

Quantum Tunneling Phenomenon

Quantum Mechanics - Brief Recap

The wavefunction $\psi(x, t)$ (considering 1D) is governed by Schrödinger's equation -

$$i\hbar \frac{\partial}{\partial t} \psi(x, t) = \frac{-\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi(x, t) + V(x, t) \psi$$

For time independent potential $V(x)$, we can use separation of variables -

$$\psi(x, t) = X(x)T(t)$$

Substituting back in Schrödinger's equation -

$$X i\hbar \frac{dT}{dt} = \frac{-\hbar^2}{2m} T \frac{d^2 X}{dx^2} + V(x) X T$$

Quantum Mechanics Brief Recap – Contd.

Dividing by $\psi(x, t)$ and setting both sides equal to constant E -

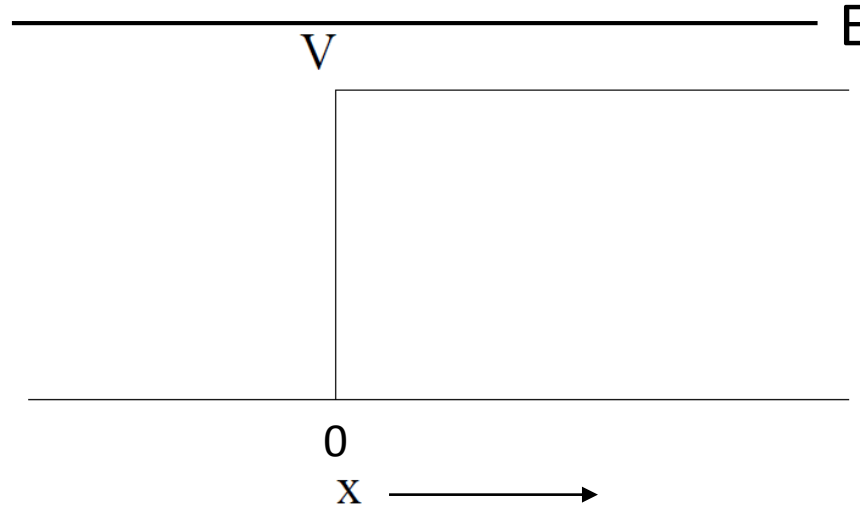
$$i\hbar \frac{1}{T} \frac{dT}{dt} = \frac{-\hbar^2}{2m} \frac{1}{X} \frac{d^2 X}{dx^2} + V(x) = E$$

$$T(t) = Ae^{-iEt/\hbar}$$
$$\psi(x, t) = Ae^{-iEt/\hbar} X(x)$$

To obtain $X(x)$, we need to solve -

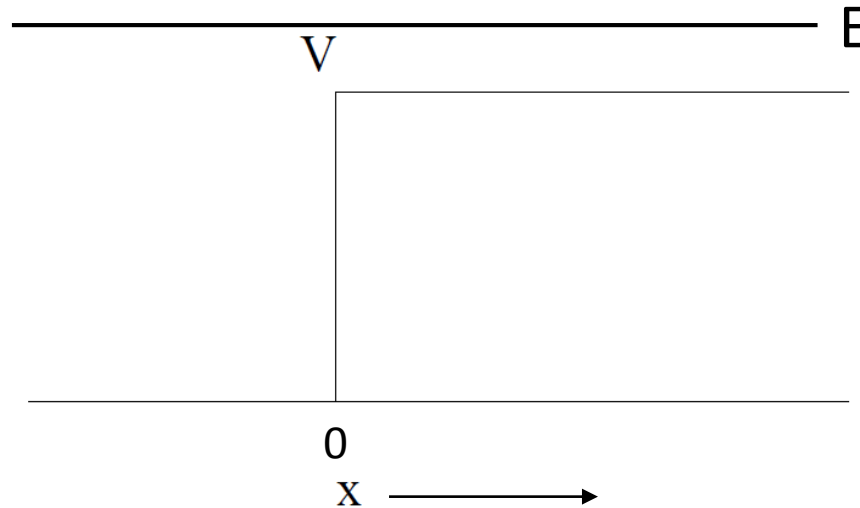
$$\frac{\hbar^2}{2m} \frac{d^2 X}{dx^2} = -[E - V(x)]X$$

Step Potential: Case I - (Energy of particle E) $> V$



We have to solve for -
$$\frac{d^2 X}{dx^2} = -\frac{2m}{\hbar^2}(E - V)X$$

Step Potential: Case I - (Energy of particle E) $> V$

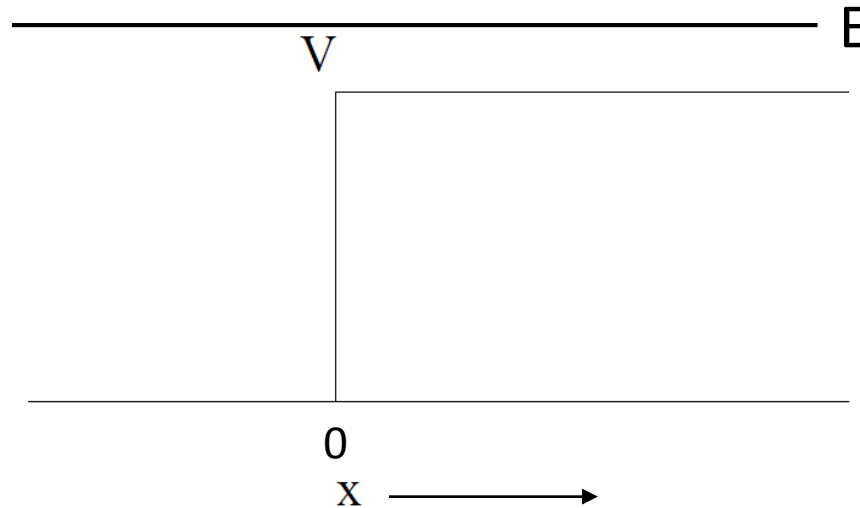


We have to solve for -
$$\frac{d^2 X}{dx^2} = -\frac{2m}{\hbar^2}(E - V)X$$

For right-side, the momentum is -
$$p' = \sqrt{2m(E - V)}$$

$$\frac{d^2 X}{dx^2} = \frac{-p'^2}{\hbar^2}X$$

Step Potential: Case I - (Energy of particle E) $> V$



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$$X(x) = A_1' e^{ip'x/\hbar} + A_2' e^{-ip'x/\hbar}$$

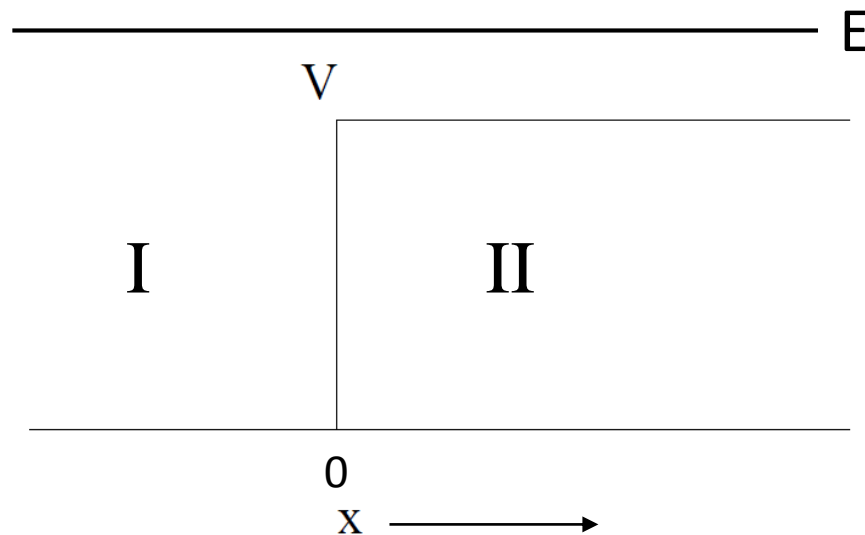
Step Potential: Case I - (Energy of particle E) $> V$

