

# solution of problem 6 kinetic theory

## Part 1, Question 1: Linearized Vlasov-Maxwell Equations

### Vlasov Equation for Electrons

The Vlasov equation describes the evolution of the electron distribution function  $f(\vec{r}, \vec{v}, t)$ . For electrons in electromagnetic fields, it is:

$$\frac{\partial f}{\partial t} + \vec{v} \cdot \nabla_r f + \frac{q}{m} (\vec{E} + \vec{v} \times \vec{B}) \cdot \nabla_v f = 0$$

where  $q = -e$  for electrons (so  $q/m = -e/m$ ),  $\vec{E}$  is the electric field, and  $\vec{B}$  is the magnetic field.

### Maxwell's Equations

Maxwell's equations govern the electromagnetic fields:

$$\nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

where  $\rho = q \int f d^3v$  is the charge density and  $\vec{J} = q \int \vec{v} f d^3v$  is the current density.

### Equilibrium and Perturbations

The equilibrium state has:

- A uniform static magnetic field  $\vec{B}_0 = B_0 \hat{z}$ .
- No electric field  $\vec{E}_0 = 0$ .
- A Maxwellian electron distribution function:

$$f_0(v) = \frac{n_0}{(2\pi v_{th}^2)^{3/2}} \exp\left(-\frac{v^2}{2v_{th}^2}\right), \quad v_{th} = \sqrt{\frac{k_B T}{m}}$$

- Ions are stationary and neutralize the plasma, so  $\rho = 0$  and  $\vec{J} = 0$  in equilibrium.

We introduce small perturbations:

- $f = f_0 + f_1$
- $\vec{E} = \vec{E}_1$  (since  $\vec{E}_0 = 0$ )
- $\vec{B} = \vec{B}_0 + \vec{B}_1$

## Linearization

We assume wave-like perturbations:  $e^{i(\vec{k} \cdot \vec{r} - \omega t)}$ , so all first-order quantities have this form. For example,  $f_1(\vec{r}, \vec{v}, t) = f_1(\vec{v})e^{i(\vec{k} \cdot \vec{r} - \omega t)}$ .

### Linearized Vlasov Equation:

Substitute into the Vlasov equation and keep only first-order terms. The linearized equation becomes:

$$\frac{\partial f_1}{\partial t} + \vec{v} \cdot \nabla_r f_1 + \frac{q}{m}(\vec{v} \times \vec{B}_0) \cdot \nabla_v f_1 = -\frac{q}{m}(\vec{E}_1 + \vec{v} \times \vec{B}_1) \cdot \nabla_v f_0$$

Using the wave-like dependence, the derivatives become:

- $\frac{\partial}{\partial t} \rightarrow -i\omega$
- $\nabla_r \rightarrow i\vec{k}$

Thus:

$$-i\omega f_1 + i\vec{k} \cdot \vec{v} f_1 + \frac{q}{m}(\vec{v} \times \vec{B}_0) \cdot \nabla_v f_1 = -\frac{q}{m}(\vec{E}_1 + \vec{v} \times \vec{B}_1) \cdot \nabla_v f_0$$

### Linearized Maxwell's Equations:

Similarly, for the fields, we linearize Maxwell's equations. Since  $\rho_1 = q \int f_1 d^3v$  and  $\vec{J}_1 = q \int \vec{v} f_1 d^3v$ , we have:

$$i\vec{k} \cdot \vec{E}_1 = \frac{\rho_1}{\epsilon_0}$$

$$i\vec{k} \cdot \vec{B}_1 = 0$$

$$i\vec{k} \times \vec{E}_1 = i\omega\vec{B}_1$$

$$i\vec{k} \times \vec{B}_1 = \mu_0\vec{J}_1 - i\omega\mu_0\epsilon_0\vec{E}_1$$

These equations describe how the perturbations  $\vec{E}_1$  and  $\vec{B}_1$  are related to  $f_1$ .

