

# Possibility of the Heavy QCD Axion

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# Strong CP Problem

- QCD should break CP symmetry

$$\theta = \theta_{\text{YM}} + \arg \det(Y_u Y_d)$$

- However, the violation looks very small

$$|\theta| \lesssim 10^{-10}$$

# Peccei-Quinn mechanism Peccei & Quinn, 1977

## $U(1)_{PQ}$ Symmetry

$$q_L \rightarrow e^{i\alpha} q_L, \quad \theta \rightarrow \theta + 2T(R)\alpha$$

- $U(1)_{PQ}$  must be broken at  $f_a$  and a pseudo NG Boson  $a$  appears

Weinberg 1978, Wilczek 1978

# Their Original Model

- The VEVs of 2HDM break EW gauge group and  $U(1)_{PQ}$  simultaneously
- It's simple and minimal, but experimentally excluded

# Which Direction?

- There are roughly two ways to achieve the PQ mechanism
  - Larger  $f_a$ , *invisible axion*
  - Heavier  $m_a$ , *heavy axion*

# Axion Mass and Decay Constant

## Axion Mass

$$m_a^2 \simeq \frac{m_q \Lambda^3}{f_a^2}$$

- Heavier  $m_a$  with sufficiently large  $f_a$  is hence difficult

# Larger $f_a$ isn't Easy, Either

- Why does no higher dim. op. exist?

$$\Delta\mathcal{L} = c \frac{\phi^5}{M_{\text{Pl}}}$$

$$\Rightarrow \Delta\theta \simeq c \frac{f_a^3}{M_{\text{Pl}} m_a^2} \gg 10^{-10},$$

even for the WW axion

# Realizing a Heavy Axion

- (Rubakov, 1997) suggested a *consistent way* to achieve a heavy axion

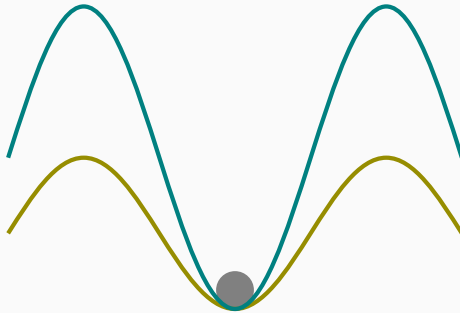
Rubakov 1997; Berezhiani, Gianfagna and Giannotti 2000

Hook 2014, HF, Harigaya, Ibe and Yanagida 2015, Albaid, Dine and Draper 2015  
(Gherghetta, Nagata and Shifman 2016)



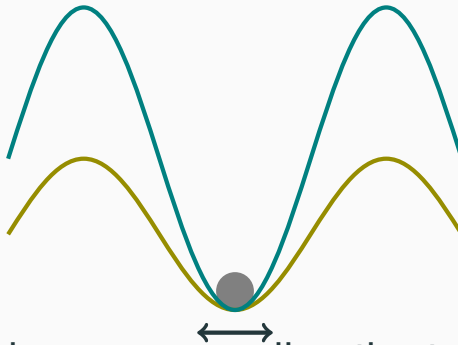
# How to Make an Axion Heavier?

- Another gauge theory is needed



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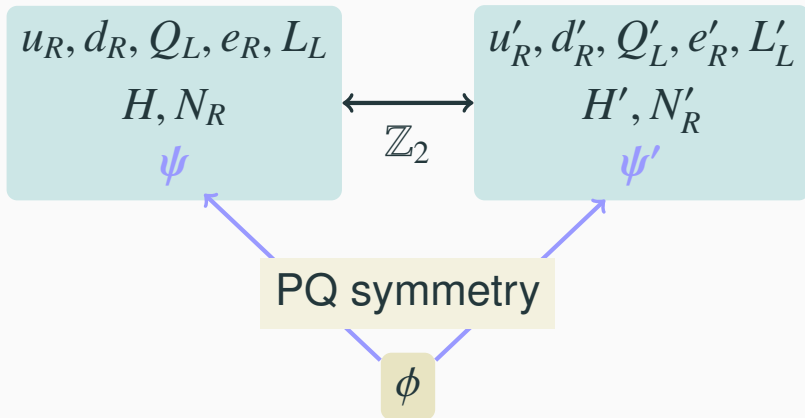
Then how can we align the two  $\theta$ s?

