

Higgs bundles twisted by a vector bundle

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CHAPTER I

Vector bundles on Riemann surfaces

§ 1. Topological classification of vector bundles

The first step towards the classification of vector bundles on Riemann surfaces is their “topological” classification. That is, we want to classify smooth complex vector bundles on a Riemann surface up to C^∞ isomorphism. This is indeed pretty easy to do, since the problem can be reduced to the classification of line bundles.

Theorem 1.1. *If E is a rank n smooth complex vector bundle over a compact Riemann surface X then it is isomorphic to $\det E \oplus (X \times \mathbb{C}^{n-1})$.*

Proof. We will proceed by induction on n . Of course, if E is a line bundle, $E \cong \det E$. Now, let $n > 1$, suppose that it is true for any vector bundle of rank $n - 1$ and let E be a rank n vector bundle.

Lemma 1.2. *E has a nowhere vanishing section.*

Proof. Let s_0 be the zero section of E and $S_0 = s_0(X) \subset E$. By the transversality theorem [Hir12], we can densely choose a section $s \in \Gamma(E)$ transversal to S_0 . Now, if s vanishes at some point $x \in X$, then $s(x) \in S_0$ and, since s is transversal to S_0 ,

$$d_x s(T_x X) + T_{s(x)} S_0 = T_{s(x)} E.$$

But $\dim_{\mathbb{R}} E = 2n + 2$, $\dim_{\mathbb{R}} S_0 = 2$ and $\dim_{\mathbb{R}} d_x s(T_x X) \leq \dim_{\mathbb{R}} X = 2$, so, if $n > 1$, dimensions do not add up to verify the above equality and therefore we get a contradiction. \square

Let us continue with the proof of the theorem. Since E has a nowhere vanishing section s , we can define the line bundle

$$L = \bigsqcup_{x \in X} \text{span}(s(x)),$$

with

$$\begin{aligned} \pi : L &\longrightarrow X \\ \lambda s(x) &\longmapsto x. \end{aligned}$$

Therefore, we can decompose $E = E' \oplus L$, with E' a rank $n - 1$ vector bundle. For example, fixing a metric on E , we can define E' to be the orthogonal complement of L . But observe now

that L is isomorphic to the trivial line bundle: the bundle morphism

$$\begin{aligned} X \times \mathbf{C} &\longrightarrow L \\ (x, \lambda) &\longmapsto \lambda s(x), \end{aligned}$$

is in fact an isomorphism. This can be proven by defining a metric on L and normalizing $s \mapsto s/\|s\|$. Then we can define the inverse $y \in L \mapsto (\pi(y), \langle s(\pi(y)), y \rangle) \in X \times \mathbf{C}$.

Thus, we have shown that $E \cong E' \oplus (X \times \mathbf{C})$. Now, applying the induction hypothesis, $E' \cong \det E' \oplus (X \times \mathbf{C}^{n-2})$, so $E \cong \det E' \oplus (X \times \mathbf{C}^{n-1})$. Finally, via transition functions it can be easily shown that $\det E \cong \det E'$. \square

This last theorem says that vector bundles can be topologically classified by their rank and their determinant, which is a line bundle. Let us proceed then with the classification of line bundles. Recall that all the data of a vector bundle can be recovered by the transition functions $\{g_{\alpha\beta} \in C^\infty(U_\alpha \cap U_\beta, \mathbf{C})\}$ defining it, where $\{U_\alpha\}$ is an open cover of X . These functions verify the cocycle condition

$$g_{\alpha\beta} = g_{\gamma\beta} \cdot g_{\alpha\gamma}.$$

Also recall that an isomorphism of vector bundles induces a coboundary on the transition functions

$$\tilde{g}_{\alpha\beta} = f_{\delta\beta}^{-1} \cdot g_{\gamma\delta} \cdot f_{\gamma\alpha}.$$

Therefore, topological (C^∞) isomorphism classes of vector bundles are parametrized by the Čech cohomology group

$$H^1(X, C_X^{\infty,*}).$$

To obtain more information about this cohomology group we are going to introduce a very powerful tool: the first Chern class of a vector bundle.

Recall from Chern-Weil theory [GH14, WGP07] that for any complex vector bundle E over X and for any connection ∇ on E with associated curvature form F , the 2-form $\text{tr}F$ is closed and its de Rham cohomology class $[\text{tr}F] \in H_{\text{dR}}^2(X)$ does not depend on the choice of the connection, so it is an invariant of the vector bundle E . If we normalize this form to get an integer cohomology class, we define the **first Chern class** of E as the cohomology class:

$$c_1(E) = \left[\frac{i}{2\pi} \text{tr}F \right] \in H_{\text{dR}}^2(X).$$

We define the **degree** of a vector bundle E as the pairing of $c_1(E)$ with the fundamental class of X , that is

$$\deg E = \int_X \frac{i}{2\pi} \text{tr}F.$$

The next proposition [WGP07] summarizes the most important properties about the degree that we are going to use:

Proposition 1.3 (Properties of the degree). *Let E and F be complex vector bundles over a compact Riemann surface X .*

1. $\deg(E)$ depends only on the isomorphism class of E .
2. $\deg(E) \in \mathbf{Z}$.
3. $\deg(E \oplus F) = \deg(E) + \deg(F)$.

4. $\deg(E \otimes F) = \text{rk} F \deg E + \text{rk} E \deg F$.
5. $\deg E = \deg(\det E)$.
6. Let $\delta : H^1(X, C_X^{\infty,*}) \rightarrow H^2(X, \mathbf{Z})$ be the connecting homomorphism of the long exact sequence in cohomology induced by the exponential sheaf exact sequence

$$0 \longrightarrow \mathbf{Z} \longrightarrow C_X^\infty \xrightarrow{\exp} C_X^{\infty,*} \longrightarrow 0,$$

where $\exp(f) = e^{2\pi i f}$. The diagram

$$\begin{array}{ccc} H^1(X, C_X^{\infty,*}) & \xrightarrow{\delta} & H^2(X, \mathbf{Z}) \\ & \searrow c_1 & \downarrow \\ & & H_{\text{dR}}^2(X), \end{array}$$

is commutative.

This last property will be crucial in the classification of line bundles. Let us consider again the exponential sheaf exact sequence

$$0 \longrightarrow \mathbf{Z} \longrightarrow C_X^\infty \xrightarrow{\exp} C_X^{\infty,*} \longrightarrow 0.$$

The existence of smooth partitions of unity implies that the sheaf C_X^∞ is fine, so $H^1(X, C_X^\infty) = H^2(X, C_X^\infty) = 0$. Therefore, the connecting map $\delta : H^1(X, C_X^{\infty,*}) \rightarrow H^2(X, \mathbf{Z})$ is an isomorphism. If we now consider the set of integer de Rham cohomology classes with $H_{\text{dR}}^2(X, \mathbf{Z})$, which is just the image of $H^2(X, \mathbf{Z})$ by the inclusion $H^2(X, \mathbf{Z}) \hookrightarrow H_{\text{dR}}^2(X)$, we get an isomorphism $c_1 : H^1(X, C_X^{\infty,*}) \rightarrow H_{\text{dR}}^2(X, \mathbf{Z})$. Now, the isomorphism $H_{\text{dR}}^2(X) \cong \mathbf{C}$ given by integration on X , descends to an isomorphism $H_{\text{dR}}^2(X, \mathbf{Z}) \cong \mathbf{Z}$. Summarizing, we have the diagram

$$\begin{array}{ccccc} H^1(X, C_X^{\infty,*}) & \xrightarrow{c_1} & H_{\text{dR}}^2(X, \mathbf{Z}) & \xrightarrow{\int_X} & \mathbf{Z} \\ & \searrow \text{deg} & & \nearrow & \end{array}$$

That is, the degree gives an isomorphism between the set isomorphism classes of smooth line bundles and \mathbf{Z} . This concludes the topological classification of vector bundles, which we can gather in the next theorem

Theorem 1.4. *Smooth complex vector bundles over a compact Riemann surfaces are classified, up to C^∞ isomorphism by their rank and their degree.*

§ 2. The problem of classification of holomorphic vector bundles

Now that we have classified vector bundles up to topological (C^∞) isomorphism, we are going to pursue the full classification of holomorphic vector bundles over Riemann surfaces. According to the results of last section, we can reduce our problem to the study of the “list” of isomorphism classes of holomorphic vector bundles of fixed rank n and degree d (regarding them in particular as complex vector bundles). From the beginning, this problem gets really involved, since this

“lists” are so big that they themselves admit a geometric structure. These are the so called **moduli spaces**. Therefore, the classification problem translates to that of investigating the geometric properties of the associated moduli spaces.

To illustrate these ideas in more detail, let us consider the case of holomorphic line bundles. The same arguments regarding Čech cocycles of the previous section also apply now to show that the isomorphism classes of holomorphic vector bundles are parametrized by the sheaf cohomology group

$$H^1(X, \mathcal{O}_X^*).$$

This cohomology group is called the **Picard group** of X and we denote it by $\text{Pic}(X)$. Now, as before, we can consider the exponential sheaf sequence

$$0 \longrightarrow \mathbf{Z} \longrightarrow \mathcal{O}_X \xrightarrow{\exp} \mathcal{O}_X^* \longrightarrow 0,$$

and the connecting operator $\delta : H^1(X, \mathcal{O}_X^*) \rightarrow H^2(X, \mathbf{Z})$ of the induced long exact sequence in cohomology. Analogously to the previous section, one can show that the diagram

$$\begin{array}{ccc} \text{Pic}(X) & \xrightarrow{\delta} & H^2(X, \mathbf{Z}) \\ & \searrow c_1 & \downarrow \\ & & H_{\text{dR}}^2(X) \end{array}$$

is commutative. However, unlike the smooth case, the sheaf \mathcal{O}_X is not fine, since there are not holomorphic partitions of unity, so in general δ is not an isomorphism anymore. Let $\text{Pic}^0(X) = \ker \delta \subset \text{Pic}(X)$ be the subgroup of degree zero line bundles. Then, we have an exact sequence

$$0 \longrightarrow H^1(X, \mathbf{Z}) \longrightarrow H^1(X, \mathcal{O}_X) \longrightarrow \text{Pic}^0(X) \xrightarrow{\delta} 0.$$

Therefore $\text{Pic}^0(X) \cong H^1(X, \mathcal{O}_X)/H^1(X, \mathbf{Z})$. Now, the Dolbeault theorem says that $H^1(X, \mathcal{O}_X) \cong H^{0,1}(X)$ and the Hodge decomposition theorem says that $H^1(X, \mathbf{C}) \cong H^{1,0}(X) \oplus H^{0,1}(X)$. Also, by Serre duality, $H^{1,0}(X) \cong H^{0,1}(X)^*$ and via Mayer-Vietoris one can easily prove that $H^1(X, \mathbf{C}) \cong \mathbf{C}^{2g}$ and that $H^1(X, \mathbf{Z}) \cong \mathbf{Z}^{2g}$, where g is the genus of X . Putting all this together we have that $H^1(X, \mathcal{O}_X) \cong \mathbf{C}^g$, so $\text{Pic}^0(X)$ is a complex torus. We call this complex torus

$$J(X) := \frac{H^1(X, \mathcal{O}_X)}{H^1(X, \mathbf{Z})}$$

the **Jacobian** of X . The other components of fixed degree of $\text{Pic}(X)$ can be also shown to be isomorphic to the Jacobian of X , via the isomorphism

$$\begin{aligned} \text{Pic}^d(X) &\longrightarrow \text{Pic}^0(X) \\ L &\longmapsto L \otimes M, \end{aligned}$$

where $M \in \text{Pic}^{-d}(X)$ is a fixed line bundle. It is of course injective, since if $L \otimes M = L' \otimes M$, then

$$L = L \otimes M \otimes M^* = L' \otimes M \otimes M^* = L',$$

and it is also surjective since for every line bundle $L \in \text{Pic}^0(X)$, $L = (L \otimes M^*) \otimes M$.

The result for line bundles already hints on the complexity of the general problem. However, for very low genus the problem can be solved relatively easily. For genus 0, Grothendieck proved

that every holomorphic vector bundle on the Riemann sphere $\mathbf{P}_{\mathbb{C}}^1$ can be decomposed as a direct sum of line bundles. A proof of this result can be found in [Har13]. For genus 1, it was Atiyah who showed that the “moduli space” of indecomposable vector bundles with fixed rank and degree over an elliptic curve is isomorphic to the curve itself. Check out [REF:Mukai?] for a construction of the moduli space and for a proof of this.

