

Mathematics and numerics for data assimilation and state estimation – Lecture 20



Summer semester 2020

Overview

- 1 The Fokker-Planck equation
- 2 Numerical integration of SDE
- 3 Filtering problems with SDE dynamics
- 4 Examples using Euler–Maruyama integration
- 5 Model error and model fitting

Summary lecture 20

- Itô integrals and theory of stochastic differential equations (SDE)

$$V_t = V_0 + \int_0^t b(V_s)ds + \int_0^t \sigma(V_s)dW_s$$

- Plan for today: Fokker-Planck equation, numerical integration of SDE and applications in filtering problems.

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The kernel density for SDE

Our plan is to study filtering problems

$$V_{j+1} = \Psi(V_j) := V_j + \int_0^1 b(V_{j+s})ds + \int_0^1 \sigma(V_{j+s})dW_s^{(j)}$$

$$Y_{j+1} = h(V_{j+1}) + \eta_{j+1}$$

where $W^{(j)}$ are independent Wiener processes.

The Bayes filter for this problem takes the form

$$\pi(v_{j+1}|y_{1:j+1}) \propto \pi(y_{j+1}|v_{j+1}) \int_{\mathbb{R}^d} \pi(v_{j+1}|v_j) \pi(v_j|y_{1:j}) dv_j$$

with $\pi_{V_{j+1}|V_j}(x|y)$ equal to the kernel density for $t \in (0, 1]$,

$$\rho(t, x|y) = \frac{\mathbb{P}(V_{j+t} \in dx | V_j \in dy)}{dx} = \frac{\mathbb{P}(V_t \in dx | V_0 \in dy)}{dx}$$

(due to the time-independent coefficients the SDE is stationary).

The density of an SDE

Consider the 1D SDE

$$dX_t = b(X_t)dt + \sigma(X_t)dW_t, \quad X_0 \sim p(0, x)$$

and assume that the density $p(t, x) = \mathbb{P}(X_t \in dx)/dx$ exists for any $t > 0$.

Recall that for any $f \in C_c^2(\mathbb{R})$ (mapping with compact support),

$$df(X_t) = f'(X_t)dX_t + \frac{1}{2}f''(X_t)(dX_t)^2 = (f'b + \sigma^2/2f'')dt + f'\sigma dW_t$$

By integration,

$$f(X_t) - f(X_0) = \int_0^t (bf' + \frac{\sigma^2}{2}f'')(X_s)ds + \int_0^t (f'\sigma)(X_s)dW_s.$$

Recalling that Itô integrals have mean-zero,

$$\mathbb{E}[f(X_t) - f(X_0)] = \int_0^t \mathbb{E}\left[(bf' + \frac{\sigma^2}{2}f'')(X_s)\right] ds$$

Note: expectation is here wrt the density $p(s, x)$

Fokker-Planck equation

$$\int_{\mathbb{R}} f(x)(p(t, x) - p(0, x))dx = \int_0^t \int_{\mathbb{R}} \left[b(x)f'(x) + \sigma^2(x)\frac{f''(x)}{2} \right] p(s, x) dx ds$$

Integration by parts, using the compact support of f (and its derivatives), we obtain

$$\begin{aligned} & \int_0^t \int_{\mathbb{R}} f(x) p_t(s, x) dx ds \\ &= \int_0^t \int_{\mathbb{R}} f(x) \left[-\partial_x (b(s)p(s, x)) + \partial_{xx} \left(\frac{\sigma^2(x)}{2} p(s, x) \right) \right] dx ds \quad \forall f \in C_c^2(\mathbb{R}) \end{aligned}$$

Conclusion: The density $p(t, x) = \mathbb{P}(X_t \in dx)/dx$ must satisfy the **Fokker-Planck PDE**

$$p_t = \partial_x(-bp) + \partial_{xx}\left(\frac{\sigma^2}{2}p\right) \quad (t, x) \in [0, T] \times \mathbb{R} \quad (1)$$

$$p(t, x)|_{t=0} = p(0, x).$$

If the SDE coefficients are sufficiently smooth and $\sigma > 0$, then (1) is well-posed and a unique classical solution exists for all $t > 0$.