Solution : $\psi_{II}(x, t) = e^{-iEt/\hbar} \left[A'_1 e^{ip'x/\hbar} + A'_2 e^{-ip'x/\hbar} \right]$

$$\psi_I(x, t) = e^{-iEt/\hbar} \left[A_1 e^{ipx/\hbar} + A_2 e^{-ipx/\hbar} \right]$$

$$p = \sqrt{2mE}$$

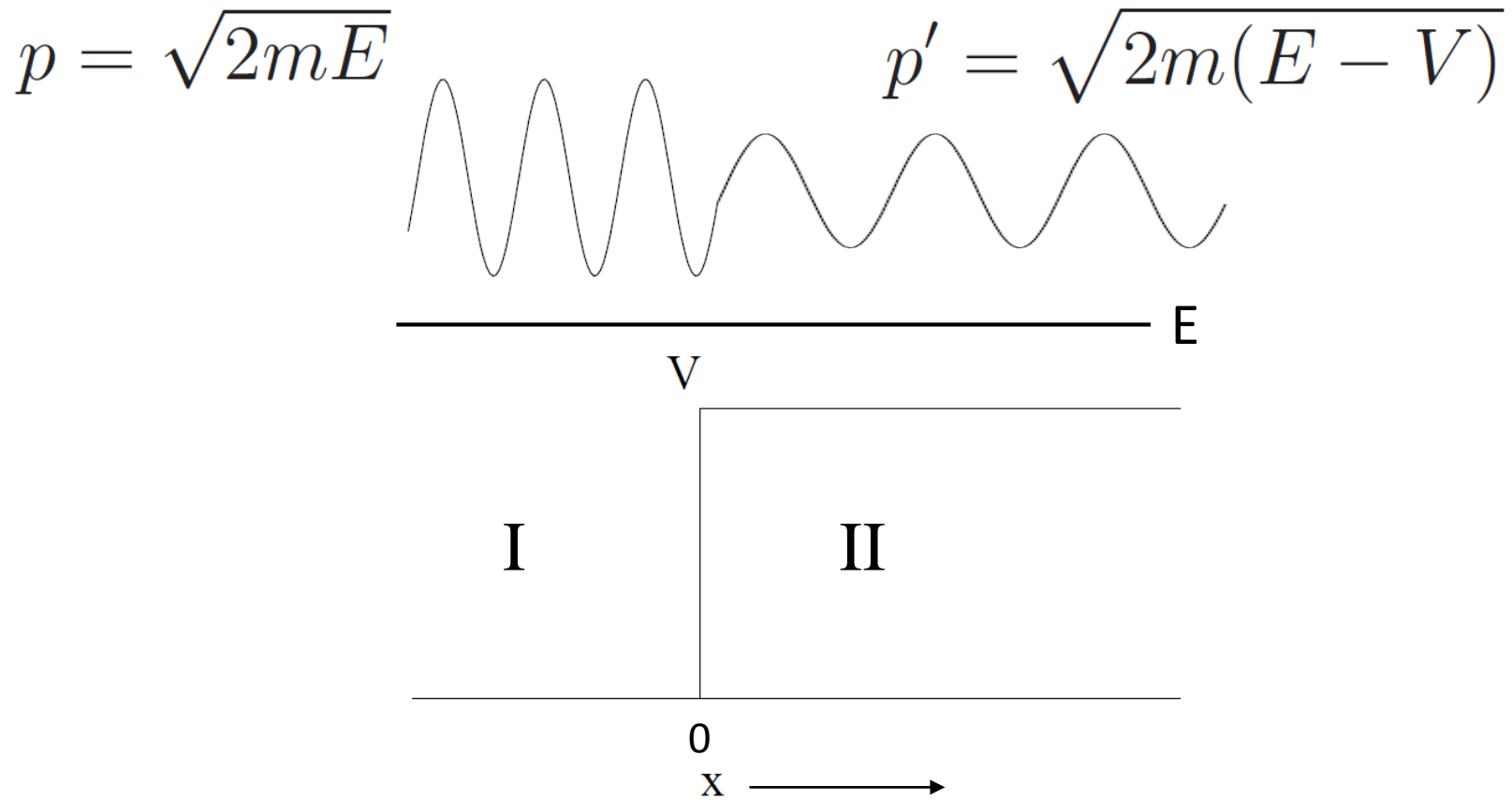
$$p' = \sqrt{2m(E - V)}$$



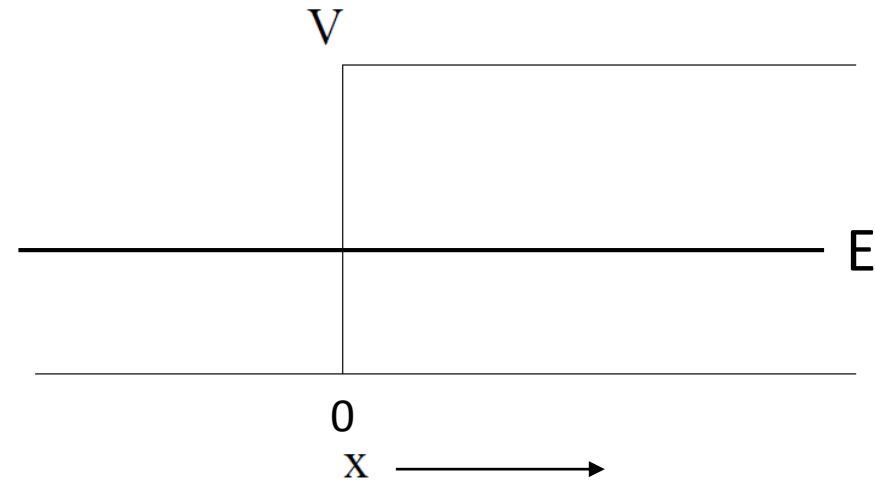
Step Potential: Case I - (Energy of particle E) $> V$

Solution : $\psi_{II}(x, t) = e^{-iEt/\hbar} \left[A'_1 e^{ip'x/\hbar} + A'_2 e^{-ip'x/\hbar} \right]$

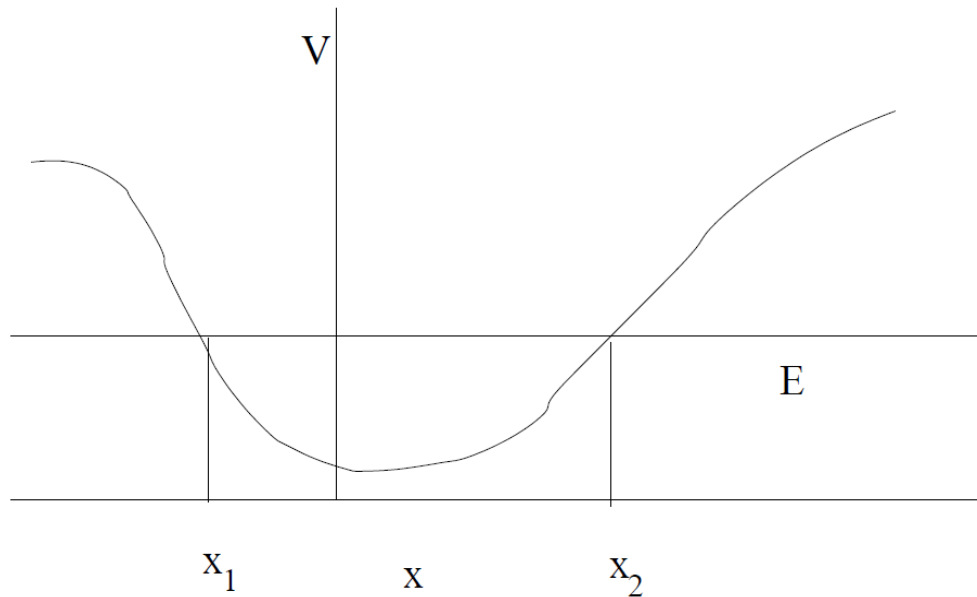
$$\psi_I(x, t) = e^{-iEt/\hbar} \left[A_1 e^{ipx/\hbar} + A_2 e^{-ipx/\hbar} \right]$$