The problem gets its full complexity in the case of general genus $g \geq 2$, for which we will dedicate the rest of the chapter. To construct a “good” moduli space (one with nice topological properties, like being Hausdorff) we need to consider the idea of stability, arising from Mumford’s Geometric Invariant Theory [MFK94]. This theory also allows to construct the moduli space of stable vector bundles, although it can also be done using analytic methods [Kob14]. This moduli space has very nice and interesting geometric properties and it has been studied from the algebraic point of view (for example in the works of Narasimhan, Ramanan or Seshadri) as well as from the analytical or gauge-theoretical point of view (by Atiyah, Bott, Donaldson or Hitchin, for example).

§ 3. Holomorphic structures as Dolbeaut operators

Let \mathbf{E} be a holomorphic vector bundle on a compact Riemann surface X . Associated to \mathbf{E} , we have the *Dolbeaut operator*

$$\bar{\partial}_{\mathbf{E}} : \Omega^{p,q}(X, \mathbf{E}) \longrightarrow \Omega^{p,q+1}(X, \mathbf{E})$$

which satisfies

$$\bar{\partial}_{\mathbf{E}}(\alpha\psi) = (\bar{\partial}\alpha)\psi + (-1)^p\alpha \wedge \bar{\partial}_{\mathbf{E}}\psi,$$

for every $\alpha \in \Omega^p(X)$, $\psi \in \Gamma(\mathbf{E})$, and

$$\bar{\partial}_{\mathbf{E}}^2 = 0.$$

Conversely, if E is a complex vector bundle, any operator $\bar{\partial}_{\mathbf{E}} : \Omega^{p,q}(X, E) \rightarrow \Omega^{p,q+1}(X, E)$ which satisfies the conditions above induces on E the structure of a holomorphic vector bundle. The idea here is that $\bar{\partial}_{\mathbf{E}}^2 = 0$ is the *integrability condition* for the PDE $\bar{\partial}_{\mathbf{E}}s = 0$ (check out [DK90] for a proof of this fact). The solutions of this PDE are the holomorphic vector functions, so the sheaf of local solutions is locally free over \mathcal{O}_X and therefore it is an holomorphic vector bundle supported on E . In conclusion, if we define $\mathcal{A}_{\bar{\partial}}$ as the set of all such $\bar{\partial}_{\mathbf{E}}$ operators satisfying the previous conditions, this set is in bijection with the set of all holomorphic structures on E .

Now let us consider the group \mathcal{G}^c of *gauge transformations* of a complex vector bundle E , that is, diffeomorphisms $g : E \rightarrow E$ such that the diagram

$$\begin{array}{ccc} E & \xrightarrow{g} & E \\ & \searrow & \swarrow \\ & X & \end{array}$$

commutes. That is, $\mathcal{G}^c = \Gamma(\text{Aut}(E))$. This group acts on $\mathcal{A}_{\bar{\partial}}$ by the rule

$$g \cdot \bar{\partial}_{\mathbf{E}} = g \bar{\partial}_{\mathbf{E}} g^{-1}.$$

Therefore we can identify the quotient set $\mathcal{A}_{\bar{\partial}}/\mathcal{G}^c$ with the set of isomorphism classes of holomorphic vector bundles of rank $\text{rk}E$ and degree $\deg E$. Even though $\mathcal{A}_{\bar{\partial}}$ and \mathcal{G}^c are infinite-dimensional spaces, using analytic techniques [Kob14] they can be given the structure of *Banach manifolds* (roughly, spaces locally modeled by Banach spaces) and this quotient space can be

precisely constructed. This has however a serious problem: the obtained space is not Hausdorff. Nevertheless, if we restrict ourselves to the class of **stable vector bundles**, with the same analytic methods we can obtain a “good” moduli space.

Definition 3.1. Let E be a complex vector bundle over a compact Riemann surface X . We define the **slope** of E as the number

$$\mu(E) = \deg E / \operatorname{rk} E.$$

We say that a holomorphic vector bundle $\mathbf{E} = (E, \bar{\partial}_{\mathbf{E}})$ is **stable** if for every proper holomorphic subbundle $\mathbf{E}' \subset \mathbf{E}$ (that is, for every subbundle $E' \subset E$ such that $\bar{\partial}$ preserves E'),

$$\mu(E') < \mu(E).$$

Let $n = \operatorname{rk} E$ and $d = \deg E$. We consider $\mathcal{A}_{\bar{\partial}}^s \subset \mathcal{A}_{\bar{\partial}}$ the subset of stable holomorphic bundles $\mathbf{E} = (E, \bar{\partial}_{\mathbf{E}})$ and define the **moduli space of stable holomorphic vector bundles** of rank n and degree d as the quotient

$$\mathbf{M}_{\text{Bun}}(n, d) = \mathcal{A}_{\bar{\partial}}^s / \mathcal{G}^c.$$

Theorem 3.2. *The moduli space $\mathbf{M}_{\text{Bun}}(n, d)$ is a complex manifold of dimension $1 + n^2(g - 1)$, where g is the genus of X .*

§ 4. Deformations of holomorphic structures

In this section we are going to compute the dimension of $\mathbf{M}_{\text{Bun}}(n, d)$, stated in Theorem 3.2, using deformation theory. The idea is to identify the tangent space of $\mathbf{M}_{\text{Bun}}(n, d)$ at some point $\bar{\partial}_{\mathbf{E}}$ and find its dimension.

In first place, observe that $\mathcal{A}_{\bar{\partial}}$ is an affine space modelled on $\Omega^{0,1}(X, \operatorname{End} E)$. Therefore, the tangent space of $\mathcal{A}_{\bar{\partial}}$ at $\bar{\partial}_{\mathbf{E}}$ is isomorphic to $\Omega^{0,1}(X, \operatorname{End} E)$. Now, let $\operatorname{Lie} \mathcal{G}^c = \Gamma(\operatorname{End} E)$ be the Lie algebra of \mathcal{G}^c and pick an element $\xi \in \Gamma(\operatorname{End} E)$. Let us compute the action of this element on the tangent space $\Omega^{0,1}(X, \operatorname{End} E)$,

$$\left. \frac{d}{dt} \right|_{t=0} \exp(\xi t) \bar{\partial}_{\mathbf{E}} \exp(-\xi t) = -\bar{\partial}_{\mathbf{E}} \xi.$$

Hence, since stability is an open condition, the tangent space $T_{\bar{\partial}_{\mathbf{E}}} \mathbf{M}_{\text{Bun}}(n, d)$ is isomorphic to

$$H^1(X, \operatorname{End} \mathbf{E}) = \frac{\Omega^{0,1}(X, \operatorname{End} E)}{\bar{\partial}_{\mathbf{E}} \Gamma(\operatorname{End} E)}.$$

To compute the dimension of this tangent space, we use the **Riemann-Roch theorem**

$$\dim H^0(X, \operatorname{End} \mathbf{E}) - \dim H^1(X, \operatorname{End} \mathbf{E}) = \deg(\operatorname{End} \mathbf{E}) + \operatorname{rk}(\operatorname{End} \mathbf{E})(1 - g).$$

Now, note that $\operatorname{End} E = \operatorname{Hom}(E, E) \cong E^* \otimes E$, so

$$\operatorname{rk}(\operatorname{End} E) = n^2$$

and

$$\deg(\operatorname{End} E) = \operatorname{rk}(E^*) \deg(E) + \operatorname{rk}(E) \deg(E^*) = n \deg E - n \deg E = 0.$$

Therefore

$$\dim H^0(X, \operatorname{End} \mathbf{E}) - \dim H^1(X, \operatorname{End} \mathbf{E}) = n^2(1 - g).$$

To find the dimension of $H^0(X, \operatorname{End} \mathbf{E})$ we use the following result:

Proposition 4.1. *If \mathbf{E} is a stable holomorphic vector bundle, then*

$$H^0(X, \text{End } \mathbf{E}) = \mathbb{C}.$$

Proof. What we will see is that every endomorphism is a scalar multiple of the identity $\mathbf{1}_{\mathbf{E}}$. Pick any endomorphism $f : \mathbf{E} \rightarrow \mathbf{E}$ and let $\lambda \in \mathbb{C}$ be an eigenvalue of $f_x : \mathbf{E}_x \rightarrow \mathbf{E}_x$. Let us define the endomorphism $g = f - \lambda \mathbf{1}_{\mathbf{E}}$. Since λ is an eigenvalue of f_x , we have $\det g_x = 0$. Assuming that g is nonzero, we will prove that g is injective and hence $\det g_x \neq 0$ arriving at a contradiction. To see this, suppose that $\text{im } g \subset \mathbf{E}$ is a holomorphic strict subbundle of \mathbf{E} . Then, since \mathbf{E} is stable, $\mu(\text{im } g) < \mu(\mathbf{E})$. But we also have that $\mu(\mathbf{E}) < \mu(\text{im } g)$, so $\mu(\mathbf{E}) < \mu(\mathbf{E})$ and we have a contradiction. Therefore g is injective. We conclude then that $g = 0$, so $f = \lambda \mathbf{1}_{\mathbf{E}}$. \square

Finally, we get

$$\dim \mathbf{M}_{\text{Bun}}(n, d) = \dim H^1(X, \text{End } \mathbf{E}) = 1 + n^2(g - 1).$$

§ 5. Holomorphic structures and unitary connections

Although we will not enter into detail, the main reason why the quotient $\mathcal{A}_{\bar{\partial}}/\mathcal{G}^c$ is not Hausdorff and why we need to introduce the stability condition is that the group $\mathcal{G}^c = \Gamma(\text{Aut}(E))$ is a complex group. Indeed, it is the complexification of the group $\mathcal{G} = \Gamma(U_h(E))$, where h is a Hermitian metric on E and $U_h(E)$ is the subgroup of $\text{Aut}(E)$ consisting on automorphisms of E preserving the metric h . The idea for this is essentially that the general linear group $\text{GL}(n, \mathbb{C})$ is the complexification of the unitary group $U(n)$. This motivates the study of unitary connections, which will give a gauge-theoretical approach to holomorphic vector bundles. This will allow us to give another analytical construction of the moduli space, this time as the space of solutions (up to gauge equivalence) to some differential equation.

Let $\mathbf{E} = (E, \bar{\partial}_{\mathbf{E}})$ be a holomorphic vector bundle on X and let h be a Hermitian metric on E . Recall that for any connection $\nabla : \Gamma(E) \rightarrow \Omega^1(X, E) = \Omega^{1,0}(X, E) \oplus \Omega^{0,1}(X, E)$ on E there is a natural splitting $\nabla = \nabla^{1,0} + \nabla^{0,1}$, where

$$\begin{aligned} \nabla^{1,0} : \Gamma(E) &\longrightarrow \Omega^{1,0}(X, E), \\ \nabla^{0,1} : \Gamma(E) &\longrightarrow \Omega^{0,1}(X, E). \end{aligned}$$

The **Chern connection** in (\mathbf{E}, h) is the unique *h -unitary* connection (that is,

$$d\langle \xi, \eta \rangle = \langle \nabla \xi, \eta \rangle + \langle \xi, \nabla \eta \rangle,$$

for ξ, η local sections of E) such that

$$\nabla^{0,1} = \bar{\partial}_{\mathbf{E}} : \Gamma(E) \rightarrow \Omega^{0,1}(X, E).$$

Go to [WGP07] for a proof of the existence and uniqueness of this connection.

Conversely, given any h -unitary connection ∇ on E , we can define a holomorphic structure by fixing $\bar{\partial}_{\mathbf{E}} = \nabla^{0,1}$ and extending to operators $\Omega^{p,q}(X, E) \rightarrow \Omega^{p,q+1}(X, E)$ by linearity. Therefore, the set holomorphic structures on E can be identified with the set \mathcal{A}_h of all h -unitary connections on E .

Note now that the curvature of any connection on E must satisfy that the cohomology class

$$[\text{tr} F] = -i2\pi c_1(E).$$

Fixing an area form ω_X on X so that $\int_X \omega_X = 1$, we can choose a representative of $c_1(E)$ of the form $k\omega_X$, where $k \in \mathbb{C}$ is a constant. Of course,

$$\deg E = \int_X c_1(E) = \int_X k\omega_X = k,$$

so $k = \deg E$. Therefore we can ask if there is a connection on E such that its curvature satisfies

$$\mathrm{tr} F = -2\pi i \deg(E) \omega_X.$$

Or, more generally, if $\mathbf{1}_E$ denotes the identity endomorphism of E , we can ask whether

$$F = -2\pi i \frac{\deg E}{\mathrm{rk} E} \mathbf{1}_E \omega_X.$$

Recall that we defined the number $\mu(E) = \deg E / \mathrm{rk} E$ as the slope of E .

Definition 5.1. We say that a connection ∇ on a complex vector bundle E has *constant central curvature* if its curvature F satisfies

$$F = -2\pi i \mu(E) \mathbf{1}_E \omega_X.$$

In particular, if $\deg E = 0$, $F = 0$ and we say that ∇ is a *flat connection*.

