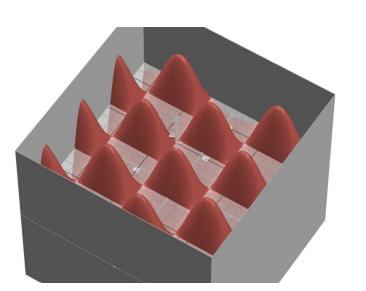
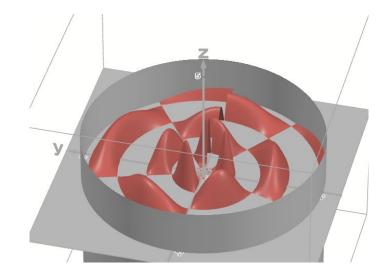
Acoustophoresis and Standing Waves



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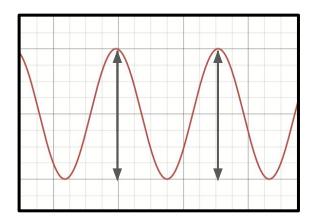


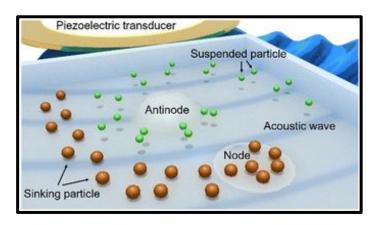
Background

Wave equation models propagation of a wave (c = propagation speed)

$$\frac{\partial^2 f}{\partial t^2} = c^2 \nabla^2 f = c^2 \left(\frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} \right)$$

Standing waves are used in acoustophoresis to separate particles





What standing wave boundaries are most helpful for acoustophoresis?

Square Case

The wave equation is the simplest form

$$\frac{\partial^2 f}{\partial t^2} = c^2 \left(\frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} \right)$$

- Assume a solution of the form |f(x,y,t) = X(x)Y(y)T(t)|
- Let $\left|\omega^2 = -\frac{1}{T}\frac{d^2T}{dt^2}\right|$ and $\left|m^2 = -\frac{1}{Y}\frac{d^2Y}{du^2}\right|$
- This reduces the wave equation to a set of 3 independent 2nd-order differential equations.

Multiplying them together,
$$\left| f(x,y,t) = C \cos \left(\sqrt{\frac{\omega^2 - m^2}{c^2}} x \right) \cos(my) \cos(\omega t) \right|$$

Circular Case

The wave equation in polar coordinates is

$$\frac{\partial^2 f}{\partial t^2} = c^2 \nabla^2 f = c^2 \left(\frac{\partial^2 f}{\partial r^2} + \frac{1}{r} \frac{\partial f}{\partial r} + \frac{1}{r^2} \frac{\partial^2 f}{\partial \theta^2} \right)$$

Assuming $|f(r, \theta, t) = R(r)\Theta(\theta)T(t)|$ and making substitutions,

$$0 = \frac{d^2R}{dr^2} + \frac{1}{r}\frac{dR}{dr} + \left(\frac{\omega^2}{c^2} - \frac{m^2}{r^2}\right)R \qquad \qquad \omega^2 = -\frac{1}{T}\frac{d^2T}{dt^2} \qquad \boxed{m^2 = -\frac{1}{\Theta}\frac{d^2\Theta}{d\theta^2}}$$

$$\omega^2 = -\frac{1}{T} \frac{d^2 T}{dt^2}$$

$$m^2 = -\frac{1}{\Theta} \frac{d^2 \Theta}{d\theta^2}$$

$$0 = u^{2} \frac{d^{2}R}{du^{2}} + u \frac{dR}{du} + (u^{2} - m^{2}) R \qquad u = \omega r/c$$

$$u = \omega r/c$$

Circular Case (cont'd)

$$0 = u^{2} \frac{d^{2}R}{du^{2}} + u \frac{dR}{du} + (u^{2} - m^{2}) R$$
 $u = \omega r/c$

• This is the Bessel equation, and its solution has the form

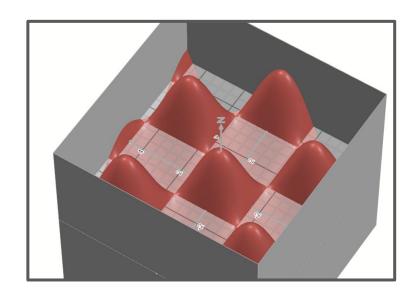
$$J_m(u) = \sum_{k=0}^{\infty} \frac{(-1)^k}{k!(k+m)!} \left(\frac{u}{2}\right)^{2k+m}$$

Now we can get our solution:

$$f(r, \theta, t) = CR(r)\Theta(\theta)T(t) = CJ_m(u)\cos(m\theta)\cos(\omega t)$$

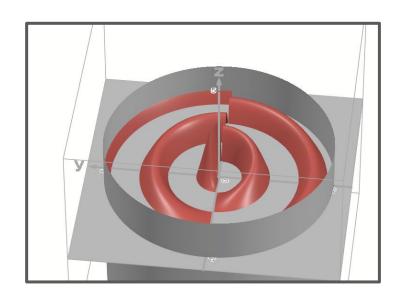
$$f(r,\theta,t) = C\cos(m\theta)\cos(\omega t)\sum_{k=0}^{\infty} \frac{(-1)^k}{k!(k+m)!} \left(\frac{\omega r}{2c}\right)^{2k+m}$$

Conclusions



Square boundary

- Uniform peak heights
- Always a peak at the origin
- Better for few number of particles



Circular boundary

- Variable heights on each "ring"
- z=0 around origin if $m >> \omega$ or m >> c
- Better for spectrum of particle weights

Sources

Images

https://pubs.rsc.org/en/content/articlelanding/2024/lc/d4lc00277f

Materials

- Circular membrane case
 https://www-eng.lbl.gov/~shuman/NEXT/MATERIALS&COMPONENTS/MISC/Standing-Waves-on-a-Circular-Membrane.pdf
- Bessel functions

 <a href="https://math.libretexts.org/Bookshelves/Differential_Equations/A_First_Course_in_Differential_Equations for Scientists and Engineers (Herman)/04%3A_Series_Solutions/4.06%3A_Bessel_Functions

 sel_Functions
- Nice video
 https://www.youtube.com/watch?v=6yV17h26llc