数学物理方法 期末考试

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1 第一题

2 第二题

$$z = \frac{1 - i \tan x}{1 + i \tan x} = \frac{\cos x - i \sin x}{\cos x + i \sin x} = \frac{-i \left[\cos \left(\frac{\pi}{2} - x\right) + i \sin \left(\frac{\pi}{2} - x\right)\right]}{\exp \left(ix\right)}$$
$$= \exp\left(i\frac{3\pi}{2}\right) \exp\left(i\frac{\pi}{2}\right) \exp\left(-2ix\right) = \exp\left(-2ix\right)$$
$$= \cos 2x - i \sin 2x$$

3 第三题

$$1 + i = \sqrt{2} \exp\left(i\frac{\pi}{4}\right)$$
$$1 - i = \sqrt{2} \exp\left(-i\frac{\pi}{4}\right)$$
$$(1 + i)^{1000} + (1 - i)^{1000} = 2^{500} \left[\exp\left(250\pi i\right) + \exp\left(-250\pi i\right)\right] = 2^{501}$$

4 第四题

二者模长必然相同,若要求二者相等,只需要辐角满足

$$\arg\left(\sqrt{3}+i\right)^n = 2k\pi + \arg\left(\sqrt{3}-i\right)^n, \quad k \in \mathbb{Z}$$

$$\therefore n\frac{\pi}{6} = -n\frac{\pi}{6} + 2k\pi$$
$$\therefore n = 6k, \quad k \in \mathbb{Z}$$

所以只需要 n 是整数且能被 6 整除即可。

5 第五题

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5 第五题

平面上圆的一般方程可以写成

$$\alpha_1 (x^2 + y^2) + 2\alpha_2 x + 2\alpha_3 y + \alpha_4 = 0$$

其中 $\alpha_{1,2,3,4} \in \mathbb{R}$, 若 $\alpha_1 = 0$, 其退化为直线 若代入

$$x = \frac{z + \bar{z}}{2}, y = \frac{z - \bar{z}}{2i}$$

设 $\beta = \alpha_2 + i\alpha_3$ 有复数表示

$$\alpha_1 z \bar{z} + \bar{\beta} z + \beta \bar{z} + \alpha_4 = 0$$

对于此题,

$$w = \frac{z-i}{z+i} \Rightarrow z = -\frac{(w+1)i}{w-1}, \bar{z} = \frac{(\bar{w}+1)i}{\bar{w}-1}$$

代入化简有

$$\gamma_1 w \bar{w} + \bar{\beta}_2 w + \beta_2 \bar{w} + \gamma_4 = 0$$

其中

$$\begin{pmatrix} \gamma_1 \\ \gamma_2 \\ \gamma_3 \\ \gamma_4 \end{pmatrix} = \begin{pmatrix} \alpha_1 - 2\alpha_3 + \alpha_4 \\ \alpha_1 - \alpha_4 \\ -2\alpha_2 \\ \alpha_1 + 2\alpha_3 + \alpha_4 \end{pmatrix} \qquad \beta_2 = \gamma_2 + i\gamma_3$$

因此广义上(直线是特殊的圆)看,w 将圆变成圆。特别的,对于此题的情况有:

$$\begin{pmatrix} \alpha_1 & \alpha_2 & \alpha_3 & \alpha_4 \end{pmatrix} = \begin{pmatrix} 0 & 1/2 & 0 & -a \end{pmatrix} \rightarrow \begin{pmatrix} \gamma_1 & \gamma_2 & \gamma_3 & \gamma_4 \end{pmatrix} = \begin{pmatrix} -a & a & -1 & -a \end{pmatrix}$$

$$\begin{pmatrix} \alpha_1 & \alpha_2 & \alpha_3 & \alpha_4 \end{pmatrix} = \begin{pmatrix} 0 & 0 & 1/2 & -b \end{pmatrix} \rightarrow \begin{pmatrix} \gamma_1 & \gamma_2 & \gamma_3 & \gamma_4 \end{pmatrix} = \begin{pmatrix} -1 - b & b & 0 & 1 - b \end{pmatrix}$$
 显然当且仅当 $-a = 0, -1 - b = 0$ 即 $a = 0, b = 1$ 时,退化为直线。

6 第六题

$$a = \sum_{k=0}^{\infty} \frac{\cos nz}{n!}$$

$$b = \sum_{n=0}^{\infty} \frac{\sin nz}{n!}$$

7 第七题

考虑做一对替换有:

$$c_{+} \equiv a + bi = \sum_{k=0}^{\infty} \frac{\cos nz + i\sin nz}{n!} = \sum_{k=0}^{\infty} \frac{\exp(inz)}{n!} = \exp(\exp(iz))$$

