting eqs. (2.11) and (2.12) into eq. (2.10) yields the following weak, element-local formulation of the original equation (2.1):

$$\begin{aligned}&\int_{t_i}^{t_i+\Delta t_i} \int_K \frac{\partial}{\partial t} [\boldsymbol{u}]_v [\boldsymbol{w}]_v \, d\boldsymbol{x}\, dt - \int_{t_i}^{t_i+\Delta t_i} \int_K \left[\boldsymbol{F}(\boldsymbol{u})\right]_{vd} \frac{\partial}{\partial x_d} [\boldsymbol{w}]_v \, d\boldsymbol{x}\, dt + \\ &\int_{t_i}^{t_i+\Delta t_i} \int_{\partial K} \left[\boldsymbol{F}(\boldsymbol{u})\right]_{vd} [\boldsymbol{w}]_v [\boldsymbol{n}]_d \, ds(\boldsymbol{x})\, dt = \int_{t_i}^{t_i+\Delta t_i} \int_K \left[\boldsymbol{s}(\boldsymbol{u})\right]_v [\boldsymbol{w}]_v \, d\boldsymbol{x}\, dt.\end{aligned} \tag{2.13}$$

Again we require the weak formulation to hold for all $v \in \{1, 2, \dots V\}$, $\boldsymbol{w} \in L^2(\boldsymbol{\Omega})^V$, $K \in \mathcal{T}_h$ and $i \in \{0, 1, \dots, I-1\}$.

2.1.5 Restriction to finite-dimensional function spaces

To discretize eq. (2.13) we need to impose the restriction that both test and ansatz functions come from a finite-dimensional space. First, let $\mathbb{Q}_N(K)^V$ and $\mathbb{Q}_N(K \times (t_i, t_i + \Delta t_i))^V$ be the space of vector-valued, multivariate polynomials of degree less or equal N in each variable on K and $K \times (t_i, t_i + \Delta t_i)$, respectively. We then make the following choices:

- For spatial functions we restrict ourselves to

$$\mathbb{W}_h = \left\{ \boldsymbol{w}_h \in L^2(\boldsymbol{\Omega})^V : \boldsymbol{w}_h|_K := \boldsymbol{w}_h^K \in \mathbb{Q}_N(K)^V \, \forall K \in \mathcal{T}_h \right\}. \tag{2.14}$$

- For space-time functions we restrict ourselves to

$$\begin{aligned}\tilde{\mathbb{W}}_h^i = \Big\{ &\tilde{\boldsymbol{w}}_h^i \in L^2 \left(\boldsymbol{\Omega} \times (t_i, t_i + \Delta t_i) \right) : \\ &\tilde{\boldsymbol{w}}_h^i \Big|_K := \tilde{\boldsymbol{w}}_h^{Ki} \in \mathbb{Q}_N \left(K \times (t_i, t_i + \Delta t_i) \right) \forall K \in \mathcal{T}_h \Big\}\end{aligned} \tag{2.15}$$

  for all $i \in \{0, 1, \dots, I-1\}$.

Replacing $\boldsymbol{w}$ by $\boldsymbol{w}_h \in \mathbb{W}_h$ and $\boldsymbol{u}$ by $\tilde{\boldsymbol{u}}_h^i \in \tilde{\mathbb{W}}_h^i$ in eq. (2.13) yields a finite-dimensional approximation of the weak formulation

$$\int_{t_i}^{t_i+\Delta t_i} \int_K \frac{\partial}{\partial t} \left[\tilde{\boldsymbol{u}}_h^{Ki}\right]_v \left[\boldsymbol{w}_h^K\right]_v dx\, dt - \int_{t_i}^{t_i+\Delta t_i} \int_{\partial K} \left[\boldsymbol{F}(\tilde{\boldsymbol{u}}_h^{Ki})\right]_{vd} \frac{\partial}{\partial x_d} \left[\boldsymbol{w}_h^K\right]_v dx\, dt +$$
$$\int_{t_i}^{t_i+\Delta t_i} \int_{\partial K} \left[\boldsymbol{\mathcal{G}}(\tilde{\boldsymbol{u}}_h^{Ki}, \tilde{\boldsymbol{u}}_h^{K^+i}, \boldsymbol{n})\right]_v \left[\boldsymbol{w}_h^K\right]_v ds(\boldsymbol{x})\, dt =$$
$$\int_{t_i}^{t_i+\Delta t_i} \int_K \left[\boldsymbol{s}(\tilde{\boldsymbol{u}}_h^{Ki})\right]_v \left[\boldsymbol{w}_h^K\right]_v dx\, dt, \tag{2.16}$$