Fokker-Planck for kernel densities

The PDE extends to kernel densities $p(t, x|y) = \mathbb{P}(X_t \in dx|y \in dy)/dx$:

$$\begin{aligned} p_t(\cdot, \cdot|y) &= \partial_x(-bp(\cdot, \cdot|y)) + \partial_{xx}\left(\frac{\sigma^2}{2}p(\cdot, \cdot|y)\right) \quad (t, x) \in [0, T] \times \mathbb{R} \\ p(0, x|y) &= \delta_y(x). \end{aligned} \quad (2)$$

Remarks: The operator

$$(\mathcal{L}^*p)(x) := \partial_x(-bp)(x) + \partial_{xx}\left(\frac{\sigma^2}{2}p\right)(x)$$

may be associated to the transition function of Markov chains (here denoted P):

$$p(t + \Delta t, x) \approx p(t, x) + \Delta t(\mathcal{L}^*p)(x),$$

vs

$$\pi_i^{n+1} = \sum_{j=1}^N P_{ji} \pi_j^n = \pi_i^n + \left((P - I)^T \pi^n \right)_i$$

And just like Markov chains, SDE may have stationary distributions:

$$\mathcal{L}^*p = 0 \iff p \text{ stationary}, \quad (P - I)^T \pi = 0 \iff \pi \text{ stationary}.$$

Application in filtering

Returning to the filtering problem

$$V_{j+1} = \Psi(V_j) := V_j + \int_0^1 b(V_{j+s})ds + \int_0^1 \sigma(V_{j+s})dW_s^{(j)}$$

$$Y_{j+1} = h(V_{j+1}) + \eta_{j+1}$$

the iterative Bayes filter equation

$$\pi(v_{j+1}|y_{1:j+1}) \propto \pi(y_{j+1}|v_{j+1})\pi(v_{j+1}|y_{1:j})$$

can be written

$$\pi(v_{j+1}|y_{1:j+1}) \propto \pi(y_{j+1}|v_{j+1})p(1, v_{j+1})$$

where p solves

$$\begin{aligned} p_t &= \mathcal{L}^* p & (t, x) &\in [0, T] \times \mathbb{R} \\ p(t, x)|_{t=0} &= \pi_{V_j|Y_{1:j}}(x|y_{1:j}) \end{aligned}$$

Conclusion: In principle we can solve these filtering problems exactly!

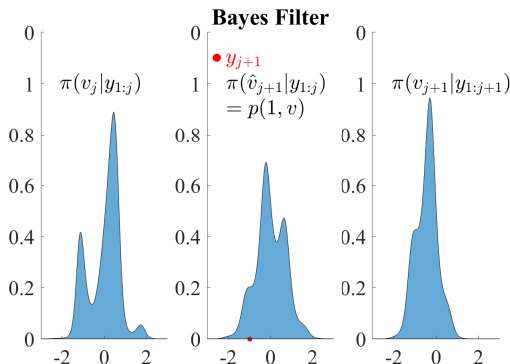
Example

Filtering problem:

$$V_{j+1} = V_j + \int_0^1 U'(V_{j+s}) ds + \int_0^1 dW_s^{(j)}$$

$$Y_{j+1} = V_{j+1} + \eta_{j+1}$$

with $U(x) = x^2/2 + 0.15 \sin(2\pi x)$ and for some j , we have set $\pi(v_j | y_{1:j}) \propto \exp(-2U(v_j) + \sin(4v_j))$.



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Euler–Maruyama scheme

For the SDE

$$dX_t = b(X_t)dt + \sigma(X_t)dW_t \quad t \in [0, T], \quad X_t|_{t=0} = X_0,$$

the Euler-Maruyama scheme on a uniform mesh $t_j = j\Delta t$

$$\bar{X}_{t_{j+1}} = \bar{X}_{t_j} + b(\bar{X}_{t_j})\Delta t + \sigma(\bar{X}_{t_j})\Delta W_j$$

where $\Delta W_j = W_{t_{j+1}} - W_{t_j}$ and $\bar{X}_0 = X_0$.

Motivation:

$$\begin{aligned} X_{t_{j+1}} - X_{t_j} &= \int_{t_j}^{t_{j+1}} b(X_t)dt + \int_{t_j}^{t_{j+1}} \sigma(X_t)dW_t \\ &\approx \int_{t_j}^{t_{j+1}} b(X_{t_j})dt + \int_{t_j}^{t_{j+1}} \sigma(X_{t_j})dW_t \end{aligned}$$

Let $\bar{X}_t := \text{LinInterp}(t; \{(t_j, \bar{X}_{t_j})\}_{j=0}^{T/\Delta t})$.

Strong convergence rate for Euler–Maruyama

Under the regularity assumptions in Thm 4, Lecture 19, most importantly

$$\begin{aligned} |b(x) - b(y)| + |\sigma(x) - \sigma(y)| &\leq K|x - y| \\ |b(x)|^2 + |\sigma(x)|^2 &\leq K(1 + |x|^2), \end{aligned}$$

the Euler–Maruyama method converges strongly with rate $1/2$.

$$\sqrt{\max_{t \in [0, T]} \mathbb{E} [|\bar{X}_t - X_t|^2]} \leq C \Delta t^{1/2}$$

for some $C > 0$.

