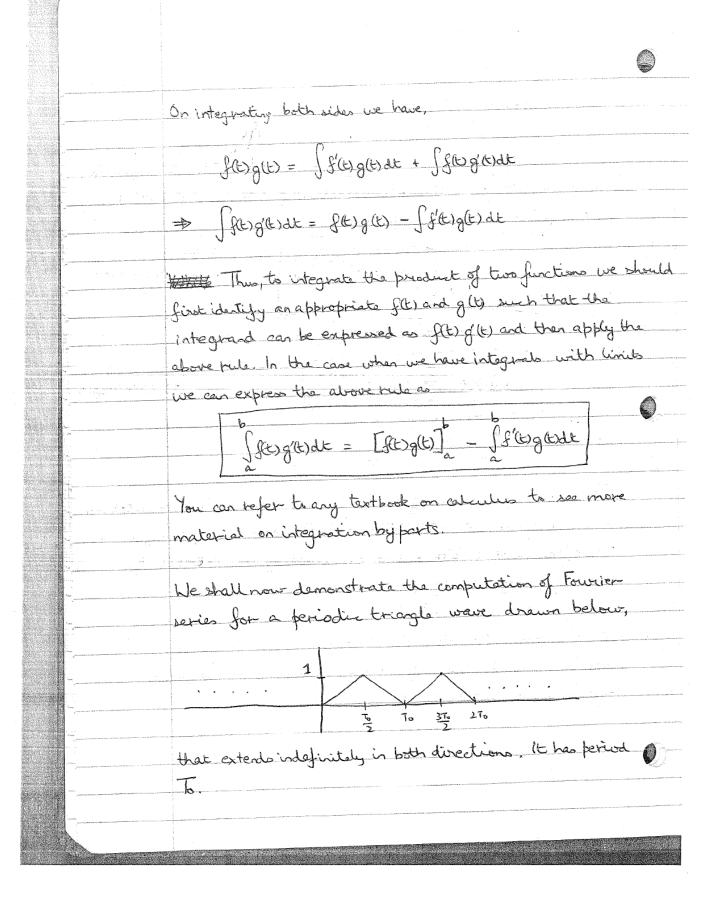
Digital Filters & Spectral Analysis Lecture 2

The Fourier Series

Derivation of the Fourier Series coefficients of a triangle wave

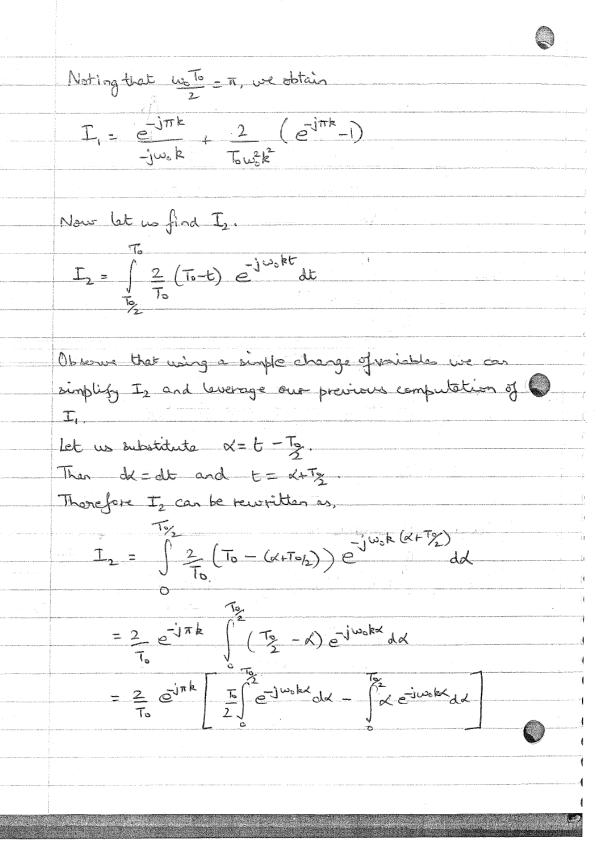
(From: J From Aditya Ramamoorthy, Review Notes, "Signals and Systems I" course, Iowa State University)

