

Theorem 3.1.6: for constant a & function $f(t)$,
 $\mathcal{L}[e^{at}f(t)](s) = \mathcal{L}[f](s-a)$.

Proof: $\mathcal{L}[e^{at}f] = \int_0^\infty e^{at}f(t)e^{-st}dt = \int_0^\infty f(t)e^{-(s-a)t}dt = \mathcal{L}[f](s-a)$.

Example 3.1.7 $f(t) = e^{at} \cos(bt)$
 $\mathcal{L}[f] = \int_0^\infty e^{at} \cos(bt) e^{-st} dt = \int_0^\infty \cos(bt) e^{-(s-a)t} dt$

So $\mathcal{L}[e^{at} \cos(bt)] = \frac{s-a}{(s-a)^2 + b^2}$

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Definition 3.1.8 The Gamma function $\Gamma(\frac{\alpha}{n}) : (-1, \infty) \rightarrow \mathbb{R}$ is
 defined by
 $\Gamma(\frac{\alpha}{n}) = \int_0^\infty t^{\alpha-1} e^{-t} dt$

We'll most often use it in the form $\Gamma(1+a) = \int_0^\infty t^a e^{-t} dt$

Observe: $\Gamma(1) = \int_0^\infty t^0 e^{-t} dt = \int_0^\infty e^{-t} dt = 1$

And for positive integers n ,

$$\Gamma(n+1) = \int_0^\infty t^n e^{-t} dt$$

$$\begin{aligned} \text{IBP} &= \left[-t^n e^{-t} \right]_0^\infty + n \int_0^\infty t^{n-1} e^{-t} dt \\ &= n \Gamma(n) \end{aligned}$$

So $\Gamma(n+1) = n \Gamma(n)$.

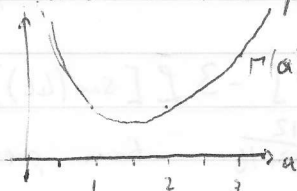
Since $\Gamma(1) = 1 \Rightarrow \Gamma(2) = 1 \cdot \Gamma(1) = 1$

$\Gamma(3) = 2 \cdot \Gamma(2) = 2$

$\Gamma(4) = 3 \cdot \Gamma(3) = 6$ etc.

So $\Gamma(n+1) = n!$

Thus the Gamma function interpolates the factorial function.
 Looks like:



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LAPLACE TRANSFORMS INTRO, PART 2 (ROYCE G.I.)

Example 3.1.9: $f(t) = t^a$ for some $a > -1$.

$$\begin{aligned} \text{Then } \mathcal{L}[f] &= \int_0^\infty t^a e^{-st} dt = \int_0^\infty \left(\frac{u}{s}\right)^a e^{-u} \cdot \frac{1}{s} du \\ &= \frac{1}{s^{a+1}} \int_0^\infty u^a e^{-u} du \\ &= \frac{\Gamma(1+a)}{s^{1+a}} \end{aligned}$$

$$\begin{aligned} u &= st & t=0 & u=0 \\ du &= s dt & t=\infty & u=\infty \\ \text{or } \frac{1}{s} du &= dt \end{aligned}$$

$$\text{So } \mathcal{L}[t^a] = \frac{\Gamma(1+a)}{s^{1+a}}$$

Suppose we know that $\int_0^\infty e^{-x^2} dx = \frac{\sqrt{\pi}}{2}$ (show if there's time)

Example 3.1.10 i) $f(t) = \frac{1}{\sqrt{t}} = t^{-\frac{1}{2}}$.

$$\text{Then } \mathcal{L}[f] = \frac{\Gamma(\frac{1}{2})}{s^{\frac{1}{2}}} = \frac{1}{\sqrt{s}} \Gamma(\frac{1}{2}) \text{ by Example 3.1.9.}$$

$$\begin{aligned} \text{And } \Gamma(\frac{1}{2}) &= \int_0^\infty t^{-\frac{1}{2}} e^{-t} dt = 2 \int_0^\infty e^{-u^2} du \\ u &= \sqrt{t} & t=0 & u=0 \\ du &= \frac{1}{2\sqrt{t}} dt & t=\infty & u=\infty \end{aligned} \quad \begin{aligned} &= 2 \cdot \frac{\sqrt{\pi}}{2} \text{ by our supposition} \\ &= \sqrt{\pi} \end{aligned}$$

$$\text{Hence } \mathcal{L}[t^{-\frac{1}{2}}] = \sqrt{\frac{\pi}{s}}.$$

$$\begin{aligned} 2) \quad f(t) &= \sqrt{t} = t^{\frac{1}{2}} \\ \text{Then } \mathcal{L}[f] &= \frac{\Gamma(\frac{3}{2})}{s^{\frac{3}{2}}} \text{ by Ex. 3.1.9} \end{aligned}$$

$$\begin{aligned} \text{But } \Gamma(\frac{3}{2}) &= \frac{1}{2} \Gamma(\frac{1}{2}) \text{ by properties of the } \Gamma\text{-function} \\ &= \frac{1}{2} \cdot \sqrt{\pi} \text{ by 1).} \end{aligned}$$

$$\text{So } \mathcal{L}[\sqrt{t}] = \frac{\sqrt{\pi}}{2s^{\frac{3}{2}}} \text{ or } \frac{\sqrt{\pi}}{2s^{3/2}}.$$

