Quantum Circuit Simulation of Charge-Spin Separation in 1D Hubbard Model

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Introduction

- Spin-charge separation is a phenomenon observed in 1D strongly correlated systems.
- ► Luttinger's Liquid theory describes the low energy excitations in 1D electron gas as of electrons into *Spinons* and *Holons*.

System

$$\mathcal{H} = -J\sum_{j=1}^{L-1}\sum_{\nu=\uparrow,\downarrow}c_{j,\nu}^{\dagger}c_{j+1,\nu} + h.c. + U\sum_{j=1}^{L}n_{j,\uparrow}n_{j,\downarrow} + \sum_{j=1}^{L}\sum_{\nu=\uparrow,\downarrow}\epsilon_{j,\nu}n_{j,\nu}$$

- ▶ 1D Fermi-Hubbard model on L lattice sites with open-boundary conditions.
 - $ightharpoonup \epsilon_{j,
 u}$ represents spin-dependent local potentials

$$\rho_j^{\pm} = \langle n_{j,\uparrow} \rangle \pm \langle n_{j,\downarrow} \rangle$$

► Charge and spin densities are defined as the sum and difference of the spin-up and down particle densities respectively.

Motivation

- ► Transport of electrons through nanowires is subject to spin-charge separation
- ► Advancements towards spintronics
- ightharpoonup High T_c superconductivity

Google Quantum Al Experiment

- Simulated Hubbard model on a programmable superconducting quantum processor.
 - In highly excited regime
 - ► Trapping potentials were abruptly removed
 - ► The on-site interactions were suddenly activated
- ► Introduced accurate gate calibration procedure fast enough to capture temporal drifts of the gate parameters.

Compare with Numerical Simulations

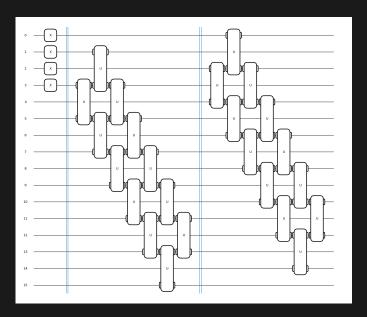
- For comparison, OpenFermion-FQE, a specialized quantum emulator for simulating Fermionic many-body problems was used.
- ► We have recreated the results of numerical simulation using another package *Pennylane*.

Initial State

- ► Exact Diagonalization
 - ► Ground state is obtained by diagonalizing the 4 particle non interacting Hamiltonian.
- Can be implemented using Givens Network
 - ► Initial state prepared is in one particle basis.
 - Converted to a state in two mode fermionic basis by performing a basis transformation using Given's rotation.

$$G = egin{pmatrix} 1 & 0 & 0 & 0 \ 0 & \cos heta & -\sin heta & 0 \ 0 & \sin heta & \cos heta & 0 \ 0 & 0 & 0 & 1 \end{pmatrix}$$

Given's Rotation Circuit



Trotterization

▶ Trotter's Product

$$|\psi(t)
anglepprox \left(\prod_{j}e^{-i\mathsf{a}_{j}\mathcal{H}_{j}dt}
ight)^{n}|\psi(0)
angle$$

Time independent Hamiltonian H can be written as a weighted sum of Pauli terms which may or may not commute.

$$\mathcal{H} = \sum_{j} a_{j} h_{j}$$

Product of matrix exponentials is a good approximation for the exponent of the sum.

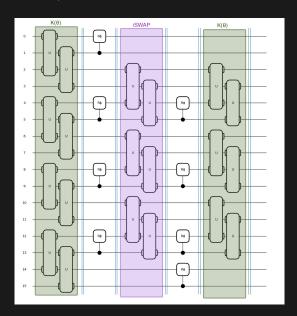
Operators Required for Trotterization

$$K = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos\theta & -i\sin\theta & 0 \\ 0 & -i\sin\theta & \cos\theta & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

$$iSWAP = K(-\frac{\pi}{2})$$

$$CPhase = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & e^{-i\phi} \end{pmatrix}$$

Trotterization Step Circuit



Measurement

► Jordan-Wigner Transformation

$$a_0^{\dagger} = \left(\frac{X_0 - iY_0}{2}\right), \dots, a_n^{\dagger} = Z_0 \otimes Z_1 \otimes \dots \otimes Z_{n-1} \otimes \left(\frac{X_n - iY_n}{2}\right)$$

$$a_0 = \left(\frac{X_0 + iY_0}{2}\right), \ldots, a_n = Z_0 \otimes Z_1 \otimes \ldots \otimes Z_{n-1} \otimes \left(\frac{X_n + iY_n}{2}\right)$$

► The number operator transforms to

$$n_j = \frac{\mathbb{1} - Z_j}{2}$$

Charge and Spin Density for U = 0

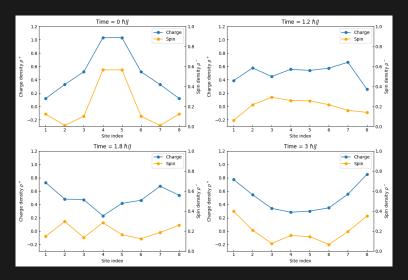


Figure: Evolution of charge and spin density at each site

Charge and Spin Density for U = 3

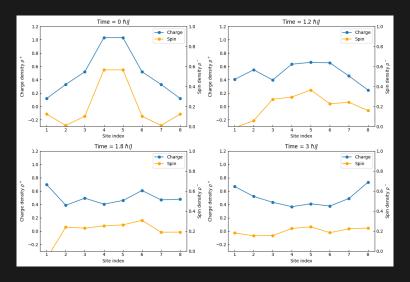
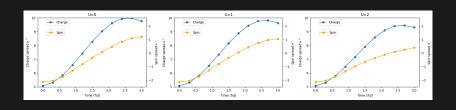