	Fourier Series
	For a given periodic signal x(t) with period To it is possible
Ŧ.	to write x(t) as a linear combination of complex exponentials. The coefficients can be computed as.
	ar = 1 (xb) e-iwok t - (1) (ANALY)
	and x(b) can be reconstructed as
	$\chi(t) = \frac{2}{2} a_k e^{j w_0 kt} \qquad \qquad (2) (SYNTHE)$
	where wo = 1 x 2 T.
	Thus, the computation of Fourier series requires the computation
	of an integral where the integrand is the product of two functions. We often need to either find or simplify this integral by
	means of integration by parts.
3	Let us review the integration by parts rule.
	Consider two functions $f(t)$ and $g(t)$, then by the differentiation rule for a product of functions, we have.
	$\frac{d}{dt}\left(f(t)g(t)\right) = f'(t)g(t) + f(t)g'(t)$
	A STATE OF THE STA



let us define the signal is nathandical terms first. Let us call the signal between [0, To], S(E). Than, $f(t) = \begin{cases} \frac{2}{T_0}t & t \in [0, \frac{T_0}{2}) \\ \frac{2}{T_0}(T_0 - t) & t \in [\frac{T_0}{2}, T_0) \end{cases}$ ak= I f(t) = jwokt dt > To ak = ∫f(t) = i wokt at = \frac{2}{76} t e^{-j w_0 kt} dt + \int_0^2 (T_0 - t) e^{-j w_0 kt} dt Let us first compute I $I_{i} = \frac{2}{T_{0}} \int_{0}^{\infty} t e^{j\omega_{0}kt} dt = \frac{2}{T_{0}} \left[\frac{1}{t} e^{j\omega_{0}kt} \int_{0}^{\infty} - \int_{0}^{\infty} e^{j\omega_{0}kt} dt \right]$ = 2 To ejwoktog | [ejwoktog]

To 2 -jwok jwok -jwok $= \frac{2}{10} \left[\frac{7}{2} - \frac{1}{1000} \frac{1}{1000} + \frac{1}{1000} \left(\frac{1}{2} \frac{1}{1000} \frac{1}{1000} \frac{1}{1000} \right) \right]$



= 2 eink To leiwokx To de iwokx dx = ejnk (e-juok = -1) ejnk I since 2 | x = j work dd = I. $I_2 = \frac{e^{j\pi k}}{e^{j\pi k}} \left(e^{j\pi k} - 1 \right) - e^{j\pi k} I,$ To ak = I, +I, = I,+ eink (eink-D eink I, $= (1 - e^{j\pi k}) I_1 + e^{j\pi k} \frac{(e^{j\pi k} - 1)}{-j\omega_0 k}$ Substituting for I, we have

To $a_k = \begin{bmatrix} e^{i\pi k} \\ -jw_0k \end{bmatrix} \begin{bmatrix} e^{j\pi k} \\ -jw_0k \end{bmatrix} \begin{bmatrix} -jw_0k \end{bmatrix}$ $= \frac{-i\pi k}{-j\omega_0 k} \left(\frac{1}{1-e^{j\pi k}} \right) + \frac{2}{70\omega_0^2 k^2} \left(\frac{e^{j\pi k}}{1-1} \right) \times (-1) + \frac{e^{-j\pi k}}{1-1} \left(\frac{e^{j\pi k}}{1-1} \right)$ $= -2 \left(e^{i^2\pi k} + 1 - 2e^{i\pi k} \right)$

Noting that e 12nk = 1 and e 1nk = (-Dk, we have

9k= -4 (1-61)x)

Substituting for us, we obtain

 $a_k = 4(-1)^k - 1$ = $(-1)^k - 1$ $T_0^2 \times \frac{4\pi^2}{T_0^2} k^2$ = $T_0^2 k^2$

It remains to determine the value of ao. Note that a substitution of k=0 is the above expression is not correct as it yields a modhematically undefined of form. In general as always needs to be computed separately. We have,

Que = 1 feb e jus x 0 xt

= 1 feb dt

= 1 feb dt

= 1 x 1. To x 1 (using the formule for the area of a triangle)

= 1/2

We can now compactly express the result as:

$$a_{k} = \begin{cases} 0 & \text{if } k = \pm 2, \pm 4... \\ \frac{-2}{\pi^{2} R} & \text{if } k = \pm 1, \pm 3... \end{cases}$$

Note that the rate of decay of as with k is ~ /2. Thus, intuitively coefficients with high values of k are not likely to contribute significantly to the "synthesis" sum. i.e. we expect a few terms in the sum to result in a good reconstruction of the original signal.

The behavior of ak for the periodic triagle wave should be compared with the behavior of ak for the periodic square wave where the rote of decay to ~ 1/k. This should provide you are insight that in the reconstruction of the periodic square wave you would probably need to sun more tomo to have the same quality of reconstruction.