

Ryder's QFT, chapter 9.

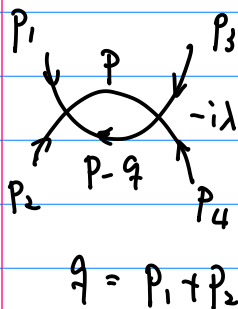
Peskin & Schroeder chapter 10 p325 - chapter 12

2pt func:  $\frac{1}{O(\lambda^0)} + \frac{\text{loop}}{O(\lambda)} = \frac{i\lambda m^2}{16\pi\epsilon} + \text{finite}$

$\epsilon = 4-d$

4pt func =  $\text{X} + \text{bubble} + \text{triangle} + \text{box}$

$O(\lambda)$   $O(\lambda^2)$   $O(\lambda^2)$   $O(\lambda^2)$



$$= \frac{1}{2} \lambda^2 (\mu^2)^{4-d} \int \frac{d^d p}{(2\pi)^d} \frac{1}{p^2 - m^2} \frac{1}{(p-q)^2 - m^2}$$

$$\text{integrand} = \frac{1}{a} \frac{1}{b} = \frac{1}{b-a} \left( \frac{1}{a} - \frac{1}{b} \right) = \frac{1}{b-a} \int_a^b \frac{dx}{x^2}$$

$$x = a z + b(1-z)$$

$$= \int_0^1 \frac{dz}{[a z + b(1-z)]^2}$$

$$a = p^2 - m^2 \quad b = (p-q)^2 - m^2$$

$$\frac{1}{p^2 - m^2} \frac{1}{(p-q)^2 - m^2} = \int_0^1 \frac{dz}{[p^2 - m^2 - 2pq(1-z) + q^2(1-z)]^2}$$

$$p = p' + q(1-z)$$

$$\text{bubble} = \frac{1}{2} \lambda^2 (\mu^2)^{4-d} \int \frac{d^d p'}{(2\pi)^d} \int_0^1 dz \frac{1}{[p'^2 - m^2 + q^2(1-z)z]^2}$$

$p'^0 \rightarrow i p'_E{}^0 \Rightarrow p'^2 \rightarrow -p_E'^2 = |p_E'|^2$

$d^d p \sim d|p| d\theta \dots (\dots)$

HW compute this integral

$$\text{Bubble} = \frac{i\lambda^2}{32\pi^2} (\mu^2)^{2-\frac{d}{2}} \Gamma(2-\frac{d}{2}) \int_0^1 dz \left[ \frac{q^2 z(1-z) - m^2}{4\pi\mu^2} \right]^{\frac{d}{2}-2}$$

$$d=4-\varepsilon \quad \Gamma(2-\frac{d}{2}) = \Gamma(\frac{\varepsilon}{2}) = \frac{2}{\varepsilon} - \gamma + O(\varepsilon)$$

pole at 0  $\Leftrightarrow$  logarithmic div.

$$= \frac{i\lambda^2 \mu^2}{32\pi^2} \left( \frac{2}{\varepsilon} - \gamma + O(\varepsilon) \right) \left( 1 - \frac{\varepsilon}{2} \int_0^1 dz \ln \left[ \frac{q^2 z(1-z) - m^2}{4\pi\mu^2} \right] + O(\varepsilon^2) \right)$$

$$\begin{array}{c} p_1 \quad p_3 \\ \text{Bubble} \\ p_2 \quad p_4 \\ \text{s-channel} \end{array} = \frac{i\lambda^2 \mu^\varepsilon}{16\pi^2 \varepsilon} - \underbrace{\frac{i\lambda \mu^\varepsilon}{32\pi^2} \left( \gamma + \int_0^1 dz \ln \left[ \frac{q^2 z(1-z) - m^2}{4\pi\mu^2} \right] \right)}_{\text{finite}} + O(\varepsilon)$$

$\uparrow$  div.  $\uparrow$   $F(s, m, \mu)$

$$q = p_1 + p_2 = p_3 + p_4$$

$$q^2 = (p_1 + p_2)^2 \quad \text{Lorentz inv.} \quad \parallel \quad s$$

$$= \int_0^1 dz \ln \left[ \frac{s z(1-z) - m^2}{4\pi\mu^2} \right]$$

Mandelstam variables:

$$s = (p_1 + p_2)^2, \quad t = (p_1 + p_3)^2, \quad u = (p_1 + p_4)^2$$

$$\text{Bubble} = \frac{i\lambda^2 \mu^\varepsilon}{16\pi^2 \varepsilon} - \frac{i\lambda \mu^\varepsilon}{32\pi^2} \left[ \gamma + F(s, m, \mu) \right]$$

