### 例 点电荷电场中试验电荷的电势能和电势

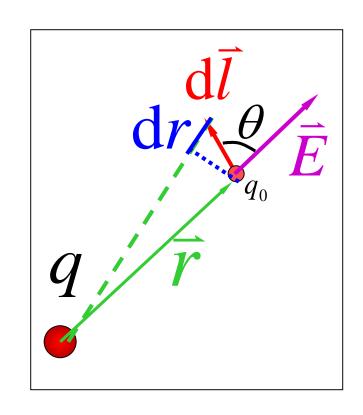
$$\vec{E} = \frac{q}{4 \pi \varepsilon_0 r^2} \hat{e}_r$$

$$\Rightarrow W_{\infty} = 0$$

$$W = \int_{r}^{\infty} \frac{q_0 q}{4 \pi \varepsilon_0 r^2} \vec{r}_0 \cdot d\vec{l}$$

$$= \int_{r}^{\infty} \frac{q_0 q \mathrm{d}r}{4 \pi \varepsilon_0 r^2}$$

$$W = \frac{q_0 q}{4 \pi \varepsilon_0 r}$$



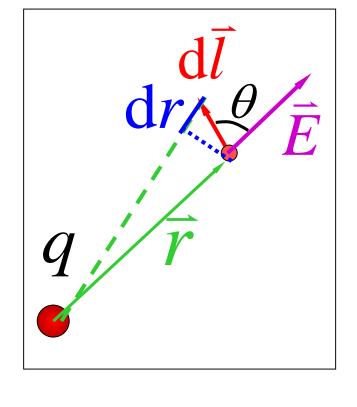
#### 点电荷的电势

$$\vec{E} = \frac{q}{4 \pi \varepsilon_0 r^2} \vec{r}_0 \quad \Leftrightarrow U_{\infty} = 0$$

$$U = \int_{r}^{\infty} \frac{q}{4 \pi \varepsilon_{0} r^{2}} \hat{e}_{r} \cdot d\vec{l}$$

$$q \quad \mathbf{r}^{\infty} d\mathbf{r}$$

$$= \frac{q}{4 \pi \varepsilon_0} \int_r^\infty \frac{\mathrm{d}r}{r^2}$$



$$U = \frac{q}{4 \pi \varepsilon_0 r}$$

——球对称性

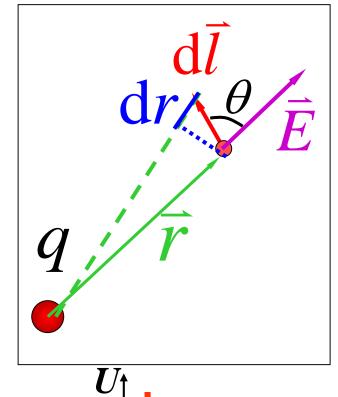
### 六、电势的计算

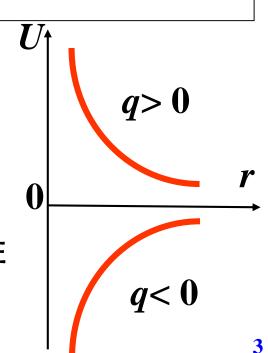
1. 点电荷的电势

$$\vec{E} = \frac{q}{4\pi\varepsilon_0 r^2} \vec{r}_0 \quad \Leftrightarrow U_{\infty} = 0$$

$$U = \int_r^{\infty} \frac{q}{4\pi\varepsilon_0 r^2} \vec{r}_0 \cdot d\vec{l}$$

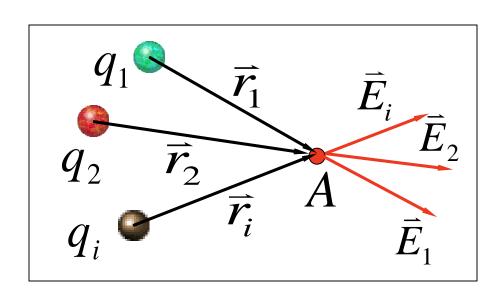
$$= \frac{q}{4\pi\varepsilon_0} \int_r^{\infty} \frac{dr}{r^2}$$





#### 2. 电势的叠加原理

$$\begin{split} U_{\infty} &= 0 \\ \mathbf{点电荷系} \; \vec{E} = \sum_{i} \vec{E}_{i} \\ U_{A} &= \int_{A}^{\infty} \vec{E} \cdot \mathrm{d}\vec{l} = \int_{A}^{\infty} \sum_{i} \vec{E}_{i} \cdot \mathrm{d}\vec{l} \\ &= \sum_{A}^{\infty} \vec{E}_{i} \cdot \mathrm{d}\vec{l} = \sum_{A}^{\infty} U_{i} \end{split}$$