Weak convergence rate Euler–Maruyama

Under more restrictive regularity conditions, the Euler–Maruyama converges weakly with rate 1.

$$\max_{t \in [0, T]} |\mathbb{E} [f(\bar{X}_t) - f(X_t)]| \leq C_f \Delta t$$

for any mapping $f \in C_p^\infty(\mathbb{R}^d, \mathbb{R})$ with $C_f > 0$ depending on f .¹

Remark: See [ELV-E 7, 8] for more on results in higher-dimensional state space, and on higher order numerical methods.

¹ C_p^∞ is set of functions with at most polynomial growth in any partial derivative: $|\partial_\alpha f(x)| \leq C_\alpha |x|^{p_\alpha}$ for any $\alpha \in \mathbb{N}^d$, some $p_\alpha \in \mathbb{N}$ and all $x \in \mathbb{R}^d$.

Example - geometric Brownian motion

Consider the SDE

$$dX_t = X_t dt + X_t dW, \quad X_0 = 1,$$

and let us approximate: $\mathbb{E}[X_1] = e^1$ [Ubung 9].

Monte Carlo strategy:

- 1 Fix $\Delta t = 1/N$ and generate M numerical solutions of the SDE $\bar{X}_1^{(i)}$ by the EM scheme

$$\bar{X}_{t_{j+1}}^{(i)} = X_{t_j} + \bar{X}_{t_j}^{(i)} \Delta t + \bar{X}_{t_j}^{(i)} \Delta W_j^{(i)}, \quad j = 0, 1, \dots, N-1,$$

with independent Wiener paths $W^{(i)}$ and $X_0^{(i)} = 1$.

- 2 And apply the Monte Carlo method:

$$\mathbb{E}[X_1] = E_M[\bar{X}_1^{(\cdot)}] = \frac{1}{M} \sum_{i=1}^M \bar{X}_1^{(i)}$$

Illustration of approximation

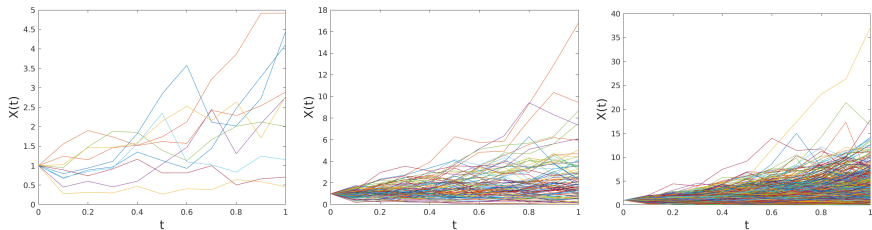


Figure: From left to right $(M, N) = (10, 10), (100, 10), (1000, 10)$.

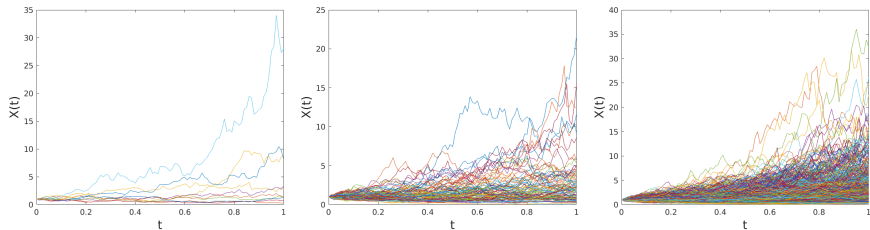


Figure: From left to right $(M, N) = (10, 100), (100, 100), (1000, 100)$

Approximation error

For any $f \in C_P(\mathbb{R}^d, \mathbb{R})$,

$$\begin{aligned}\mathbb{E} [(\mathbb{E}[f(X_1)] - E_M[f(\bar{X}_1)])^2] &= \mathbb{E} \left[\left(\mathbb{E}[f(X_1)] \pm \mathbb{E}[f(\bar{X}_1)] - E_M[f(\bar{X}_1)] \right)^2 \right] \\ &\leq \mathbb{E} [(\mathbb{E}[f(X_1)] - \mathbb{E}[f(\bar{X}_1)])^2] \\ &\quad + 2\mathbb{E} [(\mathbb{E}[f(X_1)] - \mathbb{E}[f(\bar{X}_1))](\mathbb{E}[f(\bar{X}_1)] - E_M[f(\bar{X}_1)])] \\ &\quad + \mathbb{E} [(\mathbb{E}[f(\bar{X}_1)] - E_M[f(\bar{X}_1)])^2] \\ &= (\mathbb{E}[f(X_1)] - \mathbb{E}[f(\bar{X}_1)])^2 + \mathbb{E} [(\mathbb{E}[f(\bar{X}_1)] - E_M[f(\bar{X}_1)])^2] \\ &\leq C(N^{-2} + M^{-1}).\end{aligned}$$