Finally, we want to consider connections that are *irreducible*:

Definition 5.2. A unitary connection ∇ on a complex Hermitian vector bundle (E, h) is *reducible* if $(E, h) = (E_1, h_1) \oplus (E_2, h_2)$ and $\nabla = \nabla_1 \oplus \nabla_2$. We say that ∇ is *irreducible* if it is not reducible.

Let us consider then the set \mathcal{A}_h^s of all h -unitary irreducible connections of constant central curvature on E . The gauge group \mathcal{G} acts on connections by conjugation $\nabla \mapsto g\nabla g^{-1}$, $g \in \mathcal{G}$, and this action preserves irreducibility and the equation of constant central curvature, so the group \mathcal{G} acts on \mathcal{A}_h^s . Now, the same analytic techniques mentioned in the previous section [Kob14] allow us to construct a “good” quotient:

Theorem 5.3. *The moduli space of irreducible constant central curvature unitary connections $\mathcal{A}_h^s / \mathcal{G}$ on (E, h) has the structure of a smooth real manifold of dimension $2 + 2n^2(g - 1)$, where $n = \mathrm{rk} E$ and g is the genus of X .*

Now, Donaldson’s version of the theorem of Narasimhan-Seshadri relates this moduli space with the moduli space of stable holomorphic vector bundles.

Theorem 5.4 (Narasimhan-Seshadri). *Let (E, h) be a Hermitian complex vector bundle of rank n and degree d on a compact Riemann surface X . An irreducible unitary connection ∇ has constant central curvature if and only if the associated holomorphic vector bundle $(E, \nabla^{0,1})$ is stable.*

This can be reformulated in terms of moduli spaces:

Corollary 5.5. *The map*

$$\begin{aligned} \mathcal{A}_h &\longrightarrow \mathcal{A}_{\bar{\partial}} \\ \nabla &\longmapsto \nabla^{0,1}, \end{aligned}$$

descends to a homeomorphism

$$\mathcal{A}_h^s / \mathcal{G} \cong \mathcal{A}_{\bar{\partial}}^s / \mathcal{G}^c = \mathbf{M}_{\mathrm{Bun}}(n, d).$$

We will prove the “easy” direction of the equivalence. The proof in the other direction consists in defining a Yang-Mills functional and looking for a minimum of it using analytic techniques, in particular a theorem by Uhlenbeck [REF:Uhlenbeck?]. Check [Don83] for the details.

Proof. First, let us suppose that ∇ has constant central curvature

$$F = -2\pi i \frac{d}{n} \mathbf{1}_E \omega_X,$$

and define $\bar{\partial}_E = \nabla^{0,1}$. Let $E' \subset E$ be a subbundle preserved by $\bar{\partial}_E$. The Hermitian metric gives a smooth splitting

$$E = E' \oplus E'',$$

and we can write

$$\bar{\partial}_E = \begin{pmatrix} \bar{\partial}_{E'} & \beta \\ 0 & \bar{\partial}_{E''} \end{pmatrix},$$

where $\bar{\partial}_{E'}$ and $\bar{\partial}_{E''}$ are the restrictions of $\bar{\partial}_E$ to E' and E'' and $\beta \in \Omega^{0,1}(X, \text{Hom}(E'', E'))$. Now ∇ can be written as

$$\nabla = \begin{pmatrix} \nabla_{E'} & \beta \\ -\beta^\dagger & \nabla_{E''} \end{pmatrix},$$

where $\nabla_{E'}$ and $\nabla_{E''}$ are the connections associated to $\bar{\partial}_{E'}$ and $\bar{\partial}_{E''}$ and $\beta^\dagger = \star \bar{\beta} \in \Omega^{1,0}(X, \text{Hom}(E', E''))$ is the transpose (on the matrix part) conjugate (on the form part) of β . The curvature of ∇ can be written now as

$$F = \begin{pmatrix} F_{E'} - \beta \wedge \beta^\dagger & \nabla_{\text{Hom}(E'', E')} \beta \\ -\nabla_{\text{Hom}(E', E'')} \beta^\dagger & F_{E''} - \beta^\dagger \wedge \beta \end{pmatrix} = -2\pi i \frac{d}{n} \mathbf{1}_E \omega_X.$$

The first corner of this equality now says that

$$F_{E'} - \beta \wedge \beta^\dagger = -2\pi i \frac{d}{n} \mathbf{1}_{E'} \omega_X.$$

Taking the trace and integrating we get

$$\frac{i}{2\pi} \int_X \text{tr} F_{E'} - \frac{i}{2\pi} \int_X \text{tr}(\beta \wedge \beta^\dagger) = d \frac{\text{rk} E'}{n}.$$

Therefore

$$\frac{\deg E'}{\text{rk} E'} = \frac{d}{n} + \frac{i}{2\pi} \int_X \text{tr}(\beta \wedge \beta^\dagger).$$

Now, $\text{tr}(\beta \wedge \beta^\dagger)$ is precisely the Hodge pairing $\text{tr}(\beta \wedge \star \bar{\beta}) = -\|\beta\|^2 \omega_X$. Thus we have proven that

$$\mu(E') = \mu(E) - \|\beta\|^2.$$

The connection ∇ is irreducible, so $\|\beta\| \neq 0$ and $\mu(E') < \mu(E)$. That is, E is stable. \square

CHAPTER II

Momentum maps and symplectic quotients

§ 1. Banach manifolds

For the purpose of this work it is necessary to work on an infinite-dimensional context, so we can construct the moduli spaces as quotients of infinite-dimensional manifolds. To do this, we need to consider an infinite-dimensional calculus analogous to differential calculus and the notion of an infinite-dimensional smooth manifold. Here we just state the basic definitions and results, which are a direct generalization of the classical ones, and refer to [AMR12, Lan12] for more details.

Definition 1.1. Let E and F be Banach spaces and $U \subset E$ an open set. Let $f : E \rightarrow F$ be a continuous map. We say that f is **differentiable** at a point $x_0 \in U$ if there exists a continuous linear map $d_{x_0}f : E \rightarrow F$ such that

$$\lim_{h \rightarrow 0} \frac{\|f(x_0 + h) - f(x_0) - d_{x_0}f(h)\|_F}{\|h\|_E} = 0.$$

If this map $d_{x_0}f$ exists, then it is unique and it is called the **derivative of f at x_0** . If f is differentiable at every point of U , then we have a map

$$\begin{aligned} df : U &\longrightarrow L(E, F) \\ x &\longmapsto d_x f, \end{aligned}$$

where $L(E, F)$ denotes the set of bounded linear maps $E \rightarrow F$. If df is continuous, we say that f is **of class C^1** . Inductively, we define maps of class C^p , for $p \in \mathbf{N} \cup \{\infty\}$. If f is of class C^∞ , we say that it is **smooth**. If f is a bijective map of class C^p such that its inverse is also of class C^p , we say that f is a **C^p -diffeomorphism**. If we do not specify, by a **diffeomorphism** we just mean a C^∞ -diffeomorphism.

As we said before, all the basic definitions and results of classical differential calculus (the chain rule, Taylor's formula, inverse and implicit function theorems etc.) can be generalized to the infinite-dimensional context.

Definition 1.2. Let X be a Hausdorff topological space. An **atlas of class C^p** on X is a collection of **charts** $(U_\alpha, \varphi_\alpha)$ such that

1. The $U_\alpha \subset X$ are open subsets of X , and $X = \bigcup_\alpha U_\alpha$.
2. Each $\varphi_\alpha : U_\alpha \rightarrow V_\alpha$ is a homeomorphism of U_α onto some open subset of a Banach space $V_\alpha \subset E_\alpha$ and for any α, β , $\varphi_\alpha(U_\alpha \cap U_\beta)$ is open in E_α .

3. The map

$$\varphi_\beta \circ \varphi_\alpha^{-1} : \varphi_\alpha(U_\alpha \cap U_\beta) \rightarrow \varphi_\beta(U_\alpha \cap U_\beta)$$

is a C^p -diffeomorphism for each pair of indices α, β .

Compatibility classes for this kind of atlases are defined like in the finite-dimensional case. A equivalence class of this atlases (or a maximal atlas) is what gives X the structure of a C^p -**Banach manifold**. A C^∞ -Banach manifold is what we call a **Banach smooth manifold**.

Let X be a Banach smooth manifold and $x \in X$. Consider triples (U, φ, v) , where (U, φ) is a chart at x with $\varphi : U \rightarrow E$, E a Banach space, and $v \in E$. We say that two such triples (U, φ, v) and (U', φ', v') are **equivalent** if

$$d_{\varphi(x)}(\varphi' \circ \varphi^{-1})(v) = v'.$$

The chain rule guarantees that this is an equivalence relation and an equivalence class is called a **tangent vector of X at x** . The set of these equivalence classes is denoted by $T_x X$ and it acquires the structure of a Banach space via the bijection $[(U, \varphi, v)] \leftrightarrow v$. We call this space the **tangent space of X at x** .

Remark. As one can notice, the previous definition is analogous to one of the classical definitions of the tangent space. However, not all the classical definitions coincide in the infinite-dimensional setting. For example, although in general this definition can be seen to coincide with the one given as equivalence classes of curves with the same velocity, it does not coincide with the one given using derivations. In a general case in which the model space is not reflexive there are more derivations than tangent vectors.

Now that we have a good notion of what the tangent space is in an infinite-dimensional setting one can give all the typical constructions derivated from it just like in the classical theory. In that way we can generalize to the context of Banach manifolds the notions of the tangent and cotangent bundles and all its tensor powers, vector fields, differential forms, the d operator, the induced maps on these sets, etc.

To finish the section, we generalize the notion of a Lie group. A **Banach Lie group** is a Banach smooth manifold that has a group structure consistent with its manifold structure, that is, such that the group multiplication

$$\begin{aligned} G \times G &\longrightarrow G \\ (g, h) &\longmapsto gh \end{aligned}$$

is a smooth map. As in the classical case, the **Lie algebra** \mathfrak{g} of G is just the tangent space at the identity $T_e G$, which again happens to be isomorphic to the space the set of (left) invariant vector fields, so it is naturally equipped with a **Lie bracket**.

§ 2. Symplectic manifolds and the momentum map

Definition 2.1. Let X be a Banach smooth manifold. A **symplectic form** on X is a non-degenerate closed 2-form on X , that is, a 2-form ω satisfying:

1. For each $x \in X$, $\omega_x : T_x X \times T_x X \rightarrow \mathbf{R}$ is continuous;
2. For each $x \in X$, ω_x is non-degenerate, i.e., if $\omega_x(u, v) = 0$ for all $v \in T_x X$, then $u = 0$;
3. ω_x is smooth in x ;

4. ω is closed, i.e., $d\omega = 0$.

The pair (X, ω) is called a *symplectic manifold*.

Note that the form ω_x defines a bounded linear map

$$\begin{aligned} T_x X &\longrightarrow T_x X^* \\ v &\longmapsto \omega(v, -). \end{aligned}$$

Unlike in the finite-dimensional case, the non-degeneracy condition does not imply that this map is bijective, although it does imply that it is injective.

Example 2.2. The cotangent space T^*M of any smooth manifold can be endowed with a symplectic structure in the following way. We first consider the bundle projection $\pi : T^*M \rightarrow M$ and the pull-back bundle

$$\begin{array}{ccc} \pi^*T^*M & \longrightarrow & T^*M \\ \downarrow & & \downarrow \\ T^*M & \xrightarrow{\pi} & M. \end{array}$$

This bundle has a tautological section $\theta \in \Gamma(T^*M, \pi^*T^*M)$: $\theta(y) = (y, y)$ for each $y \in T^*M$. The differential of this section $\omega = d\theta$ is a symplectic form on T^*M . Locally, if M is finite-dimensional and (x_1, \dots, x_n) are coordinates on M and cotangent vectors are parametrized by coordinates (y_1, \dots, y_n) , the form θ is defined by

$$\theta = \sum_{i=1}^n y_i dx_i,$$

and

$$\omega = \sum_{i=1}^n dy_i \wedge dx_i.$$

□

We say that a vector field v on a symplectic manifold (X, ω) is *symplectic* if the Lie derivative $L_v \omega$ vanishes, that is, if

$$L_v \omega = d(i_v \omega) + i_v(d\omega) = 0,$$

where i_v denotes the contraction. Since ω is closed, v is symplectic if and only if the 1-form $i_v \omega$ is closed. We say that v is *Hamiltonian* if $i_v \omega$ is exact. In that case, there exists a function $f : X \rightarrow \mathbf{R}$, called the *Hamiltonian* of v such that

$$df = -i_v \omega.$$

The minus sign in the last equality is just a widely used convention. Of course, if the first de Rham cohomology of X vanishes, $H_{\text{dR}}^1(X) = 0$, then every symplectic vector field is Hamiltonian. In particular, every symplectic vector field is *locally Hamiltonian*. Reciprocally, to every function $f : X \rightarrow \mathbf{R}$, we can define the *Hamiltonian vector field associated to f* as the vector field v_f determined by the equation

$$df = -i_{v_f} \omega.$$

Given two functions f and g , we define their *Poisson bracket* as the function

$$\{f, g\} = v_f g = -v_g f = -\{g, f\}.$$

Two functions are said to *Poisson commute* if $\{f, g\} = 0$.