Step Potential: Case II - (Energy of particle E) $< V$



In Classical Mechanics



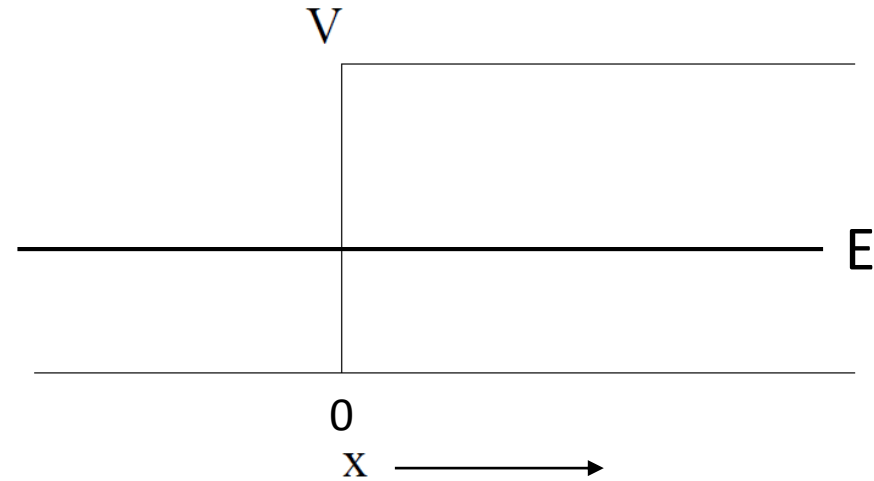
$$\boxed{\frac{p^2}{2m} + V(x) = E}$$

$$p = \pm \sqrt{2m(E - V(x))}$$

In Classical Mechanics, the particle's motion is restricted between x_1 and x_2 , as otherwise p becomes imaginary

What happens in Quantum Mechanics?

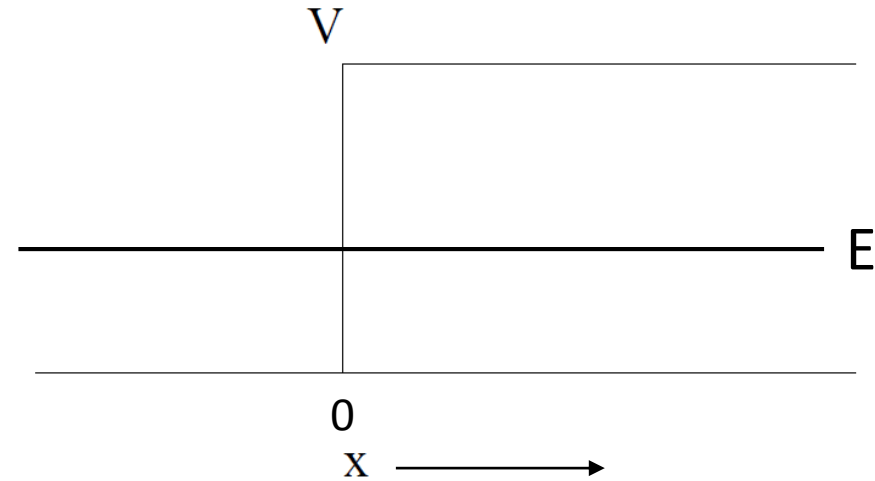
Step Potential: Case II - (Energy of particle E) $< V$



Defining - $\sqrt{2m(E - V)} = \sqrt{-1} \sqrt{2m(V - E)} = iq$

$$\frac{d^2 X}{dx^2} = \frac{q^2}{\hbar^2} X$$

Step Potential: Case II - (Energy of particle E) $< V$

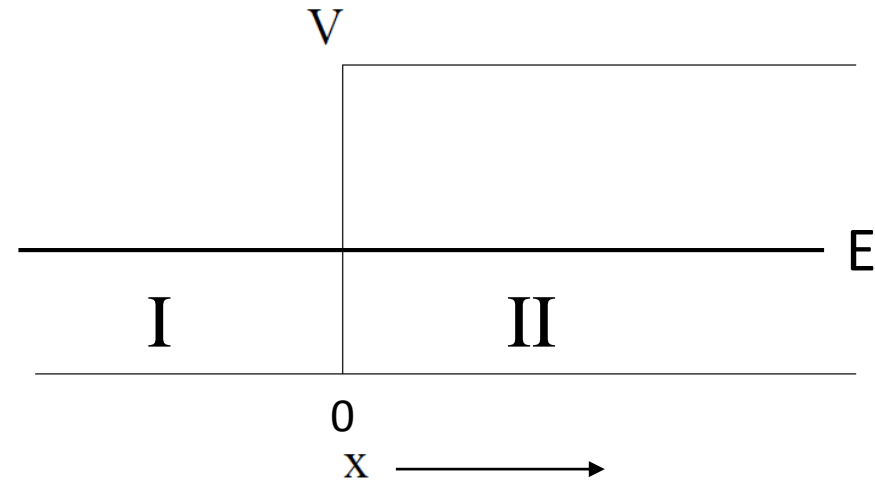


Defining - $\sqrt{2m(E - V)} = \sqrt{-1} \sqrt{2m(V - E)} = iq$

$$\frac{d^2 X}{dx^2} = \frac{q^2}{\hbar^2} X$$

$$X(x) = A_1 e^{-qx/\hbar} + A_2 e^{qx/\hbar}$$

Step Potential: Case II - (Energy of particle E) $< V$

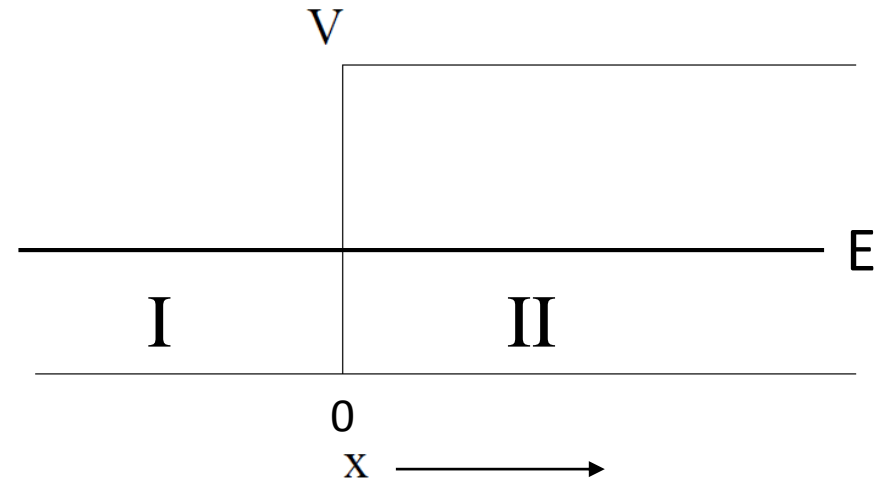


$$X(x) = A_1 e^{-qx/\hbar} + A_2 e^{qx/\hbar}$$

To ensure wavefunction remains finite as $x \rightarrow +\infty$, $A_2 = 0$

$$X_{II}(x) = A_1 e^{-qx/\hbar}$$

Step Potential: Case II - (Energy of particle E) $< V$



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To ensure wavefunction remains finite as $x \rightarrow +\infty$, $A_2 = 0$

$$X_{II}(x) = A_1 e^{-qx/\hbar}$$

Total solution for within the potential step -

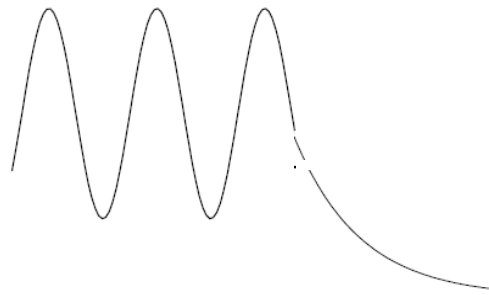
$$\psi_{II}(x, t) = A_1 e^{-iEt/\hbar} e^{-qx/\hbar}$$