Computational cost of the Euler–Maruyama+Monte Carlo approach is of the order

$$Cost = M \times N.$$

Minimization of the error as a function of the cost:

$$M = \mathcal{O}(N^2) = \mathcal{O}(\Delta t^{-2})$$

Then

$$\|\mathbb{E}[f(X_1)] - E_M[f(\bar{X}_1)]\|_{L^2(\Omega)} = \mathcal{O}(\Delta t) = Cost^{-1/3}.$$

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Filtering problem

$$V_{\tau_{j+1}} = \Psi(V_{\tau_j}) := V_{\tau_j} + \int_0^{\Delta\tau} b(V_{\tau_j+t})dt + \int_0^{\Delta\tau} \sigma(V_{\tau_j+t})dW_t^{(j)}$$

$$Y_{\tau_{j+1}} = H(V_{\tau_{j+1}}) + \eta_{j+1}$$

where $\Delta\tau$ denotes the observation time interval, and $\eta_j \stackrel{iid}{\sim} N(0, \Gamma)$, $\Gamma > 0$ and $k \times k$ matrix, $\{W^{(j)}\}$ are independent Wiener processes and

$$V_0 \perp \{\eta_j\} \perp \{W^{(j)}\}$$

Shorthand notation: To align with previous notation we write $V_j := V_{\tau_j}$ and $Y_j := Y_{\tau_j}$.

Objective: Approximate the Bayes filter $\pi_{V_{\tau_j}|Y_{\tau_{1:j}}}$.

Exact model EnKF method

- 1 Set $\tau_j = 0$ and sample initial distribution $v_j^{(i)} \stackrel{iid}{\sim} \mathbb{P}_{V_0}$.

and for $j = 0, 1, \dots$:

- 2 **Prediction:** Simulate particles

$$\hat{v}_{\tau_{j+1}}^{(i)} = \Psi(v_{\tau_j}^{(i)}) = v_{\tau_j}^{(i)} + \int_0^{\Delta\tau} b(v_{\tau_j+t}^{(i)})dt + \int_0^{\Delta\tau} \sigma(v_{\tau_j+t}^{(i)})dW_t^{(j,i)},$$

for $i = 1, 2, \dots, M$, where $\{W^{(i,j)}\}_{i,j}$ are independent Wiener processes.

- 3 **Analysis:**

$$v_{\tau_{j+1}}^{(i)} = \hat{v}_{\tau_{j+1}}^{(i)} + K_{j+1}(y_{\tau_{j+1}}^{(i)} - h(\hat{v}_{\tau_{j+1}}^{(i)})),$$

for $i = 1, 2, \dots, M$ where

$$y_{\tau_{j+1}}^{(i)} = y_{\tau_{j+1}} + \eta_{j+1}^{(i)}, \quad \eta_{j+1}^{(i)} \stackrel{iid}{\sim} N(0, \Gamma)$$

and

$$K_{j+1} = \text{Cov}_M[v_{\tau_{j+1}}^{(\cdot)}, h(v_{\tau_{j+1}}^{(\cdot)})](\text{Cov}_M[h(v_{\tau_{j+1}}^{(\cdot)})] + \Gamma)^{-1}.$$

Problem: In many cases Ψ must be approximated by numerical integration.

Artificial example

Consider the Ornstein-Uhlenbeck process

$$\Psi(V_{\tau_j}) = V_{\tau_j} - \int_0^{\Delta\tau} \theta V_{\tau_j+t} dt + \int_0^{\Delta\tau} dW_s \stackrel{D}{=} e^{-\theta\Delta\tau} V_{\tau_j}$$

with $\theta, \sigma > 0$.

