

# On level crossing in conformal field theories

1512.05362

## 1. Intro

Assume a CFT has a large- $N$  limit,

and a "coupling const"  $g$

e.g.

$$N=4: N \cdot \lambda$$

large-spin.

$$\text{twist-2 sector} = N \cdot \frac{1}{S}$$

$$QCD \text{ Veneziano lim.} = N_c \cdot \frac{N_F}{N_c}$$

$\partial_{\bar{z}}$  operator

$\partial_{\bar{z}}$



so that  $\Delta_i = \Delta_i^{(0)}(g) + \frac{1}{N^2} \Delta_i^{(2)}(g)$

(satisfy this)

$$+ \frac{1}{N^4} \Delta_i^{(4)}(g) + \dots$$

$\mathcal{O}_1, \mathcal{O}_2$ : scalar primary ops

Suppose  $\Delta_1^{(0)} = \Delta_2^{(0)}$  at some  $g = g_*$

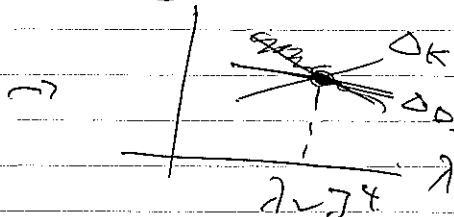
e.g.  $N=4$

$$\mathcal{O}_K = \frac{1}{N} \text{Tr}(\phi^A \phi^A)$$

$$\Delta_K = 2\lambda^{1/4} - 2 + \dots$$

$$\mathcal{O}_D = \frac{1}{N^3} \text{Tr}(\phi^{A_1} \dots \phi^{A_j}) \text{Tr}(\phi^{A_1} \dots \phi^{A_j})$$

$$\Delta_D = 2J + \dots$$



$$\sim \lambda^{1/4} \sim J$$

von Neumann - Wigner thm

It is generally hard that  $\mathcal{O}_1$  and  $\mathcal{O}_2$  have the same dilatation eigenvalue.

(generic)

Suppose that at some  $(N, g)$ ,  $\Delta_1(N, g) \neq \Delta_2(N, g)$   
 at  $(N + \delta N, g + \delta g)$ , consider dilatation matrix

$$D = \begin{pmatrix} \Delta_1 & \\ & \Delta_2 \end{pmatrix} + \begin{pmatrix} \delta_{11} & \delta_{12} \\ \delta_{21} & \delta_{22} \end{pmatrix}$$

$$D = \delta \mathbb{I}$$

$$\Rightarrow \begin{cases} \delta_{11} - \delta_{22} = -\Delta_1 + \Delta_2 \\ \delta_{12} (= \delta_{21}) = 0 \end{cases}$$

2 variables vs. ~~2 eqns~~

3 (real) equations.

→ generically, solutions are isolated points  
 in  $(N, g)$ -space

( $\exists$  点の  $N \in \mathbb{Z}^7 \subset \mathbb{Z}^{12}$ ,  $\Delta_1 = \Delta_2 \exists \delta g \neq 0$   
 $\exists$  generic  $(\tau \neq \tau_2 \in \mathbb{R}^{11})$ )

Remember that in QM, the perturbation theory  
 is valid only when  $|\Delta_1^{(0)} - \Delta_2^{(0)}|$  is much  
 larger than the leading energy correlation.

$$\delta \Delta_n \sim \sum_i \frac{|\langle n | \delta V | i \rangle|^2}{\Delta_n - \Delta_i}$$

When  $\Delta_n - \Delta_i \sim \mathcal{O}(1/N)$ ,  $1/N$ -perturbation  
 breaks down. & resummation is needed.

( $\exists \tau \neq \tau_2 \in \mathbb{R}^{11}$  to  $\Delta_1 = \Delta_2$  の  $\delta g \neq 0$  の  $\exists$  点  $N \in \mathbb{Z}^7 \subset \mathbb{Z}^{12}$ )

## 2. resummed scaling dimensions

$$\langle O_i(x) O_j(y) \rangle \propto \frac{\delta_{ij}}{|x-y|^{2\Delta_i}}$$

Assume at  $g=g_*$ ,  $\Delta_1^{(0)} = \Delta_2^{(0)}$ ,

& other scalar primaries satisfy

$$|\Delta_{\text{other}} - \Delta_{1,2}^{(0)}| \sim \mathcal{O}(1)$$

At  $N \rightarrow \infty$ ,  $\langle \mathcal{O}_1^{(0)}(x) \mathcal{O}_2^{(0)}(y) \rangle = 0$