Example 2.3. We can consider the cotangent bundle T^*M of some manifold M with the canonical symplectic structure $\omega = \sum dy_i \wedge dx_i$. If f and g are functions of (y_1, \dots, y_n) alone,

$$df = \sum_i \frac{\partial f}{\partial y_i} dy_i = -i_{v_f} \omega$$

so

$$v_f = \sum_i \frac{\partial f}{\partial y_i} \frac{\partial}{\partial x_i}.$$

Thus,

$$\{f, g\} = v_f g = \sum_i \frac{\partial f}{\partial y_i} \frac{\partial g}{\partial x_i} = 0.$$

□

Let G be a Banach Lie group acting symplectically on X , that is, if we denote the action by $\rho : G \rightarrow \text{Diff}(X)$, for every $g \in G$ we have

$$\rho(g)^* \omega = \omega.$$

Let \mathfrak{g} be the Lie algebra of G and \mathfrak{g}^* its dual Banach space. Recall that to every element $\xi \in \mathfrak{g}$ we can associate the vector field $\vec{\xi}$ defined as the infinitesimal generator of $\rho(\exp t\xi)$. Consider now a smooth map $\mu : X \rightarrow \mathfrak{g}^*$ and, for every $\xi \in \mathfrak{g}$ define the function

$$\begin{aligned} \mu_\xi : X &\longrightarrow \mathbf{R} \\ x &\longmapsto \langle \mu(x), \xi \rangle, \end{aligned}$$

where $\langle -, - \rangle$ denotes the natural pairing between \mathfrak{g} and its dual. We say that μ is a **momentum map**¹ for the action of G on X if for every $\xi \in \mathfrak{g}$, the vector field $\vec{\xi}$ is Hamiltonian with Hamiltonian μ_ξ , that is, if

$$d\mu_\xi = -i_{\vec{\xi}} \omega$$

or, equivalently, if

$$\langle d_x \mu(u), \xi \rangle = \omega(u, \vec{\xi}(x)),$$

for every $u \in T_x X$, where $d_x \mu : T_x X \rightarrow \mathfrak{g}^*$ is the derivative of μ at x .

Example 2.4. Let (X, ω) be any finite-dimensional symplectic manifold and $f_1, \dots, f_n : X \rightarrow \mathbf{R}$ be functions. If the vector fields v_{f_1}, \dots, v_{f_n} form the basis of a Lie sub-algebra \mathfrak{g} of the Lie algebra of vector fields on X , then the functions define a map

$$\begin{aligned} \mu : X &\longrightarrow \mathfrak{g}^* \\ x &\longmapsto \sum_{i=1}^n f_i(x) \xi_i, \end{aligned}$$

where $\{\xi_i\}$ is the dual basis of $\{v_{f_i}\}$. If the vector fields integrate to give the action of a Lie group G on X , μ is a momentum map for that action.

An special case of this example is that in which $n = \frac{1}{2} \dim X$ and the functions f_1, \dots, f_n pairwise Poisson commute. In that case the vector fields v_{f_1}, \dots, v_{f_n} generate an abelian Lie

¹Originally due to a bad translation by Marsden and Weinstein of the French term “*application moment*”, introduced by Souriau, it is not uncommon in the literature to call this the “moment” map. The physically correct term is “momentum” map since it is a generalization of the physical notions of linear and angular momentum.

algebra. If the functions are independent, that is, $df_1 \wedge \cdots \wedge df_n$ is generically nonzero, the momentum map, which is simply $f = (f_1, \dots, f_n) : X \rightarrow \mathbf{R}^n$, has the property that a generic fibre is an n -dimensional submanifold with n linearly independent commuting vector fields v_{f_1}, \dots, v_{f_n} . This is called a **completely integrable system**. Using the aforementioned properties, it can be easily shown that on the regular fibres the symplectic form vanishes, that the flow of any of the v_{f_i} is linear in them and that generic fibres are open sets in tori $\mathbf{R}^n/\mathbf{Z}^n$. This is the content of the theorem of Arnold-Liouville [Arn97]. \square

§ 3. Symplectic reduction

We are going now to introduce **symplectic reduction**, for which we will need to assume that we have a symplectic action of a Banach Lie group G on a symplectic manifold (X, ω) that admits a momentum map with the following technical conditions. First, assume that μ is **G -equivariant**, that is,

$$\mu(\rho(g)(x)) = (\text{ad}_g)^*(\mu(x)),$$

for every $g \in G$, where $\text{ad} : G \rightarrow \text{Aut}(\mathfrak{g})$ denotes the adjoint action

$$\text{ad}_g \xi = \left. \frac{d}{dt} \right|_{t=0} (g \exp(t\xi) g^{-1}).$$

As a consequence of this, G leaves $\mu^{-1}(0) \subset X$ invariant. The **symplectic quotient** is defined as the quotient set

$$Z = \mu^{-1}(0)/G,$$

and we have the following diagram

$$\begin{array}{ccc} \mu^{-1}(0) & \xhookrightarrow{j} & X \\ \downarrow \pi & & \\ Z = \mu^{-1}(0)/G & & \end{array}$$

where j is the inclusion and π is the natural projection. Assume now that $\mu^{-1}(0)$ is a submanifold of X and that for every $x \in \mu^{-1}(0)$, $T_x(\mu^{-1}(0)) = \ker(d_x \mu)$. In particular this is true if $0 \in \mathfrak{g}^*$ is a **regular value** of μ , that is, if $d_x \mu : T_x X \rightarrow \mathfrak{g}^*$ is surjective for every $x \in \mu^{-1}(0)$, by the implicit function theorem. Finally, assume that the action of G on $\mu^{-1}(0)$ is **free** (without fixed points) and that at each point $x \in \mu^{-1}(0)$ there is a **slice** $S_x \subset \mu^{-1}(0)$ for the action, i.e., a submanifold $S_x \subset \mu^{-1}(0)$, $x \in S_x$, transversal to the orbit Gx , that is

$$T_x(\mu^{-1}(0)) = T_x(S_x) + T_x(Gx).$$

Taking S_x small, the projection $\pi : \mu^{-1}(0) \rightarrow Z$ defines a homeomorphism of S_x onto an open set of Z , making Z into a manifold. In principle Z could be non-Hausdorff, so in order to get a Hausdorff manifold we also need to ask that the action is **proper**, that is, the map

$$\begin{aligned} G \times X &\longrightarrow X \times X \\ (g, x) &\longmapsto (\rho(g)(x), x), \end{aligned}$$

is proper. For details on why this properness condition is necessary to get a Hausdorff space, check [AMR12].

Theorem 3.1 (Marsden-Weinstein). *Under the previous conditions, there is a unique symplectic form ω_Z on the symplectic quotient Z such that*

$$\pi^* \omega_Z = j^* \omega,$$

on $\mu^{-1}(0)$.

Proof. We can easily define ω_Z by

$$\omega_Z(d_x \pi(u), d_x \pi(v)) = \omega(u, v),$$

for $u, v \in T_x(\mu^{-1}(0))$. To see that it is well defined choose another representative $u' \in T_{\rho(g)(x)}(\mu^{-1}(0))$, for some $g \in G$, such that $d_{\rho(g)(x)} \pi(u') = d_x \pi(u)$. Then, u and u' are related by

$$u' = d_x \rho(g)(u + \vec{\xi}(x)),$$

for some $\xi \in \mathfrak{g}$, since

$$T_{\rho(g)(x)}(\mu^{-1}(0)) = d_x \rho(g)(T_x(\mu^{-1}(0))) = d_x \rho(g)(T_x(S_x) + T_x(Gx)).$$

Therefore,

$$\omega_{\rho(g)x}(u', v') = (\rho(g)^* \omega)_x(u + \vec{\xi}(x), v + \vec{\eta}(x)).$$

Now, since the action is symplectic, $\rho(g)^* \omega = \omega$, so

$$\begin{aligned} \omega_{\rho(g)x}(u', v') &= \omega_x(u + \vec{\xi}(x), v + \vec{\eta}(x)) \\ &= \omega_x(u, v) + \omega_x(\vec{\xi}(x), v) \\ &\quad + \omega_x(u, \vec{\eta}(x)) + \omega_x(\vec{\xi}(x), \vec{\eta}(x)) \\ &= \omega_x(u, v) + \langle d_x \mu(u), \eta \rangle - \langle d_x \mu(v), \xi \rangle + \langle d_x \mu(\vec{\xi}(x)), \eta \rangle \\ &= \omega_x(u, v), \end{aligned}$$

since $u, v, \vec{\xi}(x) \in T_x(\mu^{-1}(0)) = \ker(d_x \mu)$.

This gives the existence of ω_Z , while the uniqueness follows from the fact that $d_x \pi : T_x(\mu^{-1}(0)) \rightarrow T_{\pi(x)}Z$ is surjective. Since $\pi : S_x \rightarrow \pi(S_x) \subset Z$ is a diffeomorphism and $\pi^* \omega_Z|_{S_x} = \omega|_{S_x}$, ω_Z is also smooth and closed.

It remains to check that ω_Z is non-degenerate. Let $u \in T_x(\mu^{-1}(0))$ be such that $\omega(u, v) = 0$ for all $v \in T_x(\mu^{-1}(0))$. We have to show that $d_x \pi(u) = 0$, that is, that $u = \vec{\xi}(x)$ for some $\xi \in \mathfrak{g}$. To see this we need the following technical lemma, we refer to [Kob14] for a proof.

Lemma 3.2. *Let E be a Banach space and $\omega : E \times E \rightarrow \mathbf{R}$ a continuous non-degenerate skew-symmetric form. For any closed subspace $F \subset E$, set*

$$F^\omega = \{v \in E : \omega(u, v) = 0 \text{ for all } u \in F\}.$$

Then $(F^\omega)^\omega = F$.

In our case $E = T_x X$ and $F = \{\vec{\xi}(x) : \xi \in \mathfrak{g}\} = \{u \in T_x(\mu^{-1}(0)) : d_x \pi(u) = 0\}$. Therefore

$$F^\omega = \{v \in T_x X : \omega(\vec{\xi}(x), v) = 0 \text{ for all } \xi \in \mathfrak{g}\} = \{v \in T_x X : d_x \mu(v) = 0\} = T_x(\mu^{-1}(0)),$$

since $\ker(d_x \mu) = T_x(\mu^{-1}(0))$. Now, applying the Lemma

$$F = (F^\omega)^\omega = \{u \in T_x X : \omega(u, v) = 0 \text{ for all } v \in T_x(\mu^{-1}(0))\},$$

which is exactly what we wanted to prove. \square

Remark. Let M be a symplectic manifold with the action of a Lie group G and momentum map $\mu : X \rightarrow \mathfrak{g}^*$ defined by functions f_1, \dots, f_n as in example 2.4. If g is a G -invariant function on M by restriction we can define a G -invariant function \tilde{g} on $\mu^{-1}(0)$ and therefore on the quotient $\mu^{-1}(0)/G$. If g, h are two such functions such that $\{g, h\} = 0$, $X_g h = 0$, so h is constant along the orbits of X_g . But the projection of these orbits to the quotients are the orbits of $X_{\tilde{g}}$, so \tilde{h} is constant on these orbits and so $\{\tilde{g}, \tilde{h}\} = 0$.

§ 4. The moduli space as a symplectic quotient

In this section we are going to construct the moduli space of stable holomorphic vector bundles $\mathbf{M}_{\text{Bun}}(n, d)$, regarded as the moduli space of irreducible constant central curvature unitary connections $\mathcal{A}_h^s/\mathcal{G}$, as a symplectic quotient.

Let X be a compact Riemann surface of genus $g \geq 2$, E a complex vector bundle on X of rank n and degree d and h a Hermitian metric on E . Consider the set \mathcal{A}_h of all h -unitary connections on E . Any two h -unitary connections can be seen to differ in a 1-form with values in $\mathfrak{u}_h(E)$, the Lie algebra of $U_h(E)$. Therefore, \mathcal{A}_h is an affine space modeled on $\Omega^1(X, \mathfrak{u}_h(E))$. This space admits a non-degenerate skew-symmetric form

$$\omega(A, B) = - \int_X \text{tr}(A \wedge B),$$

endowing \mathcal{A}_h , at least formally, with the structure of a symplectic manifold. Strictly speaking, we would need to give in \mathcal{A}_h the structure of a Banach manifold. This is a technical procedure that we will not detail here, but the essential idea is to give completions of this space with respect to **Sobolev norms**. In a similar fashion, for our purposes we also need to give Sobolev completions of the group $\mathcal{G} = \Gamma(U_h(E))$, in order to get a Banach Lie group. Go to [AB83, Kob14] for explicit constructions.