We can solve this exactly:

$$\Psi(V_{\tau_j}) \stackrel{D}{=} AV_{\tau_j} + \xi_j$$

where $\xi_j \sim N(0, \Sigma_{\Delta\tau})$ and we are in the familiar the linear-Gaussian setting. [see Übung 9]

Approximation of the stochastic integrator

Let Ψ^N be the Euler-Maruyama approximation of

$$\Psi(V_{\tau_j}) = V_{\tau_j} + \int_0^{\Delta\tau} b(V_{\tau_j+t})dt + \int_0^{\Delta\tau} \sigma(V_{\tau_j+t})dW_t^{(j)}$$

using a uniform timestep $\Delta t = \Delta\tau/N$.

$\bar{V}_{\tau_{j+1}} = \Psi^N(\bar{V}_{\tau_j})$ is computed as follows

1 Input: \bar{V}_{τ_j} .

2 For $k = 0 : N - 1$, compute

$$\begin{aligned}\bar{V}_{\tau_j+(k+1)\Delta t} &= \bar{V}_{\tau_j+k\Delta t} + b(\bar{V}_{\tau_j+k\Delta t})\Delta t \\ &\quad + \sigma(\bar{V}_{\tau_j+k\Delta t})\left(W_{\tau_j+(k+1)\Delta t} - W_{\tau_j+k\Delta t}\right)\end{aligned}$$

3 Output: $\bar{V}_{\tau_{j+1}} = \bar{V}_{\tau_j+N\Delta t}$

EnKF method using a numerical integrator

- 1 Set $\tau_j = 0$ and sample initial distribution $v_j^{(i)} \stackrel{iid}{\sim} \mathbb{P}_{V_0}$.

and for $j = 0, 1, \dots$:

- 2 **Prediction:** Simulate particles

$$\hat{v}_{\tau_{j+1}}^{(i)} = \Psi^N(v_{\tau_j}^{(i)}),$$

for $i = 1, 2, \dots, M$, where $\{W^{(i,j)}\}_{i,j}$ are independent Wiener processes used in Ψ^N .

- 3 **Analysis:**

$$v_{\tau_{j+1}}^{(i)} = \hat{v}_{\tau_{j+1}}^{(i)} + K_{j+1}(y_{\tau_{j+1}}^{(i)} - h(\hat{v}_{\tau_{j+1}}^{(i)})),$$

for $i = 1, 2, \dots, M$ where

$$y_{\tau_{j+1}}^{(i)} = y_{\tau_{j+1}} + \eta_{j+1}^{(i)}, \quad \eta_{j+1}^{(i)} \stackrel{iid}{\sim} N(0, \Gamma)$$

and

$$K_{j+1} = \text{Cov}_M[\hat{v}_{\tau_{j+1}}^{(\cdot)}, h(\hat{v}_{\tau_{j+1}}^{(\cdot)})](\text{Cov}_M[h(\hat{v}_{\tau_{j+1}}^{(\cdot)})] + \Gamma)^{-1}.$$

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Ornstein-Uhlenbeck process

$$V_{\tau_{j+1}} = \Psi(V_{\tau_j}) := V_{\tau_j} + -\frac{1}{4} \int_0^{\Delta\tau} V_{\tau_j+t} dt + \frac{1}{4} \int_0^{\Delta\tau} dW_t^{(j)}$$

$$Y_{\tau_{j+1}} = V_{\tau_{j+1}} + \eta_{j+1}$$

with $V_0 = 1$, $\Delta\tau = 1/2$ and $\eta_j \sim N(0, \Gamma)$.

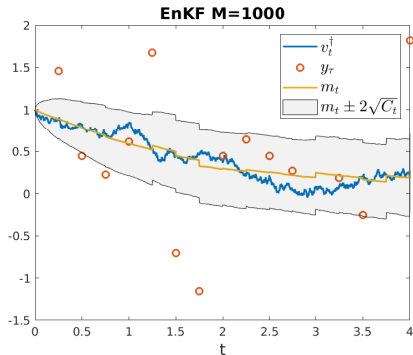
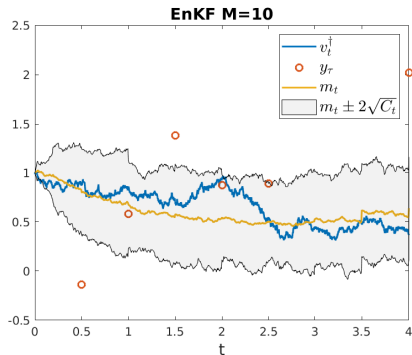
We generate an observation sequence for $y_{\tau_{1:J}}$ for $J = 100$ from synthetic data $v_{\tau_{1:J}}^\dagger$.

Approximation method: EnKF with numerical integrator Ψ^N with $N = 100$ and $\Delta t = \Delta\tau/N$.

Note: Continuum dynamics makes it possible to also estimate the filtering distribution for times between observation times.