## Appendix : The Maxwellian fontion

### 1. What is it?

The provided equation is the **Maxwell-Boltzmann distribution** (or Maxwellian) for particle velocities in a plasma (or gas) in thermal equilibrium:

$$f_0(v) = \frac{n_0}{(2\pi v_{th}^2)^{3/2}} \exp\left(-\frac{v^2}{2v_{th}^2}\right)$$

Where:

- $f_0(v)$  is the equilibrium distribution function.
- $n_0$  is the particle number density (e.g., electrons/m<sup>3</sup>).
- $v_{th} = \sqrt{\frac{k_B T}{m}}$  is the **thermal speed** (a characteristic speed of particles at temperature T).
- $v^2 = v_x^2 + v_y^2 + v_z^2$  is the square of the speed.
- $k_B$  is Boltzmann's constant.
- $T$  is the temperature.
- $m$  is the particle mass.

### 2. Where does it come from? (Physics Behind It)

The Maxwellian distribution emerges naturally from **statistical mechanics**:

- It describes the most probable distribution of particle speeds in a system that has reached **thermal equilibrium**.
- It is derived by maximizing the entropy of the system subject to the constraints of constant total energy and particle number.
- In the context of plasma, it represents a plasma where collisions (or other processes) have randomized the particle velocities, resulting in no net energy transfer in any direction—a state of maximum disorder.

### 3. When is it used? For what?

It is used as the **background** or **equilibrium distribution** function  $f_0$  in the linearization of the Vlasov equation:

- **Equilibrium Assumption:** The plasma is assumed to be in a steady state with no initial electric or magnetic fields (other than the constant  $\vec{B}_0$ ).
- **Perturbation Theory:** When studying waves or instabilities, we linearize the Vlasov equation around this equilibrium state. We write the total distribution as  $f = f_0 + f_1$ , where  $f_1$  is a small perturbation. The Maxwellian  $f_0$  serves as the reference state.
- **Zero Current/Charge in Equilibrium:** For a pure Maxwellian, the average velocity  $\langle \vec{v} \rangle = 0$ , so the equilibrium current density  $\vec{J}_0 = 0$ . Similarly, the charge density is zero if ions provide a neutralizing background.

### 4. Difference from a Gaussian

- A **Gaussian** in one dimension has the form  $e^{-x^2/(2\sigma^2)}$ .
- The Maxwellian is essentially a **product of three Gaussians** (for  $v_x, v_y, v_z$ ), normalized so that its integral over all velocity space equals the density  $n_0$ :

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f_0(\vec{v}) dv_x dv_y dv_z = n_0$$

- So, it is a **multivariate Gaussian distribution in velocity space**. The term "Maxwellian" specifically refers to this Gaussian distribution applied to particle velocities in thermal equilibrium.

## Appendix : the vlasov eqt's

### 1. Where does it come from? The Physics & Relation to

# BBGKY

let's recall the **BBGKY hierarchy** (Bogoliubov–Born–Green–Kirkwood–Yvon). This is its formal origin.

- **The Full Problem:** In a plasma with  $N$  particles, to know everything, you'd need to solve for the  $N$ -particle distribution function, which depends on  $3N$  position and  $3N$  velocity coordinates. This is impossible.
- **The BBGKY Hierarchy:** This is a set of coupled equations that describes the evolution of  $s$ -particle distribution functions ( $f_1, f_2, f_3, \dots$ ). The equation for  $f_1$  depends on  $f_2$ , which depends on  $f_3$ , and so on. It's an infinite, intractable chain.
- **The Closure Assumption (The Key Step):** To break the chain, we make a physical assumption. The Vlasov equation is derived by assuming that **the two-particle correlation function  $f_2$  is negligible**. This is known as the **"mean-field"** approximation.

## What does this mean physically?

It means an individual particle doesn't feel the specific, discrete force from every other nearby particle (a two-body collision). Instead, it only feels the **smooth, averaged collective force** generated by the total charge and current densities in the plasma. This is an excellent approximation for most plasmas because the long-range nature of the Coulomb force means a particle interacts with a vast number of others simultaneously, making the collective effect dominate over short-range, binary collisions.