Remember that the Lie group \mathcal{G} acts on \mathcal{A}_h by conjugation $\nabla \mapsto g\nabla g^{-1}$.

Proposition 4.1. *This action does a momentum map, given by*

$$\mu(\nabla) = -F - 2\pi i \mu(E) \mathbf{1}_E \omega_X,$$

where F is the curvature of ∇ .

Remark. This definition makes sense if we identify the Lie algebra of \mathcal{G} as $\text{Lie } \mathcal{G} = \Gamma(\mathfrak{u}_h(E))$ and its dual with $(\text{Lie } \mathcal{G})^* = \Omega^2(X, \mathfrak{u}_h(E))$, via the pairing

$$\langle \alpha, \xi \rangle = \int_X \text{tr}(\xi \alpha),$$

for $\xi \in \Gamma(\mathfrak{u}_h(E))$ and $\alpha \in \Omega^2(X, \mathfrak{u}_h(E))$.

Remark. In fact, $\mu(\nabla) = -F$ could be also a suitable momentum map, however, the second term is added in there in order to get a non-empty symplectic quotient. Indeed, $\mu^{-1}(0)$ would be empty unless $\deg E = 0$.

Proof. To see how $\vec{\xi}$ looks like for an element $\xi \in \Gamma(\mathfrak{u}_h(E))$, just compute

$$\vec{\xi}(\nabla) = \left. \frac{d}{dt} \right|_{t=0} \exp(t\xi) \nabla \exp(-t\xi) = -\nabla \xi.$$

Let us compute also $d_{\nabla}\mu(A)$ for some $A \in \Omega^1(X, \mathfrak{u}_h(E))$,

$$d_{\nabla}\mu(A) = \left. \frac{d}{dt} \right|_{t=0} \mu(\nabla + tA) = \left. \frac{d}{dt} \right|_{t=0} [-(\nabla + tA) \circ (\nabla + tA) - 2\pi i \mu(E) \mathbf{1}_E \omega X] = -\nabla A.$$

Therefore we have

$$\langle d_{\nabla}\mu(A), \xi \rangle = - \int_X \text{tr}(\xi \nabla A) = \int_X \text{tr}(\nabla \xi \wedge A) = \omega(A, \nabla \xi) = \omega(A, \vec{\xi}),$$

so μ is a momentum map for the action of \mathcal{G} in \mathcal{A}_h . □

The \mathcal{G} -action can be shown to verify all the technical conditions for it to define a symplectic structure on the symplectic quotient $\mu^{-1}(0)/\mathcal{G}$. In general, this quotient will classify **polystable** holomorphic vector bundles. To get stable vector bundles we consider the submanifold \mathcal{A}_h^s of irreducible h -unitary connections and the same method can be applied to construct the moduli space $\mathbf{M}_{\text{Bun}}(n, d) = \mathcal{A}_h^s/\mathcal{G}$ as a symplectic quotient $\mu^{-1}(0)/\mathcal{G}$. All the good properties of the action guarantee that $\mathbf{M}_{\text{Bun}}(n, d)$ is indeed a smooth manifold.

CHAPTER III

Higgs bundles and the Hitchin system

§ 1. The Hitchin system

Recall that the moduli space $\mathbf{M}_{\text{Bun}}(n, d)$ has the structure of a complex manifold and that its (smooth) tangent space at some point \mathbf{E} is isomorphic to $H^1(X, \text{End } \mathbf{E})$. Now, by Serre duality, the cotangent space $T_{\mathbf{E}}^* \mathbf{M}_{\text{Bun}}(n, d)$ is isomorphic to $H^0(X, \text{End } \mathbf{E} \otimes K_X)$, where K_X denotes the canonical line bundle of the Riemann surface X , that is, the cotangent bundle $K_X = (T^{1,0} X)^*$. As we shall see, the cotangent bundle $T^* \mathbf{M}_{\text{Bun}}(n, d)$ can be given the structure of a symplectic manifold and it admits a completely integrable system, the Hitchin system. More specifically, what we will prove is the following

Theorem 1.1 (Hitchin). *Let X be a Riemann surface and $\mathbf{M}_{\text{Bun}}(n, d)$ the moduli space of stable holomorphic vector bundles of rank n and degree d on X . Let $k = n^2(g - 1) + 1$ be the complex dimension of $\mathbf{M}_{\text{Bun}}(n, d)$. There is a map*

$$H : T^* \mathbf{M}_{\text{Bun}}(n, d) \rightarrow \mathcal{B},$$

where \mathcal{B} is a complex vector space of complex dimension k , such that its components Poisson commute and generic fibres of it are open sets in some k -dimensional complex tori.

Example 1.2. For the line bundle case this result is trivial. The moduli space $\mathbf{M}_{\text{Bun}}(1, d) = \text{Pic}^d(X)$ is the degree d component of the Picard group $\text{Pic}(X) = H^1(X, \mathcal{O}_X^*)$, that, as we saw, can be identified with the Jacobian of the curve, $J(X)$, which is indeed a complex torus. The tangent bundle of the Jacobian is trivial and isomorphic to $H^1(X, \mathcal{O}_X)$, so by Serre duality the cotangent bundle is

$$T^* \mathbf{M}_{\text{Bun}}(1, d) = J(X) \times H^0(X, K_X).$$

The dimension of $H^0(X, K_X)$ is $g = \dim \mathbf{M}_{\text{Bun}}(1, d)$, so we can take $\mathcal{B} = H^0(X, K_X)$ and define

$$H : T^* \mathbf{M}_{\text{Bun}}(1, d) \rightarrow \mathcal{B}$$

as the projection on the second factor $\text{pr}_2 : J(X) \times H^0(X, K_X) \rightarrow H^0(X, K_X)$. \square

First of all, let us describe the symplectic form on $T^* \mathbf{M}_{\text{Bun}}(n, d)$. In order to do this, we are going to regard $T^* \mathbf{M}_{\text{Bun}}(n, d)$ as

$$T^* \mathbf{M}_{\text{Bun}}(n, d) = T^*(\mathcal{A}_{\bar{\partial}}^s / \mathcal{G}^c).$$

Fix now a smooth vector bundle E over X , of rank n and degree d . Now we define the “complexification”

$$\mathcal{A}^c = \mathcal{A}_{\bar{\partial}}^s \times \Omega^{1,0}(X, \text{End}E),$$

which is an open subset of a complex affine space over $\Omega^{0,1}(X, \text{End}E) \oplus \Omega^{1,0}(X, \text{End}E)$ and $\mathcal{A}^c = T^*\mathcal{A}_{\bar{\partial}}^s$. This space carries a natural symplectic form

$$\omega((A_1, \varphi_1), (A_2, \varphi_2)) = 2i \int_X \text{tr}(A_1 \wedge \varphi_2 - A_2 \wedge \varphi_1),$$

where $A_i \in \Omega^{0,1}(X, \text{End}E)$ and $\varphi_i \in \Omega^{1,0}(X, \text{End}E)$. If we denote points of \mathcal{A}^c as pairs $(\bar{\partial}_E, \varphi)$, \mathcal{A}^c has a natural action of \mathcal{G}^c that we can write as

$$(\bar{\partial}_E, \varphi) \mapsto (g\bar{\partial}_E g^{-1}, g\varphi g^{-1}).$$

We now define the momentum map

$$\begin{aligned} \mu : \mathcal{A}^c = \mathcal{A}_{\bar{\partial}}^s \times \Omega^{1,0}(X, \text{End}E) &\longrightarrow (\text{Lie}\mathcal{G}^c)^* = \Omega^2(X, \text{End}E) \\ (\bar{\partial}_E, \varphi) &\longmapsto -2i\bar{\partial}_E\varphi. \end{aligned}$$

To check that this is indeed the momentum map for this action, compute

$$\begin{aligned} d_{(\bar{\partial}_E, \varphi)}\mu(A, \psi) &= \left. \frac{d}{dt} \right|_{t=0} \mu((\bar{\partial}_E, \varphi) + t(A, \psi)) \\ &= \left. \frac{d}{dt} \right|_{t=0} 2i(\bar{\partial}_E(\varphi + t\psi) + t[A, (\varphi + t\psi)]) \\ &= -2i(\bar{\partial}_E\psi + [A, \varphi]). \end{aligned}$$

Recall now that an element of the Lie algebra $\xi \in \Gamma(\text{End}E)$ acts on $\Omega^{0,1}(X, \text{End}E)$ as $-\bar{\partial}_E\xi$. On a similar way one shows that the action on $\Omega^{1,0}(X, \text{End}E)$ is given by $[\xi, \varphi]$, so

$$\vec{\xi}(\bar{\partial}_E, \varphi) = (-\bar{\partial}_E\xi, [\xi, \varphi]).$$

Therefore

$$\omega((A, \psi), \vec{\xi}) = 2i \int_X \text{tr}(A \wedge [\xi, \varphi] + \bar{\partial}_E\xi \wedge \psi) = -2i \int_X \text{tr}(\xi[A, \varphi] + \xi\bar{\partial}_E\psi) = \langle d_{\bar{\partial}_E, \varphi}\mu(A, \psi), \xi \rangle.$$

Therefore we get a symplectic quotient

$$T^*\mathbf{M}_{\text{Bun}}(n, d) = \mathcal{A}^c / \mathcal{G}^c = \mu^{-1}(0) / \mathcal{G}^c.$$

Note also that for a pair $(\bar{\partial}_E, \varphi) \in \mathcal{A}^c$ we have that $\mu(\bar{\partial}_E, \varphi) = 0$ if $\bar{\partial}_E\varphi = 0$, that is, if $\varphi \in H_{\bar{\partial}_E}^{1,0}(X, \text{End}E) = H^0(X, \text{End}E \otimes K_X)$, just as we wanted.

Let us now construct the map

$$H : T^*\mathbf{M}_{\text{Bun}}(n, d) \rightarrow \mathcal{B}.$$

First of all, consider any element $\varphi \in \Omega^{1,0}(X, \text{End}E) = \Gamma(X, \text{End}E \otimes K_X)$. Associated to this element we have the characteristic polynomial, formally written as

$$\det(T - \varphi) = T^n + \sum_{i=1}^n \sigma_i(\varphi) T^{n-i},$$

with its coefficients being sections $\sigma_i(\varphi) \in \Gamma(K_X^i)$. This defines a map

$$\begin{aligned} \mathcal{A}_{\bar{\partial}} \times \Omega^{1,0}(X, \text{End} E) &\longrightarrow \bigoplus_{i=1}^n \Gamma(K_X^i) \\ (\partial_{\mathbf{E}}, \varphi) &\longmapsto (\sigma_1(\varphi), \dots, \sigma_n(\varphi)). \end{aligned}$$

Since the components of this map are functions of $\Omega^{1,0}(X, \text{End} E)$ alone, they Poisson commute and also will the components of the map defined in the symplectic quotient

$$\begin{aligned} H : T^* \mathbf{M}_{\text{Bun}}(n, d) &\longrightarrow \bigoplus_{i=1}^n H^0(X, K_X^i) \\ (\mathbf{E}, \varphi) &\longmapsto (\sigma_1(\varphi), \dots, \sigma_n(\varphi)). \end{aligned}$$

This is the **Hitchin map**. We define now the vector space $\mathcal{B} = \bigoplus_{i=1}^n H^0(X, K_X^i)$ and let us check that it has the desired dimension. This is a straightforward computation using Riemann-Roch:

$$\begin{aligned} \dim \mathcal{B} &= \sum_{i=1}^n \dim H^0(X, K_X^i) = \dim H^0(X, K_X) + \sum_{i=2}^n \dim H^0(X, K_X^i) \\ &= g + \sum_{i=2}^n [i(2g-2) - g + 1] = 1 + \sum_{i=1}^n [i(2g-2) - g + 1] \\ &= 1 + (2g-2) \frac{n(n+1)}{2} - n(g-1) = 1 + (g-1)(n^2+n) - n(g-1) \\ &= n^2(g-1) + 1. \end{aligned}$$

To finish the proof of Hitchin's theorem it remains to check that the fibers are open sets in $n^2(g-1) + 1$ -dimensional complex tori. In order to prove this, we are going to need a very powerful tool that will be introduced in the following section.

§ 2. The spectral correspondence

In this section we are going to study the spectral data of “twisted” endomorphisms on vector bundles. We are going to work in a more general setting than the previous section, in which we were only considering pairs (\mathbf{E}, φ) where $\varphi \in \text{End } \mathbf{E} \otimes K_X$ was an endomorphism “twisted” by the canonical line bundle K_X . We are now going to allow twisting by any holomorphic line bundle L .