which now has to hold for all $\boldsymbol{w}_h \in \mathbb{W}_h$, $K \in \mathcal{T}_h$ and $i \in \{0, 1, \ldots, I-1\}$. Since for a cell $K \in \mathcal{T}_h$ and one of its Voronoi neighbors $K' \in \mathcal{V}(K)$ one has

$$\tilde{\boldsymbol{u}}_h^{Ki}(\boldsymbol{x}) \neq \tilde{\boldsymbol{u}}_h^{K'i}(\boldsymbol{x}),\ \boldsymbol{x} \in K \cap K', \tag{2.17}$$

i.e. $\tilde{\boldsymbol{u}}_h^i$ is double-valued at the interface between K and K', in order to compute the surface integral we need to introduce the numerical flux function $\boldsymbol{\mathcal{G}}(\tilde{\boldsymbol{u}}_h^{Ki}, \tilde{\boldsymbol{u}}_h^{K'i}, \boldsymbol{n})$. The numerical flux at a position $\boldsymbol{x} \in K \cap K'$ on the interface is obtained by (approximately) solving a Riemann problem in normal direction.

Integration by parts in time of the first term of eq. (2.16) and noting that $\boldsymbol{w}_h$ is constant in time yields the following one-step update scheme for the cell-local time-discrete solution $\tilde{\boldsymbol{u}}_h^{Ki}$:

$$\int_K \left[\tilde{\boldsymbol{u}}_h^{Ki}\Big|_{t_i+\Delta t_i}\right]_v \left[\boldsymbol{w}_h^K\right]_v d\boldsymbol{x} = \int_K \left[\tilde{\boldsymbol{u}}_h^{Ki}\Big|_{t_i}\right]_v \left[\boldsymbol{w}_h^K\right]_v d\boldsymbol{x} +$$
$$\int_{t_i}^{t_i+\Delta t_i} \int_K \left[\boldsymbol{F}(\tilde{\boldsymbol{u}}_h^{Ki})\right]_{vd} \frac{\partial}{\partial x_d} \left[\boldsymbol{w}_h^K\right]_v d\boldsymbol{x}\, dt -$$
$$\int_{t_i}^{t_i+\Delta t_i} \int_{\partial K} \left[\boldsymbol{\mathcal{G}}(\tilde{\boldsymbol{u}}_h^{Ki}, \tilde{\boldsymbol{u}}_h^{K^+i}, \boldsymbol{n})\right]_v \left[\boldsymbol{w}_h^K\right]_v ds(\boldsymbol{x})\, dt +$$
$$\int_{t_i}^{t_i+\Delta t_i} \int_K \left[\boldsymbol{s}(\tilde{\boldsymbol{u}}_h^{Ki})\right]_v \left[\boldsymbol{w}_h^K\right]_v d\boldsymbol{x}\, dt. \tag{2.18}$$

Again we require eq. (2.18) to hold for all $v \in \{1, 2, \ldots, V\}$, $\boldsymbol{w}_h \in \mathbb{W}_h$, $K \in \mathcal{T}_h$ and $i \in \{0, 1, \ldots, I-1\}$.

Problem: We only have $\tilde{\boldsymbol{u}}_h^i\Big|_t$ at the discrete time steps $t \in \{t_i, t_i + \Delta t_i\}$, not within the open interval, i.e. for $t \in (t_i, t_i + \Delta t_i)$.

Idea: Replace $\tilde{\boldsymbol{u}}_h$ in $K \times (t_i, t_i + \Delta t_i)$ by an approximation $\tilde{\boldsymbol{q}}_h^i \in \tilde{\mathbb{W}}_h^i$ which we call space-time predictor.

