

Private Notes (1) 10/10/04

## Coolie's guide to partial differential Equations

→ The Heat Equation →

$$\frac{dT}{dt} = \alpha \nabla^2 T$$

Imagine you have a piece of metal



→ How is the heat distributed across it at any one moment → what's the temperature at any one point

→ we want to know how this distribution will change over time → as heat flows from warmer spots to cooler ones

The 1D Example →

lets say you have 2 different rods at 2 different temperatures along the x-axis.

→ now lets bring these two rods into contact,  
→ the heat will flow from the hot one to the cold one and over time the temperature will be  $\neq$  equal.

→ but how? what will the temperature look like at each point in time → can we do the maths modelling for it? → math model

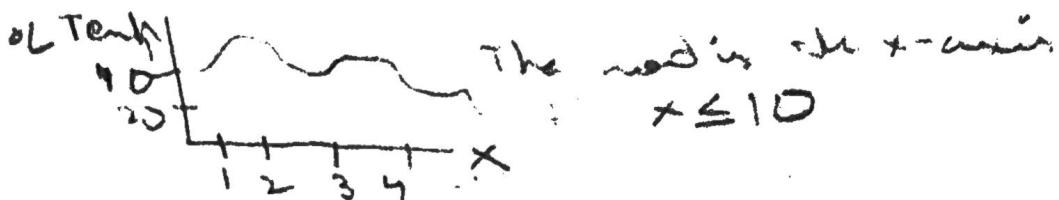
$$\frac{\partial T(x,t)}{\partial t} = \alpha \cdot \frac{\partial^2 T(x,t)}{\partial x^2}$$

# Mathematical Modelling

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## Building up the Heat Equation

Currently we are imagining a rod along the  $x$ -axis, with each point of that rod being given by the value on the  $x$ -axis



→ The temperature is a function of this rod so  $T(x)$

However the temperature is changing over time, so really we have another dimension to consider



→ since we have an entire function changing over time and the function has multiple inputs, there are more than 1 order of change happening.

$$\frac{\partial T}{\partial x} \quad \begin{cases} \text{temperature change with respect to} \\ x \end{cases}$$

$$\frac{\partial T}{\partial t} \quad \begin{cases} \text{temperature change with respect to} \\ t \end{cases}$$

→ so actually our heat transfer is controlled by

$$\frac{\partial T}{\partial t} \quad \begin{cases} \text{small change to temperature with small} \\ \text{change in time} \end{cases}$$

AND

$$\frac{\partial T}{\partial x} \quad \begin{cases} \text{small change to temperature with small} \\ \text{change along x spatial dimension} \end{cases}$$

→ Therefore, the change in the time dimension depends upon how the function changes with respect to time.

So, the modelling of how the heat of things with time affects on how it changes with space, if they are proportional.

$$\frac{\partial T}{\partial t} = k \cdot \frac{\partial^2 T}{\partial x^2} \rightarrow \text{Change with respect to } t \text{ is proportional to the second partial derivative with respect to } x$$

