

A Model of Heavy QCD Axion and the LHC Signature

arXiv:1504.06084, 1602.07909, 1607.01936

Hajime Fukuda (Kavli IPMU)

July 13, 2016

Outline

1. Heavy axion - Why and How?
2. Our model and the LHC signature

Strong CP Problem

- QCD should break CP symmetry

$$\theta = \theta_{\text{YM}} + \arg \det(Y_u Y_d)$$

- The PQ mechanism can set $\theta = 0$.
 - The original model has been excluded.

What is the Alternative?

- Roughly, two choices:
 - Larger f_a / Heavier m_a

Axion Mass

$$m_a^2 \simeq \frac{m_q \Lambda^3}{f_a^2}$$

indicates heavier m_a is difficult.

Higher Dimensional Operator

$$\Delta\mathcal{L} = c \frac{\phi^5}{M_{\text{Pl}}}$$

$$\Rightarrow \Delta\theta \simeq c \frac{f_a^3}{M_{\text{Pl}} m_a^2} \gg 10^{-10},$$

indicates heavier m_a is preferred!

Realizing a Heavy Axion

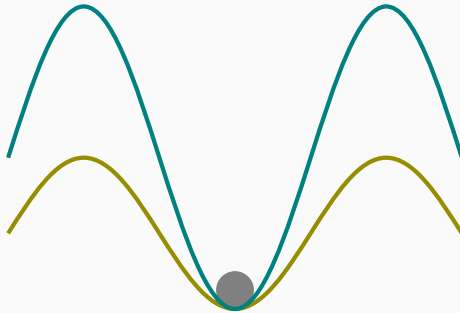
- (Rubakov, 1997) suggested a *consistent way* to achieve a heavy axion

Rubakov 1997; Berezhiani, Gianfagna and Giannotti 2000

Hook 2014, HF, Harigaya, Ibe and Yanagida 2015, Albaid, Dine and Draper 2015
(Kobakhidze 2016), (Gherghetta, Nagata and Shifman 2016)

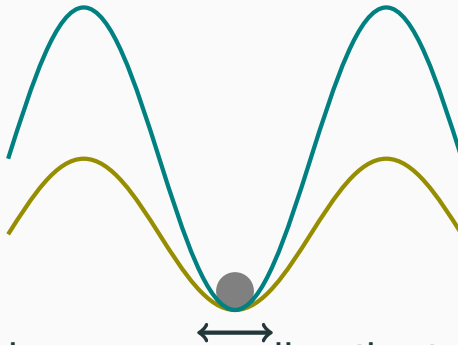
How to Make an Axion Heavier?

- Another gauge theory is needed



How to Make an Axion Heavier?

- Another gauge theory is needed



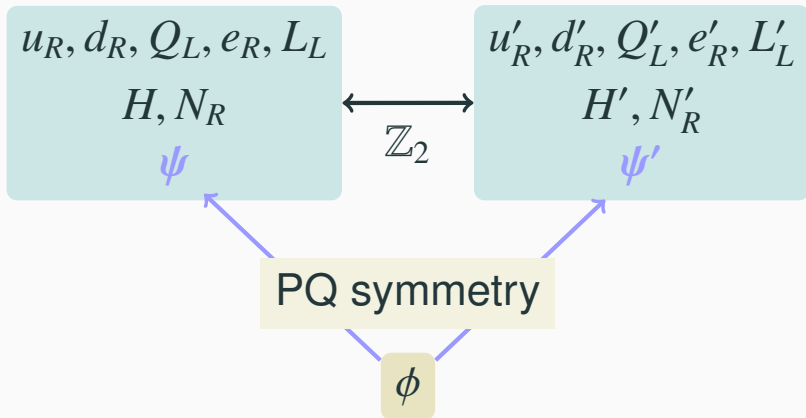
Then how can we align the two θ s?

Copy of SM

$$\theta = \theta_{\text{YM}} + \arg \det(Y_u Y_d)$$

- θ' must also have Yukawa sector
- Thus, we need a complete copy of SM
 - We assume \mathbb{Z}_2 parity, which is spontaneously broken

Our Model



HF, Harigaya, Ibe and Yanagida,
arXiv:1504.06084

Low Energy Spectrum

Axion a

$$m_a \gtrsim 400 \text{ MeV}$$

Vector like quark ψ, ψ'

$$m_\psi = \frac{1}{\sqrt{2}} g f_a \gtrsim 900 \text{ GeV}$$

Dilaton s

$$m_s = \sqrt{2\lambda} f_a \gtrsim O(100) \text{ GeV}$$

Low Energy Spectrum

Axion a

$$m_a \gtrsim 400 \text{ MeV}$$

Vector like quark ψ, ψ'

$$m_\psi = \frac{1}{\sqrt{2}} g f_a \gtrsim 900 \text{ GeV}$$

Dilaton s

$$m_s = \sqrt{2\lambda} f_a \simeq 750 \text{ GeV??}$$

or, "let us consider the LHC signal of the dilaton"...