$$\text{4-pt amplitude: } \Gamma^{(4)}(p_1 \dots p_4) = \begin{array}{c} p_1 \quad p_3 \\ \text{Cross} \\ p_2 \quad p_4 \end{array} + \begin{array}{c} p_1 \quad p_3 \\ \text{Bubble} \\ p_2 \quad p_4 \\ \text{s-channel} \end{array} + \begin{array}{c} p_1 \quad p_3 \\ \text{Bubble} \\ p_2 \quad p_4 \\ \text{t-channel} \end{array} + \begin{array}{c} p_1 \quad p_3 \\ \text{Bubble} \\ p_2 \quad p_4 \\ \text{u-channel} \end{array}$$

$$= -i\lambda \mu^\varepsilon + \frac{3i\lambda^2 \mu^\varepsilon}{16\pi^2 \varepsilon} - \frac{i\lambda \mu^\varepsilon}{32\pi^2} \left[ 3\gamma + F(s, m, \mu) + F(t, m, \mu) + F(u, m, \mu) \right] + O(\varepsilon)$$

$$\text{Cross} = \text{Bubble} + \text{Bubble}$$

$$\text{Cross} = \text{Bubble} + \text{Bubble}$$

$$\Gamma^{(4)} = -i\lambda \mu^\varepsilon \left( 1 - \frac{3\lambda}{16\pi^2 \varepsilon} \right) + \text{finite}$$

"pre-baby" renormalization:  $\lambda = \lambda_0$  bare coupling (div.)

$$\lambda_0 \mu^\epsilon \left(1 - \frac{3\lambda_0}{16\pi^2 \epsilon}\right) \equiv \lambda_{\text{phys}} \text{ (finite) } \text{renormalized coupling}$$

$$\Gamma^{(4)} = -i \lambda_{\text{phys}} + \text{finite}$$

Perturbative renormalization of  $\lambda \phi^4$  ( $d=4-\epsilon$ ) (renormalization for babies)

the interacting Lagrangians dep. on bare quantities, instead of physical quantities

$$\mathcal{L} = \frac{1}{2} \partial_\mu \phi_0 \partial^\mu \phi_0 - \frac{1}{2} m_0^2 \phi_0^2 - \frac{\lambda_0 \mu^{4-d}}{4!} \phi_0^4$$

$\phi_0$ : bare field (wave function)

$m_0$ : bare mass

$\lambda_0$ : bare coupling

Renormalized field  $\phi(x) = Z^{\frac{1}{2}} \phi_0(x)$

$$\mathcal{L} = \frac{1}{2} Z (\partial_\mu \phi)^2 - \frac{1}{2} Z m_0^2 \phi^2 - \frac{\lambda_0 \mu^{4-d}}{4!} \phi^4 Z^2$$

Define renormalized mass & coupling


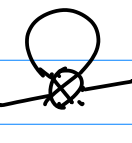
$$\begin{aligned} \delta_m &= m_0^2 Z - m^2, & \delta_\lambda &= \lambda_0 Z^2 - \lambda \\ \delta_Z &= Z - 1 \end{aligned} \quad \begin{array}{c} \uparrow \\ \text{renormalized} \\ \text{mass} \end{array} \quad \begin{array}{c} \uparrow \\ \text{renormalized} \\ \text{coupling} \end{array}$$

$$\begin{aligned} \mathcal{L} &= \frac{1}{2} (\partial_\mu \phi)^2 - \frac{1}{2} m^2 \phi^2 - \frac{\lambda \mu^{4-d}}{4!} \phi^4 \leftarrow \mathcal{L}_r \text{ renormalized Lagrangian} \\ &+ \frac{1}{2} \delta_Z (\partial_\mu \phi)^2 - \frac{1}{2} \delta_m \phi^2 - \frac{\delta_\lambda \mu^{4-d}}{4!} \phi^4 \leftarrow \text{local counter terms } \delta \mathcal{L} \\ &= \mathcal{L}_r + \delta \mathcal{L} \end{aligned}$$

$$Z[J] = e^{\int \left(\frac{\delta}{\delta J}\right)^4 Z_0[J]}$$

$$\text{nonlocal term: } \int d^4x d^4y \phi(x) G(x-y) \phi(y)$$

$$\begin{array}{ll}
 \phi_0 \rightarrow \phi & \text{wavefunction renormalization: } \delta_2 \\
 m_0 \rightarrow m & \text{mass renormalization: } \delta_m \\
 \lambda_0 \rightarrow \lambda & \text{coupling const. renormalization: } \delta_1
 \end{array}
 \left. \vphantom{\begin{array}{l} \phi_0 \rightarrow \phi \\ m_0 \rightarrow m \\ \lambda_0 \rightarrow \lambda \end{array}} \right\} \text{renormalization const.}$$