Step Potential: Case II - (Energy of particle E) $< V$

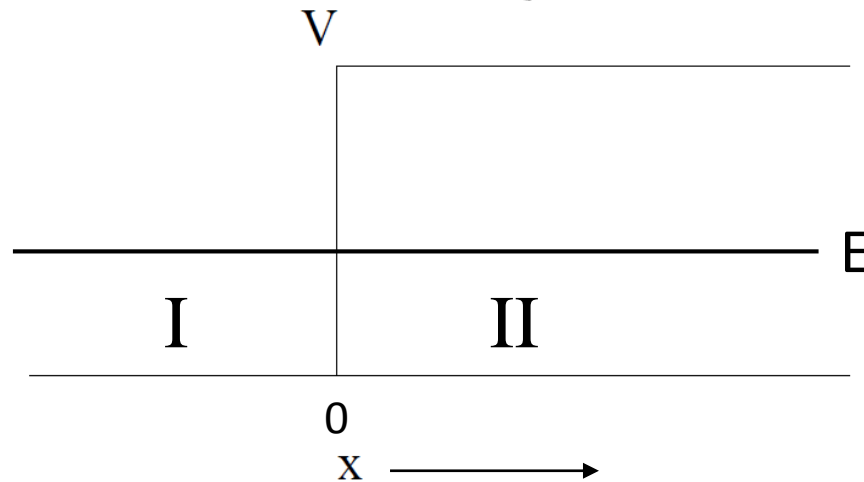
$$\psi_{II}(x, t) = A_1 e^{-iEt/\hbar} e^{-qx/\hbar}$$

$$\psi_I(x, t) = e^{-iEt/\hbar} [A_1 e^{ipx/\hbar} + A_2 e^{-ipx/\hbar}]$$

$$p = \sqrt{2mE}$$



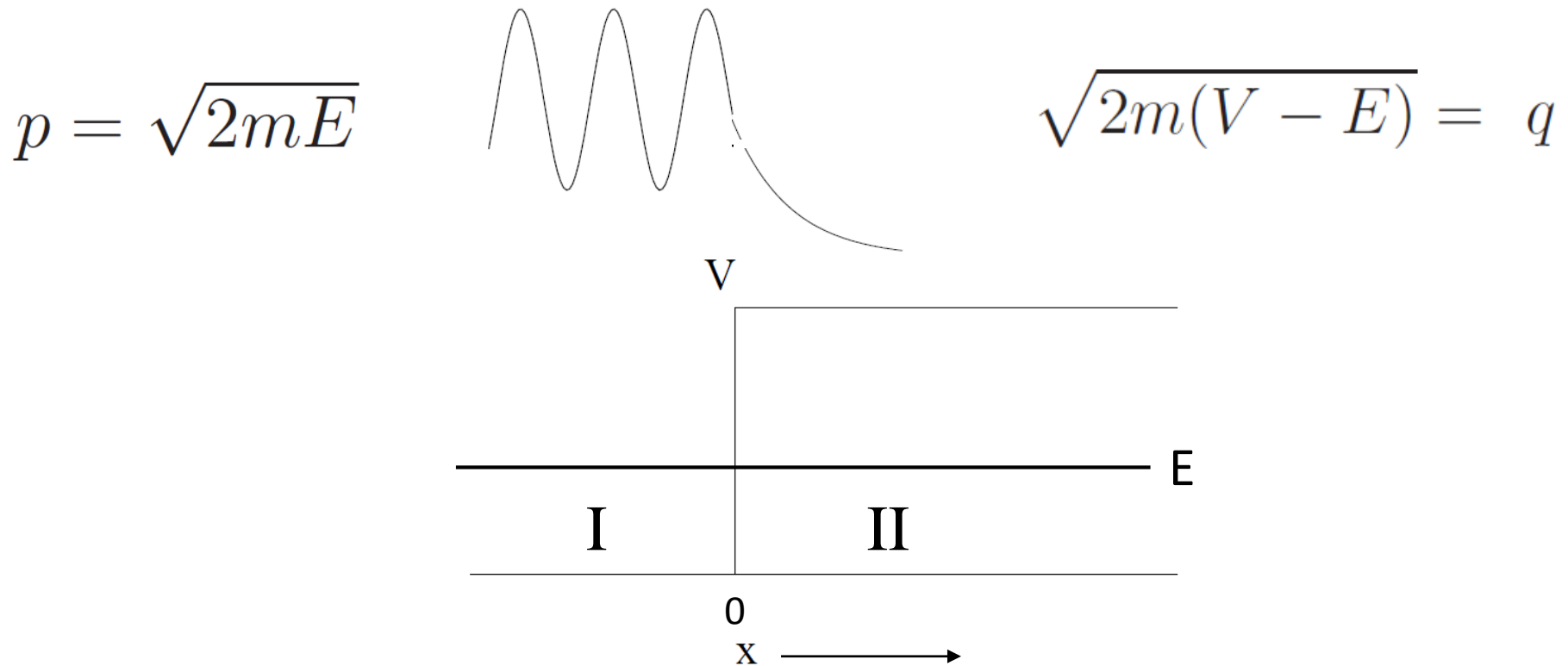
$$\sqrt{2m(V - E)} = q$$



Step Potential: Case II - (Energy of particle E) $< V$

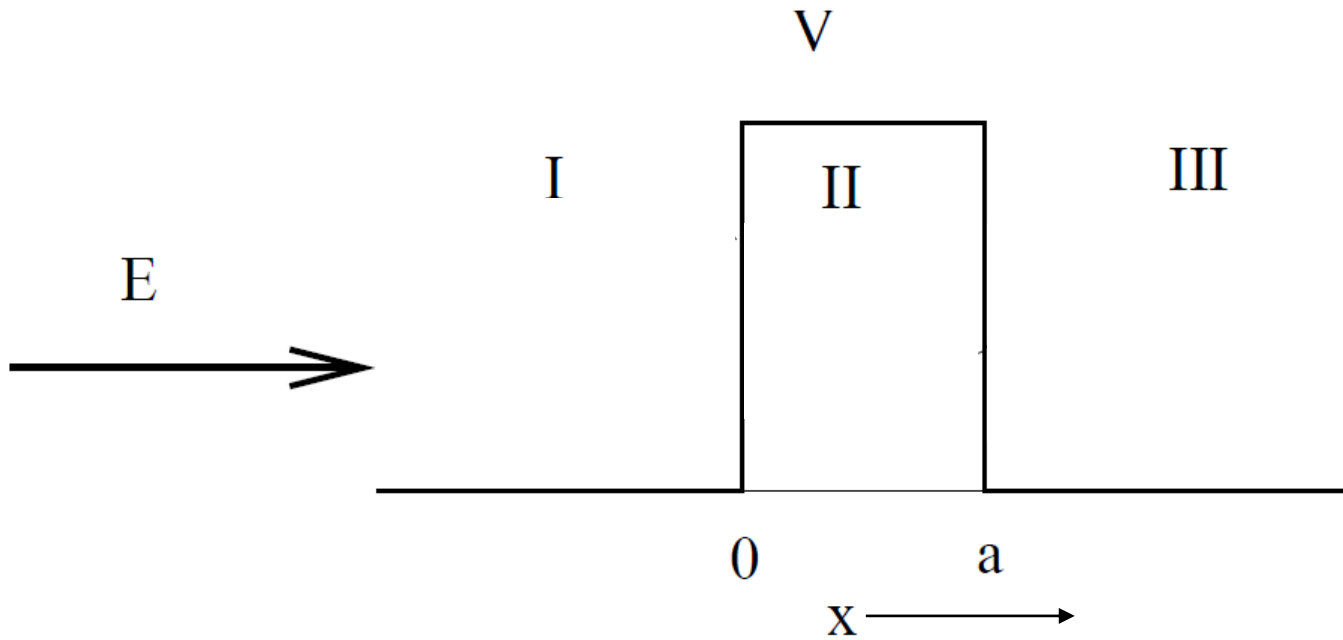
$$\psi_{II}(x, t) = A_1 e^{-iEt/\hbar} e^{-qx/\hbar}$$