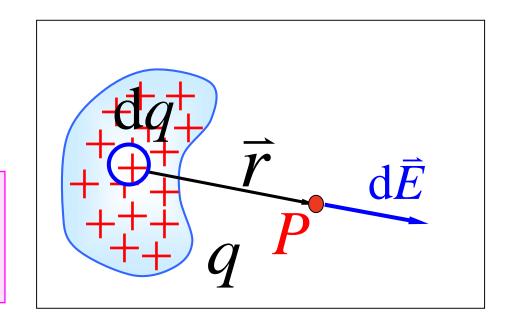
对于点电荷—— 
$$U_i = \frac{q_i}{4\pi \varepsilon_0 r_i}$$

对于点电荷系——
$$U_{A}=\sum_{i}U_{iA}=\sum_{i}rac{q_{i}}{4\pi\,arepsilon_{0}r_{i}}$$

#### 3. 连续分布电荷的电势

$$dU = \frac{dq}{4 \pi \varepsilon_0 r}$$

$$U_P = \int dU = \int_q \frac{dq}{4 \pi \varepsilon_0 r}$$



求电势 的方法

$$ightarrow$$
利用  $U_P = \int dU = \int_q \frac{dq}{4 \pi \varepsilon_0 r}$ 

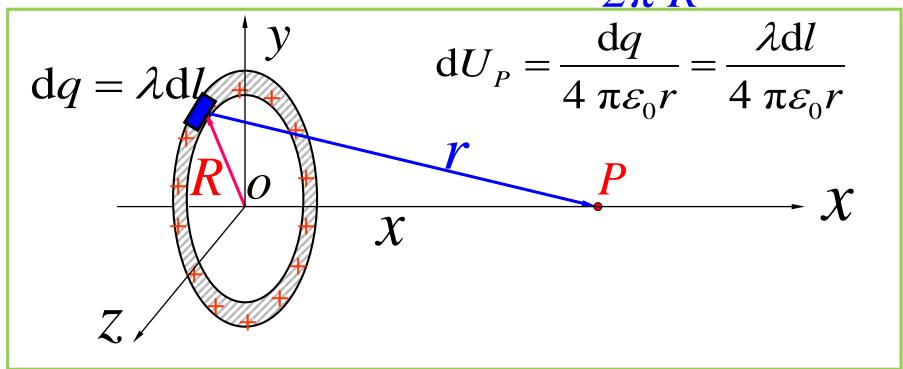
——这一结果已选无限远处为电势零点。

> 若已知在积分路径上 产的函数表达式,

则 
$$U_A = \int_A^{U=0} ec{E} \cdot \mathrm{d}ec{l}$$

例1 正电荷q均匀分布在半径为R的细圆环上。 求圆环轴线上距环心为x处点 P的电势。

$$\beta = \frac{q}{2\pi R} \qquad dq = \lambda dl = \frac{q dl}{2\pi R}$$

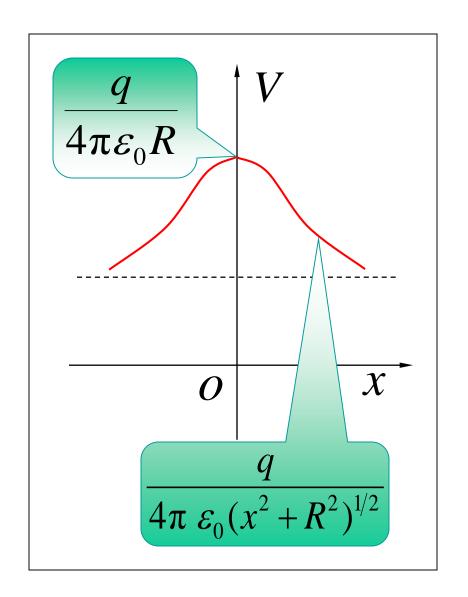


$$U_{P} = \int_{(q)} dU_{p} = \frac{\lambda}{4 \pi \varepsilon_{0} r} \int_{0}^{2 \pi R} dl = \frac{q}{4 \pi \varepsilon_{0} r} = \frac{q}{4 \pi \varepsilon_{0} \sqrt{x^{2} + R^{2}}}$$