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$$c_{-} \equiv a - bi = \sum_{k=0}^{\infty} \frac{\cos nz - i \sin nz}{n!} = \sum_{k=0}^{\infty} \frac{\exp(-inz)}{n!} = \exp(\exp(-iz))$$

其中

$$\exp(\exp(iz)) = \exp(\cos z + i\sin z) = e^{\cos z}(\cos\sin z + \sin\sin z)$$

$$\exp(\exp(-iz)) = \exp(\cos z - i\sin z) = e^{\cos z}(\cos\sin z - \sin\sin z)$$

于是由替换的关系可以得到:

$$a = \frac{c_+ + c_-}{2} = e^{\cos z} \cos \sin z$$

$$b = \frac{c_+ - c_-}{2i} = -ie^{\cos z} \sin \sin z$$

7 第七题

求

$$f\left(z\right) = \frac{1}{e^{z} - 1}$$

在 z=0 的洛朗级数, 很明显此处不解析, 考察 $n \in N^*$ 时,

$$\lim_{z \to 0} \frac{z^n}{e^z - 1} = \begin{cases} 1 & n = 1 \\ 0 & n > 1 \end{cases}$$

因此这是一阶极点,考虑求解析函数 a(z)=zf(z) 在 z=0 的泰勒展开,设 其为

$$a\left(z\right) = \sum_{n=0}^{+\infty} a_n z^n$$

这不容易计算, 考虑其倒数 b(z) = 1/a, 可以展开为

$$b(z) = \sum_{n=0}^{+\infty} \frac{z^n}{(n+1)!}$$

又因为恒等关系 a(z)b(z) = 1

$$\left(\sum_{n=0}^{+\infty} a_n z^n\right) \cdot \left(\sum_{n=0}^{+\infty} \frac{z^n}{(n+1)!}\right) = 1$$

比较 z 的各阶系数可以得到 $a_0=1, a_1=-\frac{1}{2}...$

8 第八题 4

在 $n \ge 1$ 时,其满足线性齐次方程组方程为:

$$\sum_{k=0}^{n} \frac{a_k}{(n-k+1)!} = 0$$

理论上可以解出任意的 a_n

我试图找到通项公式但是失败了,查阅资料发现这个数列与伯努利数相关, 关系为 $a_n = B_n/n!$ 其被定义为

$$\frac{z}{e^z - 1} = \sum_{n=0}^{\infty} B_n \frac{z^n}{n!}$$

递推关系为

$$B_m = [m = 0] - \sum_{k=0}^{m-1} \binom{m}{k} \frac{B_k}{m-k+1}$$

 $B_0 = 1$, 其中 [m = 0] 表示当 m = 0 时,取 1,其余取 0。 综上可得,

$$f(z) = \sum_{n=-1}^{+\infty} \frac{z^n}{(n+1)!} B_{n+1}$$

8 第八题

由于所有极点都在围道内部,直接考察无穷远的留数,根据引理有:

$$\int_{|z|=200} f(z) dz = -2\pi i \mathop{\rm Res}_{z \to \infty} f(z) = 2\pi i \mathop{\rm Res}_{t \to 0} \left[f\left(\frac{1}{t}\right) \frac{1}{t^2} \right]$$

其中

$$f\left(\frac{1}{t}\right)\frac{1}{t^2} = \frac{1}{t^2} \prod_{k=1}^{100} \frac{1}{1 - kt}$$
$$\because \lim_{t \to 0} t^2 \frac{1}{t^2} \prod_{k=1}^{100} \frac{1}{1 - kt} = 1$$
$$\therefore 二阶极点$$

所以可以计算出留数的值

$$\operatorname{Res}_{t \to 0} \left[f\left(\frac{1}{t}\right) \frac{1}{t^2} \right] = \lim_{t \to 0} \frac{d}{dt} \prod_{k=1}^{100} \frac{1}{1 - kt}$$

$$= \lim_{t \to 0} \sum_{n=1}^{100} \prod_{n=1 \land k \neq n}^{100} \frac{n}{1 - kt}$$

$$= \sum_{n=1}^{100} n = 5050$$

$$\therefore \int_{|z| = 200} f(z) \, dz = 10100\pi i$$

9 第九题

9 第九题

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考虑配对,设

$$I_{1} = \int_{0}^{2\pi} e^{\cos \theta} \cos (n\theta - \sin \theta) d\theta$$

$$I_{2} = \int_{0}^{2\pi} e^{\cos \theta} \sin (n\theta - \sin \theta) d\theta$$

$$I_{1} + iI_{2} = \int_{0}^{2\pi} \exp (\cos \theta - i \sin \theta) \exp (in\theta) d\theta$$

$$I_{1} - iI_{2} = \int_{0}^{2\pi} \exp (\cos \theta + i \sin \theta) \exp (-in\theta) d\theta$$