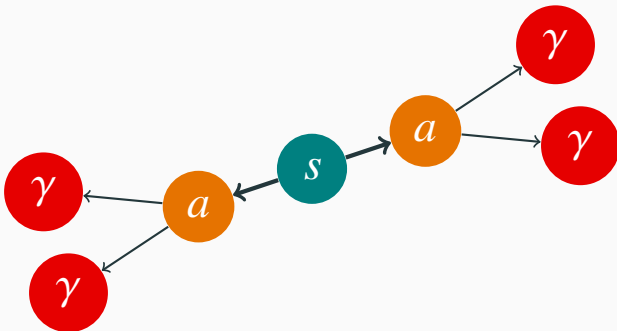
Effective Lagrangian

$$\begin{aligned}\mathcal{L} = & \frac{s}{f_a} \partial a \partial a + N_1 \frac{\alpha_s}{8\pi} \frac{s}{f_a} G G + N_2 \frac{\alpha}{8\pi} \frac{s}{f_a} F^{(\prime)} F^{(\prime)} \\ & + N_1 \frac{\alpha_s}{8\pi} \frac{a}{f_a} G \tilde{G} + N_2 \frac{\alpha}{8\pi} \frac{a}{f_a} F^{(\prime)} \tilde{F}^{(\prime)}\end{aligned}$$

- $s\partial a\partial a$ is the strongest
- How does it look?

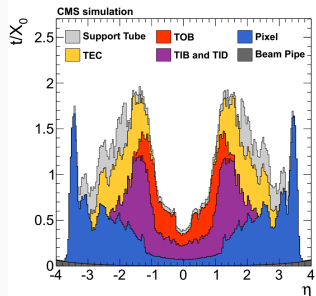
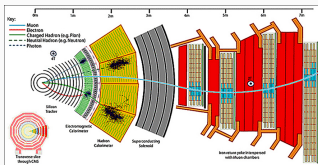
Photons and Photon Jets

- ECAL can't count the number of γ
- The decay looks like “diphoton”!



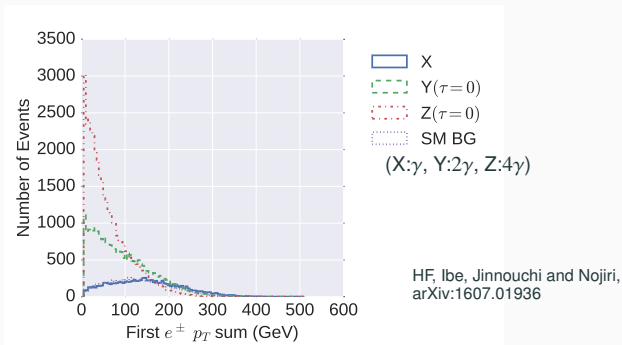
Difference b/w γ s and γ -jets

- Some γ s are “converted”



How to Distinguish the Jet

- We simulate trackers in CMS.
 - Conversion, bremsstrahlung, ...
- p_T^{track} has greater discrimination power!



Axion Decay

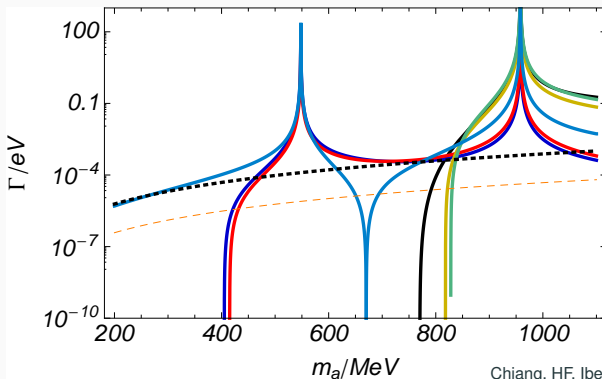
Lagrangian

$$\mathcal{L}_a = N_1 \frac{\alpha_s}{8\pi} \frac{a}{f_a} G^{(\prime)} \tilde{G}^{(\prime)} + N_2 \frac{\alpha}{8\pi} \frac{a}{f_a} F^{(\prime)} \tilde{F}^{(\prime)}$$

- We need large BR
 - $\text{BR}(s \rightarrow 4\gamma) = \text{BR}(a \rightarrow 2\gamma)^2$
- a - G - G coupling looks too strong

Mixings with Mesons

- The phase space suppresses $a \rightarrow 3\pi$



Chiang, HF, Ibe and Yanagida,
arXiv:1602.07909

Summary

- The heavy axion is plausible
- The dilaton may appear at the LHC as a “di-photon-jet” signal.
- “Diphoton” and “Di-photon-jet” is distinguishable using p_T