2.1.6 Space-time predictor

To derive a procedure to compute the space-time predictor $\tilde{\boldsymbol{q}}_h^i \in \tilde{\mathbb{W}}_h^i$ we again start from the original PDE (2.1), but this time we do not use a spatial test function $\boldsymbol{w}_h \in \mathbb{W}_h$, but a space-time test function $\tilde{\boldsymbol{w}}_h^i \in \tilde{\mathbb{W}}_h^i$. If we furthermore replace the solution $\boldsymbol{u}$ by the the space-time predictor $\tilde{\boldsymbol{q}}_h^i \in \tilde{\mathbb{W}}_h^i$, integrate over the space-time element $K \times (t_i, t_i + \Delta t_i)$ and apply the divergence theorem analogously to eq. (2.12) we obtain the following relation:

$$
\begin{aligned}
&\int_{t_i}^{t_i+\Delta t_i} \int_K \frac{\partial}{\partial t} \left[\tilde{\boldsymbol{q}}_h^{Ki}\right]_v \left[\tilde{\boldsymbol{w}}_h^{Ki}\right]_v \, d\boldsymbol{x}\, dt - \\
&\int_{t_i}^{t_i+\Delta t_i} \int_K \left[\boldsymbol{F}(\tilde{\boldsymbol{q}}_h^{Ki})\right]_{vd} \frac{\partial}{\partial x_d} \left[\tilde{\boldsymbol{w}}_h^{Ki}\right]_v \, d\boldsymbol{x}\, dt + \\
&\int_{t_i}^{t_i+\Delta t_i} \int_{\partial K} \left[\boldsymbol{\mathcal{G}}\left(\tilde{\boldsymbol{q}}_h^{Ki}, \tilde{\boldsymbol{q}}_h^{K^+i}, \boldsymbol{n}\right)\right]_v \left[\tilde{\boldsymbol{w}}_h^{Ki}\right]_v \, ds(\boldsymbol{x})\, dt = \\
&\int_{t_i}^{t_i+\Delta t_i} \int_K \left[\boldsymbol{s}\left(\tilde{\boldsymbol{q}}_h^{Ki}\right)\right]_v \left[\tilde{\boldsymbol{w}}_h^{Ki}\right]_v \, d\boldsymbol{x}\, dt.
\end{aligned}
\tag{2.19}
$$

We require eq. (2.19) to hold for all $v \in \{1, 2, \ldots, V\}$, $\tilde{\boldsymbol{w}}_h^i \in \tilde{\mathbb{W}}_h^i$, $K \in \mathcal{T}_h$ and $i \in \{0, 1, \ldots, I-1\}$.

The assumption that the solution is balanced, i.e. that there is no net inflow or outflow for cell $K \in \mathcal{T}_h$ allows us to drop the third term. Together with integration by parts in time of the first term this yields

$$
\begin{aligned}
&\int_K \left[\tilde{\boldsymbol{q}}_h^{Ki}\Big|_{t_i+\Delta t_i}\right]_v \left[\tilde{\boldsymbol{w}}_h^{Ki}\Big|_{t_i+\Delta t_i}\right]_v d\boldsymbol{x} - \int_{t_i}^{t_i+\Delta t_i} \int_K \left[\tilde{\boldsymbol{q}}_h^{Ki}\right]_v \frac{\partial}{\partial t} \left[\tilde{\boldsymbol{w}}_h^{Ki}\right]_v d\boldsymbol{x}\, dt = \\
&\int_K \left[\tilde{\boldsymbol{q}}_h^{Ki}\Big|_{t_i}\right]_v \left[\tilde{\boldsymbol{w}}_h^{Ki}\Big|_{t_i}\right]_v d\boldsymbol{x} + \int_{t_i}^{t_i+\Delta t_i} \int_K \left[\boldsymbol{s}\left(\tilde{\boldsymbol{q}}_h^{Ki}\right)\right]_v \left[\tilde{\boldsymbol{w}}_h^{Ki}\right]_v d\boldsymbol{x}\, dt,
\end{aligned}
\tag{2.20}
$$