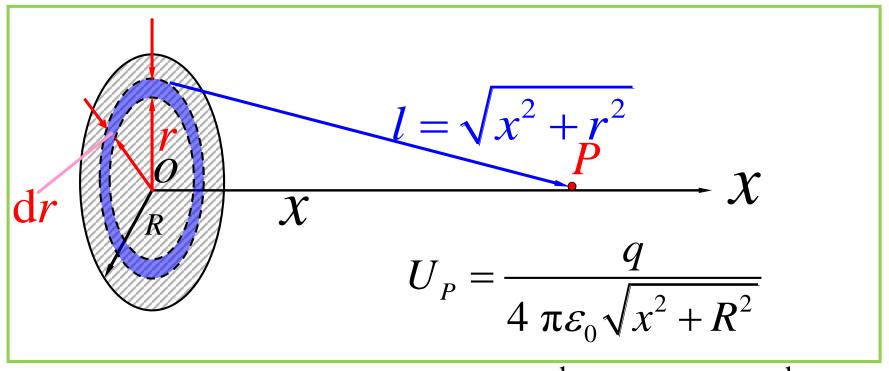
$$U_P = \frac{q}{4 \pi \varepsilon_0 \sqrt{x^2 + R^2}}$$

### 讨论:

若
$$x >> R$$
, $U_P = \frac{q}{4 \pi \varepsilon_0 x}$ 



### 例2 求均匀带电薄圆盘轴线上的电势。



解: 
$$dq = \sigma 2 \pi r dr$$

$$dq = \sigma 2 \pi r dr \qquad dU_p = \frac{dq}{4\pi \varepsilon_0 \sqrt{x^2 + r^2}} = \frac{\sigma \pi r dr}{4\pi \varepsilon_0 \sqrt{x^2 + r^2}}$$

$$U_{P} = \int_{(q)} dU = \frac{1}{4 \pi \varepsilon_{0}} \int_{0}^{R} \frac{\sigma 2 \pi r dr}{\sqrt{x^{2} + r^{2}}} = \frac{\sigma}{2\varepsilon_{0}} (\sqrt{x^{2} + R^{2}} - x)$$

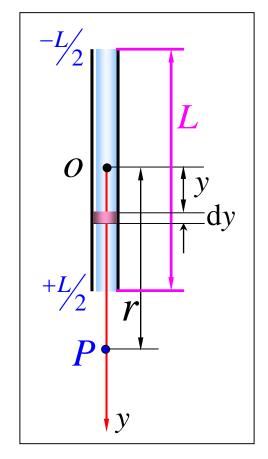
### 例3 求长为L的均匀带电q直线延长上一点P的电势。

解: 
$$\lambda = \frac{q}{L}$$
  $dq = \lambda dy$   $\Rightarrow U_{\infty} = 0$ 

$$dU = \frac{dq}{4\pi\varepsilon_0(r-y)} = \frac{\lambda dy}{4\pi\varepsilon_0(r-y)}$$

$$U = \int dU = \int_{-L/2}^{L/2} \frac{\lambda dy}{4\pi\varepsilon_0 (r - y)}$$

$$= \frac{\lambda}{4\pi\varepsilon_0} \int_{-L/2}^{L/2} \frac{\mathrm{d}y}{r - y} = \frac{\lambda}{4\pi\varepsilon_0} \ln \frac{r + \frac{L}{2}}{r - \frac{L}{2}}$$



### 例4 真空中,有一带均匀带电球壳,带电量为9 ,半径为8。

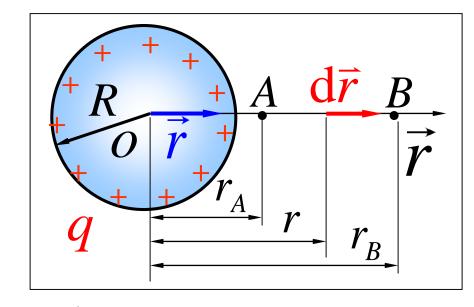
- 试求(1) 球壳外任意点的电势;(2) 球壳内任意点的电势;
  - (3) 球壳外两点间的电势差; (4) 球壳内两点间的电势差。

#### 解:

$$\begin{cases} r < R, \quad \vec{E}_1 = 0 \\ r > R, \quad \vec{E}_2 = \frac{q}{4 \pi \varepsilon_0 r^2} \hat{e}_r \end{cases}$$

以无限远为参考点。

(1) r > R 时



$$U_P(r) = \int_r^\infty \vec{E}_2 \cdot d\vec{r} = \frac{q}{4\pi\varepsilon_0} \int_r^\infty \frac{1}{r^2} \hat{e}_r \cdot d\vec{r}$$

$$= \frac{q}{4\pi\varepsilon_0} \int_r^\infty \frac{\mathrm{d}r}{r^2} = \frac{q}{4\pi\varepsilon_0 r}$$