At  $N < \infty$ ,  $\langle \mathcal{O}_1^{(0)}(x) \mathcal{O}_2^{(0)}(y) \rangle = \frac{1}{N} \mathcal{O}(|x-y|^2)$

↑  
op. mixing

e.g.  $N < \infty$  so  $\Delta_K \neq \Delta_{\bar{K}}$   
dilatation eigenstate  $2^{\text{nd}}$  or  $1^{\text{st}}$   
of  $\mathbb{Z}_2 \times \mathbb{Z}_2$

Assume  $\langle \mathcal{O}_1^{(0)}(x) \mathcal{O}_2^{(0)}(y) \rangle \Big|_{\text{non-planar}} = \mathcal{O}(1/N^2)$

At  $N < \infty$ , we expect

$$\left. \begin{aligned} \mathcal{O}_+ &= \mathcal{O}_1^{(0)} + c_2 \mathcal{O}_2^{(0)} \\ \mathcal{O}_- &= \mathcal{O}_2^{(0)} - c_1 \mathcal{O}_1^{(0)} \end{aligned} \right\} \begin{array}{l} \text{dilatation} \\ \text{eigenstates} \end{array}$$

↓  
 $\mathbb{Z}_2 \times \mathbb{Z}_2$  24 24 24 24

$$\langle \mathcal{O}_+ \mathcal{O}_+ \rangle \neq 0, \langle \mathcal{O}_- \mathcal{O}_- \rangle \neq 0, \langle \mathcal{O}_+ \mathcal{O}_- \rangle = 0$$

$$0 = \langle \mathcal{O}_+(x) \mathcal{O}_-(y) \rangle$$

$$\begin{aligned} \rightarrow \frac{1+c_1c_2}{N} \mathcal{O}(|x|) &= -|x|^{-\Delta_1^{(0)}} \left[ c_1 + c_2 + c_2 \mathcal{O}(\log|x|^2) \right] \\ \epsilon &= \Delta_1^{(0)} - \Delta_2^{(0)} \\ &= -|x|^{-2\Delta_2} (1 + \dots \log|x|^2) \end{aligned}$$

$$A+ |e| \sim 1/N.$$

$$C_1 + C_2 = 0 + \frac{\beta}{N} + O(1/N^2)$$

$$\rightarrow \phi(x) = (x^2)^{-\Delta_1} \frac{\beta + \epsilon N \log x^2}{1 - C_1^2}$$

$$\langle \phi_+(t) \phi_-(0) \rangle = (\text{const.}) \frac{1}{|x|^{2\Delta_-}}$$

$$\sim \Delta_- = \Delta_2 + \frac{\beta}{N} C_1 \quad \gamma = \frac{\epsilon N C_1}{1 - C_1^2}$$

$$\Delta_1 = \frac{\Delta_1 + \Delta_2}{2} \pm \sqrt{\frac{\epsilon^2}{4} + \frac{\beta^2}{N^2}} = O(1/N^2)$$

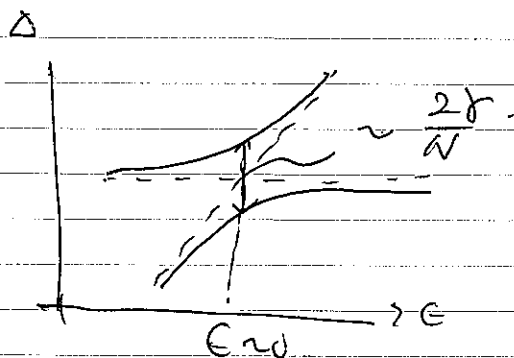
Expanding in  $\frac{1}{N}$  leads to

$$\Delta_- = \Delta_2 - \frac{\beta^2}{N^2 \epsilon} + \frac{\beta^4}{N^4 \epsilon^3} - \dots$$

$\rightarrow$  each term diverges at  $\epsilon \sim 0$

around  $\epsilon \sim 0$ .

