

1 Basic of Fourier Transform

Fourier Series

If $x(t) = x(t + T)$ then $x(t)$ can be written as

$$x(t) = \sum_{-\infty}^{+\infty} c_k e^{\frac{2\pi i k t}{T}}$$

i is the imaginary unit, and k is an integer. The above expression is eligible because $e^{\frac{2\pi i k t}{T}}$ is a periodic function

$$e^{\frac{2\pi i k t}{T}} = e^{\frac{2\pi i k (t+T)}{T}}$$

Each basis $e^{\frac{2\pi i k t}{T}}$ represents a signal with frequency $f_k = \frac{k}{T}$. So the interval between each adjacent frequency $\Delta f = \frac{1}{T}$. Based on orthogonality, we can get c_k

$$c_k = \frac{1}{T} \int_0^T x(t) e^{-i \frac{2\pi k t}{T}} dt$$

Fourier Series: Example

$$x(t) = \cos(2\pi f_0 t) = \frac{1}{2}(e^{i2\pi f_0 t} + e^{-i2\pi f_0 t})$$

where $f_0 = \frac{1}{T}$

$$\begin{aligned} c_k &= \frac{1}{T} \int_0^T x(t) e^{-i \frac{2\pi k t}{T}} dt \\ &= \frac{1}{T} \int_0^T \frac{1}{2} (e^{\frac{2i\pi t}{T}} + e^{\frac{-2i\pi t}{T}}) e^{-i \frac{2\pi k t}{T}} dt \end{aligned}$$

Only terms with $k = \pm 1$ in the above expression can survive, so

$$c_1 = \frac{1}{T} \int_0^T \frac{1}{2} dt = \frac{1}{2}$$

Similarly, $c_{-1} = \frac{1}{2}$.

Fourier Transform

We can generalize the Fourier series to non-periodic functions. We define the Fourier transform as

$$\mathcal{F}(f) = \int_{-\infty}^{\infty} x(t) e^{-2\pi i f t} dt$$

With the inverse Fourier transform defined as

$$x(t) = \int_{-\infty}^{\infty} \mathcal{F}(f) e^{2\pi i f t} df$$

Fourier Transform: Example

Take the same $x(t)$ as above

$$x(t) = \cos(2\pi f_0 t) = \frac{1}{2}(e^{2i\pi f_0 t} + e^{-2i\pi f_0 t})$$

$$\mathcal{F}(f) = \int_{-\infty}^{\infty} \frac{1}{2}(e^{2i\pi f_0 t} + e^{-2i\pi f_0 t}) e^{-2i\pi f t} dt = \frac{1}{2}\delta(f - f_0) + \frac{1}{2}\delta(f + f_0)$$

Discrete Fourier Transform

The above is the Fourier transform in continuous case, in discrete case If $x = n\Delta t$, where $n = 1 \dots N$, and $T = N\Delta t$, then the Fourier series can be written as

$$\begin{aligned} f(n) &= \sum_{-\infty}^{+\infty} c_k e^{\frac{2\pi i k n \Delta t}{N \Delta t}} \\ &= \sum_{-\infty}^{+\infty} c_k e^{\frac{2\pi i k n}{N}} \end{aligned}$$

$$c_k = \frac{1}{N\Delta t} \sum_{n=1}^N f(n\Delta t) e^{-i2\pi k \frac{1}{N\Delta t} n\Delta t} d(n\Delta t) = \frac{1}{N} \sum_{n=1}^N f(n) e^{-i2\pi k \frac{n}{N}}$$

This is the discrete Fourier transform.

$$\Delta F = \frac{1}{T} = \frac{1}{N\Delta t}$$

Example

Let $N = 4$, and

$$f(n) = \cos(2\pi \frac{n}{4}) = \frac{1}{2}(e^{i2\pi \frac{n}{4}} + e^{-i2\pi \frac{n}{4}})$$

$$c_k = \frac{1}{4} \sum_{n=1}^4 \frac{1}{2}(e^{i2\pi \frac{n}{4}} + e^{-i2\pi \frac{n}{4}}) e^{-i \frac{2\pi k n}{4}}$$