$$\psi_I(x, t) = e^{-iEt/\hbar} [A_1 e^{ipx/\hbar} + A_2 e^{-ipx/\hbar}]$$

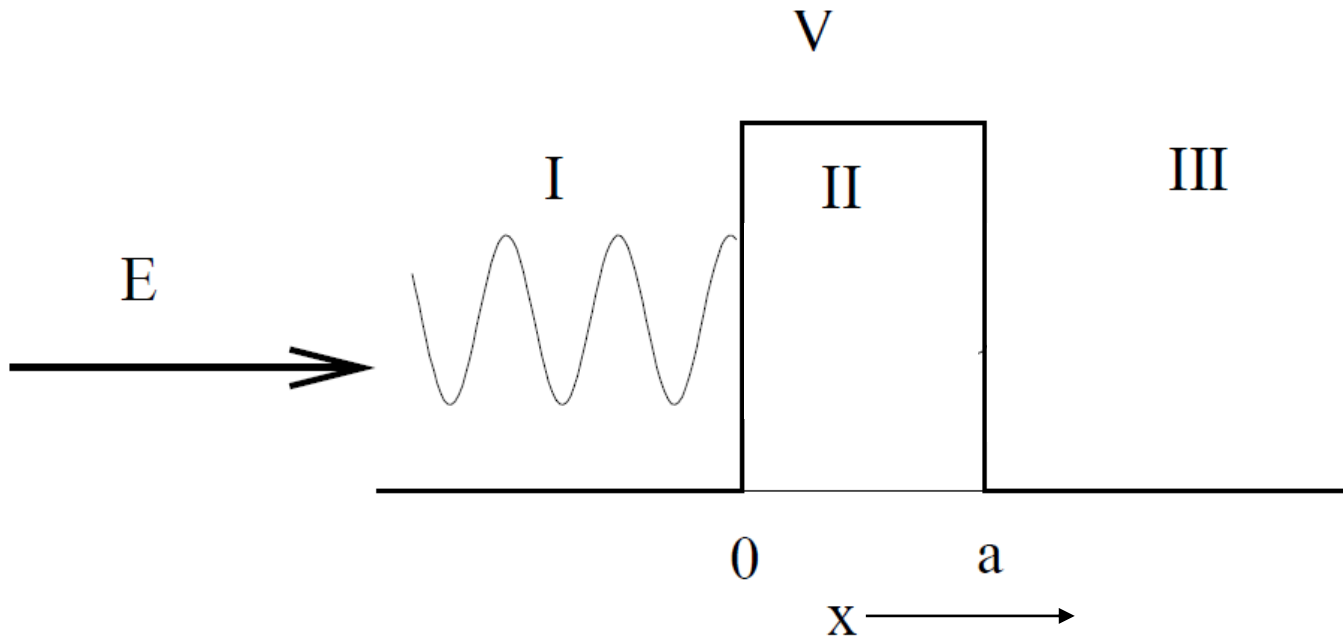


So finite probability of finding particle within barrier for $E < V$!

Quantum Tunneling

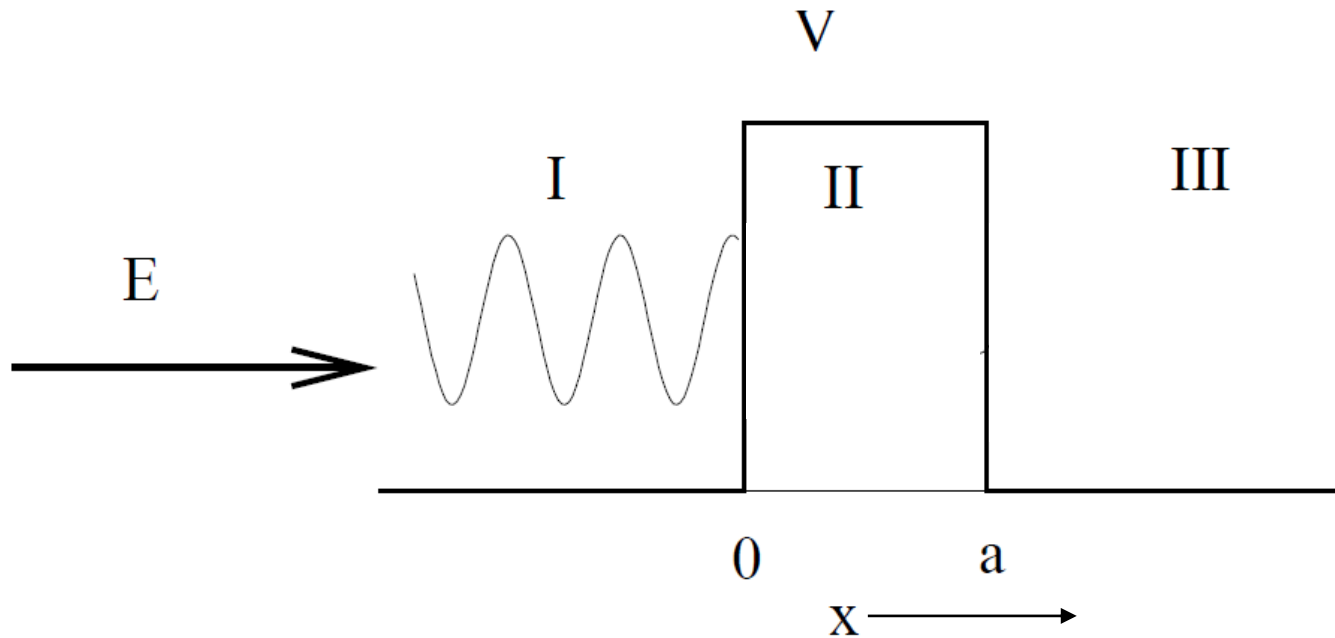


Quantum Tunneling



$$\psi_I(x, t) = e^{-iEt/\hbar} [A_I e^{ipx/\hbar} + B_I e^{-ipx/\hbar}]$$

Quantum Tunneling

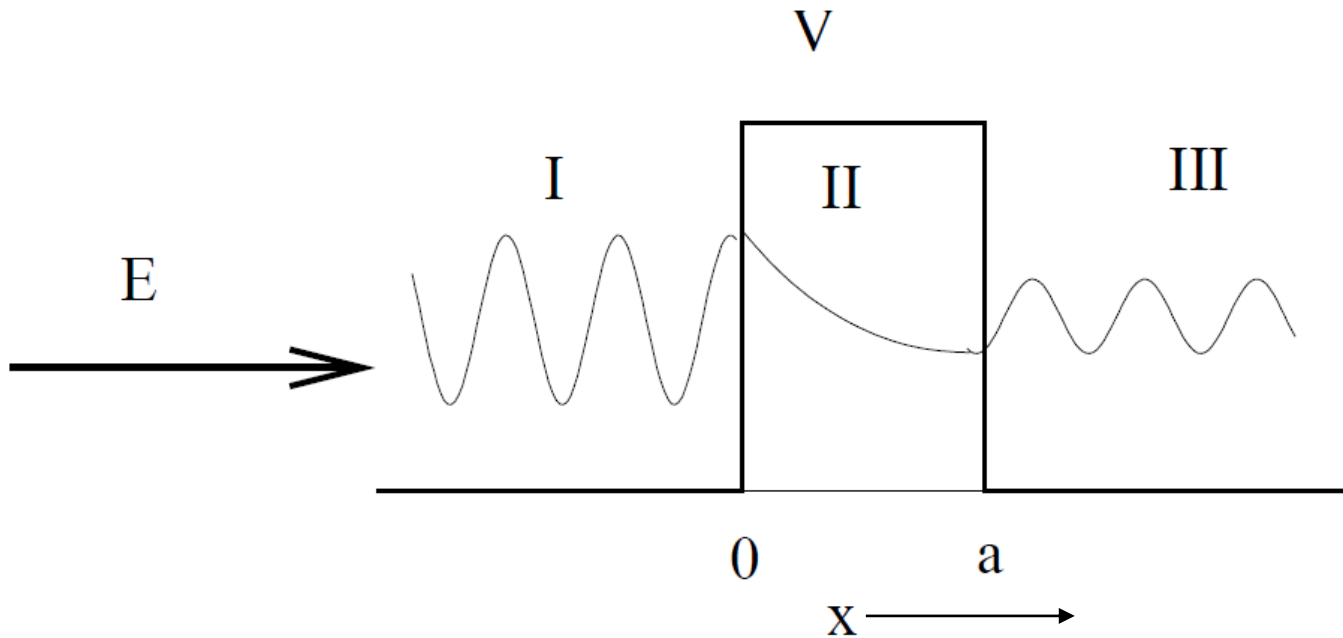


$$\psi_I(x, t) = e^{-iEt/\hbar} \left[\underbrace{A_I e^{ipx/\hbar}}_{\text{Incident particle travelling along } +x \text{ direction}} + \underbrace{B_I e^{-ipx/\hbar}}_{\text{Reflected particle travelling along } -x \text{ direction}} \right]$$