**In short: The Vlasov equation is the BBGKY hierarchy truncated at the first level, using the mean-field approximation. It's a collisionless kinetic equation.**

## 2. The Vlasov Equation and its Components

The Vlasov equation for a species of charged particles (e.g., electrons) is:

$$\frac{\partial f}{\partial t} + \vec{v} \cdot \nabla_r f + \frac{q}{m} (\vec{E} + \vec{v} \times \vec{B}) \cdot \nabla_v f = 0$$

This is a equation in **6-dimensional phase space**  $(\vec{r}, \vec{v})$ . The distribution function  $f(\vec{r}, \vec{v}, t)$  tells you the density of particles in a small volume  $d^3r d^3v$  around the point  $(\vec{r}, \vec{v})$  at time  $t$ .

Let's give each term a simple, physical meaning. Imagine you are riding along on a tiny spaceship through phase space, following a group of particles with a specific position and

velocity.

### Term 1: $\frac{\partial f}{\partial t}$ (The "Time Derivative")

- **Physical Meaning:** The local rate of change of the particle density *at a fixed point* in phase space.
- **Easy to Remember:** It's like a photographer standing in one spot and taking pictures of the crowd density at that spot over time.

### Term 2: $\vec{v} \cdot \nabla_r f$ (The "Spatial Flow" or "Convection" Term)

- **Physical Meaning:** This term accounts for particles with a velocity  $\vec{v}$  *flowing into or out of* the spatial volume  $d^3r$  due to their motion.
- **Easy to Remember:** This is the  $\vec{v}$  part. If particles are moving, they naturally carry their distribution function with them. A gradient ( $\nabla_r f$ ) means the density is different from one point to the next, so flow will change the local density.

### Term 3: $\frac{q}{m}(\vec{E} + \vec{v} \times \vec{B}) \cdot \nabla_v f$ (The "Acceleration" or "Force" Term)

- **Physical Meaning:** This is the most crucial term for plasmas. It accounts for particles with a position  $\vec{r}$  *flowing into or out of* the velocity volume  $d^3v$  due to forces accelerating them.
- **Easy to Remember:** This is the  $\vec{a}$  part.
- $\frac{q}{m}(\vec{E} + \vec{v} \times \vec{B})$  is the acceleration  $\vec{a}$  of a particle due to the Lorentz force.
- $\nabla_v f$  is the gradient of the distribution in *velocity space*. If  $f$  changes with velocity (e.g., more slow particles than fast ones), acceleration will move particles along the velocity axis, changing the density at a specific velocity  $\vec{v}$ .

### The Sum (= 0): The "Total Derivative"

The whole equation states that the **total derivative** of  $f$  is zero:

$$\frac{Df}{Dt} = 0$$

**Profound Physical Meaning:** As you move along with a group of particles in 6D phase space, the density of that group *does not change*. Particles are neither created nor destroyed (it's a continuity equation in phase space). They just move around in phase space due to their velocity (Term 2) and due to forces (Term 3). This is a beautiful and powerful result.

## Summary and Connection to Your Problem

In your problem, you **linearize** this equation. You assume:

1. **Equilibrium ( $f_0$ ):** The smooth, unchanging background state (the Maxwellian you asked about earlier).
2. **Perturbation ( $f_1$ ):** A very small wave-like ripple on top of that background ( $f = f_0 + f_1$ ).
3. **Perturbed Fields ( $\vec{E}_1, \vec{B}_1$ ):** The electromagnetic fields created by the perturbations in charge and current density ( $\rho_1, \vec{J}_1$ ).

## Appendix: Maxwell's\_and\_vlasov\_eqt

the concept of **self-consistency**.

The Vlasov equation and Maxwell's equations are not just used together; they are **coupled**. They form a closed, self-consistent system that describes how particles create fields and how fields, in turn, govern the motion of particles.

