

Massachusetts Institute of Technology  
 Department of Electrical Engineering and Computer Science

6.453 QUANTUM OPTICAL COMMUNICATION

**Mid-Term Examination**

Fall 2016

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This examination is closed book except for one  $8\frac{1}{2} \times 11$  handwritten formula sheet (both sides) of your own devising.

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**Problem 1** (20 points)

For each statement below, indicate whether it is True or whether it is False, and provide a brief explanation of your reasoning.

- (a) (10 points) Consider a pair of single-mode electromagnetic fields, with annihilation operators  $\hat{a}_A$  and  $\hat{a}_B$ , whose joint state  $|\psi\rangle_{AB}$  is a pure state. Suppose that the  $\hat{N}_A = \hat{a}_A^\dagger \hat{a}_A$  and  $\hat{N}_B = \hat{a}_B^\dagger \hat{a}_B$  measurements are made on these modes and that the resulting classical outcomes,  $N_A$  and  $N_B$ , have measurement statistics which satisfy

$$\text{Var}(N_A - N_B) < \text{Var}(N_A) + \text{Var}(N_B),$$

where  $\text{Var}(\cdot)$  denotes variance.

**True or False:** The joint state of the  $\hat{a}_A$  and  $\hat{a}_B$  modes *must* be non-classical.

- (b) (10 points) Consider a single-mode electromagnetic field with photon annihilation operator  $\hat{a}$  whose Wigner distribution is  $W(\alpha^*, \alpha)$ .

**True or False:** The function  $F(\alpha_1) \equiv \int_{-\infty}^{\infty} d\alpha_2 W(\alpha^*, \alpha)$ , where  $\alpha_1$  and  $\alpha_2$  are the real and imaginary parts of  $\alpha$ , is non-negative for *all* values of  $\alpha_1$ .

**In case you've forgotten:**

$$W(\alpha^*, \alpha) = \int \frac{d^2\zeta}{\pi^2} \chi_W(\zeta^*, \zeta) e^{\zeta^* \alpha - \zeta \alpha^*},$$

where  $\chi_W(\zeta^*, \zeta) \equiv \langle e^{-\zeta^* \hat{a} + \zeta \hat{a}^\dagger} \rangle$ , and

$$\int \frac{d\alpha_2}{\pi} e^{2j\zeta_1 \alpha_2} = \delta(\zeta_1)$$

where  $\delta(\cdot)$  is the impulse function.

**Problem 2** (40 points)

Consider the asymmetric beam-splitter setup shown in Fig. 1. In this setup, the beam splitter is illuminated by a signal mode (with annihilation operator  $\hat{a}_S$ ) and a local-oscillator (LO) mode (with annihilation operator  $\hat{a}_{\text{LO}}$ ). We will be interested in the output mode from that beam splitter whose annihilation operator is  $\hat{a}_{\text{out}} = \sqrt{\epsilon} \hat{a}_S + \sqrt{1-\epsilon} \hat{a}_{\text{LO}}$ , where  $0 < \epsilon < 1$  and the  $\hat{a}_{\text{LO}}$  mode is in the coherent state  $|\beta \sqrt{\epsilon/(1-\epsilon)}\rangle_{\text{LO}}$ .

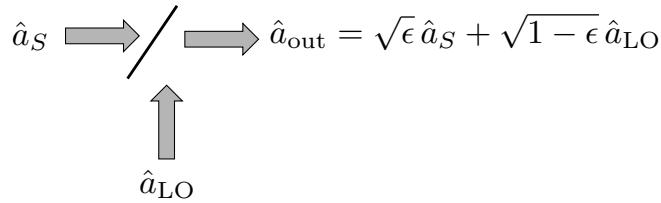


Figure 1: Asymmetric beam-splitter setup.

- (a) (10 points) Suppose that the  $\hat{a}_S$  mode is in the coherent state  $|\gamma\rangle_S$ .
  - (i) With only a *simple* statement of justification, find the state of the  $\hat{a}_{\text{out}}$  mode.
  - (ii) Use your result from (i) to find  $\rho_{a_{\text{out}}}^{(n)}(\alpha^*, \alpha) \equiv {}_{\text{out}}\langle \alpha | \hat{\rho}_{a_{\text{out}}} | \alpha \rangle_{\text{out}}$  in the limit  $\epsilon \rightarrow 1$ .
- (b) (10 points) Figure 2 uses the beam-splitter setup in a photon-counting communication receiver with the following characteristics.
  - The binary message  $b$  being communicated is equally likely to be 0 or 1.
  - When  $b = 0$ , the  $\hat{a}_S$  mode is in the coherent state  $|-\sqrt{N_S}\rangle_S$ . When  $b = 1$ , the  $\hat{a}_S$  mode is in the coherent state  $|\sqrt{N_S}\rangle_S$ .
  - The beam-splitter setup has  $0 < \epsilon < 1$  and  $\beta = \sqrt{N_S}$ .
  - The receiver's output is  $\tilde{b} = 1$  when the  $\hat{N}_{\text{out}} = \hat{a}_{\text{out}}^\dagger \hat{a}_{\text{out}}$  measurement's outcome  $N_{\text{out}}$  is non-zero. The receiver's output is  $\tilde{b} = 0$  when  $N_{\text{out}} = 0$ .
  - (i) Use your result from (a) to find the states that the  $\hat{a}_{\text{out}}$  mode is in when  $b = 0$  and  $b = 1$ .
  - (ii) Use your results from (i) to find the receiver's error probability,  $\Pr(\tilde{b} \neq b)$ .