which we require to hold for all $v \in \{1, 2, \ldots, V\}$, $\tilde{\boldsymbol{w}}_h^i \in \tilde{\mathbb{W}}_h^i$, $K \in \mathcal{T}_h$ and $i \in \{0, 1, \ldots, I-1\}$. Together with the initial condition

$$
\tilde{\boldsymbol{q}}_h^{Ki}\Big|_{t_i} = \tilde{\boldsymbol{u}}_h^{K}\Big|_{t_i} \tag{2.21}
$$

and an initial guess

$$
\tilde{\boldsymbol{q}}_h^{Ki}\Big|_{t} = \tilde{\boldsymbol{u}}_h^{K}\Big|_{t_i} \; \forall t \in (t_i, t_i + \Delta t_i) \tag{2.22}
$$

this relation can be used as a fixed-point iteration to find $\tilde{\boldsymbol{q}}_h^{Ki}\Big|_t \; \forall t \in (t_i, t_i + \Delta t_i)$.

In the following two sections we will introduce mappings from space-time elements $K \times (t_i, t_i + \Delta t_i)$ to reference space-time cells and orthogonal bases for the spaces $\mathbb{W}_h$ and $\tilde{\mathbb{W}}_h^i$. We will then insert these results into eq. (2.20) and derive a fully-discrete iterative method to compute the space-time predictor $\tilde{\boldsymbol{q}}_h^{Ki}$.

2.1.7 Mappings

Let $\hat{K} = (0,1)^D$ be the spatial reference element and $\boldsymbol{\xi} \in \hat{K}$ be a point in the reference element. Let $(0,1)$ be the reference time interval and $\tau \in (0,1)$ be a point in time in reference time.

We can then introduce the following mappings:

Spatial mappings: Let $K \in \mathcal{T}_h$ be a cell in global coordinates with extent $\Delta\boldsymbol{x}^K$ and "lower-left corner" $\boldsymbol{P}_K$, more precisely that is

$$\left[\Delta\boldsymbol{x}^K\right]_d = \max_{x\in K}\,[\boldsymbol{x}]_d - \min_{x\in K}\,[\boldsymbol{x}]_d \tag{2.23}$$

and

$$[\boldsymbol{P}_K]_d = \min_{x\in K}\,[\boldsymbol{x}]_d \tag{2.24}$$

for $d \in \{1,2,\dots,D\}$. We can then define a mapping

$$\boldsymbol{\mathcal{X}}_K : \hat{K} \to K, \boldsymbol{\xi} \mapsto \boldsymbol{\mathcal{X}}_K(\boldsymbol{\xi}) = \boldsymbol{x} \tag{2.25}$$

via the relation

$$[\boldsymbol{x}]_d = \left[\boldsymbol{\mathcal{X}}_K(\boldsymbol{\xi})\right]_d = [\boldsymbol{P}_K]_d + [\Delta\boldsymbol{x}]_d\,\left[\boldsymbol{\xi}\right]_d \tag{2.26}$$

for $v \in \{1,2,\dots,V\}$ (i.e. no summation on v) and for all $\boldsymbol{x} \in K$, $\boldsymbol{\xi} \in \hat{K}$ and $K \in \mathcal{T}_h$.

Temporal mappings: Let $(t_i, t_i + \Delta t_i), i \in \{0,1,\dots,I-1\}$ be an interval in global time. The mapping

$$\mathcal{T}_i : (0,1) \to (t_i, t_i + \Delta t_i), \tau \mapsto \mathcal{T}_i(\tau) = t_i + \Delta t_i\tau = t \tag{2.27}$$

maps a point in reference time $\tau \in (0,1)$ to a point in global time $t \in (t_i, t_i + \Delta t_i)$ for all $i \in \{0,1,\dots,I-1\}$.