Here's the purpose and the physics behind their relationship, broken down into a simple cause-and-effect loop.

### The Self-Consistency Loop

The core idea is that **particles generate fields** and **fields move particles**. This creates a feedback loop:

#### 1. Particles → Fields (Maxwell's Equations):

- The charge density ( $\rho$ ) and current density ( $\vec{J}$ ) of the particles are the **sources** for the electromagnetic fields in Maxwell's equations.
- $\nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0}$ : Electric charges are the source of electric field divergence.
- $\nabla \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$ : Electric currents are the source of magnetic field curl.

#### 2. Fields → Particles (Vlasov Equation):

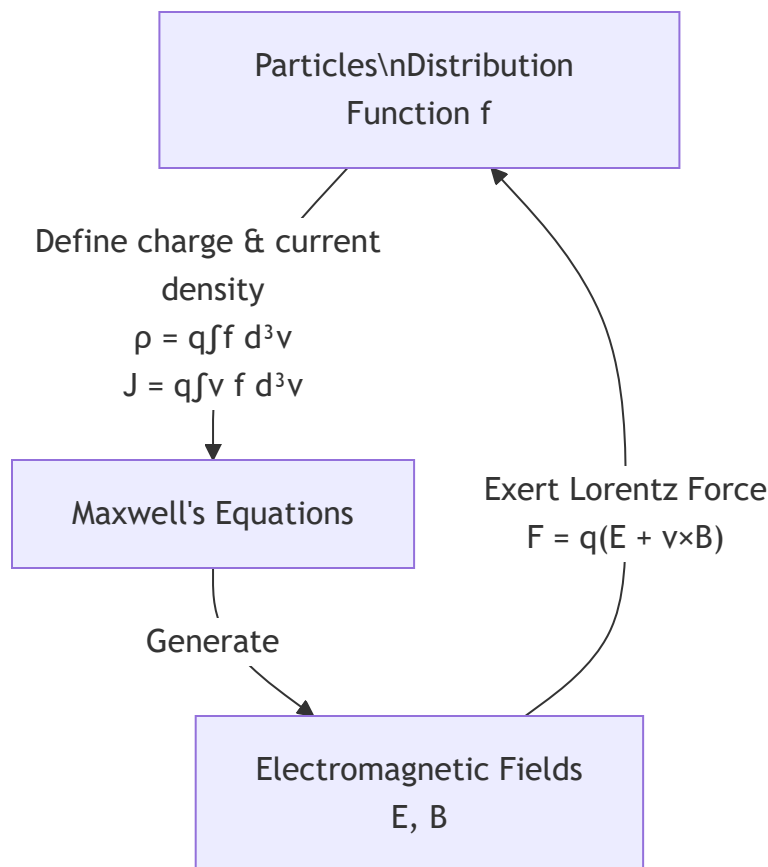
- The electromagnetic fields ( $\vec{E}$  and  $\vec{B}$ ) determine the force on every particle via the Lorentz Force Law:  $\vec{F} = q(\vec{E} + \vec{v} \times \vec{B})$ .

- This force dictates how particles **move through phase space** (change their position and velocity), which is exactly what the Vlasov equation describes:  $\frac{\partial f}{\partial t} + \vec{v} \cdot \nabla_r f + \frac{q}{m} (\vec{E} + \vec{v} \times \vec{B}) \cdot \nabla_v f = 0$ .

### 3. The Loop Closes:

- The motion of the particles (from step 2) changes their spatial distribution and flow, which **alters the charge density  $\rho$  and current density  $\vec{J}$** .
- These new  $\rho$  and  $\vec{J}$  become the new sources for Maxwell's equations (back to step 1), which modify the  $\vec{E}$  and  $\vec{B}$  fields.
- The modified fields then again alter the motion of the particles, and so on.

This loop happens continuously in time. To describe a plasma, you **must** solve both sets of equations simultaneously because each one depends on the solution of the other.