$$\Delta |\Delta_+ - \Delta_-| \geq \frac{2\beta}{N}$$



### 3. Resummed OPE coefficients

$\phi$ : another scalar primary

$$|\Delta_\phi - \Delta_\pm| = \mathcal{O}(1)$$

$$\langle \phi(x_1) \phi(x_2) \dots \phi(x_n) \rangle = \frac{1}{(x_{12}^2 x_{34}^2)^{\Delta_\phi}}$$

$$C_{\pm} \sim \underbrace{\sum_{\Delta} F_{\Delta}(u,v)}_{\text{conf. block}} \times \left[ F_{\Delta_+}(u,v) + F_{\Delta_-}(u,v) + \sum_i F_{\Delta_i}(u,v) \right]$$

$\lambda, 2, 3 \in \mathbb{Z}$

$F_0(u,v)$   $u, v$ : cross ratios.

$$u = \left( \frac{|x_{12}| |x_{34}|}{|x_{13}| |x_{24}|} \right)^2, \quad v = u \Big|_{2 \leftrightarrow 4}$$

In the limit  $x_{12} \rightarrow 0$ .

$$F_{\Delta_\pm}(u,v) \sim C_{\phi\phi\Delta_\pm}^2 u^{\frac{1}{2}\Delta_\pm} + (x_{12}\text{-suppressed})$$

$$C_{\phi\phi\Delta_\pm}^2 \sim \frac{\langle \phi\phi\Delta_\pm \rangle \langle \Delta_\pm\phi\phi \rangle}{\langle \phi\phi \rangle^2 \langle \Delta_\pm\Delta_\pm \rangle}$$

$$C_{\phi\phi\Delta_\pm} \sim \frac{(C_{\phi\phi\Delta_{1,2}} \mp C_{\phi\phi\Delta_{2,1}})^2}{1 + C_i^2}$$

$\langle \phi\phi\phi\phi \rangle$  is a well-def func. in CFT

$\hookrightarrow$  must be non-singular in the limit  $\epsilon \rightarrow 0$

$F_{\phi_i}$ : non-singular.

$F_{\Delta_\pm}$ : the large- $N$  perturbation becomes invalid at  $\epsilon \rightarrow 0$ .

$$F_{0+} + F_{0-} \underset{x_{12} \rightarrow 0}{\sim} u^{\frac{\Delta_1 + \Delta_2}{2}} \left[ C_{\phi\phi 0_1}^2 u^{\frac{\epsilon}{4}} + C_{\phi\phi 0_2}^2 u^{-\frac{\epsilon}{4}} - 2\gamma C_{\phi\phi 0_1} C_{\phi\phi 0_2} \frac{u^{\frac{\epsilon}{4}} - u^{-\frac{\epsilon}{4}}}{N\epsilon} \right] + O(1/N^2)$$

$O(\epsilon^0)$

By the assumption,  $O(1/N)$  term does not exist

$$\rightarrow \frac{C_{\phi\phi 0_1}}{C_{\phi\phi 0_2}} = O\left(\frac{1}{N}\right) \text{ or } O(N)$$

$$C_{\phi\phi 0_2}^2 = \frac{1}{2} C_{\phi\phi 0_1}^2 \left( 1 \mp \frac{\epsilon}{\sqrt{\epsilon^2 + \frac{4\gamma^2}{N^2}}} \right)$$

$$\sim \begin{cases} \frac{C_{\phi\phi 0_+}}{C_{\phi\phi 0_1}} = 1 - \frac{\gamma^2}{2N^2 \epsilon^2} + \dots \\ \frac{C_{\phi\phi 0_-}}{C_{\phi\phi 0_1}} = \frac{\gamma}{N\epsilon} \left( 1 - \frac{3\gamma^2}{2N^2 \epsilon^2} + \dots \right) \end{cases}$$

resum  $\gamma \epsilon \rightarrow 0$  is ok.

$$\downarrow$$

$$|C_{\phi\phi 0_+}| \sim 1$$

$$C_{\phi\phi 0_-} \sim 0, (1 \gg \epsilon \gg 1/N)$$

$$F_{0+} + F_{0-} \sim \frac{1}{2} \left( u^{\frac{1}{2}(\Delta_1 + \frac{\gamma}{N})} + u^{\frac{1}{2}(\Delta_1 - \frac{\gamma}{N})} \right)$$

$$\rightarrow (\Delta_+ - \Delta_-) \geq \frac{4\gamma}{N}$$