The inverse mappings, the Jacobian matrices and the Jacobi determinants of the mappings are given in the following:

Spatial mappings: The inverse spatial mappings

$$\boldsymbol{\mathcal{X}}_K^{-1} : K \to \hat{K}, \boldsymbol{x} \mapsto \boldsymbol{\mathcal{X}}_K^{-1}(\boldsymbol{x}) = \boldsymbol{\xi} \tag{2.28}$$

are defined via the relation

$$[\boldsymbol{\xi}]_d = \left[\boldsymbol{\mathcal{X}}_K^{-1}(\boldsymbol{x})\right]_d = \frac{1}{\left[\Delta \boldsymbol{x}^K\right]_d}\left([\boldsymbol{x}]_d - [\boldsymbol{P}_K]_d\right) \tag{2.29}$$

for $v \in \{1, 2, \ldots, V\}$ and for all $\boldsymbol{\xi} \in \hat{K}$, $\boldsymbol{x} \in K$ and $K \in \mathcal{T}_h$. The Jacobian of $\boldsymbol{\mathcal{X}}_K$ is found to be

$$\left[\frac{\partial \boldsymbol{\mathcal{X}}_K}{\partial \boldsymbol{\xi}}\right]_{dd'} = \frac{\partial\left[\boldsymbol{\mathcal{X}}_K\right]_d}{\partial \xi_{d'}} = \left[\Delta \boldsymbol{x}^K\right]_d \delta_{dd'}, \tag{2.30}$$

where $d, d' \in \{1, 2, \ldots D\}$ (i.e. no summation on d) and for all $K \in \mathcal{T}_h$. As usual $\delta_{dd'}$ denotes the Kronecker delta defined as

$$\delta_{dd'} = \begin{cases} 0 & \text{if } d \neq d' \\ 1 & \text{if } d = d'. \end{cases} \tag{2.31}$$

The Jacobi determinant of $\boldsymbol{\mathcal{X}}_K$ for $K \in \mathcal{T}_h$ then simply is

$$J_{\boldsymbol{\mathcal{X}}_K} = \left\|\frac{\partial \boldsymbol{\mathcal{X}}_K}{\partial \boldsymbol{\xi}}\right\| = \prod_{d=1}^{D}\left[\Delta \boldsymbol{x}^K\right]_d, \tag{2.32}$$

i.e. the determinant is constant for all $\boldsymbol{x} \in K$.

Temporal mappings: The inverse temporal mappings are given as

$$\mathcal{T}_i^{-1} : (t_i, t_i + \Delta t_i) \to (0,1), t \mapsto \mathcal{T}_i^{-1}(t) = \frac{t - t_i}{\Delta t_i} = \tau \tag{2.33}$$

for all $\tau \in (0, 1)$, $t \in (t_i, t_i + \Delta t_i)$ and $i \in \{1, 2, \ldots, I - 1\}$. In the trivial case of a one-dimensional mapping the Jacobian of $\mathcal{T}_i$ is a scalar which in turn is its own determinant. One finds

$$\frac{d\mathcal{T}_i}{\partial \tau} = \Delta t_i = J_{\mathcal{T}_i} \tag{2.34}$$

which again is constant for all $t \in (t_i, t_i + \Delta t_i)$ for a fixed $i \in \{0, 1, \ldots, I - 1\}$.

2.1.8 Orthogonal bases for the finite-dimensional spatial and space-time function spaces

Lagrange interpolation

Let $f \in \mathbb{Q}_N\left((0,1)\right)$ be a polynomial of degree N and let $\{\hat{x}_n\}_{n \in \{0,1,\ldots,N\}}$ be a set of distinct nodes in $(0, 1)$. The the Lagrange interpolation of f,

$$\hat{f}(x) = \sum_{n=0}^{N} L_n(x) f(x_n) \tag{2.35}$$

with Lagrange functions

$$L_n(x) = \prod_{m=0, m\neq n}^{N} \frac{x - \hat{x}_m}{\hat{x}_n - \hat{x}_m} \tag{2.36}$$

is exact, i.e.

$$f(x) = \hat{f}(x) \, \forall x \in (0,1). \tag{2.37}$$

Since every polynomial $f \in \mathbb{Q}_N\left((0,1)\right)$ can be represented as a linear combination of the Legendre polynomials L_n the set of functions $\{L_n\}_{n\in\{0,1,\ldots,N\}}$ is a basis of $Q_N\left((0,1)\right)$.

The following observation is an important property of the Lagrange polynomials:

$$L_n(\hat{x}_{n'}) = \delta_{nn'}, \tag{2.38}$$

i.e. at each node $\hat{x}_n$ only L_n has value 1 and all other polynomials evaluate to 0.