Incident particle travelling
along $+x$ direction

Reflected particle travelling
along $-x$ direction

Quantum Tunneling

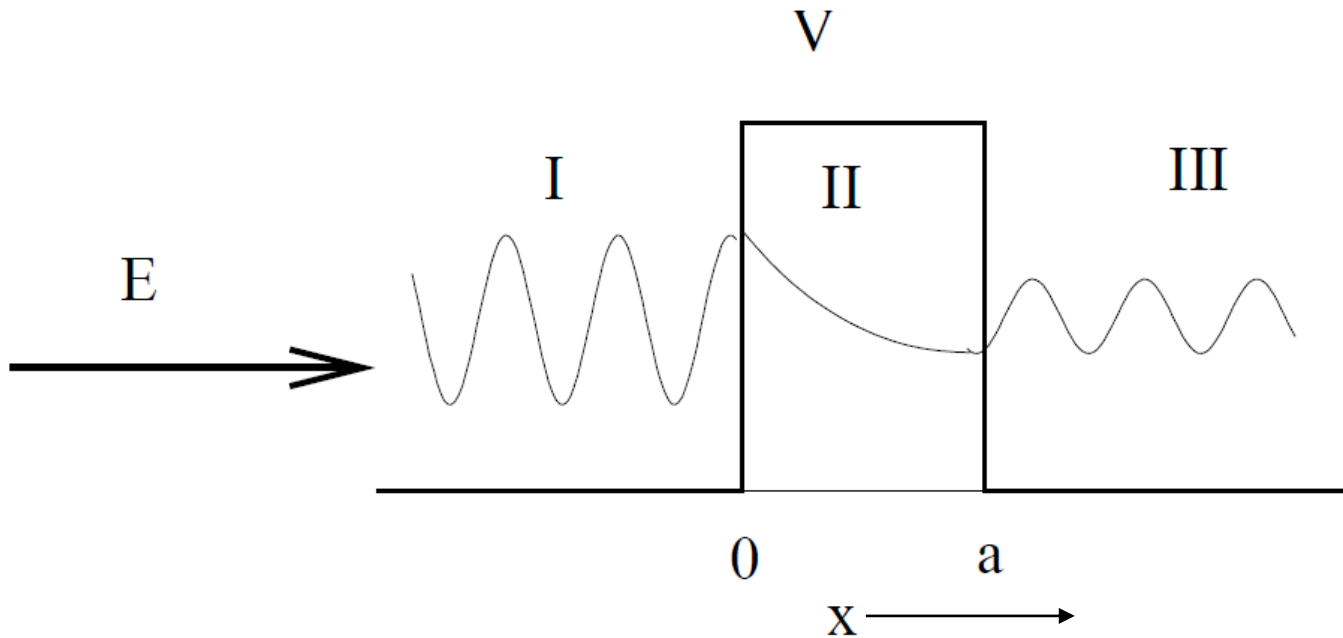


$$\psi_I(x, t) = e^{-iEt/\hbar} [A_I e^{ipx/\hbar} + B_I e^{-ipx/\hbar}]$$

$$\psi_{II}(x, t) = e^{-iEt/\hbar} [A_{II} e^{-qx/\hbar} + B_{II} e^{qx/\hbar}]$$

$$\psi_{III}(x, t) = e^{-iEt/\hbar} [A_{III} e^{ipx/\hbar} + B_{III} e^{-ipx/\hbar}]$$

Quantum Tunneling

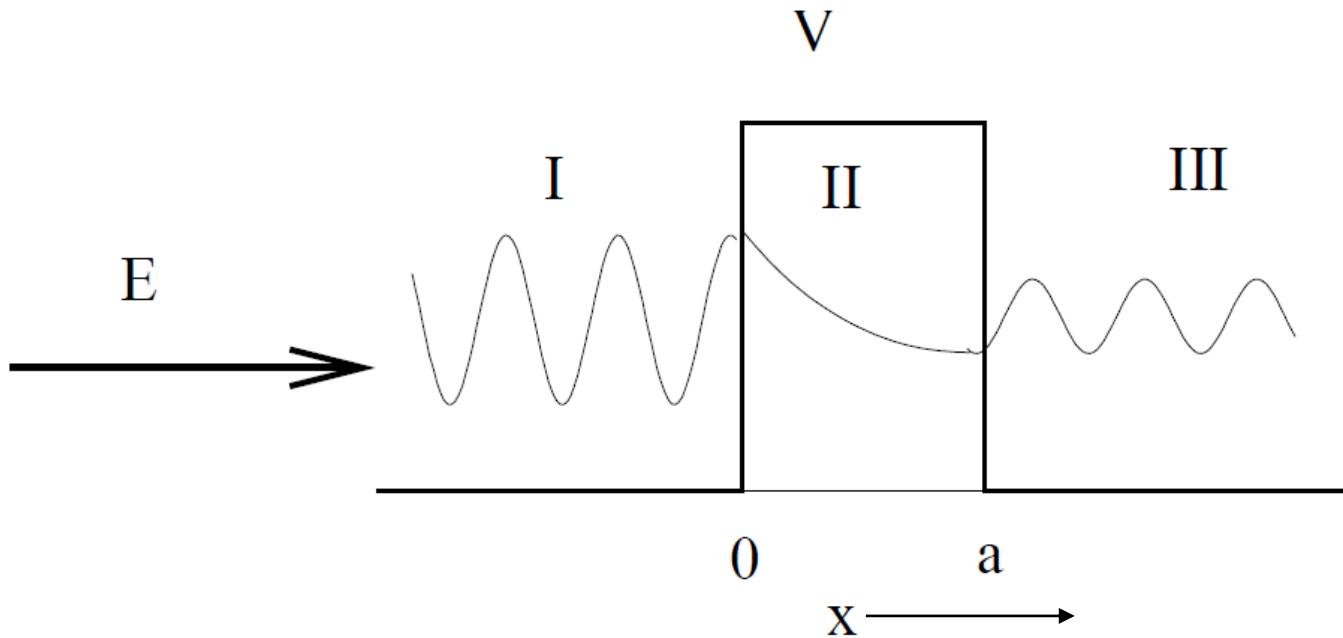


$$\psi_{III}(x, t) = e^{-iEt/\hbar} \left[\underbrace{A_{III}e^{ipx/\hbar}}_{\text{Particle travelling through the barrier and along } +x \text{ direction}} + \underbrace{B_{III}e^{-ipx/\hbar}}_{\text{Particle incident from the right along } -x \text{ direction. Here } B_{III} = 0} \right]$$