对于 $I_1 + iI_2$ 设 $z = e^{i\theta}$, 有

$$I_1 + iI_2 = \int_{|z|=1} e^{\frac{1}{z}} z^n \frac{dz}{zi} = 2\pi \underset{z \to 0}{\text{Res}} e^{\frac{1}{z}} z^{n-1}$$

包裹的奇点只在 z=0, 我们计算其在 z=0 洛朗展开为

$$e^{\frac{1}{z}}z^{n-1} = \sum_{k=-\infty}^{0} \frac{e^{n+k-1}}{(-k)!}$$

为求其留数, 取 z^{-1} 项要求 k = -n 所以有

$$\underset{z\to 0}{\text{Res}} e^{\frac{1}{z}} z^{n-1} = \frac{1}{n!}$$

对于 $I_1 - iI_2$ 设 $z = e^{-i\theta}$, 注意积分方向相反需要取出一个负号, 化简有

$$I_1 - iI_2 = \int_{|z|=1} e^{\frac{1}{z}} z^n \frac{dz}{zi} = I_1 + iI_2 = \frac{2\pi}{n!}$$

于是要求的量即为

$$I_1 = \frac{2\pi}{n!}$$

10 第十题

直接考察积分, 试图利用柯西公式

$$D_{F}(x-y) = \int_{C_{F}} \frac{d^{4}p}{(2\pi)^{4}} \frac{ie^{-ip\cdot(x-y)}}{p^{2} - m^{2}}$$

$$= \int_{C_{F}} \frac{d^{3}\vec{p}dp_{0}}{(2\pi)^{4}} \frac{ie^{i\vec{p}\cdot(\vec{x}-\vec{y})}e^{-ip_{0}(x_{0}-y_{0})}}{p_{0}^{2} - (\vec{p}^{2} + m^{2})}$$

$$= \int \frac{d^{3}\vec{p}\left[e^{i\vec{p}\cdot(\vec{x}-\vec{y})}\right]}{(2\pi)^{4}} \int_{C_{F}} \frac{ie^{-ip_{0}(x_{0}-y_{0})}dp_{0}}{p_{0}^{2} - E_{\vec{p}}^{2}}$$

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考察其中的

$$\int_{C_F} \frac{ie^{-ip_0(x_0-y_0)}dp_0}{p_0^2 - E_{\vec{p}}^2} = \int_{C_F} \frac{ie^{-ip_0(x_0-y_0)}dp_0}{(p_0 - E_{\vec{p}})(p_0 + E_{\vec{p}})}$$

注意到在 C_F 的围道内只有一个一阶极点为

$$\mathop{\rm Res}_{p_0 \to E_{\vec{p}}} \frac{i e^{-i p_0 (x_0 - y_0)}}{(p_0 - E_{\vec{p}}) \left(p_0 + E_{\vec{p}}\right)} = \frac{i e^{-i p_0 (x_0 - y_0)}}{2 E_{\vec{p}}}$$

考察其在无穷远的性质有

$$\lim_{|p_0| \to \infty} \left| \frac{ip_0 e^{-ip_0(x_0 - y_0)}}{p_0^2 - E_{\vec{p}}^2} \right| = 0$$

实际上是一致收敛的

注意到 C_F 的围道方向为顺时针,因此有

$$\int_{C_F} \frac{ie^{-ip_0(x_0 - y_0)} dp_0}{p_0^2 - E_{\vec{p}}^2} = 2\pi \frac{e^{-ip_0(x_0 - y_0)}}{2E_{\vec{p}}}$$

代入 Feynman 传播子表示可得到:

$$D_F(x-y) = \int_{C_F} \frac{d^4 p}{(2\pi)^4} \frac{ie^{-ip\cdot(x-y)}}{p^2 - m^2} = \int \frac{d^3 \vec{p}}{(2\pi)^3 2E_{\vec{p}}} e^{-ip\cdot(x-y)} \bigg|_{p_0 = E_{\vec{p}}}$$

