

# Numerical Simulation Methods in Geophysics, Exercise 13: **1D/2D EM**

## 1. MGPY+MGIN

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# 1D/2D (FD) EM modelling

The induction equation for perpendicular (E or H) fields

$$-\nabla^2 u + \omega\mu\sigma u = f$$

is discretized by stiffness matrix  $\mathbf{A}$  and mass matrix  $\mathbf{M}$

$$(\mathbf{A} + \omega\mathbf{M})\mathbf{u} = \mathbf{f}$$

See [Lecture in Theory EM](#)

# TM polarization

$$\nabla \times \sigma^{-1} \nabla \times \mathbf{H} + \omega \mu \mathbf{H} = \nabla \times \sigma^{-1} \mathbf{j}_s$$

## Transverse magnetic (TM) mode

Assume the source field is oscillating perpendicular to the modelling plane, i.e.

$$\mathbf{H} = [H_x, 0, 0]^T e^{i\omega t}.$$

Then the PDE holds for the scalar  $H_x$  (now only  $H$ )

$$-\nabla \cdot \sigma^{-1} \nabla H_x(y, z) + \omega \mu H_x(y, z) = 0$$

# TE polarization

$$\nabla \times \mu^{-1} \nabla \times \mathbf{E} + \imath\omega\sigma\mathbf{E} = \nabla \times \mathbf{j}_s$$

## Transverse electric (TE) mode

Assume the source field is oscillating perpendicular to the modelling plane, i.e.

$$\mathbf{E} = [E_x, 0, 0]^T e^{\imath\omega t}.$$

Then the PDE holds for the scalar  $E_x$  (now only  $E$ )