Particle travelling through the barrier and along $+x$ direction

Particle incident from the right along $-x$ direction.
Here $B_{III} = 0$

Quantum Tunneling

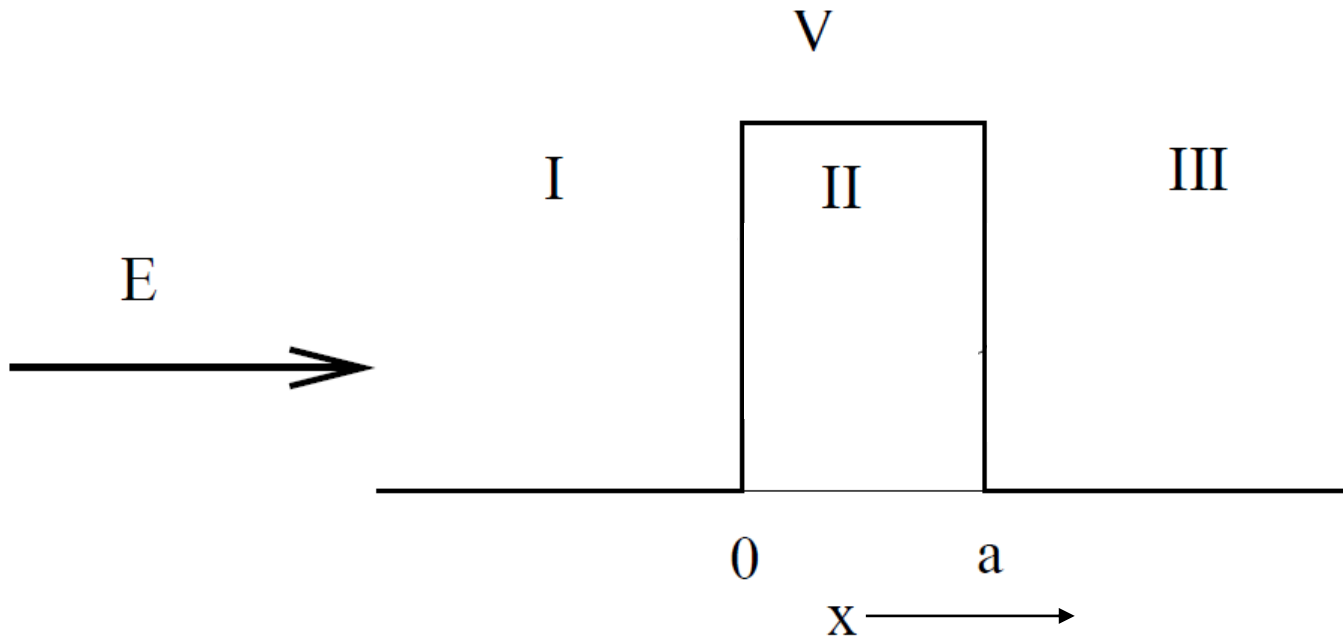


$$\psi_I(x, t) = e^{-iEt/\hbar} [A_I e^{ipx/\hbar} + B_I e^{-ipx/\hbar}]$$

$$\psi_{II}(x, t) = e^{-iEt/\hbar} [A_{II} e^{-qx/\hbar} + B_{II} e^{qx/\hbar}]$$

$$\psi_{III}(x, t) = e^{-iEt/\hbar} [A_{III} e^{ipx/\hbar}]$$

Quantum Tunneling



Boundary conditions to be satisfied at $x = 0$ –

$$\psi_I(0, t) = \psi_{II}(0, t)$$

$$\left(\frac{\partial \psi_I}{\partial x} \right)_{x=0} = \left(\frac{\partial \psi_{II}}{\partial x} \right)_{x=0}$$

Origin of boundary conditions

The wavefunction $\psi(x, t)$ (considering 1D) is governed by Schrödinger's equation -

$$i\hbar \frac{\partial}{\partial t} \psi(x, t) = \frac{-\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi(x, t) + V(x, t) \psi$$

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$$i\hbar \frac{\partial}{\partial t} \psi(x, t) = \frac{-\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi(x, t) + V(x, t) \psi$$

Further assumption for ease of calculation

Let us assume that the potential step is very high, i.e. $V \gg E$

$$q = \sqrt{2m(V - E)} \approx \sqrt{2mV}$$

$$p = \sqrt{2mE} \qquad \frac{p}{q} = \sqrt{\frac{E}{V}} \ll 1$$

Quantum Tunneling

$$\psi_I(x, t) = e^{-iEt/\hbar} [A_I e^{ipx/\hbar} + B_I e^{-ipx/\hbar}]$$

$$\psi_{II}(x, t) = e^{-iEt/\hbar} [A_{II} e^{-qx/\hbar} + B_{II} e^{qx/\hbar}]$$

Applying boundary conditions at $x = 0$

$$\psi_I(0, t) = \psi_{II}(0, t) \Rightarrow A_I + B_I = A_{II} + B_{II}$$

$$\left(\frac{\partial \psi_I}{\partial x} \right)_{x=0} = \left(\frac{\partial \psi_{II}}{\partial x} \right)_{x=0} \Rightarrow ip(A_I - B_I) = -q(A_{II} - B_{II})$$

$$A_I - B_I = \frac{iq}{p} (A_{II} - B_{II})$$

Quantum Tunneling

$$\psi_{II}(x, t) = e^{-iEt/\hbar} [A_{II}e^{-qx/\hbar} + B_{II}e^{qx/\hbar}]$$

$$\psi_{III}(x, t) = e^{-iEt/\hbar} [A_{III}e^{ipx/\hbar}]$$

Applying boundary conditions at $x = a$

Wavefunction continuity -

$$A_{II}e^{-qa/\hbar} + B_{II}e^{qa/\hbar} = A_{III}e^{ipa/\hbar}$$

First derivative continuity -

$$-q [A_{II}e^{-qa/\hbar} - B_{II}e^{qa/\hbar}] = ipA_{III}e^{ipa/\hbar}$$

$$A_{II}e^{-qa/\hbar} - B_{II}e^{qa/\hbar} = \frac{-ip}{q} A_{III}e^{ipa/\hbar}$$

Quantum Tunneling

Applying boundary conditions at $x = a$

First derivative continuity -

$$A_{II}e^{-qa/\hbar} - B_{II}e^{qa/\hbar} = \frac{-ip}{q}A_{III}e^{ipa/\hbar}$$

Since - $p/q \ll 1$

$$A_{II}e^{-qa/\hbar} - B_{II}e^{qa/\hbar} = 0$$

$$B_{II} = e^{-2qa/\hbar}A_{II} \ll A_{II}$$

Small

$$\psi_{II}(x, t) = e^{-iEt/\hbar} [A_{II}e^{-qx/\hbar} + B_{II}e^{qx/\hbar}]$$

Quantum Tunneling

Wavefunction continuity at $x = a$

$$A_{II}e^{-qa/\hbar} + B_{II}e^{qa/\hbar} = A_{III}e^{ipa/\hbar}$$

Also from last slide - $A_{II}e^{-qa/\hbar} - B_{II}e^{qa/\hbar} = 0$

$$A_{III} = 2e^{-ipa/\hbar}e^{-qa/\hbar}A_{II}$$

Quantum Tunneling

Wavefunction continuity at $x = a$

$$A_{II}e^{-qa/\hbar} + B_{II}e^{qa/\hbar} = A_{III}e^{ipa/\hbar}$$

Also from last slide - $A_{II}e^{-qa/\hbar} - B_{II}e^{qa/\hbar} = 0$

$$A_{III} = 2e^{-ipa/\hbar}e^{-qa/\hbar}A_{II}$$

Recalling boundary conditions at $x = 0$

$$A_I + B_I = A_{II} + \cancel{B_{II}} \quad \& \quad A_I - B_I = \frac{iq}{p} (A_{II} - \cancel{B_{II}})$$
$$\left(1 + \frac{iq}{p}\right) A_{II} = 2A_I$$