11 第十一题

以下都建立在 D=2 上

引理 11.1. 导数定理

$$f(x) = \mathcal{F}(g(k)) \Rightarrow \partial_{\mu} f(x) = \mathcal{F}(-ik_{\mu}g(k))$$

证明.

$$\mathcal{F}^{-1}(\partial_{\mu}f(x)) = \frac{1}{2\pi} \int d^2x \exp(ik \cdot x) \partial_{\mu}f(x)$$
$$= -\frac{1}{2\pi} \int d^2x (ik_{\mu}) \exp(ik \cdot x) f(x) = -ik_{\mu}g(k)$$
$$\therefore \partial_{\mu}f(x) = \mathcal{F}(-ik_{\mu}g(k))$$

同理,

$$\partial^{\mu} f(x) = \mathcal{F}(-ik^{\mu}g(k))$$

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引理 11.2. 卷积定理

如果定义

$$f * g(x) \equiv \int d^2\xi f(\xi)g(x - \xi)$$

则有

$$\mathcal{F}\left(f * g\left(x\right)\right) = 2\pi \mathcal{F}\left(f\left(x\right)\right) \mathcal{F}\left(g\left(x\right)\right)$$

证明.

$$\begin{split} \mathcal{F}\left(f*g\left(x\right)\right) &= \frac{1}{2\pi} \int d^2x e^{-ik\cdot x} \int d^2\xi f\left(\xi\right) g\left(x-\xi\right) \\ &= \int d^2\xi f\left(\xi\right) \left[\frac{1}{2\pi} \int d^2x e^{-ik\cdot x} g\left(x-\xi\right)\right] \\ &= \mathcal{F}\left(g\left(x\right)\right) \int d^2\xi f\left(\xi\right) e^{-ik\cdot \xi} \\ &= 2\pi \mathcal{F}\left(f\left(x\right)\right) \mathcal{F}\left(g\left(x\right)\right) \end{split}$$

推论:对于三个函数的情形,只需要依次计算,即为

$$\mathcal{F}\left(f*g*h\left(x\right)\right) = 4\pi^{2}\mathcal{F}\left(f\left(x\right)\right)\mathcal{F}\left(g\left(x\right)\right)\mathcal{F}\left(h\left(x\right)\right)$$

引理 11.3. 设对于 f(x) 有 $f(x) = \mathcal{F}(g(k))$ 则,

$$\int d^2x \mathcal{F}(g(k)) = (2\pi)^2 g(0)$$

证明.

$$\int d^2x \mathcal{F}(g(k)) = \int d^2x \int d^2k \exp(-ik \cdot x) g(k)$$
$$= \int d^2k g(k) \int d^2x \exp(-ik \cdot x)$$
$$= (2\pi)^2 \int d^2k g(k) \delta(-k) = (2\pi)^2 g(0)$$

我们考虑原始表达式

$$S = \int d^{2}x \left\{ -\frac{1}{2} \partial_{\mu} \phi \left(x \right) \partial^{\mu} \phi \left(x \right) + \frac{1}{2} \lambda_{\mu v} \left(x \right) \left[\partial^{\mu} \phi \left(x \right) - \epsilon^{\mu \sigma} \partial_{\sigma} \phi \left(x \right) \right] \left[\partial^{v} \phi \left(x \right) - \epsilon^{v \rho} \partial_{\rho} \phi \left(x \right) \right] \right\}$$

先分别计算其中的积分

$$S_{1} = \int d^{2}x \left[\partial_{\mu}\phi \left(x \right) \partial^{\mu}\phi \left(x \right) \right]$$

$$S_{2} = \int d^{2}x \left\{ \lambda_{\mu\nu} \left(x \right) \left[\partial^{\mu}\phi \left(x \right) - \epsilon^{\mu\sigma}\partial_{\sigma}\phi \left(x \right) \right] \left[\partial^{\nu}\phi \left(x \right) - \epsilon^{\nu\rho}\partial_{\rho}\phi \left(x \right) \right] \right\}$$

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$$\therefore \phi(x) = \mathcal{F}(\phi(k)), \lambda_{\mu\nu}(x) = \mathcal{F}(\lambda_{\mu\nu}(k))$$

换元,令

$$g_{\mu} = -ik_{\mu}\phi\left(k\right), g^{\mu} = -ik^{\mu}\phi\left(k\right)$$

利用导数定理 11.1

$$\therefore \mathcal{F}\left(-ik_{\mu}\phi\left(k\right)\right) = \partial_{\mu}\phi\left(x\right), \mathcal{F}\left(-ik^{\mu}\phi\left(k\right)\right) = \partial^{\mu}\phi\left(x\right)$$