Definition 2.1. Let X be a compact Riemann surface and let $\mathbf{E} \rightarrow X$ be a holomorphic vector bundle of rank n and degree d . Let us fix a holomorphic line bundle $L \rightarrow X$. A **L -twisted endomorphism** is a bundle homomorphism $\varphi : \mathbf{E} \rightarrow \mathbf{E} \otimes L$, or equivalently, a holomorphic section $\varphi \in H^0(X, \text{End } \mathbf{E} \otimes L)$.

Any L -twisted endomorphism induces, for each $i = 1, \dots, n$ a homomorphism

$$\wedge^i \varphi : \wedge^i E \longrightarrow \wedge^i (E \otimes L) = \wedge^i E \otimes L^i.$$

We can then take the traces of these homomorphisms and get sections $\text{tr } \wedge^i \varphi \in H^0(X, L^i)$. With these sections we can construct the characteristic polynomial of φ ,

$$P_{\varphi}(T) = \det(T - \varphi) = T^n + \sum_{i=1}^n (-1)^i \sigma_i(\varphi) T^{n-i},$$

where the coefficients are precisely $\sigma_i(\varphi) = (-1)^i \text{tr} \wedge^i \varphi \in H^0(X, L^i)$. For example,

$$\sigma_1(\varphi) = -\text{tr} \varphi, \quad \sigma_n(\varphi) = (-1)^n \det \varphi.$$

Therefore, to any L -twisted endomorphism φ we can associate an element

$$(\sigma_1(\varphi), \dots, \sigma_n(\varphi)) \in \bigoplus_{i=1}^n H^0(X, L^i).$$

Let us now consider any element $b = (b_1, \dots, b_n) \in \bigoplus_{i=1}^n H^0(X, L^i)$. If we take the pullback bundle of L , p^*L given by the diagram

$$\begin{array}{ccc} p^*L & \longrightarrow & L \\ \downarrow & & \downarrow p \\ L & \xrightarrow{p} & X, \end{array}$$

where $p : L \rightarrow X$ is just the canonical projection of the bundle, it is easy to check that it has a tautological section $\lambda \in H^0(L, p^*L)$. Locally, λ can be seen as a coordinate on the total space of the bundle L . Let us define then the section

$$s_b = \lambda^n + \sum_{i=1}^n p^* b_i \lambda^{n-i} \in H^0(L, p^*L^n).$$

Definition 2.2. The *spectral curve* S_b associated to an element $b \in \bigoplus_{i=1}^n H^0(X, L^i)$ is the zero locus of the section s_b ,

$$S_b = (s_b)_0 \subset L.$$

Note that generic values of $b \in \bigoplus_{i=1}^n H^0(X, L^i)$ define an “irreducible polynomial” and therefore the spectral curve will be generically irreducible.

Near some point of X we take some neighbourhood U and think of the spectral curve as the set

$$\left\{ (x, \lambda) \in U \times \mathbf{C} : \lambda^n + \sum_{i=1}^n b_i \lambda^{n-i} = 0 \right\}.$$

Therefore, if $b = (\sigma_1(\varphi), \dots, \sigma_n(\varphi))$ for some L -twisted endomorphism $\varphi : \mathbf{E} \rightarrow \mathbf{E} \otimes L$, then locally the spectral curve can be thought as the set

$$\{(x, \lambda) \in U \times \mathbf{C} : \det(\lambda \mathbf{1} - \varphi) = 0\}$$

That is, fiberwise, over some point $x \in X$ the points of the spectral curve are precisely the “eigenvalues” $\lambda_1(x), \dots, \lambda_n(x) \in L_x$ of the twisted endomorphism $\varphi_x : E_x \rightarrow E_x \otimes L_x$; hence the name *spectral curve*.

[dibujo]

Proposition 2.3. Assume that L^n is base point free. Then, for generic elements $b \in \bigoplus_{i=1}^n H^0(X, L^i)$, the spectral curve S_b is smooth.

Proof. The spectral curve S_b is an irreducible hypersurface of L , so it can be thought as a divisor on it and on its compactification $\mathbf{P}(L \oplus \mathcal{O}_X)$. Moreover, as the b varies, s_b defines a linear subspace of $\mathbf{P}(H^0(L, p^*L^n))$, so

$$\mathfrak{d} = \left\{ S_b : b \in \bigoplus_{i=1}^n H^0(X, L^i) \right\}$$

is a linear system of divisors on L . Bertini's theorem says that, away from the base locus of \mathfrak{d} , the generic divisor of the system is smooth. Let us check then that the base locus of \mathfrak{d} is empty.

Suppose that $y \in L$ is a base point of \mathfrak{d} . Then, since $(\lambda^n)_0 \in \mathfrak{d}$, $\lambda(y) = 0$. But then $s_b(y) = p^*b_n(y) = b_n(p(y))$ for every b and, since y is a base point, $s_b(y) = 0$. Therefore $b_n(p(y)) = 0$ for every $b_n \in H^0(X, L^n)$, so $p(y)$ is a base point of L^n . But by hypothesis L^n is base point free, so we reach a contradiction. \square

Note now that the restriction of the natural projection $p : L \rightarrow X$ defines a map,

$$\begin{array}{ccc} S_b & \hookrightarrow & L \\ & \searrow \pi & \downarrow p \\ & & X. \end{array}$$

Since $\pi : S_b \rightarrow X$ is a morphism of Riemann surfaces it defines a branched covering of S_b over X . Note that at a point x where this morphism is étale the fiber has exactly n points, the n eigenvalues of φ_x . Therefore the degree of the map π is $\deg \pi = n$. Similarly, the **branch locus** will be given precisely by those points $x \in X$ where some eigenvalues of φ_x have algebraic multiplicity > 1 . By the same reason, the **ramification divisor** R on S_b is defined as the zero divisor of $\text{Disc}(s_b) \in H^0(X, p^*L^{n(n-1)})$, which is the discriminant of s_b , or equivalently, the resultant of s_b and its derivative

$$\text{Disc}(s_b) = \text{Res} \left(s_b, \frac{\partial}{\partial \lambda} s_b \right).$$

R is the zero locus of a section of $p^*L^{n(n-1)}$, so its degree is precisely

$$\deg R = \deg(L^{n(n-1)}) = n(n-1) \deg L.$$

The Riemann-Hurwitz formula yields the genus of the spectral curve,

$$2g_{S_b} - 2 = n(2g - 2) + \deg R = n(2g - 2) + n(n-1) \deg L,$$

$$g_{S_b} = 1 + n(g - 1) + \frac{n(n-1)}{2} \deg L.$$

We are now in position to prove the spectral correspondence:

Theorem 2.4 (Beauville-Narasimhan-Ramanan, [BNR89]). *Let $b \in \bigoplus_{i=1}^n H^0(X, L^i)$ such that the spectral curve S_b is irreducible and smooth. We have a bijective correspondence, up to isomorphism, between holomorphic line bundles over S_b of degree δ and pairs (\mathbf{E}, φ) , where \mathbf{E} is a holomorphic vector bundle over X of rank n and degree d , and φ is a L -twisted endomorphism with characteristic polynomial*

$$P_\varphi(T) = P_b(T) := T^n + \sum_{i=1}^n b_i T^{n-i}.$$

The degrees d and δ are related by

$$d = \delta - \frac{n(n-1)}{2} \deg L.$$

Proof. Let M be a holomorphic line bundle over S_b of degree δ . We can consider the direct image bundle $\pi_* M$, which is a rank $\deg \pi = n$ vector bundle over X and its degree is given by the formula

$$\deg(\pi_* M) = \deg M + (1 - g_{S_b}) - \deg \pi(1 - g) = \delta - \frac{n(n-1)}{2} \deg L.$$

Let now $U \subset X$ be any open subset. If we take tensor product by the tautological section restricted to $\pi^{-1}(U)$, $\lambda|_{\pi^{-1}(U)} \in H^0(\pi^{-1}(U), \pi^* L)$ we can construct a homomorphism

$$\otimes \lambda|_{\pi^{-1}(U)} : H^0(\pi^{-1}(U), M) \rightarrow H^0(\pi^{-1}(U), M \otimes \pi^* L).$$

But, by definition of the direct image sheaf,

$$H^0(\pi^{-1}(U), M) = H^0(U, \pi_* M),$$

$$H^0(\pi^{-1}(U), M \otimes \pi^* L) = H^0(U, \pi_*(M \otimes \pi^* L)) = H^0(U, \pi_* M \otimes L).$$

Therefore, we have a homomorphism of locally free sheaves

$$\otimes \lambda|_{\pi^{-1}(U)} : \mathcal{O}(\pi_* M) \rightarrow \mathcal{O}(\pi_* M \otimes L),$$

that induces a L -twisted endomorphism

$$\varphi : \pi_* M \rightarrow \pi_* M \otimes L.$$

Moreover, by construction $P_b(\varphi) = 0$, and since P_b is irreducible, the Cayley-Hamilton theorem guarantees that P_b is the characteristic polynomial of φ .

On the other hand, take a pair (\mathbf{E}, φ) , with \mathbf{E} a rank n holomorphic vector bundle and φ a L -twisted endomorphism such that $b = (\sigma_1(\varphi), \dots, \sigma_n(\varphi))$. Consider the pullback bundle $\pi^* \mathbf{E}$,

$$\begin{array}{ccc} \pi^* \mathbf{E} & \xrightarrow{\quad} & \mathbf{E} \\ \downarrow & & \downarrow \\ S_b & \xrightarrow{\quad \pi \quad} & X. \end{array}$$

The pullback section $\pi^* \varphi \in H^0(S_b, \text{End}(\pi^* \mathbf{E}) \otimes \pi^* L)$ must satisfy

$$P_{\pi^* \varphi}(\lambda|_{S_b}) = (\lambda|_{S_b})^n + \sum_{i=1}^n \pi^* b_i (\lambda|_{S_b})^{n-i} = s_b = 0,$$

since $S_b = (s_b)_0$. Therefore, locally we can see λ as an eigenvalue of φ and one could construct the line bundle M , at least away from the ramification divisor, as the bundle defined at each point x by the eigenspace of φ_x associated to λ_x . Globally what we do is to consider $\lambda|_{S_b} \mathbf{1} - \pi^* \varphi$ as a morphism of locally free sheaves and the sheaf $\ker(\lambda|_{S_b} \mathbf{1} - \pi^* \varphi) \subset \pi^* \mathbf{E}$, which is a rank one locally free sheaf, so it defines a line bundle M over S_b such that $\pi_*(M) = \mathbf{E}$. Another direct way to do it is to define M by the exact sequence

$$0 \longrightarrow M(-R) \longrightarrow \pi^* \mathbf{E} \xrightarrow{\lambda|_{S_b} \mathbf{1} - \pi^* \varphi} \pi^*(\mathbf{E} \otimes L) \longrightarrow M \otimes \pi^* L \longrightarrow 0.$$

The naturality of the constructions implies that the correspondence is given up to isomorphism. It is also clear from the constructions that they are inverse to each other. \square

§ 3. Hitchin's equations

Let us now apply the results of the previous section to the case of the Hitchin system

$$H : T^*\mathbf{M}_{\text{Bun}}(n, d) \longrightarrow \mathcal{B} = \bigoplus_{i=1}^n H^0(X, K_X^i)$$

$$(\mathbf{E}, \varphi) \longmapsto (\sigma_1(\varphi), \dots, \sigma_n(\varphi)).$$

The fiber of some point $b = (b_1, \dots, b_n) \in \mathcal{B}$ by this map is formed precisely by (isomorphism classes of) that pairs (\mathbf{E}, φ) such that \mathbf{E} is a holomorphic *stable* vector bundle of rank n and degree d and $\varphi : \mathbf{E} \rightarrow \mathbf{E} \otimes K_X$ is a K_X -twisted endomorphism with $P_\varphi(T) = P_b(T)$. Moreover, K_X^n is base point free, so for generic $b \in \mathcal{B}$ the spectral curve S_b will be irreducible and smooth. Stability is an open condition, so in this case the Beauville-Narasimhan-Ramanan theorem says that the fiber of a generic b is an open subset of the set of isomorphism classes of holomorphic line bundles over S_b of degree

$$\delta = d + \frac{n(n-1)}{2} \deg K_X = d + n(n-1)(g-1).$$

But this set of isomorphism classes is precisely the component $\text{Pic}^\delta(S_b)$ of the Picard group, which is isomorphic to the Jacobian $J(S_b)$; so it is indeed a complex torus! Moreover, the dimension of this complex torus is the genus of the spectral curve, and if we compute it the miracle occurs:

$$g_{S_b} = 1 + n(g-1) + \frac{n(n-1)}{2} \deg K_X = 1 + n^2(g-1).$$

Summing up, we have proven that for some point $b \in \mathcal{B}$, the fiber of the Hitchin map $H^{-1}(b)$ is an open set of $\text{Pic}^\delta(S_b)$, a complex torus of dimension $n^2(g-1) + 1$. This completes the proof of Hitchin's theorem 1.1.