# Copy of SM

$$\theta = \theta_{\text{YM}} + \arg \det(Y_u Y_d)$$

- $\theta'$  must also have Yukawa sector
- Thus, we need a complete copy of SM
  - We assume  $\mathbb{Z}_2$  parity, which is spontaneously broken

# Our Model



# Use of spontaneous $\mathbb{Z}_2$ breaking

- Recall

$$m_a^2 \simeq \frac{m'_q \Lambda'^3}{f_a^2}$$

- We have to increase  $m'_q \propto v'$  and  $\Lambda'$ 
  - For  $\Lambda'$ , we introduce color charged particles and change their masses.

# Cosmological Properties

- $\gamma'$  is massless
  - The axion must decouple before QCD PT
- Seesaw mechanism in  $\nu'$  is forbidden
  - $\nu'$ s have large Dirac mass
  - No fine-tuning:  $\sigma_{\mathbb{Z}_2} = \sigma_{B'-L'}^2 / M_{\text{Pl}}$

# Low Energy Spectrum

**Axion  $a$**

$$m_a \gtrsim 400 \text{ MeV}$$

**Vector like quark  $\psi, \psi'$**

$$m_\psi = \frac{1}{\sqrt{2}} g f_a \gtrsim 900 \text{ GeV}$$

**Dilaton  $s$**

$$m_s = \sqrt{2\lambda} f_a \gtrsim O(100) \text{ GeV}$$

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$$m_s = \sqrt{2\lambda} f_a \simeq 750 \text{ GeV??}$$



# Dilaton Decay

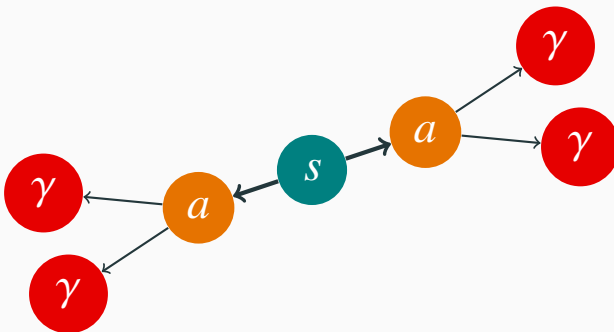
- Obviously,  $\frac{s}{f_a} \partial a \partial a$  is the strongest
- Almost no  $s \rightarrow 2\gamma^{(\prime)}$  decay
- Does it fail?

# Dilaton Decay

- Obviously,  $\frac{s}{f_a} \partial a \partial a$  is the strongest
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- Does it fail? - **No!**

# Photons and Photon Jets

- ECAL can't count the number of  $\gamma$ 
  - Use “ $s \rightarrow 2a, a \rightarrow 2$  collinear  $\gamma$ ” mode



# Axion Decay

## Lagrangian

$$\mathcal{L}_a = N_1 \frac{\alpha_s}{8\pi} \frac{a}{f_a} G^{(\prime)} \tilde{G}^{(\prime)} + N_2 \frac{\alpha}{8\pi} \frac{a}{f_a} F^{(\prime)} \tilde{F}^{(\prime)}$$

- We need large BR
  - $\text{BR}(s \rightarrow 4\gamma) = \text{BR}(a \rightarrow 2\gamma)^2$
- $a$ - $G$ - $G$  coupling looks too strong

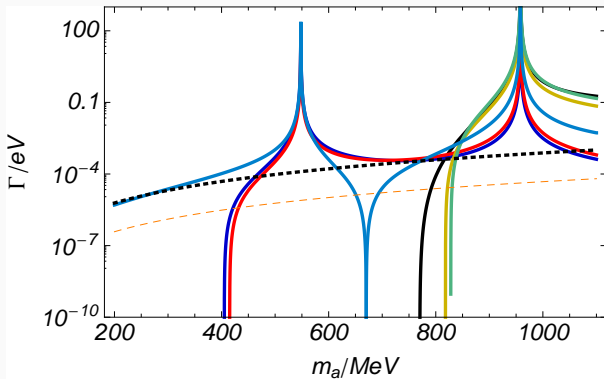
# Is Large BR Possible?

## Two possibility

- $m_a < 3m_\pi$ , the threshold of  $a \rightarrow 2g$
- Use the mixings with mesons

# Mixings with Mesons

- The phase space suppresses  $a \rightarrow 3\pi$



# Summary

- The heavy axion is possible
- We need a complete copy of SM
- The diphoton excess can be explained as the dilaton using our model