利用卷积定理 11.2

$$S_{1} = \int d^{2}x \mathcal{F}(g_{\mu}) \mathcal{F}(g^{\mu}) = \frac{1}{2\pi} \int d^{2}x \mathcal{F}(g_{\mu} * g^{\mu})$$

最后利用 11.3

$$\therefore S_{1} = (g_{\mu} * g^{\mu})|_{k=0} = \int d^{2}\xi g_{\mu}(\xi) g^{\mu}(0 - \xi) = \int d^{2}k (-ik_{\mu}) \phi(k) (-ik^{\mu}) \phi(k)$$

同理我们设

$$\eta^{\mu} = ik^{\mu}\phi(k) - \epsilon^{\mu\sigma}ik_{\sigma}\phi(k), \eta^{\nu} = ik^{\nu}\phi(k) - \epsilon^{\nu\rho}ik_{\rho}\phi(k)$$

有

$$S_{2} = \int d^{2}x \mathcal{F}(\lambda_{\mu\nu}(k)) \mathcal{F}(\eta^{\mu}) \mathcal{F}(\eta^{\nu}) = \frac{1}{4\pi^{2}} \int d^{2}x \mathcal{F}(\lambda_{\mu\nu}(k) * \eta^{\mu} * \eta^{\nu})$$

$$S_{2} = \frac{1}{2\pi} \lambda_{\mu\nu}(k) * \eta^{\mu} * \eta^{\nu}|_{k=0} = \frac{1}{2\pi} \int d^{2}\xi \int d^{2}\zeta \lambda_{\mu\nu}(\zeta + \xi) \eta^{\mu}(0 - \xi) \eta^{\nu}(0 - \zeta)$$
即为

$$S_{2} = \frac{1}{2\pi} \int d^{2}k \int d^{2}k' \lambda_{\mu\nu} \left(-k-k'\right) \left[ik^{\mu}\phi\left(k\right) - \epsilon^{\mu\sigma}ik_{\sigma}\phi\left(k\right)\right] \left[ik^{\nu}\phi\left(k'\right) - \epsilon^{\nu\rho}ik_{\rho}\phi\left(k'\right)\right]$$

现在我们终于可以代回到原表达式

$$S_{m} = -\frac{1}{2} \int d^{2}k \left(-ik_{\mu}\right) (ik^{\mu}) \phi(k) \phi(-k)$$

$$+ \frac{1}{4\pi} \int d^{2}k d^{2}k' \lambda_{\mu\nu} \left(-k - k'\right) (ik^{\mu} - \epsilon^{\mu\sigma}ik_{\sigma}) \left(ik'^{\nu} - \epsilon^{\nu\rho}ik'_{\rho}\right) \phi(k) \phi(k')$$

这就是结果

12 第十二题

12.1

$$g_1(t) = \begin{cases} 2t/T & 0 \le t < T/2\\ 2(1 - t/T) & T/2 \le t \le T \end{cases}$$

12 第十二题

9

根据周期函数

$$\mathcal{L}\left(g\left(t\right)\right) = \frac{1}{1 - e^{-pT}} \int_{0}^{T} e^{-pt} g(t) dt$$

计算得

$$\int_{0}^{T/2} \exp(-pt) \frac{2t}{T} dt = \frac{2}{T} \left[\frac{1}{p^2} - e^{-Tp/2} \left(\frac{T}{2p} + \frac{1}{p^2} \right) \right]$$
$$- \int_{T/2}^{T} \exp(-pt) \frac{2t}{T} dt = \frac{2}{T} e^{-Tp/2} \left[\left(1 - e^{-Tp/2} \right) \left(\frac{T}{2p} + \frac{1}{p^2} \right) \right]$$
$$\int_{T/2}^{T} 2 \exp(-pt) dt = \frac{2}{p} \left(e^{-Tp/2} - e^{-Tp} \right)$$

代入化简得

$$\mathcal{L}\left(g_{1}\left(t\right)\right) = \frac{2}{Tp^{2}} \tan h\left(\frac{Tp}{4}\right)$$

12.2

$$g_{2}(t) = \frac{2}{T} \left[\text{tH}(t) + 2 \sum_{n=1}^{\infty} (-1)^{n} \left(t - \frac{1}{2} nT \right) H \left(t - \frac{1}{2} nT \right) \right]$$
$$\mathcal{L}(g_{2}(t)) = \frac{2}{Tp^{2}} \left[1 + 2 \sum_{n=1}^{\infty} (-1)^{n} \exp \left(-\frac{1}{2} nT \right) \right]$$

如果 $g_1(t) = g_2(t)$, 比照系数令 x = Tp/4 换元有

$$\tanh x = 1 + 2\sum_{n=0}^{\infty} (-1)^n e^{-2nx}$$

这正是我们想证明的。因此只需证明 $g_1(t) = g_2(t)$.