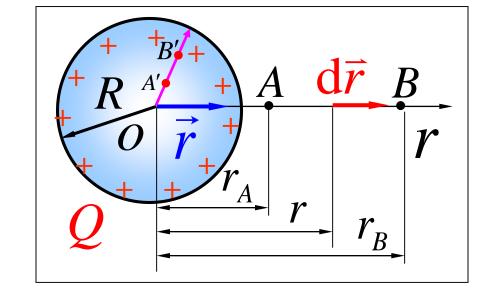
(2) 
$$r < R$$
 时

$$U_{\bowtie}(r) = \int_{r}^{R} \vec{E}_{1} \cdot d\vec{r} + \int_{R}^{\infty} \vec{E}_{2} \cdot d\vec{r} = \frac{q}{4 \pi \varepsilon_{0} R}$$

(3) r > R

$$U_A - U_B = \int_{r_A}^{r_B} \vec{E}_2 \cdot d\vec{r}$$

$$= \frac{q}{4 \pi \varepsilon_0} \int_{r_A}^{r_B} \frac{\vec{r}_0 \cdot d\vec{r}}{r^2}$$



$$= \frac{q}{4 \pi \varepsilon_0} \int_{r_A}^{r_B} \frac{\mathrm{d}r}{r^2} = \frac{q}{4 \pi \varepsilon_0} \left( \frac{1}{r_A} - \frac{1}{r_B} \right)$$

(4) 
$$r < R$$
  $U_{A'} - U_{B'} = \int_{r_{A'}}^{r_{B'}} \vec{E}_1 \cdot d\vec{r} = 0$ 

### 电场强度与电势的关系概括如下:

积分关系: 
$$U = \int_{(P)}^{(P_0)} \vec{E} \cdot d\vec{l}$$
 ---- 由场强求电势

微分关系: 
$$ar{E} = -
abla U$$
 ---- 由电势求场强

# 真空中静电场小结

## 一. 线索(基本定律、定理):

库仑定律
$$\vec{E} = \vec{F} / q_0$$

$$\vec{E} = \sum_{i} \frac{q_i \hat{e}_{r_i}}{4\pi \varepsilon_0 r_i^2} \rightarrow \begin{bmatrix} \oint_{S} \vec{E} \cdot d\vec{s} = \frac{\sum q_h}{\varepsilon_0} \\ \oint_{L} \vec{E} \cdot d\vec{l} = 0 \end{bmatrix}$$

## 二. 求静电场的方法:

静电场可以用电场强度来描述,静电场也可以用电势来描述。

# 真空中静电场小结(两两歌)

1.两个物理量

$$ec{E}$$
  $U$ 

2.两个基本性质方程 
$$\int_{S} \vec{E} \cdot d\vec{S} = \frac{\sum_{i} q_{i}}{\varepsilon_{0}} \int_{L} \vec{E} \cdot d\vec{l} = 0$$

$$\vec{E} = \int_{(Q)} d\vec{E} \qquad U = \int_{(Q)} dU \qquad \text{叠加}$$

$$\oint_{S} \vec{E} \cdot d\vec{S} = \frac{\sum_{i} q_{i}}{\varepsilon_{0}} \qquad U = \int_{(P)} \vec{E} \cdot d\vec{l} \quad \text{高斯}$$

## 4.强调两句话

## 注重典型场

# 注重叠加原理

$$\vec{E} = \frac{Q}{4\pi\varepsilon_0 r^2} \hat{r} \qquad U = \frac{Q}{4\pi\varepsilon_0 r}$$

$$U = \frac{Q}{4\pi\varepsilon_0 r}$$

$$r < R \quad E = 0$$

$$U = \frac{Q}{4\pi\varepsilon_0 R}$$

$$r > R \vec{E} = \frac{Q}{4\pi\varepsilon_0 r^2} \hat{r} \quad U = \frac{Q}{4\pi\varepsilon_0 r}$$

$$E = \frac{\lambda}{2\pi\varepsilon_0 r}$$

$$E = \frac{\sigma}{2\varepsilon_0}$$

# 三、几种典型电荷分布的 $\bar{E}$ 和 U

```
点电荷(?)
均匀带电球面(?)
均匀带电球体(?)
均匀带电无限长直线(?)
均匀带电无限大平面(?)
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