$$-\nabla \cdot \mu^{-1} \nabla E_x(y, z) + \imath\omega\sigma E_x(y, z) = 0$$

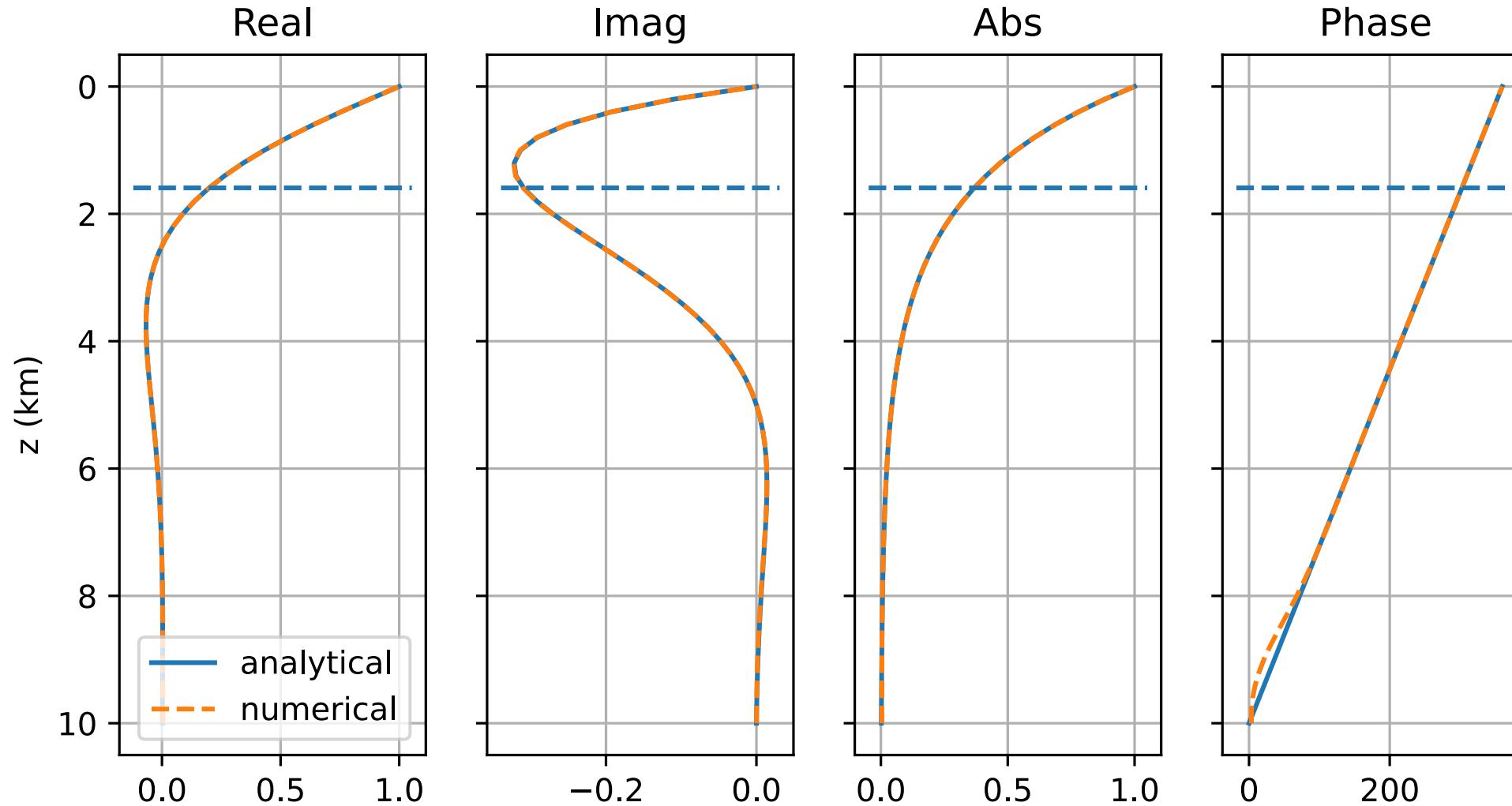
# The analytical solution

$$H = H_0 e^{-\imath kz} \quad \text{with} \quad k^2 = -\omega\mu\sigma$$

```
1 import numpy as np
2 import matplotlib.pyplot as plt
3 from poisson1d import stiffnessMatrix1DFE, massMatrix1DFE
4 T = 0.1 # 1, 10
5 w = 2 * np.pi / T
6 sigma0 = 1/100 # 100 Ohmm, 10, 1000
7 mu = np.pi * 4e-7
8 d = np.sqrt(2/(mu*w*sigma0))
9 print(d)
10 z = np.arange(0, 100001, 100)
11
12 def Hanalytical(z, sigma, w):
13     k = np.sqrt(-1j*w*mu*sigma0)
14
15 #def plotResults(z, u)
16 fig, ax = plt.subplots(ncols=4, sharey=True)
```

1591.5494309189535

# The analytical solution



# The problem in 1D

```
1 import numpy as np
2 import matplotlib.pyplot as plt
3 from poisson1d import stiffnessMatrix1DFE, massMatrix1DFE
4 T = 0.1 # 0.1
5 w = 2 * np.pi / T
6 sigma0 = 1/100 # 1/100
7 mu = np.pi * 4e-7
8 d = np.sqrt(2/(mu*w*sigma0))
9 z = np.arange(-10000, 0.1, 200)
10
11 A, b = stiffnessMatrix1DFE(x=z, uL=1)
12 M = massMatrix1DFE(x=z, a=mu*sigma0)
13 AM = A + M * 1j * w
14 u = np.linalg.solve(AM, b)
15 # plotResult(u)
```

# Complex or real-valued?

The complex-valued system

$$(\mathbf{A} + \imath\omega\mathbf{M})\mathbf{u} = (\mathbf{A} + \imath\omega\mathbf{M})(\mathbf{u}_r + \imath\mathbf{u}_i) = \mathbf{b}_r + \imath\mathbf{b}_i$$

can be transferred into a doubled real-valued system

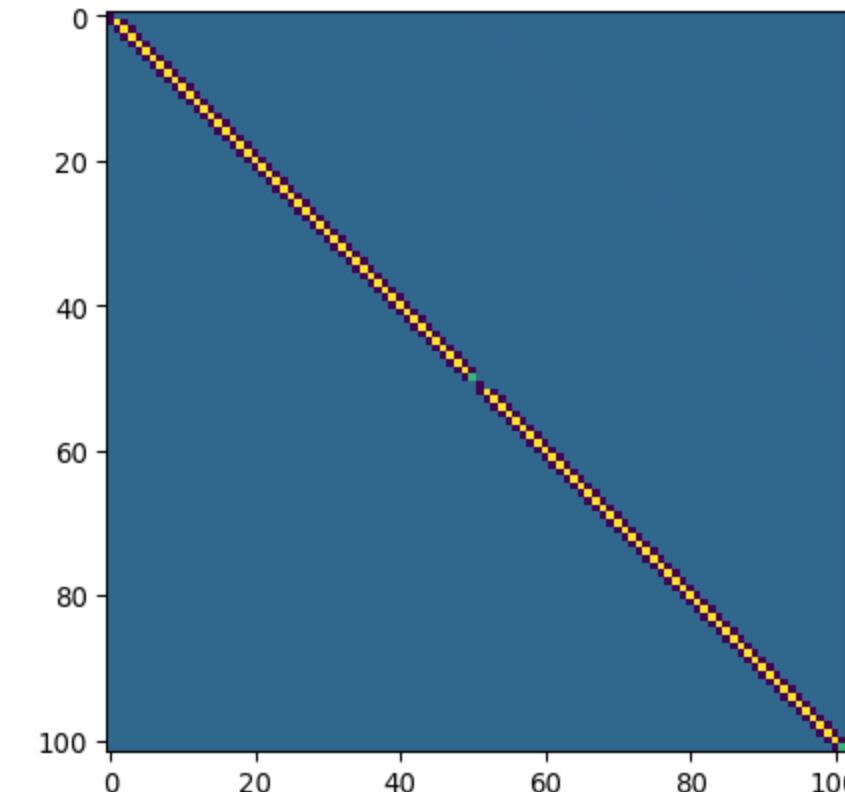
$$\mathbf{A}\mathbf{u}_r + \imath\mathbf{A}\mathbf{u}_i + \imath\omega\mathbf{M}\mathbf{u}_r - \omega\mathbf{M}\mathbf{u}_i = \mathbf{b}_r + \imath\mathbf{b}_i$$