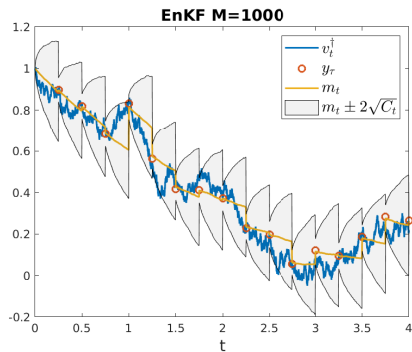
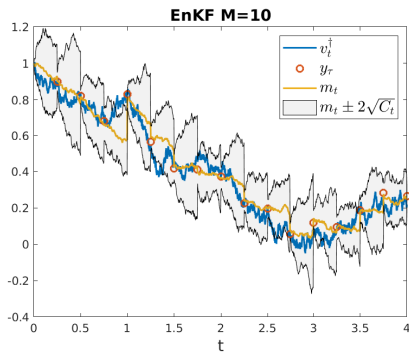
Numerical results

Large uncertainty in observations, $\Gamma = 1$, yields small correction at observation times:



Numerical results

Small uncertainty in observations, $\Gamma = 1/1000$, yields small correction at observation times:



Langevin equation

$$dX_t = V_t dt$$

$$dV_t = (-0.25V_t - U'(X_t))dt + 0.5dW_t$$

with $(X_0, V_0) = (0, 1)$.

Observations

$$Y_{\tau_k} = V_{\tau_k} + \eta_k, \quad k = 1, 2, \dots$$

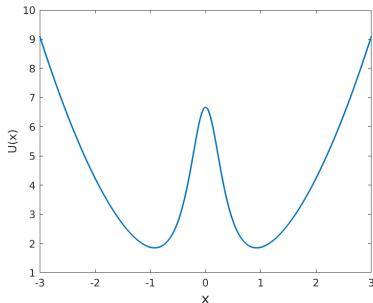
with $\eta \sim N(0, \Gamma)$.

The state X_t will oscillate
between local minima of $U(x)$.

Can we infer the pseudo-stable
state of X_t from observing V_t ?

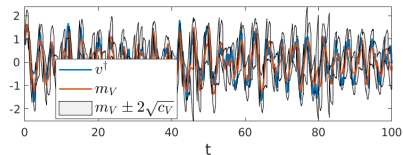
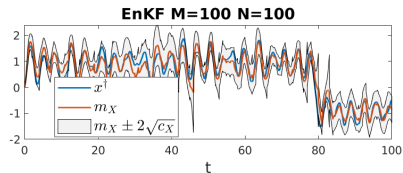
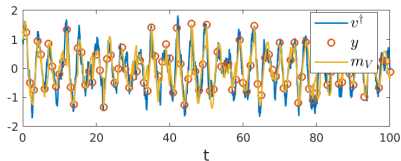
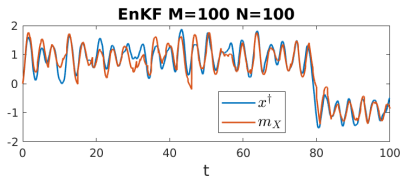
Potential:

$$U(X) = X^2 + 1/(0.15 + X^2)$$



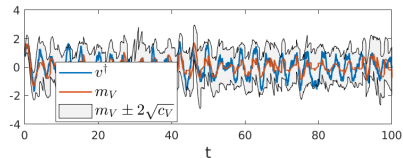
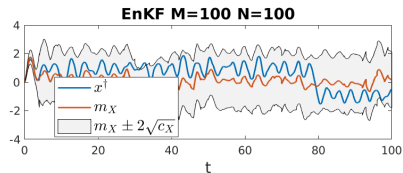
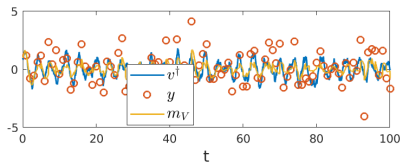
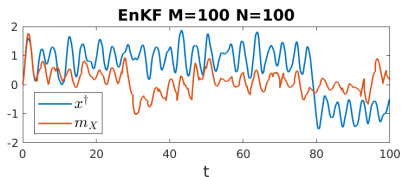
Numerical results

Small observation noise $\Gamma = 1/100$ and infrequent observations $\Delta\tau = 1$,