下面证明 $g_1(t) = g_2(t)$

我们首先考察区间 [0,T/2), 对于 $t \in [0,T/2]$ 由于全部的 H(t-1/2nT)=0 因此二者显然相等。

然后考察区间 [T/2,T), 对于 $t \in [0,T/2]$ 求和只保留到 n=1,有

$$g_2 = \frac{2}{T} \left[t - 2 \left(t - \frac{1}{2} T \right) \right] = 2 - \frac{2t}{T} = g_1$$

综上二者在 $t \in [0,T)$ 上相等。

一般地,对于 n > 0 情形,仅当 $t \ge 1/2nT$ 时,H(t - 1/2nT) = 1,其余情况为 0。如果我们只考察 t 非负半轴情况,求和可以写成:

$$g_2(t) = \frac{2}{T} \left[t + 2 \sum_{n=1}^{\left[\frac{2t}{T}\right]} (-1)^n \left(t - \frac{1}{2} nT \right) \right]$$

其中 [2] 表示括号内值向下取整。

接下来证明 $g_2(t)$ 也以 T 为周期,既 $g_2(t) = g_2(t+T)$

$$g_{2}(t+T) = \frac{2}{T} \left[t + T + 2 \sum_{n=1}^{\left[\frac{2t}{T}+2\right]} (-1)^{n} \left(t - \frac{1}{2}nT + T \right) \right]$$

$$= \frac{2}{T} \left[t + T - (2t+T) + 2t + 2 \sum_{n=3}^{\left[\frac{2t}{T}+2\right]} (-1)^{n} \left(t - \frac{1}{2}nT + T \right) \right]$$

$$= \frac{2}{T} \left[t + 2 \sum_{n=1}^{\left[\frac{2t}{T}\right]} (-1)^{n} \left(t - \frac{1}{2}nT \right) \right]$$

$$= g_{2}(t)$$

由于在 [0,T) 区间相等并都具有以 T 为周期的周期性,因此 $g_1(t)=g_2(t)$,这正是我们需要的。于是

$$\tanh x = 1 + 2\sum_{n=0}^{\infty} (-1)^n e^{-2nx}$$

得证

13 第十三题

以下的拉普拉斯方程的表达式都基于

$$\frac{1}{\sqrt{|g|}}\partial_{\mu}\left(g^{\mu\nu}\sqrt{|g|}\partial_{\nu}\psi\right) = 0$$

在正交曲线坐标系下为

$$\sum_{\mu=1}^{3}\frac{1}{|h_1h_2h_3|}\partial_{\mu}\left(\frac{1}{h_{\mu}}\sqrt{|h_1h_2h_3|}\partial_{\mu}\psi\right)=0$$

13.1 球坐标

$$h_r = 1, h_\theta = r, h_\omega = r \sin \theta$$

拉普拉斯方程为:

$$\frac{1}{r^2}\frac{\partial}{\partial r}\left(r^2\frac{\partial u}{\partial r}\right) + \frac{1}{r^2\sin\theta}\frac{\partial}{\partial \theta}\left(\sin\theta\frac{\partial u}{\partial \theta}\right) + \frac{1}{r^2\sin^2\theta}\frac{\partial^2 u}{\partial \varphi^2} = 0$$

首先,设

$$u(r, \theta, \varphi) = R(r)Y(\theta, \varphi)$$

$$\frac{1}{R}\frac{\mathrm{d}}{\mathrm{d}r}\left(r^2\frac{\mathrm{d}R}{\mathrm{d}r}\right) = -\frac{1}{\mathrm{Y}\sin\theta}\frac{\partial}{\partial\theta}\left(\sin\theta\frac{\partial\mathrm{Y}}{\partial\theta}\right) - \frac{1}{\mathrm{Y}\sin^2\theta}\frac{\partial^2\mathrm{Y}}{\partial\varphi^2} = l(l+1)$$

分解为两个方程

$$\frac{d}{dr}\left(r^2\frac{dR}{dr}\right) - l(l+1)R \equiv 0$$

$$\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial Y}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2 Y}{\partial \varphi^2} + l(l+1)Y = 0$$