Similar to the continuous case, only terms with $k = +/ -1$ in the above expression can survive, when $k = 1$

$$\begin{aligned} c_1 &= \frac{1}{4} \sum_{n=1}^4 \frac{1}{2} e^{i2\pi \frac{n}{4}} e^{-i \frac{2\pi n}{4}} \\ &= \frac{1}{4} \frac{1}{2} 4 \\ &= \frac{1}{2} \end{aligned}$$

What about case for $k = -1$? We define $k = 1, 2, 3, 4$ so $k = -1$ is not defined. However, in discrete case we note $c_{-1} = c_3$ due to the periodicity. Similarly, we can calculate $c_3 = \frac{1}{2}$.

N is the total sample within time T .

Properties

1) To be eligible, $f(x)$ has to be a period function with time T (with frequency $F = \frac{1}{T}$) in both continuous case and discrete case. The requirement in discrete case leads to uniform sampling theorem used in signal processing. The total sampling time $T_{sampling}$ has to be an integer multiple of T .

$$T_{sampling} = MT$$

while $T = \frac{N}{F_s}$ So

$$MT = N\Delta t$$

if we let $\Delta t = \frac{1}{F_s}$, where F_s is the sampling frequency, and $T = \frac{1}{F}$, we have

$$\frac{M}{F} = \frac{N}{F_s}$$

2) If $f(x)$ is real, which means $f(x) = f^*(x)$. We then substitute Fourier series for both $f(x)$ and $f^*(x)$,

$$\sum_{-\infty}^{+\infty} c_k e^{2\pi i \frac{1}{T} kx} = \sum_{-\infty}^{+\infty} c_k^* e^{-2\pi i \frac{1}{T} kx} \quad (1)$$

Since the summation on the right hand side is from $-\infty$ to ∞ , it is eligible to replace k with $-k$.

$$\sum_{-\infty}^{+\infty} c_k^* e^{-2\pi i \frac{1}{T} kx} = \sum_{\infty}^{-\infty} c_{-k}^* e^{2\pi i \frac{1}{T} kx} \quad (2)$$

Combine the above two equations 1 and 2, we can see $c_k = c_{-k}^*$. This means they are complex conjugate: their magnitude are equal, their phase are opposite. Namely $|c_k| = |c_{-k}|$, $\phi(c_k) = \phi(c_{-k})$.

3) Connection between complex representation and real representation.

We have shown that for real signal $c_k = c_{-k}^*$ and $c_k = |c_k|e^{j\theta_k}$, $c_{-k} = |c_k|e^{-j\theta_k}$. And in complex representation, we can combine the term with index k and $-k$,

$$c_k e^{j2\pi k F_0 t} + c_{-k} e^{-j2\pi k F_0 t} = 2|c_k| \cos(2\pi k F_0 t + \theta_k)$$

$$\begin{aligned} f(x) &= \sum_{-\infty}^{+\infty} c_k e^{\frac{2\pi i k x}{T}} \\ &= c_0 + 2 \sum_{k=1}^{\infty} |c_k| \cos(2\pi k F_0 t + \theta_k) \\ &= a_0 + \sum_{k=1}^{\infty} (a_k \cos(2\pi k F_0 t) - b_k \sin(2\pi k F_0 t)) \end{aligned}$$

where $a_0 = c_0$, $a_k = 2|c_k|\cos\theta_k$, $b_k = 2|c_k|\sin\theta_k$. 4) $c_k = c_{k+N}$. So when a signal contains frequency component no larger than B , in other words, the bandwidth of the signal is $2B(-B \text{ to } B)$, then in order to capture the whole bandwidth of the signal, $N\Delta f > 2B$. This leads to Nyquist sampling theorem $F_s > 2B(\text{bandwidth})$.