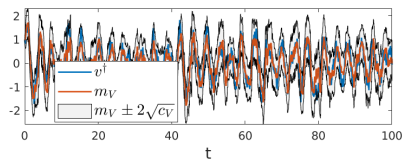
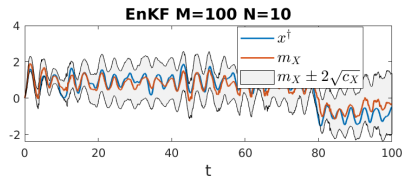
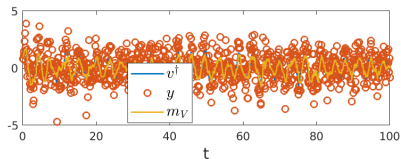
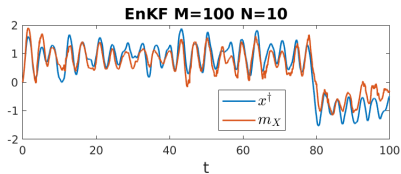
Numerical results

Large observation noise $\Gamma = 1$ and infrequent observations $\Delta\tau = 1$,



Numerical results

Large observation noise $\Gamma = 1$ and frequent observations $\Delta\tau = 0.1$,



Model and approximation error for EnKF

Let $\pi_j^{M,N}$ denote the EnKF empirical measure at time τ_j with ensemble size M and timestep $\Delta t = \Delta\tau/N$ in the Euler–Maruyama integrator.

Then, under sufficient regularity it holds for Qol f that

$$\|\pi_j^{M,N}[f] - \pi_j^{\infty,\infty}[f]\|_{L^p(\Omega)} \leq C_{p,j,f}(M^{-1/2} + N^{-1}).$$

$\pi_j^{\infty,\infty}$ – mean-field large-ensemble limit with $N = \infty$ exact-model integration. [Hoel, Law, Tempone (2016)].

Rule of thumb configuration of degrees of freedom in EnKF with Euler–Maruyama: $M = \mathcal{O}(N^2)$.

The error may be split into/bounded from above by

$$\|\pi_j^{M,N}[f] - \pi_j^{\infty,\infty}[f]\|_p \leq \underbrace{\|\pi_j^{M,N}[f] - \pi_j^{\infty,N}[f]\|_p}_{\text{bias error}} + \underbrace{\|\pi_j^{M,N}[f] - \pi_j^{\infty,N}[f]\|_p}_{\text{statistical error}}$$

Bias error is a particular kind of model error, using Ψ^N rather than the exact model Ψ as solver.

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Model uncertainty

Assume that we are given a sequence of observations $y_{1:j}$, or a collection of such, that satisfy

$$Y_j = h(V_j) + \eta_j.$$

The exact dynamics for V_j , which we denote Ψ , is unknown, but we can sample from a set of approximate dynamics $\{\Psi_\alpha\}_{\alpha \in \mathcal{M}_\Theta}$. That is

Unknown dyn: $V_{j+1} = \Psi(V_j)$, **known approx dyn** $V_{j+1}^\alpha = \Psi_\alpha(V_j^\alpha)$.

Question: given the collection of observations $y_{1:j}$ and the true observation model, how can we estimate model errors and compare models?

Strategy: Estimate error in the data space rather than in the state space.

Non-Bayesian approach

Assume the setting of exact observations

$$Y_j = h(V_j).$$

Given a collection of M_0 observation sequences $\{y_{1:J}^{(i)}\}_{i=1}^{M_0}$, we associate it to an empirical measure $\pi_{Y_{1:J}}(y_{1:J})$.

Computing the error for Ψ_α :

- Generate M_D path realizations of the dynamics $\{v_{1:J}^{\alpha,(i)}\}_{i=1}^{M_D}$.
- Associate each of these paths to observation sequences $y_{1:J}^{\alpha,(i)} = h(v_{1:J}^{\alpha,(i)})$.
- Approximate the error/divergence etc with the relevant measure in the data space. For instance, root-mean-square error,

$$RMSE(\alpha) = \|Y_{1:J}^\alpha - \mathbb{E}[Y_{1:J}]\|_{L^2(\Omega)} \approx \sqrt{\frac{1}{M_D} \sum_{i=1}^{M_D} |y_{1:J}^{\alpha,(i)} - E_{M_0}[y_{1:J}^{(\cdot)}]|^2}$$

- Best model: $\alpha^* = \arg \min_{\alpha \in \mathcal{M}_0} RMSE(\alpha)$.

[See RC 4.4] for more on scoring rules.

Bayesian approach to model selection

Assume we are given one observation sequence $Y_{1:J} = y_{1:J}$ from the noisy observation model

$$Y_{1:J} = h(V_{1:J}) + \eta_{1:J}$$

where we assume the “truth” $V_{1:J}^\dagger$ that produced the observation was generated from a model Ψ_α for some $\alpha \in \mathcal{M}\mathcal{O}$.

