

SU/PG Implementation

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1 Introduction

Implementing the SU/PG method is a variant of the continuous Galerkin method, which enforces stability by adding a term to the test function. This method is written as

$$\begin{aligned} & \sum_{e \in \mathcal{T}_h} \int_{\Omega_k} \phi_i \frac{\partial \mathbf{u}_h}{\partial t} d\Omega_e - \sum_{e \in \mathcal{T}_h} \int_{\Omega_k} \nabla \phi_i \cdot \left(\vec{\mathbf{F}}_c(\mathbf{u}_h) - \vec{\mathbf{F}}_v(\mathbf{u}_h, \nabla \mathbf{u}_h) \right) + \phi_i \mathbf{S}(\mathbf{u}_h, \nabla \mathbf{u}_h) d\Omega_e + \\ & \sum_{e \in \mathcal{T}_h} \int_{\Omega_k} \nabla \phi_i \cdot \frac{\partial \vec{\mathbf{F}}_c(\mathbf{u}_h)}{\partial \mathbf{u}_h} [\tau] \left(\frac{\partial \mathbf{u}_h}{\partial t} + \nabla \cdot \left(\vec{\mathbf{F}}_c(\mathbf{u}_h) - \vec{\mathbf{F}}_v(\mathbf{u}_h, \nabla \mathbf{u}_h) \right) + \mathbf{S}(\mathbf{u}_h, \nabla \mathbf{u}_h) \right) d\Omega_e + \\ & \sum_{b \in \mathcal{B}_h} \int_{\Gamma^b} \vec{n} \cdot \left(\vec{\mathbf{F}}_c(\mathbf{u}_h) - \vec{\mathbf{F}}_v(\mathbf{u}_h, \nabla \mathbf{u}_h) \right) ds \end{aligned} \quad (1)$$

The above equation involves gradients of test functions ϕ_i in the physical space $\vec{x} \in \mathbb{R}^d$ where d is the number of physical dimensions. The basis functions are normally defined in a reference space with $\vec{\xi} \in [-1, 1]^d$. One can define the physical coordinates \vec{x} using a mapping function $\vec{x}(\vec{\xi}) : [-1, 1]^d \mapsto \mathbb{R}^d$. This definition allows one to write the gradient in the physical space as

$$\nabla = \frac{\partial \vec{\xi}}{\partial \vec{x}} \cdot \nabla_{\vec{\xi}} \quad (2)$$

However, the relation $\vec{\xi}(\vec{x})$ is unknown. Therefore, one normally uses the following relationship

$$\begin{aligned} \frac{\partial \vec{x}}{\partial \vec{\xi}} &:= [J] \\ \frac{\partial \vec{x}}{\partial \vec{\xi}} &:= [J]^{-1} \end{aligned} \quad (3)$$

Therefore the gradient in the physical space can be written as

$$\nabla = [J]^{-1} \nabla_{\vec{\xi}} \quad (4)$$

Now consider the dot product of the gradient (i.e. divergence) with a vector $\vec{\mathbf{F}}$ is written as

$$\begin{aligned}\nabla \cdot \vec{\mathbf{F}} &= [J]^{-1} \nabla_{\xi} \cdot \vec{\mathbf{F}} \\ \nabla \cdot \vec{\mathbf{F}} &= \frac{\partial \xi_j}{\partial x_i} \frac{\partial}{\partial \xi_j} \mathbf{F}_i\end{aligned}\tag{5}$$

which can be re-arranged as

$$\nabla \cdot \vec{\mathbf{F}} = \frac{\partial}{\partial \xi_j} \frac{\partial \xi_j}{\partial x_i} \mathbf{F}_i\tag{6}$$

If one defines a new flux-vector $\mathbf{E}_j = \frac{\partial \xi_j}{\partial x_i} \mathbf{F}_i$. The equation simply can be written as

$$\mathbf{E}_j = n_i \mathbf{F}_i\tag{7}$$

When programing this method one uses the above equation and sets $n_i = \frac{\partial \xi_j}{\partial x_i}$ for a particular \mathbf{E}_j . Now considering that the divergence operator is just the dot product of the gradient with a vector one can re-write many of the operation in Eq. (1) in the spirit of the above manipulations. So Eq. (1) can be re-written as

$$\begin{aligned}& \sum_{e \in \mathcal{T}_h} \int_{\Omega_k} \phi_i \frac{\partial \mathbf{u}_h}{\partial t} d\Omega_e - \sum_{e \in \mathcal{T}_h} \int_{\Omega_k} \nabla_{\xi} \phi_i \cdot \left(\vec{\mathbf{E}}_c(\mathbf{u}_h) - \vec{\mathbf{E}}_v(\mathbf{u}_h, \nabla \mathbf{u}_h) \right) + \phi_i \mathbf{S}(\mathbf{u}_h, \nabla \mathbf{u}_h) d\Omega_e + \\& \sum_{e \in \mathcal{T}_h} \int_{\Omega_k} \nabla_{\xi} \phi_i \cdot \frac{\partial \vec{\mathbf{E}}_c(\mathbf{u}_h)}{\partial \mathbf{u}_h} [\tau] \left(\frac{\partial \mathbf{u}_h}{\partial t} + \nabla_{\xi} \cdot \left(\vec{\mathbf{E}}_c(\mathbf{u}_h) - \vec{\mathbf{E}}_v(\mathbf{u}_h, \nabla \mathbf{u}_h) \right) + \mathbf{S}(\mathbf{u}_h, \nabla \mathbf{u}_h) \right) d\Omega_e + \\& \sum_{b \in \mathcal{B}_h} \int_{\Gamma^b} \vec{n} \cdot \left(\vec{\mathbf{F}}_c(\mathbf{u}_h) - \vec{\mathbf{F}}_v(\mathbf{u}_h, \nabla \mathbf{u}_h) \right) ds\end{aligned}\tag{8}$$