$$\begin{pmatrix} A & -\omega M \\ \omega M & A \end{pmatrix} \begin{pmatrix} u_r \\ u_i \end{pmatrix} = \begin{pmatrix} b_r \\ b_i \end{pmatrix}$$

# Complex-to-real conversion (numpy)

$$\mathbf{B} = \begin{pmatrix} \mathbf{A} & -\omega\mathbf{M} \\ \omega\mathbf{M} & \mathbf{A} \end{pmatrix}$$

```
1 D = np.vstack([np.hstack([A, -M*w]),
2                 np.hstack([M*w, A])])
3 plt.imshow(D)
4 d = np.hstack([b, b*0])
5 uri = np.linalg.solve(D, d)
6 u = uri[:len(z)] + uri[len(z):] * 1j
```



# Complex-to-real conversion (pyGIMLi)

$$\mathbf{B} = \begin{pmatrix} \mathbf{A} & -\omega \mathbf{M} \\ \omega \mathbf{M} & \mathbf{A} \end{pmatrix}$$

# Secondary field approach

Consider the field to consist of a primary (background) and an secondary (anomalous) field  $F = F_0 + F_a$

solution for  $F_0$  known, e.g. analytically or 1D (semi-analytically)

⇒ form equations for  $F_a$ , because

- $F_a$  is weaker or smoother (e.g.  $F_0 \propto 1/r$  at sources)
- boundary conditions easier to set (e.g. homogeneous Dirichlet)

# Secondary field Helmholtz equation

The equation  $-\nabla^2 F - k^2 F = 0$  is solved by the primary field for  $k_0$ :

$-\nabla^2 F_0 - k_0^2 F_0 = 0$  and the total field for  $k_0 + \delta k$ :

$$-\nabla^2(F_0 + F_a) - (k_0^2 + \delta k^2)(F_0 + F_a) = 0$$

$$-\nabla^2 F_a - k^2 F_a = \delta k^2 F_0$$

 **Note**

Source terms only arise at anomalous terms, weighted by the primary field.

# Secondary field for E field

$$-\nabla^2 \mathbf{E}_0 + i\omega\mu\sigma_0 \mathbf{E}_0 = 0$$

leads to

$$-\nabla^2 \mathbf{E}_a + i\omega\mu\sigma \mathbf{E}_a = -i\omega\mu\delta\sigma \mathbf{E}_0$$

 **Note**

Source terms only arise at anomalous conductivities and increase with primary field

# Secondary field in matrix form

$$-\nabla^2 \mathbf{E}_a + \omega\mu\sigma \mathbf{E}_a = -\omega\mu\delta\sigma \mathbf{E}_0$$

leads to the discretized form (**A**-stiffness, **M**-mass)

$$\mathbf{A}\mathbf{u}_a + \omega\mathbf{M}_\sigma\mathbf{u}_a = (\mathbf{A} + \omega\mathbf{M}_\sigma)\mathbf{u}_a = -\omega\mathbf{M}_{\delta\sigma}\mathbf{u}_0$$

```
1 A = stiffnessMatrix1DFE(x=z)
2 M = massMatrix1DFE(x=z, a=w*mu*sigma)
3 dM = massMatrix1DFE(x=z, a=w*mu*(sigma-sigma0))
4 u = uAna + solve(A+M*w*1j, dM@uAna * w*1j)
```

# Secondary field for H field

$$-\nabla \cdot \rho_0 \nabla \mathbf{H}_0 + \omega \mu \mathbf{H}_0 = 0$$

leads to

$$-\nabla \cdot \rho \nabla \mathbf{H}_a + \omega \mu \mathbf{H}_a = \nabla \cdot \delta \rho_0 \nabla \mathbf{H}_0$$

$$\mathbf{A}_\rho \mathbf{u}_a + \omega \mathbf{M} \mathbf{u}_a = (\mathbf{A}_\rho + \omega \mathbf{M}) \mathbf{u}_a = -\mathbf{A}_{\delta \rho} \mathbf{u}_0$$

# Task

1. Create a 1D discretization vector (once with air and once without)
2. Use your own functions to generate stiffness and mass matrix
3. Implement Dirichlet BC on the upper and Neumann BC on the lower boundary
4. Choose a period and solve the equation for homogeneous conductivity
5. Compare with the analytical solution
6. Solve for both E and H, play with conductivity and period
7. Add inhomogeneity and repeat computations