A natural question to ask now is what happens to the remaining elements of the Jacobian: if they do not yield elements of $T^*\mathbf{M}_{\text{Bun}}(n, d)$, what do they correspond to? This suggests the existence of a wider symplectic manifold, of which $T^*\mathbf{M}_{\text{Bun}}(n, d)$ is a subset.

Recall that we constructed the space $T^*\mathbf{M}_{\text{Bun}}(n, d)$ by fixing a smooth vector bundle E of rank n and degree d over X and considering the action of \mathcal{G}^c on the space $\mathcal{A}^c = \mathcal{A}_{\bar{\partial}}^s \times \Omega^{1,0}(X, \text{End}E)$. To get this wider space we will now consider the action of the real group \mathcal{G} on the space $\mathcal{A}_h \times \Omega^{1,0}(X, \text{End}E)$. Of course, to do this we first have to fix an hermitian metric h on E . We then get a symplectic structure on $\mathcal{A}_h \times \Omega^{1,0}(X, \text{End}E)$ given by

$$\omega((A_1, \varphi_1), (A_2, \varphi_2)) = - \int_X \text{tr}(A_1 \wedge A_2) - 2i \text{Im} \text{tr}(\varphi_1 \wedge \varphi_2^\dagger),$$

where $A_i \in \Omega^1(X, \mathfrak{u}_h(E))$ and $\varphi_i \in \Omega^{1,0}(X, \text{End}E)$.

Remark. Let us see explicitly what we mean by $\varphi_1 \wedge \varphi_2^\dagger$. φ is an element of $\Omega^{1,0}(X, \text{End}E) = \Gamma(\Lambda^{1,0}(X) \otimes \text{End}E)$, thus, locally it can be seen as $\varphi = \phi \otimes \alpha$, with $\phi : X \rightarrow \text{End}E$ and $\alpha : X \rightarrow \Lambda^{1,0}(X)$. By φ^\dagger then we mean an element of $\Omega^{0,1}(X, \text{End}E)$ which locally looks like $\varphi^\dagger = \phi^\dagger \otimes \bar{\alpha}$ with ϕ^\dagger the transpose conjugate (globally the adjoint, given by h) of ϕ and $\bar{\alpha} : X \rightarrow \Lambda^{0,1}(X)$ the conjugate of α . Therefore by $\varphi_1 \wedge \varphi_2^\dagger$ we mean an element of $\Omega^{1,1}(X, \text{End}E)$, which locally looks like

$$\varphi_1 \wedge \varphi_2^\dagger = (\phi_1 \phi_2^\dagger) \otimes \alpha_1 \wedge \bar{\alpha}_2.$$

We can also consider elements of the form

$$[\varphi_1, \varphi_2] := \varphi_1 \wedge \varphi_2 + \varphi_2 \wedge \varphi_1 = (\phi_1 \phi_2) \otimes \alpha_1 \wedge \alpha_2 + (\phi_2 \phi_1) \otimes \alpha_2 \wedge \alpha_1 = [\phi_1, \phi_2] \otimes \alpha_1 \wedge \alpha_2.$$

Remember that the elements g of the group act by conjugation $(\nabla, \varphi) \mapsto (g\nabla g^{-1}, g\varphi g^{-1})$.

Proposition 3.1. *The action of \mathcal{G} on $\mathcal{A}_h \times \Omega^{1,0}(X, \text{End} E)$ admits a momentum map,*

$$\mu(\nabla, \varphi) = -F - [\varphi, \varphi^\dagger] - 2\pi i \mu(E) \mathbf{1}_E \omega_X.$$

Proof. The part of the momentum map corresponding to $-F - 2\pi i \mu(E) \mathbf{1}_E \omega_X$ comes from the action on \mathcal{A}_h and the proof is identical to the case of the previous chapter.

On the other hand, taking into account the previous remark we can focus in computing the momentum map for a much simpler case:

Lemma 3.2. *Consider the space $\text{End } \mathbf{C}^n$ equipped with the symplectic form*

$$\omega(A, B) = -\text{Im } \text{tr}(AB^\dagger),$$

and consider the action of $U(n)$ on it by conjugation. This action admits a momentum map

$$\mu(A) = \frac{i}{2} [A, A^\dagger].$$

Proof. First of all, notice that we can identify $\mathfrak{u}(n)$ with its dual via the pairing

$$\langle H, K \rangle = \text{tr}(HK^\dagger).$$

The Lie algebra $\mathfrak{u}(n)$ is formed by skew-hermitian matrices, so

$$\overline{\langle H, K \rangle} = \text{tr}(KH^\dagger) = \text{tr}((-K^\dagger)(-H)) = \text{tr}(K^\dagger H) = \text{tr}(HK^\dagger) = \langle H, K \rangle.$$

That is, $\langle H, K \rangle$ is indeed a real number. Also notice that

$$\mu(A)^\dagger = -\frac{i}{2} [A, A^\dagger] = -\mu(A),$$

so $\mu(A) \in \mathfrak{u}(n)$; the momentum map is well defined.

To see how \vec{H} looks like for an element $H \in \mathfrak{u}(n)$, just compute

$$\vec{H}(A) = \left. \frac{d}{dt} \right|_{t=0} (e^{-Ht} A e^{Ht}) = -HA + AH = [A, H].$$

Let us compute also $d_A \mu(B)$ for some $B \in \text{End } \mathbf{C}^n$

$$d_A \mu(B) = \left. \frac{d}{dt} \right|_{t=0} \mu(A + tB) = \frac{i}{2} \left. \frac{d}{dt} \right|_{t=0} [A + tB, A^\dagger + tB^\dagger] = \frac{i}{2} ([B, A^\dagger] + [A, B^\dagger]).$$

Now

$$\langle d_A \mu(B), H \rangle = \frac{i}{2} \text{tr}([B, A^\dagger]H^\dagger + [A, B^\dagger]H^\dagger) = -\frac{i}{2} \text{tr}([B, A^\dagger]H + [A, B^\dagger]H)$$

Using the cyclic property of the trace one gets

$$\langle d_A \mu(B), H \rangle = -\frac{i}{2} \text{tr}(B[A^\dagger, H] - [H, A]B^\dagger) = -\text{Im } \text{tr}((B[A, H])).$$

On the other hand

$$\omega(B, \vec{H}(A)) = -\text{Im } \text{tr}(B[A, H]) = \langle d_A \mu(B), H \rangle.$$

This finishes the proof of the lemma. □

In the general case, the symplectic structure locally has the form

$$\omega(\varphi_1, \varphi_2) = 2i \int_X \operatorname{Im} \operatorname{tr}(\varphi_1 \wedge \varphi_2) = 2i \int_X \operatorname{Im} [\operatorname{tr}(\phi_1 \phi_2^\dagger)] \otimes \alpha_1 \wedge \bar{\alpha}_2.$$

Therefore, the momentum map must have the form

$$\mu(\varphi) = -[\varphi, \varphi^\dagger].$$

This finishes the proof of the proposition. \square

Now that we have a momentum map, and assuming that the \mathcal{G} -action verifies all the technical requirements, we can take the symplectic quotient $\mu^{-1}(0)/\mathcal{G}$. Let us consider now the subvariety

$$\{(\nabla, \varphi) \in \mathcal{A}_h \times \Omega^{1,0}(X, \operatorname{End} E) : \nabla^{0,1} \varphi = 0\}.$$

Avoiding possible singularities, this variety inherits the symplectic structure from $\mathcal{A}_h \times \Omega^{1,0}(X, \operatorname{End} E)$ and, since it is \mathcal{G} -invariant, the restriction of the momentum map. This allows us to take another symplectic quotient. The space we obtain is the **moduli space of solutions to the Hitchin equations** on E , that is, the set of equivalence classes of pairs $(\nabla, \varphi) \in \mathcal{A}_h \times \Omega^{1,0}(X, \operatorname{End} E)$ that satisfy **Hitchin's equations**:

$$\begin{cases} F + [\varphi, \varphi^\dagger] = -2\pi i \mu(E) \mathbf{1}_{E\omega_X}, \\ \nabla^{0,1} \varphi = 0. \end{cases}$$

In the next section we are going to see how this space corresponds to a broader notion of stability for pairs (\mathbf{E}, φ) , with \mathbf{E} a holomorphic vector bundle and $\varphi : \mathbf{E} \rightarrow \mathbf{E} \otimes K_X$ a K_X -twisted endomorphism and we will find that this moduli space is precisely the wider symplectic manifold containing $T^*\mathbf{M}_{\text{Bun}}(n, d)$ and the pairs associated to the remaining points of the Jacobian in the spectral correspondence.

§ 4. Higgs bundles

Definition 4.1. A pair $(\nabla, \varphi) \in \mathcal{A}_h \times \Omega^{1,0}(X, \operatorname{End} E)$ is called **irreducible** if there exist no subbundles $E' \subset E$ which are preserved by both ∇ and φ .

Theorem 4.2 (Hitchin correspondence). *Let (E, h) be a Hermitian complex vector bundle of rank n and degree d on a compact Riemann surface X . An irreducible pair $(\nabla, \varphi) \in \mathcal{A}_h \times \Omega^{1,0}(X, \operatorname{End} E)$ is a solution to the Hitchin equations if and only if for every proper subbundle $E' \subset E$ which is invariant by $\nabla^{0,1}$ and φ ,*

$$\mu(E') < \mu(E).$$

The proof of this theorem uses the same techniques and is very similar to Donaldson's proof of the theorem of Narasimhan-Seshadri (Theorem 5.4). As in that case, we will only prove the “easy” direction of the equivalence.

Proof. Let (∇, φ) be a solution of the Hitchin equations

$$\begin{cases} F + [\varphi, \varphi^\dagger] = -2\pi i \mu(E) \mathbf{1}_{E\omega_X}, \\ \nabla^{0,1} \varphi = 0. \end{cases}$$

and define $\bar{\partial}_{\mathbf{E}} = \nabla^{0,1}$. Let $E' \subset E$ be a subbundle preserved by $\bar{\partial}_{\mathbf{E}}$ such that $\varphi(E') \subset E' \otimes K_X$. The Hermitian metric gives a smooth splitting $E = E' \oplus E''$, and we can write

$$\varphi = \begin{pmatrix} \varphi_{E'} & \theta \\ 0 & \varphi_{E''} \end{pmatrix},$$

for $\theta \in \Omega^{1,0}(X, \text{Hom}(E'', E'))$. The top left element of $[\varphi, \varphi^\dagger]$ is

$$\varphi_{E'} \wedge \varphi_{E'}^\dagger + \theta \wedge \theta^\dagger - \varphi_{E'}^\dagger \wedge \varphi_{E'} = [\varphi_{E'}, \varphi_{E'}^\dagger] + \theta \wedge \theta^\dagger.$$

Using the equation

$$F = \begin{pmatrix} F_{E'} - \beta \wedge \beta^\dagger & \nabla_{\text{Hom}(E'', E')} \beta \\ -\nabla_{\text{Hom}(E', E'')} \beta^\dagger & F_{E''} - \beta^\dagger \wedge \beta \end{pmatrix},$$

from the proof of Narasimhan-Seshadri (Theorem 5.4) we have

$$F_{E'} - \beta \wedge \beta^\dagger + [\varphi_{E'}, \varphi_{E'}^\dagger] + \theta \wedge \theta^\dagger = -2\pi i \frac{d}{n} \mathbf{1}_{E'} \omega_X.$$

Now, since $\text{tr}[\varphi, \varphi^\dagger] = 0$, taking the trace and integrating we get

$$\frac{i}{2\pi} \int_X \text{tr} F_{E'} + \|\beta\|^2 + \|\theta\|^2 = \mu(E) \text{rk} E',$$

and hence

$$\mu(E') = \deg E' / \text{rk} E' = \mu(E) - \|\beta\|^2 - \|\theta\|^2.$$

Since the pair (∇, φ) is irreducible, $\|\beta\|$ and $\|\theta\|$ are nonzero, so $\mu(E') < \mu(E)$. \square

Definition 4.3. Let X be a compact Riemann surface. A **Higgs bundle** (or **Hitchin pair**) over X is a pair (\mathbf{E}, φ) , where \mathbf{E} is a holomorphic vector bundle over X and $\varphi : \mathbf{E} \rightarrow \mathbf{E} \otimes K_X$ is a K_X -twisted endomorphism. A Higgs bundle (\mathbf{E}, φ) is said to be **stable** if for every proper holomorphic subbundle $\mathbf{E}' \subset \mathbf{E}$ such that $\varphi(\mathbf{E}') \subset \mathbf{E}' \otimes K_X$

$$\mu(\mathbf{E}') < \mu(\mathbf{E}).$$

The **moduli space of stable Higgs bundles of rank n and degree d** , $\mathbf{M}_{\text{Higgs}}(n, d)$, is defined as the set of isomorphism classes of stable Higgs bundles (\mathbf{E}, φ) with $\text{rk } \mathbf{E} = n$ and $\deg \mathbf{E} = d$.