进一步设

$$\begin{split} \mathbf{Y}(\theta,\varphi) &= \Theta(\theta)\Phi(\varphi) \\ \frac{\sin\theta}{\Theta}\frac{\mathrm{d}}{\mathrm{d}\theta}\left(\sin\theta\frac{\mathrm{d}\Theta}{\mathrm{d}\theta}\right) + l(l+1)\sin^2\theta = -\frac{1}{\Phi}\frac{\mathrm{d}^2\Phi}{\mathrm{d}\varphi^2} = \lambda \end{split}$$

分解为两个常微分方程:

$$\Phi'' + \lambda \Phi = 0$$

$$\sin \theta \frac{\mathrm{d}}{\mathrm{d}\theta} \left(\sin \theta \frac{\mathrm{d}\Theta}{\mathrm{d}\theta} \right) + \left[l(l+1)\sin^2 \theta - \lambda \right] \theta = 0$$

13.2 柱坐标

$$h_{\rho} = 1, h_{\varphi} = \rho, h_z = 1$$

柱坐标的拉普拉斯方程表示为

$$\frac{1}{\rho}\frac{\partial}{\partial\rho}\left(\rho\frac{\partial u}{\partial\rho}\right) + \frac{1}{\rho^2}\frac{\partial^2 u}{\partial\varphi^2} + \frac{\partial^2 u}{\partial z^2} = 0$$

设

$$u(\rho, \varphi, z) = R(\rho)\Phi(\varphi)Z(z)$$

代入得到

$$\Phi'' + \lambda \Phi = 0$$

以及

$$\frac{\rho^2}{R}\frac{\mathrm{d}^2R}{\mathrm{d}\rho^2} + \frac{\rho}{R}\frac{\mathrm{d}R}{\mathrm{d}\rho} + \rho^2\frac{Z''}{Z} = \lambda$$

分离有:

$$\frac{1}{R}\frac{\mathrm{d}^2R}{\mathrm{d}\rho^2} + \frac{1}{\rho R}\frac{\mathrm{d}R}{\mathrm{d}\rho} - \frac{m^2}{\rho^2} = -\frac{Z^{\prime\prime}}{Z} = -\mu$$

两个常微分方程

$$Z'' - \mu Z = 0$$
$$\frac{\mathrm{d}^2 R}{\mathrm{d}\rho^2} + \frac{1}{\rho} \frac{\mathrm{d}R}{\mathrm{d}\rho} + \left(\mu - \frac{m^2}{\rho^2}\right) R = 0$$

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13.3 椭圆柱坐标

标准表示为

$$x = a\xi\eta, \quad y = a\sqrt{(\xi^2 - 1)(1 - \eta^2)}, \quad z = z$$

换元,设

$$\xi = \operatorname{ch} u, \quad \eta = \cos v$$

有

$$h_u^2 = h_v^2 = a^2 \left(\cosh^2 u - \cos^2 v \right), \quad h_z = 1$$

拉普拉斯方程表示为

$$\nabla^2 \Phi = \frac{1}{a^2 \left(\operatorname{ch}^2 u - \cos^2 v \right)} \left\{ \frac{\partial^2 \Phi}{\partial u^2} + \frac{\partial^2 \Phi}{\partial v^2} \right\} + \frac{\partial^2 \Phi}{\partial z^2} = 0$$

设

$$\Phi\left(u,v,z\right) = U\left(u\right)V\left(v\right)Z\left(z\right)$$

代入有

$$\frac{1}{a^2 \left(\operatorname{ch}^2 u - \cos^2 v \right)} \left\{ \frac{1}{U} \frac{d^2 U}{du^2} + \frac{1}{V} \frac{d^2 V}{dv^2} \right\} + \frac{1}{Z} \frac{d^2 Z}{dz^2} = 0$$

分离出

$$\frac{d^2Z}{dz^2} + \lambda Z = 0$$

另有

$$\frac{1}{U}\frac{d^2U}{du^2} + \frac{1}{V}\frac{d^2V}{dv^2} - \lambda a^2 \left(\cosh^2 u - \cos^2 v\right) = 0$$

分离出

$$\frac{d^2U}{du^2} - \left(\lambda a^2 c h^2 u + \mu\right) U = 0$$
$$\frac{d^2V}{dv^2} + \left(\lambda a^2 \cos^2 v + \mu\right) V = 0$$

