Postulatus of BM

- 1) Every state is disoubted by a wave function Y(t)
- 2) Evening obserable is discribed by a Hermitian oferator A A = A[†]
- 3) The regenvalues of the operators correspond to the ratures out comes of experimental measure ment.
- 4) Time evolution of the wave femlen is desurbed by the Sets of deeper quatern

T (2)/-Proof. 1) All Hermeelean matrix has real ugen values. 2) If two oferators committee They have set of agen voutors 3) The eager rectors of a Hermtian operator form an complete orthonormal basis- $A|\phi\rangle = a_1|\phi_1\rangle$

 $A|\phi_{1}\rangle = a_{1}|\phi_{1}\rangle$ $A|\phi_{2}\rangle = a_{2}|\phi_{2}\rangle$ $= \lambda \left(\phi_{1}|A|\phi_{2}\right) = a_{2}\langle\phi_{1}|\phi_{2}\rangle$ $= \lambda \left(\phi_{1}|\phi_{2}\right) = a_{2}\langle\phi_{1}|\phi_{2}\rangle$ $= \lambda \left(\phi_{1}|\phi_{2}|\phi_{2}\right)$ $= \lambda \left(\phi_{1}|\phi_{2}|\phi_{2}|\phi_{2}\right)$ $= \lambda \left(\phi_{1}|\phi_{2}|\phi_{2}|\phi_{2}\right)$ $= \lambda \left(\phi_{1}|\phi_{2}|\phi_{2}|\phi_{2}\right)$ $= \lambda \left(\phi_{1}|\phi_{2}|\phi_{2}|\phi_{2}\right)$ $= \lambda \left(\phi_{1}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}\right)$ $= \lambda \left(\phi_{1}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}|\phi_{2}$

a / & 1 d \

 $= \left\langle \frac{1}{4}, \frac{1}{4}, \frac{4}{2} \right\rangle = \frac{1}{4} \left\langle \frac{1}{4}, \frac{1}{4}, \frac{1}{4}, \frac{1}{4} \right\rangle$ $(a_2 - a_1) \langle \theta_1 | \theta_2 \rangle = 0$ Hen if $a_2 \neq a_p \langle \theta_1 | \theta_2 \rangle = 0$ => (d, 1d2) are normal (pn 1 dm) = Snm. or the normal bassis J(Pn) dx < \ > Quare inlegrable Gran-Schmedl orthonormalization 4) Table I f the shall Y is expanded as a linear combination of the cegurectors of A Pn y= 2 an on hanh at getter value In: [an]

$$= \langle A \rangle = \sum_{n} (a_n)^n \lambda_n$$

$$= \langle A \rangle = \langle \Psi | A \Psi \rangle$$

Two level atom

Eigen values
$$E_1$$
: μM

$$E_1 = \mu M$$

$$E_1 = \mu M$$

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$$E_4 = \mu M$$

$$E_1 = \mu M$$

$$E_2 = \mu M$$

$$E_3 = \mu M$$

$$E_4 = \mu M$$

$$E_5 = \mu M$$

$$E_7 = \mu M$$

$$E_$$

(4/H/41 =

$$\langle H \rangle = \frac{1}{2} \mu M - \frac{1}{2} \mu M = 0$$

$$\langle H^{\gamma} \rangle = \frac{1}{2} \mu^{\gamma} M^{\gamma} + \frac{1}{2} \mu^{\gamma} M^{\gamma}$$

(Y(t)) = a, e | E, t | P)

Lence the vetor remains fine

Same. Hence, the eigenstalis of

Hamiltonian are called stationary

State, since the probabily doesnot

evolve with time.

Momentum oferator

$$\hat{p}$$
 ψ : \hat{p} ψ = Ae^{tlkx}

$$k = \frac{1}{t}$$

Commitation with x

$$\left[\hat{\chi},\hat{\rho}\right]\Psi = -\chi + \ln\frac{2\Psi}{2\chi} + \ln\frac{2}{2\chi}(\chi)$$

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. 2Y

Time evolution of an operator

$$\frac{d\langle A \rangle}{dt} = \frac{d}{dt} \left(\frac{1}{t} \left[A, H \right] \right)$$

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$$\frac{d\langle A \rangle}{dt} = \frac$$

dt in Lind III $[x^{\nu}, b] = 2ih \mathcal{R}$ =) d(b) = - K(2e). =) Buntaen to along average values follow the chasical qualen of motion $\frac{d\langle x\rangle}{dt} = \left(\frac{1}{t} \left[x, H\right] = \left(\frac{1}{t} \left[x, \frac{b^{r}}{2m}\right]\right)$ $= \frac{1}{(t+2m)} \left[\left(x, b^{\gamma} \right) \right]$ $\left[\begin{array}{c} \left[\times, \phi^{\gamma} \right] \end{array} \right] = \left[\begin{array}{c} \left[\times, \phi^{\gamma} \right] \\ \left[\times, \phi^{\gamma} \right] \end{array} \right] = \left[\begin{array}{c} \left[\times, \phi^{\gamma} \right] \\ \left[\times, \phi^{\gamma} \right] \end{array} \right] = \left[\begin{array}{c} \left[\times, \phi^{\gamma} \right] \\ \left[\times, \phi^{\gamma} \right] \end{array} \right] = \left[\begin{array}{c} \left[\times, \phi^{\gamma} \right] \\ \left[\times, \phi^{\gamma} \right] \end{array} \right] = \left[\begin{array}{c} \left[\times, \phi^{\gamma} \right] \\ \left[\times, \phi^{\gamma} \right] 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$$= \left[\begin{array}{c} X, p \end{array} \right] p + \left[\begin{array}{c} x, p \end{array} \right]$$

$$= \left[\begin{array}{c} 1 \\ 2 \end{array} \right] p$$

$$= \left[\begin{array}{c} 1 \\ 1 \end{array} \right]$$

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