§3.2 SOLVING IVPs WITH LAPLACE TRANSFORMS (BOYCE 6.2)

How do we use Laplace transforms to solve initial value problems? The fundamental insight comes from the following theorem:

Theorem 3.2.1 Let $f(t)$, $t \geq 0$ be a continuous, differentiable function ~~etc.~~ whose Laplace transform exists, s.t. f' is piecewise continuous. Then

$$\boxed{\mathcal{L}[f'] = s \cdot \mathcal{L}[f] - f(0)}$$

for $s > \text{some } a$.

Proof sketch: $\mathcal{L}[f'] = \int_0^\infty f'(t) e^{-st} dt$
 $\stackrel{\text{IBP}}{=} f(t) e^{-st} \Big|_0^\infty - (-s) \int_0^\infty f(t) e^{-st} dt$
 $= -f(0) + s \mathcal{L}[f]$

Corollary 3.2.2: $\mathcal{L}[f''] = s \cdot \mathcal{L}[f'] - f'(0)$
 $= s(s \mathcal{L}[f] - f(0)) - f'(0)$

So $\boxed{\mathcal{L}[f''] = s^2 \mathcal{L}[f] - s f(0) - f'(0)}$

when f is doubly differentiable & well behaved (i.e. its derivatives are at most exponentially growing etc.)

We can of course use the above methodology to obtain formulae for the Laplace transform of higher derivatives of f (if they exist).

General approach 3.2.3 The idea to solving IVPs using Laplace transforms is therefore as follows:
 Given a constant coefficient linear ODE with known initial values,

• Suppose $\phi(t)$ solves the ~~equation~~ IVP. Let $\Phi(s)$ be the Laplace transform of $\phi(t)$, i.e. $\mathcal{L}[\phi] = \Phi(s)$.

- Take the Laplace transform of the DE, to obtain an equation in Φ , s , $\phi(0)$, $\phi'(0)$ etc.
- Solve for Φ as a function of s , all other quantities being known
- Look up in a table to see which $\phi(t)$ have Laplace transform $\Phi(s)$.

Example 3.2.4 We start with an easy homogeneous equation:

$$y'' - y' - 2y = 0, \quad y(0) = 1, \quad y'(0) = 0$$

We know how to solve ~~this~~ conventionally: CE is $r^2 - r - 2 = 0$,
 so solution has the form $y = c_1 e^{-t} + c_2 e^{2t}$
 IEs: $y(0) = 1 \Rightarrow c_1 + c_2 = 1$
 $y'(0) = 0 \Rightarrow -c_1 + 2c_2 = 0$
 Solving yields $c_1 = \frac{2}{3}$, $c_2 = \frac{1}{3}$, so $y = \frac{2}{3}e^{-t} + \frac{1}{3}e^{2t}$.

Laplace transform way: Suppose $y = \phi(t)$ is the solution to this IVP.
 Then $\phi'' - \phi' - 2\phi = 0$, $\phi(0) = 1$, $\phi'(0) = 0$.

• Hit the DE with \mathcal{L} :

$$\mathcal{L}[\phi'' - \phi' - 2\phi] = \mathcal{L}[0]$$

$$\Rightarrow \mathcal{L}[\phi''] - \mathcal{L}[\phi'] - 2\mathcal{L}[\phi] = 0, \text{ as } \mathcal{L}[0] = 0, \text{ \& thereby } -2\mathcal{L}[\phi] = 0$$

$$\Rightarrow (s^2 \mathcal{L}[\phi] - s\phi(0) - \phi'(0)) - (s\mathcal{L}[\phi] - \phi(0)) - 2\mathcal{L}[\phi] = 0 \text{ using Theorems 3.2.1 \& 3.2.2.}$$

Suppose $\mathcal{L}[\phi] = \Phi$

$$\text{Then } (s^2 - s - 2)\Phi - s - 0 + \frac{1}{s} = 0 \text{ using } \phi'(0) = 0, \phi(0) = 1$$

$$\text{So } \Phi(s) = \frac{s-1}{s^2-s-2} = \frac{s-1}{(s-2)(s+1)}$$

Now we use partial fractions to get $\frac{s-1}{(s+1)(s-2)} = \frac{\frac{2}{3}}{s+1} + \frac{\frac{1}{3}}{s-2}$.

$$\text{Hence } \Phi(s) = \frac{2}{3} \cdot \frac{1}{s+1} + \frac{1}{3} \cdot \frac{1}{s-2}.$$

Finally, we know $\mathcal{L}[e^{at}] = \frac{1}{s-a}$ from previously, so we conclude that we must have

$$\phi(t) = \frac{2}{3}e^{-t} + \frac{1}{3}e^{2t}.$$

□.