Quantum Tunneling

$$\left(1 + \frac{iq}{p}\right) A_{II} = 2A_I$$

Since $q/p \gg 1$

$$A_{II} = -\frac{ip}{q} 2A_I$$

Quantum Tunneling

$$\left(1 + \frac{iq}{p}\right) A_{II} = 2A_I$$

Since $q/p \gg 1$

$$A_{II} = -\frac{ip}{q} 2A_I$$

Also from last slide -

$$A_{III} = 2e^{-ipa/\hbar} e^{-qa/\hbar} A_{II}$$

$$A_{III} = -4i\frac{p}{q} e^{-ipa/\hbar} e^{-qa/\hbar} A_I$$

Quantum Tunneling

$$A_{III} = -4i \frac{p}{q} e^{-ipa/\hbar} e^{-qa/\hbar} A_I$$

The transmission coefficient T is -

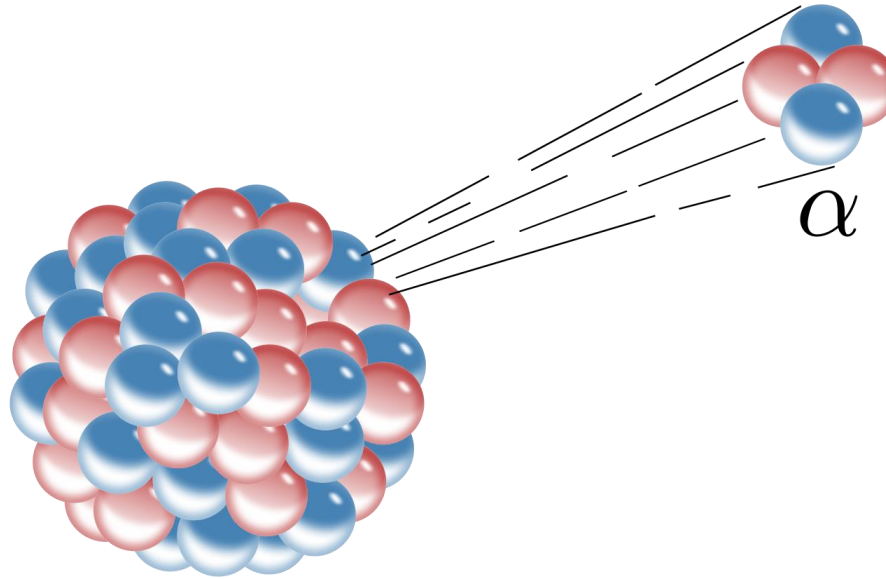
$$T = \frac{|A_{III}|^2}{|A_I|^2} = 16 \frac{p^2}{q^2} e^{-2qa/\hbar}$$

Since - $\frac{p}{q} = \sqrt{\frac{E}{V}} \ll 1$

$$T = 16 \frac{E}{V} e^{-2a\sqrt{2mV}/\hbar}$$

Some applications of Quantum Tunneling effect

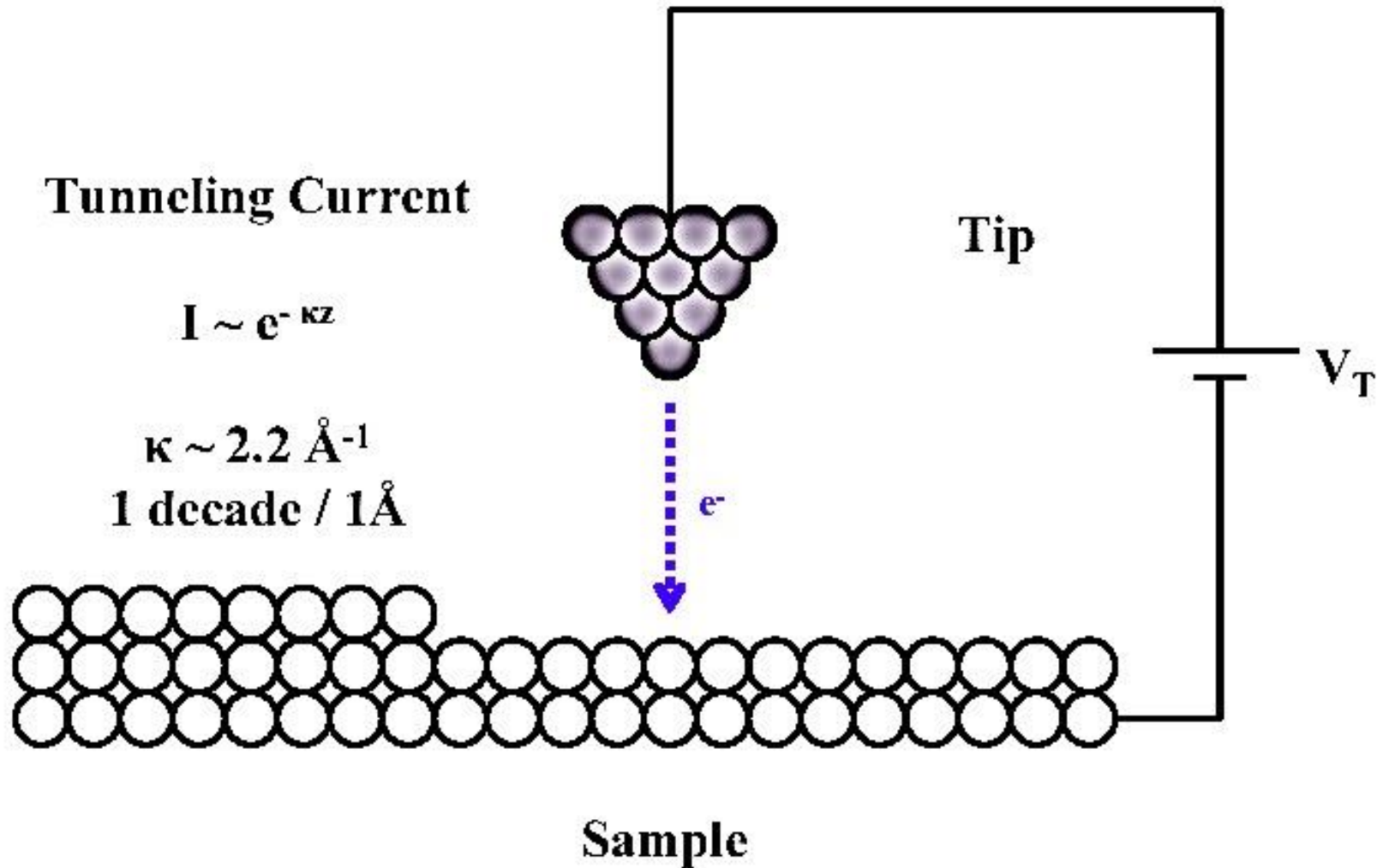
Alpha Decay



**Alpha particle (a Helium-4 nucleus) is held within nucleus by nuclear forces.
It experiences a potential barrier to come outside the nucleus.
Still Alpha particles can come outside nucleus because of tunneling**

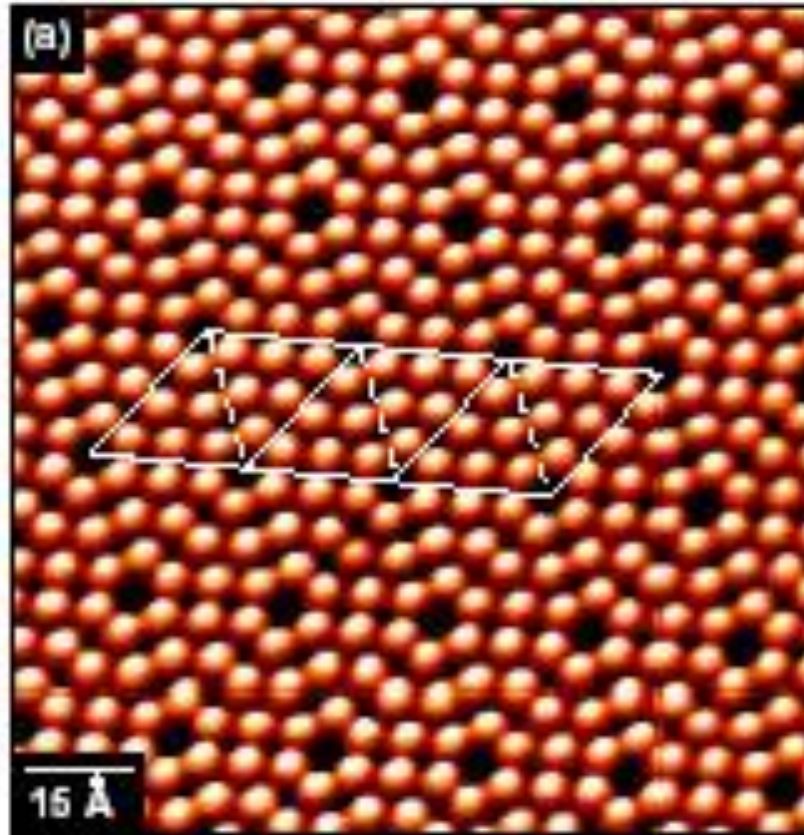
Source of image: Wikipedia

Scanning Tunneling Microscope



Electron tunneling across the air gap

Scanning Tunneling Microscope



STM scan of a Si sample
Each bright spot indicates a Si atom