In [Hit87a], Hitchin gave an analytic construction of a particular case of rank 2 of this moduli space and a general GIT construction was given by Nitsure in [REF:Nitsure].

Theorem 4.4. The moduli space $\mathbf{M}_{\text{Higgs}}(n, d)$ is a complex manifold of dimension $2 + 2n^2(g - 1)$, where g is the genus of X .

The Hitchin correspondence 4.2 can be now reinterpreted as follows

Corollary 4.5. Let (E, h) be a smooth Hermitian vector bundle of rank n and degree d . The moduli space of solutions to the Hitchin equations on E is homeomorphic to the moduli space of stable Higgs bundles $\mathbf{M}_{\text{Higgs}}(n, d)$.

We can now extend the Hitchin map to the bigger space $\mathbf{M}_{\text{Higgs}}(n, d) \supset T^* \mathbf{M}_{\text{Bun}}(n, d)$,

$$H : \mathbf{M}_{\text{Higgs}}(n, d) \longrightarrow \mathcal{B} = \bigoplus_{i=1}^n H^0(X, K_X^i)$$

$$(\mathbf{E}, \varphi) \longmapsto (\sigma_1(\varphi), \dots, \sigma_n(\varphi)).$$

Take now a point $b \in \mathcal{B}$ such that the spectral curve S_b is irreducible and smooth. For every holomorphic line bundle M over S_b the spectral correspondence gives us a Higgs bundle $(\mathbf{E} = \pi_* M, \varphi)$.

Proposition 4.6. *The pair (\mathbf{E}, φ) induced by a line bundle M on an irreducible spectral curve S_b is stable.*

Proof. Suppose that $\mathbf{E}' \subset \mathbf{E}$ is a proper holomorphic subbundle invariant by φ . Then we can take the characteristic polynomial of $\varphi|_{\mathbf{E}'}$ and it must divide the characteristic polynomial of φ , thus contradicting that the spectral curve is irreducible. \square

As a conclusion, we get that generic fibers of the Hitchin map $H : \mathbf{M}_{\text{Higgs}}(n, d) \rightarrow \mathcal{B}$ are whole Jacobians of the spectral curve

§ 5. Deformations of Higgs bundles

In this section we are going to compute the dimension of $\mathbf{M}_{\text{Higgs}}(n, d)$, stated in Theorem 4.4, using deformation theory. As in the case for vector bundles, the idea is to identify the tangent space of $\mathbf{M}_{\text{Higgs}}(n, d)$ at some point $(\bar{\partial}_{\mathbf{E}}, \varphi)$ and find its dimension.

First of all, note that an analytic construction of $\mathbf{M}_{\text{Higgs}}(n, d)$ can be given if we consider the variety

$$\mathcal{H} = \left\{ (\bar{\partial}_{\mathbf{E}}, \varphi) \in \mathcal{A}_{\partial}^s \times \Omega^{1,0}(X, \text{End } E) : \bar{\partial}_{\mathbf{E}}\varphi = 0 \right\}$$

and its subset

$$\mathcal{H}^s = \left\{ (\bar{\partial}_{\mathbf{E}}, \varphi) \in \mathcal{H} : (\bar{\partial}_{\mathbf{E}}, \varphi) \text{ is stable} \right\}.$$

As usual, an element g of the complex group $\mathcal{G}^c = \Gamma(\text{Aut } E)$ acts on \mathcal{H} by conjugation,

$$(\bar{\partial}_{\mathbf{E}}, \varphi) \mapsto (g\bar{\partial}_{\mathbf{E}}g^{-1}, g\varphi g^{-1}).$$

We can then obtain the moduli space of Higgs bundles as the quotient $\mathbf{M}_{\text{Higgs}}(n, d) = \mathcal{H}^s / \mathcal{G}^c$, where $n = \text{rk } E$ and $d = \deg E$.

The differential of action of the complex group on \mathcal{H} is mapping from the Lie algebra $\text{Lie } \mathcal{G}^c = \Gamma(\text{End } E)$ to the tangent space of \mathcal{H} ,

$$\begin{aligned} f : \Gamma(\text{End } E) &\longrightarrow T_{(E, \varphi)}\mathcal{H} \\ \xi &\longmapsto (\bar{\partial}_{\mathbf{E}}\xi, [\varphi, \xi]). \end{aligned}$$

Since stability is an open condition, we can identify the tangent space to the moduli space with the quotient of the tangent space of \mathcal{H} by the image of this mapping,

$$T_{(\bar{\partial}_{\mathbf{E}}, \varphi)}\mathbf{M}_{\text{Higgs}}(n, d) = T_{(\bar{\partial}_{\mathbf{E}}, \varphi)}\mathcal{H} / \text{im } f.$$

On the other hand, to obtain how an element of the tangent space $T_{(\bar{\partial}_{\mathbf{E}}, \varphi)}\mathcal{H}$ looks like, consider a small perturbation

$$(\bar{\partial}_{\mathbf{E}}(\epsilon), \varphi(\epsilon)) = (\bar{\partial}_{\mathbf{E}} + \epsilon A, \varphi + \epsilon \psi),$$

for $(A, \psi) \in \Omega^{0,1}(X, \text{End } E) \oplus \Omega^{1,0}(X, \text{End } E)$ and impose that the condition $\bar{\partial}_{\mathbf{E}}(\epsilon)\varphi(\epsilon) = 0$ is infinitesimally satisfied. That is,

$$0 = \left. \frac{d}{d\epsilon} \right|_{\epsilon=0} \bar{\partial}_{\mathbf{E}}(\epsilon)\varphi(\epsilon) = \left. \frac{d}{d\epsilon} \right|_{\epsilon=0} \left[\bar{\partial}_{\mathbf{E}}\varphi + \epsilon \left([A, \varphi] + \bar{\partial}_{\mathbf{E}}\psi \right) + o(\epsilon^2) \right] = [A, \varphi] + \bar{\partial}_{\mathbf{E}}\psi.$$

Therefore, we can think of $T_{(\bar{\partial}_{\mathbf{E}}, \varphi)}\mathcal{H}$ as the kernel of the map

$$\begin{aligned} g : \Omega^{0,1}(X, \text{End } E) \oplus \Omega^{1,0}(X, \text{End } E) &\longrightarrow \Omega^{1,1}(X, \text{End } E) \\ (A, \psi) &\longmapsto [A, \varphi] + \bar{\partial}_{\mathbf{E}}\psi. \end{aligned}$$

Note then that f is well defined,

$$g \circ f(\xi) = [\bar{\partial}_E \xi, \varphi] + \bar{\partial}_E[\varphi, \xi] = [\bar{\partial}_E \xi, \varphi] + [\bar{\partial}_E \varphi, \xi] - [\bar{\partial}_E \xi, \varphi] = 0,$$

since $\bar{\partial}_E \varphi = 0$. In other words, what we are saying is that we have a chain complex

$$C^\bullet : \Gamma(\text{End } E) \xrightarrow{f} \Omega^{0,1}(X, \text{End } E) \oplus \Omega^{1,0}(X, \text{End } E) \xrightarrow{g} \Omega^{1,1}(X, \text{End } E)$$

and the tangent space $T_{(\bar{\partial}_E, \varphi)} \mathbf{M}_{\text{Higgs}}(n, d)$ is the first cohomology of this complex, $H^1(C^\bullet)$.

Consider now the double complex

$$\begin{array}{ccccccc} 0 & \xrightarrow{\quad} & \Gamma(\text{End } E) & \xrightarrow{\quad} & \Gamma(\text{End } E) \\ \downarrow & & \downarrow f & & \downarrow \bar{\partial}_E \\ \Omega^0(X, \text{End } E \otimes K_X) & \xhookrightarrow{\quad} & \Omega^{0,1}(X, \text{End } E) \oplus \Omega^{1,0}(X, \text{End } E) & \xrightarrow{\text{pr}_1} & \Omega^{0,1}(X, \text{End } E) \\ \downarrow \bar{\partial}_E & & \downarrow g & & \downarrow \bar{\partial}_E \\ \Omega^{0,1}(X, \text{End } E \otimes K_X) & \xrightarrow{\quad} & \Omega^{1,1}(X, \text{End } E) & \xrightarrow{\quad} & 0 \\ \downarrow \bar{\partial}_E & & & & \\ 0 & & & & \end{array}$$

The rows of this complex are short exact sequences, so then we get a short exact sequence of complexes

$$0 \longrightarrow \Omega^{0,\bullet}(X, \text{End } E \otimes K_X)[1] \longrightarrow C^\bullet \longrightarrow \Omega^{0,\bullet}(X, \text{End } E) \longrightarrow 0,$$

where the $[1]$ after a complex denotes shifting one place and adding zero in the first place. This induces a long exact sequence in cohomology

$$\begin{aligned} 0 \rightarrow H^0(C^\bullet) \rightarrow H^0(X, \text{End } E) \rightarrow H^0(X, \text{End } E \otimes K_X) \rightarrow \\ \rightarrow H^1(C^\bullet) \rightarrow H^1(X, \text{End } E) \rightarrow H^1(X, \text{End } E \otimes K_X) \rightarrow H^2(C^\bullet) \rightarrow 0. \end{aligned}$$

The Euler characteristic of any exact sequence vanishes, so

$$H^0(C^\bullet) - H^1(C^\bullet) + H^2(C^\bullet) - \chi(\text{End } E) + \chi(\text{End } E \otimes K_X) = 0.$$

Now, note that $\chi(\text{End } E) = -\chi(\text{End } E \otimes K_X)$ and recall that we computed that $\chi(\text{End } E) = n^2(1 - g)$. Therefore

$$\dim H^1(C^\bullet) = \dim H^0(C^\bullet) + \dim H^2(C^\bullet) + 2n^2(g - 1).$$

Noting that, by Serre duality

$$\begin{aligned} H^1(X, \text{End } E)^* &\cong H^0(X, (\text{End } E)^* \otimes K_X) \cong H^0(X, \text{End } E \otimes K_X) \\ H^1(X, \text{End } E \otimes K_X)^* &\cong H^0(X, (\text{End } E \otimes K_X)^* \otimes K_X) \cong H^0(X, \text{End } E), \end{aligned}$$

we can dualize in the above exact sequence and we get that $H^2(C^\bullet)^* \cong H^0(C^\bullet)$, so

$$\dim H^2(C^\bullet) = \dim H^0(C^\bullet).$$

To find this dimension we use the following result:

Proposition 5.1. *If (\mathbf{E}, φ) is a stable Higgs bundle, then*

$$H^0(C^\bullet) = \mathbf{C}.$$

Proof. Note that

$$H^0(C^\bullet) = \ker f = \left\{ \xi \in H^0(X, \text{End } \mathbf{E}) : [\varphi, \xi] = 0 \right\}.$$

What we are going to prove is that every endomorphism ξ that satisfies that condition is a multiple of the identity $\mathbf{1}_{\mathbf{E}}$. Indeed, let $\lambda \in \mathbf{C}$ be an eigenvalue of $\xi_x : \mathbf{E}_x \rightarrow \mathbf{E}_x$ and define the endomorphism $\eta = \xi - \lambda \mathbf{1}_{\mathbf{E}}$. Since λ is an eigenvalue of ξ_x , we have $\det \eta_x = 0$. Assuming that η is nonzero, we will prove that η is injective and hence $\det \eta_x \neq 0$ arriving at a contradiction. To see this, suppose that $\text{im } \eta \subset \mathbf{E}$ is a holomorphic strict subbundle of \mathbf{E} . Note too that $\text{im } \eta$ is φ -invariant, since, if $w = \eta(v)$, then

$$\varphi(w) = \varphi((\xi - \lambda \mathbf{1}_{\mathbf{E}})(v)) = \varphi(\xi(v)) - \varphi(\lambda v) = \xi(\varphi(v)) - \lambda \varphi(v) = \eta(\varphi(v)),$$

since $[\varphi, \xi] = 0$. Therefore, since (\mathbf{E}, φ) is stable, $\mu(\text{im } \eta) < \mu(\mathbf{E})$. But we also have that $\mu(\mathbf{E}) < \mu(\text{im } \eta)$, so $\mu(\mathbf{E}) < \mu(\mathbf{E})$ and we have a contradiction. Therefore η is injective. We conclude then that $\eta = 0$, so $\xi = \lambda \mathbf{1}_{\mathbf{E}}$. \square

In conclusion, we get that $\dim H^2(C^\bullet) = \dim H^0(C^\bullet) = 1$. We then arrive at the result we were looking for:

$$\dim \mathbf{M}_{\text{Higgs}}(n, d) = \dim H^1(C^\bullet) = 2 + 2n^2(g - 1).$$

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