Examining equation indicates how many functions are required to compute the SU/PG residual. In practice one needs only a few functions that when used in the proper sequence can generate the SU/PG residual. To simplify the concepts we will write the discretization in terms a general fluxes \mathbf{F} and \mathbf{E} both of which can in general depend on \mathbf{u}_h and $\nabla \mathbf{u}_h$.

$$\begin{aligned}& \sum_{e \in \mathcal{T}_h} \int_{\Omega_k} \phi_i \frac{\partial \mathbf{u}_h}{\partial t} d\Omega_e - \sum_{e \in \mathcal{T}_h} \int_{\Omega_k} \nabla_{\xi} \phi_i \cdot \vec{\mathbf{E}} + \phi_i \mathbf{S}(\mathbf{u}_h, \nabla \mathbf{u}_h) d\Omega_e + \\& \sum_{e \in \mathcal{T}_h} \int_{\Omega_k} \nabla_{\xi} \phi_i \cdot \frac{\partial \vec{\mathbf{E}}_c(\mathbf{u}_h)}{\partial \mathbf{u}_h} [\tau] \left(\frac{\partial \mathbf{u}_h}{\partial t} + \nabla_{\xi} \cdot \vec{\mathbf{E}} + \mathbf{S}(\mathbf{u}_h, \nabla \mathbf{u}_h) \right) d\Omega_e + \\& \sum_{b \in \mathcal{B}_h} \int_{\Gamma^b} \vec{n} \cdot \vec{\mathbf{F}} ds\end{aligned}\tag{9}$$

Due to the non-linearity of the fluxes the methods of computing the divergence is as follows

$$\nabla_{\xi} \cdot \vec{\mathbf{E}} = \frac{\partial \vec{\mathbf{E}}}{\partial \mathbf{u}_h} \cdot \nabla_{\xi} \mathbf{u}_h +\tag{10}$$

If one carefully examines these equations one can deduce that all operations required by the SU/PG method effectively become computations of the form

$$\mathbf{E}_j = n_i \mathbf{F}_i \quad (11)$$

and operations of the flux jacobian of \mathbf{E} on a vector \mathbf{V} over the number of equations.

$$\frac{\partial \mathbf{E}}{\partial \mathbf{u}_h} \cdot \mathbf{V} = n_i \frac{\partial \mathbf{F}_i}{\partial \mathbf{u}_h} \cdot \mathbf{V} \quad (12)$$

Finally a method is required to compute the product $[\tau] \mathbf{V}$. However, since an explicit expression is only available for $[\tau]^{-1}$. Therefore, $[\tau]$ is never explicitly formed rather its product onto a vector is formed by recalling

$$[\tau] \mathbf{V} = \left([\tau]^{-1}\right)^{-1} \mathbf{V} = \mathbf{W} \quad (13)$$

which is simply the solution of

$$[\tau]^{-1} \mathbf{W} = \mathbf{V}; \quad (14)$$

Thus the product of $[\tau] \mathbf{V}$ is implemented as a linear solve operation. The SU/PG residual is formed using the following algorithm

Algorithm 1 :SU/PG Residual Formation Algorithm

```

R(:) = 0
for qp = 0; qp < nqp; qp++ do
   $\nabla \cdot \mathbf{F} = 0$ 
  for j = 0; j < d; j++ do
    Form  $\vec{n} = [J(:, d)]^{-1}$ 
    Compute  $\mathbf{E}_j = \vec{\mathbf{F}} \cdot \vec{n}$ 
    for i = 0; i < NDOF; i++ do
       $\mathbf{R}_i + = \frac{\partial \phi_i}{\partial \xi_j} \mathbf{E}_j w_q(qp) \text{Det}(J)$ 
    end for
    Compute  $\nabla \cdot \vec{\mathbf{F}} + = \frac{\partial \mathbf{E}_j}{\partial \mathbf{u}_h} \frac{\partial \mathbf{u}_h}{\partial \xi_j}$ 
  end for
  Compute  $[\tau]^{-1}$ 
  Compute  $[\tau] \nabla \cdot \vec{\mathbf{F}}$  via solving  $[\tau]^{-1} \mathbf{x} = \nabla \cdot \vec{\mathbf{F}}$ 
  for j = 0; j < d; j++ do
    Form  $\vec{n} = [J(:, d)]^{-1}$ 
    Compute  $D = \frac{\partial \mathbf{E}_j}{\partial \mathbf{u}_h} \mathbf{x}$ 
    for i = 0; i < NDOF; i++ do
       $\mathbf{R}_i + = \frac{\partial \phi_i}{\partial \xi_j} D w_q(qp) \text{Det}(J)$ 
    end for
  end for
end for

```

Examination of the algorithm shows that one requires only 3 functions:

1. Compute \mathbf{E}_j
2. Compute $\frac{\partial \mathbf{E}_j}{\partial \mathbf{u}_h} \cdot \mathbf{y}$
3. Compute $[\tau]$.

While it may seem that there is a missing function for inverting the τ matrix, this functionality has been provided as part of the square matrix class.