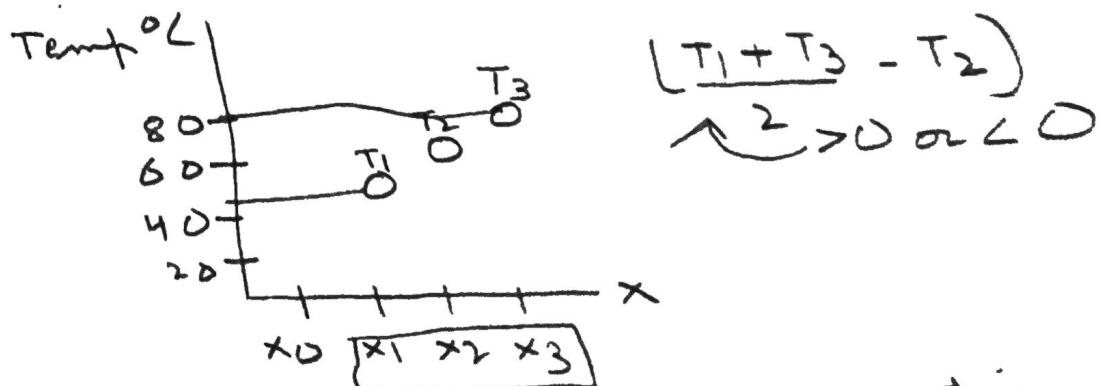
Partial differential equation

→ with just 2 inputs, 1 for  $x$ , and 1 for  $t$  in the 1D PDE

$$\frac{\partial T}{\partial t}(x, t) = k \cdot \frac{\partial^2 T}{\partial x^2}(x, t)$$

→ Where does this come from?

Let's look at a discrete number of things



→ to model how  $T_2$  changes, we can average out its neighbouring particles  $\rightarrow (T_1 + T_3 - 2T_2) \rightarrow$

$> 0 \rightarrow T_2$  will heat up Proportional to ~~difference~~

$< 0 \rightarrow T_2$  will cool down Proportional to difference

$\therefore \frac{\partial T_2}{\partial t} = \alpha \left( \frac{T_1 + T_3}{2} - T_2 \right)$

$$\begin{aligned}
 &= \alpha \left( \frac{T_1 + T_3}{2} - \frac{2T_2}{2} \right) \xrightarrow{(4)} \frac{\alpha}{2} [T_1 + T_3 - 2T_2] \\
 &\rightarrow \frac{\alpha}{2} [(T_3 - T_2) + (T_1 - T_2)] \\
 \frac{dT_2}{dt} &\rightarrow \frac{\alpha}{2} [T_3 - T_2] - [T_2 - T_1] \\
 &\quad \Delta T_2 - \Delta T_1
 \end{aligned}$$

$\Rightarrow \Delta T_2 - \Delta T_1 > 0 \Rightarrow T_2$  heats up  
 $\Rightarrow \Delta T_2 - \Delta T_1 < 0 \Rightarrow T_2$  cools down

$\Delta T_2$  → since the difference of avg's  
 in greater,  $T_2$  will change

$$\begin{aligned}
 &\Delta T_1 \xrightarrow{x_1+x_2+x_3 \text{ therefore}} \\
 &\rightarrow \Delta T_2 - \Delta T_1 = \Delta \Delta T_1 \rightarrow \frac{dT_2}{dt} = \frac{\alpha}{2} (\Delta \Delta T_1) \\
 &\qquad\qquad\qquad \text{second difference} \downarrow \qquad \text{now this is} \\
 &\rightarrow \frac{dT_2}{dt} = \frac{\alpha}{2} \Delta \Delta T_1 \rightarrow \text{discrete very similar to the} \\
 &\text{so when we go from discrete to continuous} \\
 &\text{this turns into } \frac{dT}{dt} = k \cdot \frac{d^2T}{dx^2}
 \end{aligned}$$

where  $k$  could be const  $\xrightarrow{\text{the second difference}}$   
 any proportional constant is defined by the  
 between any particles second derivative

$\rightarrow$  so instead of calculating temperature avg's  
 between two fixed points  $x_1, x_2 \rightarrow$  we are  
 now shrinking that step to a "very" small char

$\rightarrow$  so how an arbitrary point in  $x'$ 's dx  
 will change depending on the temperature of  
 the infinitely small  $dx$ 's adjacent to it.

Therefore we have now built up (5)

$$\boxed{\frac{dT}{dt} = \alpha \cdot \frac{\partial^2 T}{\partial x^2}}$$

The heat equation in 1D

$$\text{Representation: } P.5 \text{ for } \frac{\partial T}{\partial t} + \frac{\partial^2 T}{\partial x^2} = \alpha \nabla^2 T$$

3 blue 1 brown is "But what is a Partial Differential Equation?" YouTube.com

ROUGH NOTES, ONLY FOR UNDERSTANDING

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→ How does the temperature distribution from a 1D rod change over time? → From the Equations

$$\frac{\partial T}{\partial t}(x, t) = k \cdot \frac{\partial^2 T}{\partial x^2}(x, t)$$

What this means: the rate at which temperature at a given point changes over time, depends on the second derivative of that temperature at that point with respect to space.