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Note that ~~what~~ in the last step, where we go from $\Phi(s)$ to obtaining the $\phi(t)$ whose Laplace transform is $\Phi(s)$, we have made some implicit assumptions, which are laid out below.

Definition 3.2.5 The Inverse Laplace Operator \mathcal{L}^{-1} is a linear (integral) operator that takes as input functions in s -space & returns functions in t -space. It is (for all practical purposes) the unique inverse operator of \mathcal{L} , i.e.

$$\mathcal{L}^{-1}[\mathcal{L}[f(t)]] = f(t)$$

$$\mathcal{L}[\mathcal{L}^{-1}[F(s)]] = F(s)$$

for all $f(t)$ & $F(s)$ suitably well-behaved.

The Inverse Laplace Operator can be defined using a complex integral, and hence is a bit ~~outside~~ outside the scope of the course. Thus what we usually do is consult a lookup table to see which functions $f(t)$ have Laplace transform $F(s)$.

Note: \mathcal{L}^{-1} is linear, so $\mathcal{L}^{-1}[c_1 F_1(s) + c_2 F_2(s)] = c_1 \mathcal{L}^{-1}[F_1(s)] + c_2 \mathcal{L}^{-1}[F_2(s)]$.

The Laplace Transform is particularly effective when it comes to solving nonhomogeneous DEs.

Example 3.2.6: Find the solution to $y'' + y = \sin(2t)$,
 $y(0) = 0$, $y'(0) = 1$

So let $\phi(t)$ be the solution to the IVP, & let $\Phi(s) = \mathcal{L}[\phi(t)]$.

Then $\phi'' + \phi = \sin(2t)$

So $\mathcal{L}[\phi''] + \mathcal{L}[\phi] = \mathcal{L}[\sin(2t)]$

$\Rightarrow (s^2 \Phi - s\phi(0) - \phi'(0)) + \Phi = \frac{2}{s^2 + 4}$ using $\mathcal{L}[\sin(2t)] = \frac{2}{s^2 + 4}$

$\Rightarrow (s^2 + 1)\Phi - 2s - 1 = \frac{2}{s^2 + 4}$

$\Rightarrow (s^2 + 1)\Phi = \frac{2}{s^2 + 4} + \frac{(2s+1)(s^2+4)}{s^2+4}$

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$$\Rightarrow \underline{\Phi(s)} = \frac{2s^3 + s^2 + 8s + 6}{(s^2+1)(s^2+4)}$$

We can use partial fractions again to decompose the RHS:

$$\frac{2s^3 + s^2 + 8s + 6}{(s^2+1)(s^2+4)} = \frac{as+b}{s^2+1} + \frac{cs+d}{s^2+4}$$

$$\begin{aligned} \text{So } 2s^3 + s^2 + 8s + 6 &= (as+b)(s^2+4) + (cs+d)(s^2+1) \\ &= (a+c)s^3 + (b+d)s^2 + (4a+c)s + (4b+d) \end{aligned}$$

$$\text{Hence we must have } \begin{aligned} a+c &= 2, & b+d &= 1 \\ 4a+c &= 8, & 4b+d &= 6 \end{aligned}$$

These are 2 independent systems of linear equations in 2 variables each, which we know how to solve. We get:

$$a=2, \quad c=0, \quad b=\frac{5}{3} \quad \text{and} \quad d=-\frac{2}{3},$$

$$\text{So } \underline{\Phi(s)} = \frac{2s}{s^2+1} + \frac{\frac{5}{3}}{s^2+1} - \frac{\frac{2}{3}}{s^2+4}$$

$$\text{Hence } \phi(t) = 2 \mathcal{L}^{-1}\left[\frac{s}{s^2+1}\right] + \frac{5}{3} \mathcal{L}^{-1}\left[\frac{1}{s^2+1}\right] - \frac{1}{3} \mathcal{L}^{-1}\left[\frac{2}{s^2+4}\right]$$

$$\Rightarrow \underline{\phi(t) = 2 \cos(t) + \frac{5}{3} \sin(t) - \frac{1}{3} \sin(2t)}. \quad \square$$

We see then that the Laplace Transform method has advantages & disadvantages associated with it:

Advantages:

- Converts differential problem into purely algebraic problem.
- Can be used to solve higher-order IVPs.
- Works on NH equations with a large range of forcing functions.

Disadvantages:

- Algebra can get messy
- Need a lookup table of Laplace (inverse) transforms
- Constant coefficient IVPs only
(can be extended to non-constant coeff. equation, but the algebra can become unsolvably messy).