5) Power density

$$\begin{aligned} P_x &= \frac{1}{T} \int |x(t)|^2 dt \\ &= \frac{1}{T} \int x(t) \sum_{-\infty}^{\infty} c_k^* e^{-j2\pi k F_0 t} \\ &= \sum_{-\infty}^{\infty} c_k^* \left[\frac{1}{T} \int x(t) e^{-j2\pi k F_0 t} \right] \\ &= \sum_{-\infty}^{\infty} |c_k|^2 \end{aligned}$$

When signal is real, then

$$\begin{aligned} P_x &= \sum_{-\infty}^{\infty} |c_k|^2 \\ &= a_0^2 + \frac{1}{2} \sum_{k=1}^{\infty} (a_k^2 + b_k^2) \end{aligned}$$

2 Fast Fourier Transform

$$X_k = \sum_{n=0}^{N-1} x_n e^{-i2\pi k \frac{n}{N}}$$

let

$$u_k = e^{-i2\pi k \frac{n}{N}}$$

then we have the basis orthogonality

$$u_{k1}^T u_{k2} = N \delta_{k1, k2}$$

We recognize we can write X_k with even index terms and odd index terms

$X_k = \text{Even index parts} + \text{Odd index parts}$

$$= \sum_{m=0}^{N/2-1} x_{2m} e^{-\frac{2\pi i}{N} 2mk} + \sum_{m=0}^{N/2-1} x_{2m+1} e^{-\frac{2\pi i}{N} (2m+1)k}$$

$$= \sum_{m=0}^{N/2-1} x_{2m} e^{-\frac{2\pi i}{N/2} mk}$$

(We can view this as Fourier Transform of $N/2$ even indexed points, where k is $0, 1N/2$)

$$+ e^{-\frac{2\pi i}{N} k}$$

$$\sum_{m=0}^{N/2-1} x_{2m+1} e^{-\frac{2\pi i}{N/2} mk}$$

(We can view this as Fourier Transform of $N/2$ odd indexed points, where k is $0, 1N/2$)

(Since each part is a Fourier transform of $N/2$ points, k has to be smaller than $N/2$)

$$= E_k + e^{-\frac{2\pi i}{N} k} O_k$$

As noted, the above derivation is for $k < N/2$, a very similar derivation for $N/2 < k < N$ leads to

$$X_{k+N/2} = E_k - e^{-\frac{2\pi i}{N} k} O_k$$

Now we have divided the FFT of N points to two FFT with $N/2$ points. Keep going till we reach the size to one, then combine together recursively.

3 Fourier Transform of Useful Functions

The Fourier Transform of Gaussian

$$f(t) = \frac{1}{\sqrt{2\pi\sigma^2}} e^{-\frac{t^2}{2\sigma^2}}$$

$$\mathcal{F}(f) = e^{-2\pi^2\sigma^2 f^2}$$

So the Fourier transform of a Gaussian function is another Gaussian function but with different width.

The Fourier Transform of Dirac Comb

$$x(t) = \sum_{n=-\infty}^{\infty} \delta(t - nT)$$

It is clearly that $x(t)$ is periodic with period T . So we can expand that into Fourier series

$$x(t) = \sum_{k=-\infty}^{\infty} c_k e^{2\pi i k t / T}$$

Where

$$\begin{aligned} c_k &= \frac{1}{T} \int_{-T/2}^{T/2} x(t) e^{-i \frac{2\pi k t}{T}} dt \\ &= \frac{1}{T} \int_{-T/2}^{T/2} \delta(0) e^{-i \frac{2\pi k t}{T}} \\ &= \frac{1}{T} \end{aligned}$$

So

$$x(t) = \sum_{k=-\infty}^{\infty} \frac{1}{T} e^{2\pi i k t / T}$$

On the other hand, based on the formula of Fourier transform

$$\mathcal{F}(f) = \int \sum_{n=-\infty}^{\infty} \delta(t - nT) e^{-2\pi i f t} dt = \sum_{n=-\infty}^{\infty} e^{-2\pi i n T f} = \sum_{n=-\infty}^{\infty} e^{-2\pi i n f / f_0}$$