## Why is this so crucial for your problem? (Waves and Instabilities)

You are studying waves. In a plasma, waves are **self-consistent perturbations** that propagate through this coupled system.



1. **The Perturbation:** You start with an equilibrium (uniform  $f_0$ , constant  $\vec{B}_0$ ,  $\vec{E}_0 = 0$ ). Then you "wiggle" the particles a little bit, creating a small perturbed distribution  $f_1$ .
2. **Particles  $\rightarrow$  Fields:** This "wiggle" in the particles creates oscillating, wave-like perturbations in the charge and current density ( $\rho_1$ ,  $\vec{J}_1$ ).
3. **Fields  $\rightarrow$  Particles:** These oscillating sources, via Maxwell's equations, generate wave-like electromagnetic fields ( $\vec{E}_1$ ,  $\vec{B}_1$ ).
4. **The Key Question:** Do these newly created fields  $\vec{E}_1$  and  $\vec{B}_1$  reinforce the original "wiggle" ( $f_1$ ), causing it to grow (an **instability**), or do they oppose it, causing it to oscillate (a **wave**) or damp out?
5. **The Dispersion Relation:** The condition for a self-sustaining wave (where the fields created by the particle motion are just the right strength and phase to sustain that motion) leads to a mathematical equation called the **dispersion relation**, which relates the wave's frequency ( $\omega$ ) to its wavelength ( $k$ ). Solving this equation tells you everything about the wave: its existence, its propagation, and its damping.

## Part 1, Question 2: Method of Characteristics and Unperturbed Orbits

The linearized Vlasov equation can be solved using the **method of characteristics**. This involves integrating along the **unperturbed orbits** of particles—the trajectories they follow under only the static magnetic field  $\vec{B}_0$ .

### Unperturbed Orbits

For a particle in  $\vec{B}_0 = B_0 \hat{z}$ , the equation of motion is:

$$\dot{\vec{v}} = \frac{q}{m} \vec{v} \times \vec{B}_0$$

This results in **helical motion**: circular gyration in the plane perpendicular to  $\vec{B}_0$  and uniform motion along  $\vec{B}_0$ . The gyrofrequency is  $\omega_c = \frac{eB_0}{m}$  (for electrons,  $\omega_c$  is positive since  $q = -e$ , but the sense of rotation is opposite to that for ions).

### Expressing $f_1$ as an Integral

The left side of the linearized Vlasov equation represents the derivative of  $f_1$  along the unperturbed orbits. Thus, we can write:

$$\left. \frac{df_1}{dt} \right|_{\text{orbit}} = -\frac{q}{m} (\vec{E}_1 + \vec{v} \times \vec{B}_1) \cdot \nabla_v f_0$$

Integrating along the orbit from the infinite past to time  $t$ , we obtain:

$$f_1(\vec{r}, \vec{v}, t) = -\frac{q}{m} \int_{-\infty}^t dt' \left[ (\vec{E}_1 + \vec{v}' \times \vec{B}_1) \cdot \nabla_v f_0 \right]_{(\vec{r}'(t'), \vec{v}'(t'), t')}$$

Here:

- $\vec{r}'(t')$  and  $\vec{v}'(t')$  are the position and velocity along the unperturbed orbit that reaches  $\vec{r}$  and  $\vec{v}$  at time  $t$ .
- The integral sums up the effects of the perturbed fields along the particle's history.

## Effect of Magnetic Field

The magnetic field  $\vec{B}_0$  affects the orbits by causing gyration, which makes the integral more complex. Specifically:

- The helical motion introduces **cyclotron harmonics** into the response, as the particle experiences the wave fields at multiples of the gyrofrequency.
- This leads to resonances when  $\omega - k_{\parallel} v_{\parallel} - n\omega_c = 0$  for integer  $n$ , which are crucial for wave-particle interactions.

This integral form is key for deriving the dielectric tensor in Part 2.