Figure: Evolution of charge and spin densities at each site

Charge and Spin Spread

$$\kappa_{\mu}^{\pm} = \sum_{j} |j - \frac{L+1}{2}| \rho_{\mu,j}^{\pm}$$

Charge and Spin Spreads



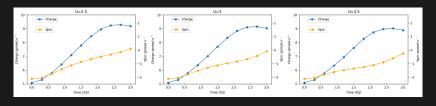


Figure: Time evolution of charge and spin spreads with varying U

More on Charge & Spin Density Waves

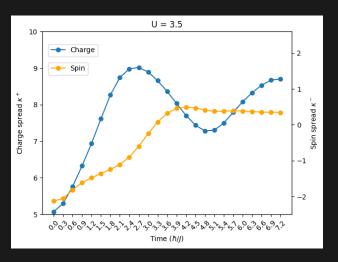
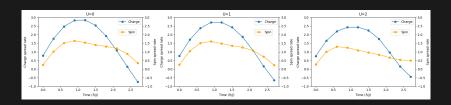


Figure: Damped oscillation like behaviour in spin and charge density waves

Rate of Change of Spreads



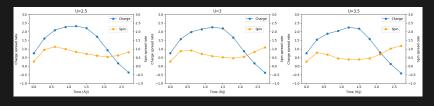


Figure: Numerical derivative of charge and spin spread with respect to evolution time

Mott Insulator Regime: 8 electron system

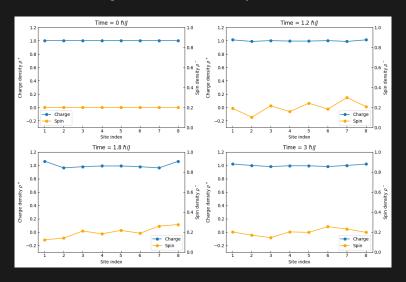


Figure: Evolution of charge and spin densities for 8 electron system

Sources of Errors

- ▶ During the calculation of the initial state some terms were truncated to 0 due to machine precision.
- ▶ First order Trotter formula is used with error $\mathcal{O}(t^2/n)$.
- CPhase gate is not a Clifford gate and hence not simulable to a good precision.
- ▶ Did not include $V \sum_{j=1}^{L-1} \sum_{\nu=\uparrow,\downarrow} n_{j,\nu} n_{j+1,\nu}$ in \mathcal{H} .

Conclusion¹

- Separation between charge and spin spreads increases as U increases.
- Maximum value of charge spread rate almost remains the same as U increases but the maximum of spin spread rate decreases.
- ► Charge spread follows a damped oscillation-like behavior.
- ► In the Mott insulator regime the charge degree of freedom is frozen while the spin degree of freedom fluctuates.

All code can be found here:

https://github.com/hammingheavy/QPM_TermPaper

References

- arXiv:2010.07965 [quant-ph], Observation of separated dynamics of charge and spin in the Fermi-Hubbard model, Google Al Quantum and collaborators.
- ► A. M. Childs and Y. Su, "Nearly Optimal Lattice Simulation by Product Formulas," Physical Review Letters 123, 050503 (2019).
- ► P.-L. Dallaire-Demers, M. Stęchły, J. F. Gonthier, N. T. Bashige, J. Romero, and Y. Cao, "An application benchmark for fermionic quantum simulations," arXiv:2003.01862 (2020).
- Y. Jompol, C. J. B. Ford, J. P. Griffiths, I. Farrer, G. a. C. Jones, D. Anderson, D. A. Ritchie, T. W. Silk, and A. J. Schofield, "Probing Spin-Charge Separation in a Tomonaga-Luttinger Liquid," Science 325, 597 (2009)

Appendix

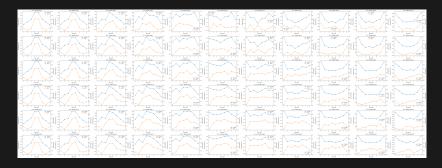


Figure: Evolution of system at different values of U for 10 trotter steps

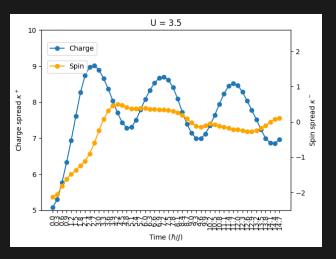


Figure: Damped oscillation like behaviour in spin and charge density waves

Results From Paper

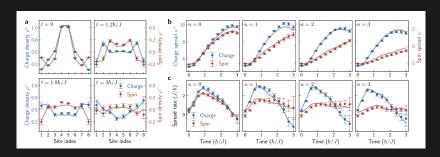


Figure: Plots in the referenced paper used to compare our results