Comparing the Fourier series of $x(t)$ and the expression of $F(f)$, they are the same except T being changed to f_0 . Therefore we can conclude that $F(f)$ itself is also a Dirac comb, which is

$$\mathcal{F}(f) = f_0 \sum_{n=-\infty}^{\infty} \delta(f - n f_0)$$

The Fourier Transform of White Noise

Assuming noise we sample in time is $n[m]$, where $m = 0, \dots, M-1$. $n[m]$ is a Gaussian random variable with zero mean and variance σ^2 . The the FFT of $n[m]$ is

$$\begin{aligned} N[k] &= \frac{1}{M} \sum_{m=0}^{M-1} n[m] e^{-i 2\pi m k / M} \\ &= \frac{1}{M} \sum_{m=0}^{M-1} n[m] (\cos(2\pi m k / M) - i n[m] \sin(2\pi m k / M)) \end{aligned}$$

The expected value is

$$\begin{aligned} E[N[k]] &= E\left[\frac{1}{M} \sum_0^{M-1} n[m] e^{-i 2\pi m k / M}\right] \\ &= \frac{1}{M} \sum_0^{M-1} E[n[m]] e^{-i 2\pi m k / M} \\ &= 0 (\text{because } E[n[m]] = 0) \end{aligned}$$

The variance of the real part is

$$\begin{aligned}
Var[R[N[k]]] &= E[(\frac{1}{M} \sum_{m=0}^{M-1} n[m](\cos(2\pi mk/M)) * (\frac{1}{M} \sum_{p=0}^{M-1} n[p](\cos(2\pi pk/M)))] \\
&= \frac{1}{M^2} E[\sum_{m=0}^{M-1} n[m]n[p]\delta(n-p)\cos(2\pi mk/M) * \cos(2\pi pk/M)] \\
&= \frac{1}{M^2} \sum_{m=0}^{M-1} E[n[m]^2]\cos^2(2\pi mk/M) \\
&= \frac{1}{M^2} \sigma^2 (\sum_{m=0}^{M-1} \cos^2(2\pi mk/M)) \\
&= \frac{1}{M^2} \sigma^2 (\frac{M}{2} + \frac{\cos((M+1)2\pi k/M)\sin(2\pi Mk/M)}{2\sin(2\pi k/M)}) \\
&= \frac{1}{M} \frac{\sigma^2}{2}
\end{aligned}$$

The same derivation applies for the imaginary part. So the FFT is Gaussian noise with mean zero and variance σ^2 .

4 Connection with Uncertainty Principle

Relationship between time length and frequency bandwidth

We consider a few examples

1) We consider a function $g(t)$ which is infinitely long in time domain

$$g(t) = \cos(2\pi f_0 t)$$

Its Fourier transform is

$$\begin{aligned}
F(f) &= \int \frac{e^{i2\pi f_0 t} + e^{-i2\pi f_0 t}}{2} e^{i2\pi f t} dt \\
&= \int \frac{1}{2} e^{i2\pi t(f_0+f)} dt + \int \frac{1}{2} e^{i2\pi t(f-f_0)} dt \\
&= \frac{1}{2} \delta(f+f_0) + \frac{1}{2} \delta(f-f_0)
\end{aligned}$$

The last line is based on $\int_{-\infty}^{\infty} e^{i2\pi f t} dt = \delta(f)$.

Since the delta function has width zero, so the bandwidth in frequency domain is zero. We see a signal which is infinitely long in time domain has zero bandwidth in frequency domain.

2) We consider a function $g(t)$ which has zero width in time, namely an impulse function.

$$g(t) = \delta(t)$$

Since this function is not a periodic function, we assume its period is infinity. Its Fourier transform is

$$F(f) = \int_{-\infty}^{\infty} \delta(t) e^{-2\pi f t} = 1$$

Now we see a signal which has zero width in time has infinitely long frequency bandwidth. This leads to the uncertainty principle.

Uncertainty Principle In quantum mechanics, if there is a particle with position x and momentum p , then uncertainty principle states

$$\Delta x \Delta p \geq \frac{\hbar}{2}$$

Similar relationship holds for time t and Energy.

$$\Delta t \Delta E \geq \frac{\hbar}{2}$$

We can modify this expression to get the time and frequency relationship in our Fourier transform. Since $E = \hbar\omega$. Then

$$\Delta t \Delta \omega \geq \frac{1}{2}$$