• 3 constraints  $T(x, t)$  must satisfy

1) PDE  $\rightarrow \frac{\partial T}{\partial t}(x, t) = k \cdot \frac{\partial^2 T}{\partial x^2}(x, t)$

2) Boundary condition

3) Initial condition

→ Let's consider  $T(x, 0) = \sin(x) \rightarrow \frac{\partial T}{\partial x} = \cos(x) \rightarrow \frac{\partial^2 T}{\partial x^2} = -\sin(x)$

$\therefore \frac{\partial T}{\partial t}(x, 0) = -k \cdot \underbrace{T(x, 0)}_{\sin(x)} \rightarrow \frac{\partial T}{\partial t}(x, 0) = -k \cdot \sin(x)$

$$T(x, 5\Delta t) = C \cdot \sin(x)$$

$$T' = -kT$$

$$\int \frac{T'}{T} dt = \int -k dt$$

$$\ln T = -kt \rightarrow T = \underbrace{C e^{-kt}}_{\text{initial temperature}} \rightarrow \text{This is Newton law of cooling}$$

→ Then extrapolating this logic to our PDE

$$T(x, t) = \underbrace{\sin(x)}_{\text{initial distribution}} e^{-kt}$$

$$\frac{\partial T}{\partial t} = \underbrace{-k \sin(x) e^{-kt}}_{\frac{\partial T}{\partial t}} = k \cdot \underbrace{-\sin(x) e^{-kt}}_{\frac{\partial^2 T}{\partial x^2}} \rightarrow \text{satisfies the PDE}$$

The function must also satisfy the boundary condition  
 $\rightarrow$  since there is no heat transfer from the edges (1)  
of the rod i.e. the rate of change is 0 at  $t > 0$ .

$$(2) \rightarrow \frac{d\bar{T}}{dx}(0, t) = \frac{d\bar{T}}{dx}(L, t) = 0 \quad \forall t > 0$$

our function which satisfies the PDE  
 $T(x, t) = \sin(\omega t)$

the new change our function is  $\sin(\omega t)$  and it will  
still satisfy the PDE.

$$-\kappa \omega \cos(\omega t) \bar{c}^{'''} = \kappa \omega \sin(\omega t)$$

$$\text{and now } \frac{d\bar{T}}{dx} = -\sin(\omega t) \bar{c}''' = 0 \text{ at } x = 0$$

However for it to satisfy the equation on  
the right hand side we need to change the  
frequency  $\omega$ . i.e. the function is flat at  $x = L$   
means

$\rightarrow$  This frequency will be  $\frac{\pi}{L}$  (or the first harmonic)

$\rightarrow$  the boundary is satisfied by  $\kappa \frac{\pi}{L}$ .

$\rightarrow$  therefore now we have an infinite set  
of functions  $T(x, t)$  that satisfy both (1)  
and (2)

$$T(x, t) = \cos\left(\frac{\pi}{L}(\frac{x}{L})\right) \bar{c}''\left(\frac{\pi}{L}t\right)$$

$$\forall n \in \mathbb{N}, n \neq 0 \quad \forall t > 0$$

$\rightarrow$  The heat equation is also linear i.e.  
for any two solutions  $T_1, T_2 \rightarrow T_3 = T_1 + T_2$  is  
also a solution

$$T(x, t) = \sum_{n=1}^{\infty} A_n \cos\left(\frac{n\pi x}{L}\right) \bar{c}''\left(\frac{n\pi}{L}t\right)$$

solves the PDE //