$\delta\mathcal{L}$ : new Feynman rules (in momentum space)  + 

$$\text{tadpole with cross} = -i\delta_1 \mu^{4-d} \quad \text{propagator with cross} = i(p^2\delta_2 - \delta_m)$$

$$\text{2pt function} = \text{renormalized quantities}$$

$$= \text{tadpole with cross} + \text{tadpole with circle} + \text{tadpole with two circles} + \dots$$

$$= \frac{i}{p^2 - m^2 + i\varepsilon} + \frac{i}{p^2 - m^2 + i\varepsilon} (-i\Sigma^2(p^2)) \frac{i}{p^2 - m^2 + i\varepsilon} + \dots$$

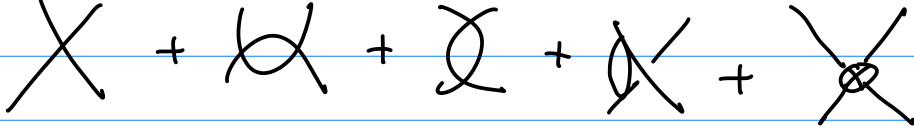
$$= \frac{i}{p^2 - m^2 - \Sigma^2(p^2) + i\varepsilon}$$

$$\frac{i}{p^2 - m^2 - \Sigma^2} = \frac{i}{(p^2 - m^2)(1 - \frac{\Sigma^2}{p^2 - m^2})} = \frac{i}{p^2 - m^2} \left(1 + \frac{\Sigma^2}{p^2 - m^2}\right)$$

$$= \frac{i}{p^2 - m^2} + \frac{i}{p^2 - m^2} (-i\Sigma^2) \frac{i}{p^2 - m^2}$$

$$-i\Sigma^2(p^2) = \frac{i\lambda m^2}{16\pi^2\varepsilon} + \frac{i\lambda m^2}{32\pi^2} \left[1 - \gamma + \ln\left(\frac{4\pi\mu^2}{-m^2}\right)\right] + \mathcal{O}(\varepsilon) + i p^2 \delta_2 - \delta_m$$

idea: we can set  $\delta_Z, \delta_m$  to cancel  $\frac{1}{\epsilon}$ -pole  $\rightarrow$  everything is finite

4pt func =  (we only consider 1-loop renormalization)

$$\epsilon = 4-d$$

$$= -i\lambda\mu^\epsilon + \frac{3i\lambda^2\mu^\epsilon}{16\pi^2\epsilon} - \frac{i\lambda^2\mu^\epsilon}{32\pi^2} [3\gamma + F(s, m, \mu) + F(t, m, \mu) + F(u, m, \mu)] + O(\epsilon) - i\delta_\lambda\mu^\epsilon$$

How to determine renormalization const.  $\delta_m, \delta_Z, \delta_\lambda$ ?

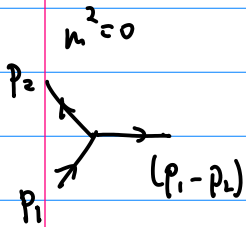
Renormalization conditions: we set an energy scale  $M$  (the energy scale of scattering experiments)

(1) 2-pt func =  $\frac{i}{p^2 - m^2 + i\epsilon}$  physical propagation

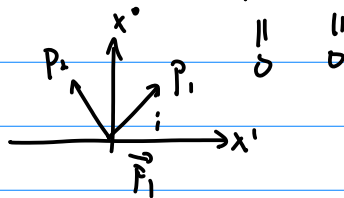
when massive ( $m^2 > 0$ ):  $p^2 \rightarrow m^2$  (propagator has the pole at physical mass)  
 $\uparrow$   
 mass is a preferred energy scale

massless ( $m^2 = 0$ ):  $p^2 \rightarrow -M^2$

$\uparrow$  arbitrary scale



$$(p_1 - p_2)^2 = p_1^2 + p_2^2 - 2p_1 \cdot p_2 < 0$$



$$p_1 = |\vec{p}_1| (1, 1, 0, 0)$$

$$p_2 = |\vec{p}_2| (1, -1, 0, 0)$$

$$\eta = \begin{pmatrix} 1 & & & \\ & -1 & & \\ & & 1 & \\ & & & -1 \end{pmatrix}$$

$$p_1 \cdot p_2 \propto 1 + 1 = 2 > 0$$

(2) 4-pt function =  $-i\lambda\mu^\epsilon$

$$p_1^2 = p_2^2 = m^2$$

when massive ( $m^2 > 0$ ):  $s = 4m^2, t = u = 0$

$$(p_1 + p_2)^2 = p_1^2 + p_2^2 + 2p_1 \cdot p_2 = 4m^2 \text{ rest particle}$$

$$m_{\text{GUTless}}(M^2=0) : s = t = u = -M^2$$

( $\lambda$  is the physical coupling at the energy scale  $\mu$ )

! the perturbative QFT is well-defined only when specifying an energy scale !

i.e. perturbative QFTs depend on a choice of energy scale.

2pt-function =

HW: derive 1-loop renormalization for QED

derive  $\beta$ -function

Due before 9/30