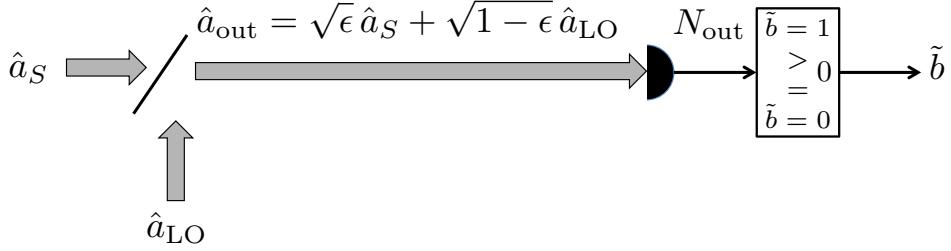


Figure 2: Photon-counting communication receiver.

- (c) (10 points) Now, let the  $\hat{a}_S$  mode be in an *arbitrary* state specified by the density operator  $\hat{\rho}_S$ .
- Find  $\chi_A^{\rho_{a_{\text{out}}}}(\zeta^*, \zeta)$ , the anti-normally ordered characteristic function of the  $\hat{a}_{\text{out}}$  mode. Your answer should be expressed in terms of the  $\hat{a}_S$  mode's anti-normally ordered characteristic function,  $\beta$ , and  $\epsilon$ .
  - Specialize your result from (i) to the limit  $\epsilon \rightarrow 1$ .

- (d) (10 points) For your  $\chi_A^{\rho_{a_{\text{out}}}}(\zeta^*, \zeta)$  from (c), use the operator-valued inverse transform relation,

$$\hat{\rho}_{a_{\text{out}}} = \int \frac{d^2\zeta}{\pi} \chi_A^{\rho_{a_{\text{out}}}(\zeta^*, \zeta)} e^{-\zeta \hat{a}_{\text{out}}^\dagger} e^{\zeta^* \hat{a}_{\text{out}}},$$

to obtain  $\rho_{a_{\text{out}}}^{(n)}(\alpha^*, \alpha) \equiv \langle \alpha | \hat{\rho}_{a_{\text{out}}} | \alpha \rangle_{\text{out}}$  in the  $\epsilon \rightarrow 1$  limit. Your answer should be expressed in terms of  $\rho_S^{(n)}(\alpha^*, \alpha) \equiv \langle \alpha | \hat{\rho}_S | \alpha \rangle_S$ , and  $\beta$ .

**Problem 3** (40 points)

The system shown in Fig. 3 is a quantum non-demolition (QND) setup for measuring the photon number of an optical mode with annihilation operator  $\hat{a}$ . The cross-Kerr-effect box has the following input-output relation:

$$\hat{c} = e^{j\kappa\hat{a}^\dagger\hat{a}}\hat{b}$$

$$\hat{d} = e^{j\kappa\hat{b}^\dagger\hat{b}}\hat{a},$$

where  $\kappa > 0$  is a constant. The homodyne detector is set up to measure the  $\hat{c}_2 = \text{Im}(\hat{c})$  observable.

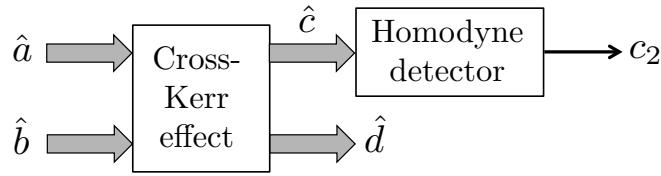


Figure 3: Quantum non-demolition detection setup.

- (a) (10 points) Evaluate the number-ket matrix elements,

$${}_b\langle n_b|_a\langle n_a|e^{j\kappa\hat{a}^\dagger\hat{a}}\hat{b}e^{j\kappa\hat{b}^\dagger\hat{b}}\hat{a}|m_a\rangle_a|m_b\rangle_b$$

and

$${}_b\langle n_b|_a\langle n_a|e^{j\kappa\hat{b}^\dagger\hat{b}}\hat{a}e^{j\kappa\hat{a}^\dagger\hat{a}}\hat{b}|m_a\rangle_a|m_b\rangle_b.$$

- (b) (10 points) Assume that the  $\hat{a}$  mode is in the number state  $|m_a\rangle_a$ . Let  $N_d$  be the outcome of the  $\hat{N}_d = \hat{d}^\dagger\hat{d}$  measurement. Find the probability mass function  $\Pr(N_d = n)$ . **Hint:** You do *not* need to know the state of the  $\hat{b}$  mode.

- (c) (10 points) Assume that the  $\hat{a}$  mode is in the number state  $|m_a\rangle_a$  and the  $\hat{b}$  mode is in the coherent state  $|\sqrt{N_b}\rangle$ . Find  $\langle\hat{c}_2\rangle$  and  $\langle\Delta\hat{c}_2^2\rangle$ , the mean and variance of the  $\hat{c}_2$  measurement.

- (d) (10 points) Assume that the states of the  $\hat{a}$  and  $\hat{b}$  modes are as given in (c), and that  $\kappa m_a \ll 1$ . Let  $c_2$  denote the outcome of the  $\hat{c}_2$  measurement and define  $\tilde{N}_a = c_2/\sqrt{N_b}\kappa$  to be the QND estimate of the  $\hat{a}$  mode's photon number. Find the mean-squared error of this estimate, i.e.,  $\langle(\tilde{N}_a - m_a)^2\rangle$ .

**I'm sure you know these things:**

$$\cos(x) = \frac{e^{jx} + e^{-jx}}{2}, \quad \sin(x) = \frac{e^{jx} - e^{-jx}}{2j}$$

$$\frac{1 - \cos(2x)}{2} = \sin^2(x), \quad \sin(x) \approx x, \text{ for } |x| \ll 1$$

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