Legendre polynomials and Gauss-Legendre integration

Let $P_0 : (-1,1) \to \mathbb{R}, x \mapsto 1$ and $P_1 : (-1,1) \to \mathbb{R}, x \mapsto x$ be the zeroth and the first Legendre polynomial, respectively. Then the $N+1$-st Legendre polynomial can be defined via the following recurrence relation:

$$P_{N+1}(x) = \frac{1}{N+1}\left((2N+1)P_N(x) - nP_{N-1}(x)\right). \tag{2.39}$$

Let $\{\tilde{x}_n\}_{n\in\{0,1,\ldots,N\}}$ be the roots of the $N+1$-st Legendre polynomial $L_{N+1}(x)$. Then $\{\hat{x}_n\}_{n\in\{0,1,\ldots,N\}}$ with

$$\hat{x}_n = \frac{1}{2}(\tilde{x}_n + 1) \tag{2.40}$$

are the roots of the $N+1$-st Legendre polynomial linearly mapped to the interval $(0,1)$. In conjunction with a set of suitable weights $\{\hat{w}_n\}_{n\in\{0,1,\ldots N\}}$ Gauss-Legendre integration can be used to integrate polynomials of degree up to $2N+1$ over the integral $[0,1]$ exactly, i.e.

$$\int_0^1 f(x)\, d\boldsymbol{x} = \sum_{n=0}^{N} \hat{w}_n f(\hat{x_n}) \, \forall f \in \mathbb{Q}_{2N+1}\left([0,1]\right). \tag{2.41}$$

A script on how to find the weights $\{\hat{w}_n\}_{n\in\{0,1,\ldots,N\}}$ can be found in appendix XXX.

1d basis functions

Let $\left\{\psi_n\right\}_{n\in\{0,1,\ldots,N\}}$ be the set of $N+1$ Lagrange polynomials with nodes at the roots of the $N+1$-st Legendre polynomial linearly mapped to the interval $(0,1)$, i.e.

$$\psi_n(x) = \sum_{n'=0}^{N} \frac{x - \hat{x}_{n'}}{\hat{x}_n - \hat{x}_{n'}} \tag{2.42}$$

for $n \in \{0,1,\ldots,N\}$. Since $\left\{\psi_n\right\}_{n\in\{0,1,\ldots,N\}}$ are Lagrange polynomials and the roots $\{\hat{x}_n\}_{n\in\{0,1,\ldots,N\}}$ are distinct the set is a basis of $\mathbb{Q}_N\left([0,1]\right)$. Since furthermore

$$\left\langle\psi_n,\psi_m\right\rangle_{L^2\left((0,1)\right)} = \int_0^1 \psi_n(x)\psi_m(x)\,d\boldsymbol{x} = \sum_{n'=0}^{N} \hat{w}'_n\psi_n(\hat{x}_{n'})\psi_m(\hat{x}_{n'}) = \hat{w}_n\delta_{mn} \tag{2.43}$$

for all $m,n \in \{0,1,\ldots,N\}$ (i.e. no summation over n) the set is even an orthogonal basis of $\mathbb{Q}_N\left([0,1]\right))$ with respect to the L^2-scalar product defined as above. In the derivation above we used the fact that $\phi_n\phi_m$ has degree $2N$ and that Gauss-Legendre integration with $N+1$ nodes is exact for polynomials up to degree $2N+1$.

2.2 Profiling and Energy-aware Computing

Chapter 3

A profiling infrastructure for ExaHyPE

- General architecture
- Architecture profiling
- Functionality

Chapter 4

Preliminary profiling results, case studies

- Analytic benchmark: Introduction, derivation
- Pie-chart per kernel
- Case-study: Cache-misses, compile-time ($\rightarrow$ Toolkit philosophy)
- Degree $\rightarrow$ Wallclock, Energy (AMR)
- Static mesh Δx $\rightarrow$ Error for polynomials (convergence tables)

Chapter 5

Conclusion and Outlook

- PA is important
- ExaHyPE as an answer to exascale challenges
- Applications

Chapter 6

Acknowledgment