Bayesian framework:

- 1 Assign a prior pdf π_α to the model space.
- 2 and Bayesian inversion yields

$$\pi_{\alpha|Y_{1:J}}(\alpha|y_{1:J}) \propto \pi_{Y_{1:J}|\alpha}(y_{1:J}|\alpha)\pi_\alpha(\alpha)$$

- 3 Select model for instance by

$$\alpha^* = \text{MAP}(\pi_{\alpha|Y_{1:J}}(\cdot|y_{1:J})).$$

Problem: evaluating $\pi_{Y_{1:J}|\alpha}(y_{1:J}|\alpha)$ may not be straightforward.

Approximating the likelihood

Note that

$$\begin{aligned}\pi_{Y_{1:J}|\alpha}(y_{1:J}|\alpha) &= \int \pi_{Y_{1:J}, V_{1:J}|\alpha}(y_{1:J}, v_{1:J}|\alpha) dv_{1:J} \\ &= \int \pi_{Y_{1:J}|V_{1:J}, \alpha}(y_{1:J}|v_{1:J}, \alpha) \pi_{V_{1:J}|\alpha}(v_{1:J}|\alpha) dv_{1:J} \\ &= \int \pi_{Y_{1:J}|V_{1:J}}(y_{1:J}|v_{1:J}) \pi_{V_{1:J}|\alpha}(v_{1:J}|\alpha) dv_{1:J} \\ &= \mathbb{E} [\pi_{Y_{1:J}|V_{1:J}}(y_{1:J}|V_{1:J})|\alpha]\end{aligned}$$

Hence, the likelihood can be approximated by the Monte Carlo method:

$$\pi_{Y_{1:J}|\alpha}(y_{1:J}|\alpha) \approx \sum_{i=1}^M \frac{\pi_{Y_{1:J}|V_{1:J}}(y_{1:J}|V_{1:J}^{(i)})}{M}$$

where $V_{1:J}^{(i)} \stackrel{iid}{\sim} \pi_{V_{1:J}|\alpha}(\cdot|\alpha)$.

Toy problem

Dynamics

$$V_{j+1} = \alpha V_j,$$

with $V_0 = 1$ and prior $\pi_\alpha(\alpha) = \mathbb{1}_{[-1,1]}(\alpha)$ Observations

$$Y_{j+1} = V_{j+1} + \eta_{j+1}, \quad \eta_j \stackrel{iid}{\sim} N(0, 1).$$

Observation sequence $y_j = (-1)^j$ for $j = 1, 2, \dots, J$.

Since $V_j = \alpha^j$ (each α leads to a unique dynamics), we derive that

$$\pi_{\alpha|Y_{1:J}}(\alpha|y_{1:J}) \propto \pi_{Y_{1:J}|\alpha}(y_{1:J}|\alpha)\pi_\alpha(\alpha) \propto \mathbb{1}_{[-1,1]} \exp\left(-\frac{1}{2} \sum_{j=1}^J ((-1)^j - \alpha^j)^2\right)$$

We conclude that

$$MAP\left(\pi_{\alpha|Y_{1:J}}(\cdot|y_{1:J})\right) = -1.$$

Model parameter estimation/selection through filtering

Consider the parameter dependent dynamics

$$V_{\tau_{j+1}} = \Psi_{\alpha}(V_{\tau_j})$$

and a sequence of observations

$$Y_{\tau_{j+1}} = h(V_{\tau_{j+1}}) + \eta_{j+1}$$

Filtering strategy to parameter estimation: Augment the state space with α . New dynamics $(V_{\tau_j}, \alpha_{\tau_j})$:

$$V_{\tau_{j+1}} = \Psi_{\alpha_{\tau_j}}(V_{\tau_j})$$

$$\alpha_{\tau_{j+1}} = \alpha_{\tau_j} + \nu_j$$

where ν_j is noise. (Adding noise may improve the exploration of possible α but, unless careful, it may also render the dynamics unstable!)

Can be implemented using for instance EnKF or particle filtering with the goal that $\alpha_{\tau_j} \rightarrow \alpha_{true}$. [See ubung 9].

Summary

- The density of SDE is described by the Fokker-Planck equation.
- Have introduced the Euler–Maruyama numerical scheme for SDE
studied applications of EnKF+Euler–Maruyama model approximation.
- Similarly, one may combine particle filtering/3DVar/ExKF and
Euler–Maruyama (and more).
- Next time: continuous-time filtering for linear-coefficient SDEs.