13.4 椭球坐标

$$x^{2} = \frac{(a^{2} + \lambda) (a^{2} + \mu) (a^{2} + v)}{(a^{2} - b^{2}) (a^{2} - c^{2})}$$
$$y^{2} = \frac{(b^{2} + \lambda) (b^{2} + \mu) (b^{2} + v)}{(b^{2} - c^{2}) (b^{2} - a^{2})}$$
$$z^{2} = \frac{(c^{2} + \lambda) (c^{2} + \mu) (c^{2} + v)}{(c^{2} - a^{2}) (c^{2} - b^{2})}$$

 λ, μ, ν 的变化范围为

$$\lambda > -c^2 > \mu > -b^2 > \nu > -a^2$$

若设

$$\varphi(\theta) = \left(a^2 + \theta\right) \left(b^2 + \theta\right) \left(c^2 + \theta\right)$$

有

$$\begin{split} h_{\lambda}^2 &= \frac{(\lambda - \mu)(\lambda - \nu)}{4\varphi(\lambda)} \\ h_{\mu}^2 &= \frac{(\mu - \lambda)(\mu - \nu)}{4\varphi(\mu)} \\ h_{\nu}^2 &= \frac{(\nu - \lambda)(\nu - \mu)}{4\varphi(\nu)} \\ \nabla^2 \Phi &= \frac{2}{(\lambda - \mu)(\lambda - \nu)(\mu - \nu)} \left[(\mu - \nu)\sqrt{\varphi(\lambda)} \frac{\partial}{\partial \lambda} \left(\sqrt{\varphi(\lambda)} \frac{\partial \Phi}{\partial \lambda} \right) \right. \\ &+ (\lambda - \nu)\sqrt{-\varphi(\mu)} \frac{\partial}{\partial \mu} \left(\sqrt{-\varphi(\mu)} \frac{\partial \Phi}{\partial \mu} \right) \\ &+ (\lambda - \mu)\sqrt{\varphi(\nu)} \frac{\partial}{\partial \nu} \left(\sqrt{\varphi(\nu)} \frac{\partial \Phi}{\partial \nu} \right) \right] = 0 \end{split}$$

设 $\Phi = \Lambda(\lambda)M(\mu)N(\nu)$ 代入有

$$\frac{\mu - \nu}{\Lambda} \sqrt{\varphi(\lambda)} \frac{\mathrm{d}}{\mathrm{d}\lambda} \left(\sqrt{\varphi(\lambda)} \frac{\mathrm{d}\Lambda}{\mathrm{d}\lambda} \right) + \frac{\lambda - \nu}{M} \sqrt{-\varphi(\mu)} \frac{\mathrm{d}}{\mathrm{d}\mu} \left(\sqrt{-\varphi(\mu)} \frac{\mathrm{d}M}{\mathrm{d}\mu} \right) + \frac{\lambda - \mu}{N} \sqrt{\varphi(\nu)} \frac{\mathrm{d}}{\mathrm{d}\nu} \left(\sqrt{\varphi(\nu)} \frac{\mathrm{d}N}{\mathrm{d}\nu} \right) = 0$$

注意到有恒等式

$$(\mu - \nu)(K\lambda + C) + (\nu - \lambda)(K\mu + C) + (\lambda - \mu)(K\nu + C) \equiv 0$$

其中 K 和 C 为常数, 比较系数得

$$4\sqrt{\varphi(\lambda)}\frac{\mathrm{d}}{\mathrm{d}\lambda}\left(\sqrt{\varphi(\lambda)}\frac{\mathrm{d}\Lambda}{\mathrm{d}\lambda}\right) = (K\lambda + C)\Lambda$$

设 K = n(n+1), 分离变量的三个方程为:

$$4\sqrt{\varphi(\lambda)} \frac{\mathrm{d}}{\mathrm{d}\lambda} \left(\sqrt{\varphi(\lambda)} \frac{\mathrm{d}\Lambda}{\mathrm{d}\lambda} \right) = \left[(n(n+1)) \lambda + C \right] \Lambda$$
$$4\sqrt{\varphi(\mu)} \frac{\mathrm{d}}{\mathrm{d}\mu} \left(\sqrt{\varphi(\mu)} \frac{\mathrm{d}M}{\mathrm{d}\mu} \right) = \left[(n(n+1)) \mu + C \right] M$$
$$4\sqrt{\varphi(\nu)} \frac{\mathrm{d}}{\mathrm{d}\nu} \left(\sqrt{\varphi(\nu)} \frac{\mathrm{d}N}{\mathrm{d}\nu} \right) = \left[(n(n+1)) \nu + C \right] N$$

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- 13.5 锥面坐标
- 13.6 抛物线柱坐标

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