

nonabelian gauge group: $SU(2)$

$$D_\mu = \partial_\mu - ig A_\mu \frac{\sigma_i}{2}$$

$$D_\mu = \partial_\mu - ig A_\mu^i \frac{\sigma_i}{2} \quad i=1,2,3$$

↑
coupling const

$$\psi^a \rightarrow u^a_b(x) \psi^b(x)$$

$$a=1,2$$

$$\bar{\psi}^a \rightarrow \bar{\psi}^b(x) (u^\dagger)_b^a(x)$$

$$i=1,2,3$$

$$\underline{A_\mu = A_\mu^i \frac{\sigma_i}{2}} \rightarrow u(x) A_\mu u(x)^{-1} + \frac{i}{g} u(x) \partial_\mu u(x)^{-1}$$

$$D_\mu \psi \rightarrow u D_\mu \psi$$

$$\mathcal{L}_D = \bar{\psi} (i \gamma^\mu D_\mu^{(A)} - m) \psi \quad \text{gauge inv.}$$

make A dynamical

$$\mathcal{L}_A = -\frac{1}{4} F_{\mu\nu}^i F^{\mu\nu i}$$

$$F_{\mu\nu} = F_{\mu\nu}^i \frac{\sigma_i}{2} = \partial_\mu A_\nu - \partial_\nu A_\mu - ig [A_\mu, A_\nu]$$

$$F_{\mu\nu}^i = \partial_\mu A_\nu^i - \partial_\nu A_\mu^i + g \underbrace{\epsilon^{ijk} A_\mu^j A_\nu^k}$$

gauge transf.

$$[D_\mu, D_\nu] \psi^b = -ig F_{\mu\nu}^{ab} \psi^b$$

$$\psi \in \mathbb{C}^2 \otimes \mathbb{C}^4$$

$$\psi = \psi_{\alpha=1\dots 4}^{a=1,2}$$

HW prove this ↗

$$\text{left} \rightarrow [D_\mu, D_\nu] u \psi = u [D_\mu, D_\nu] \psi = -ig u F_{\mu\nu} \psi$$

$$\text{right} \rightarrow -ig \underline{F'_{\mu\nu} u \psi}$$

$$\Rightarrow F'_{\mu\nu} = u F_{\mu\nu} u^{-1}$$

$$\mathcal{L}_A = -\frac{1}{2} \text{tr} (F_{\mu\nu} F^{\mu\nu}) \rightarrow -\frac{1}{2} \text{tr} (u F_{\mu\nu} u^{-1} u F^{\mu\nu} u^{-1})$$

$$= -\frac{1}{2} F_{\mu\nu}^i F^{\mu\nu j} \text{tr} \left(\frac{\sigma_i}{2} \frac{\sigma_j}{2} \right) = -\frac{1}{2} \text{tr} (F_{\mu\nu} F^{\mu\nu}) \quad \text{gauge inv.}$$

$$= -\frac{1}{4} F_{\mu\nu}^i F^{\mu\nu i}$$

ψ^a_α

$SU(2)$ gauge theory: $\mathcal{L}_{SU(2)} = \underbrace{-\frac{1}{2} \text{tr}(F_{\mu\nu} F^{\mu\nu})}_{\text{free gluon + self interaction of gluons "g"}} + \bar{\psi} (i \gamma^\mu D_\mu - m) \psi$

$$m A_\mu A^\mu$$

$$A_\mu \rightarrow A_\mu + \partial_\mu \chi$$

free quark

+ gluon-quark interaction

$$(ig \bar{\psi} \gamma^\mu \psi A_\mu)$$

$SU(N)$ gauge group.

$$u \in SU(N)$$

$N \times N$ unitary matrix $u^\dagger = u^{-1}$

generators $u = e^{i \alpha^i t^i}$ $i = 1 \dots \dim(SU(N))$

$$u = e$$

$$\alpha^i \in \mathbb{R}$$

$$\det(u) = 1$$

$$\dim(SU(N)) = N^2 - 1$$

$$SU(2): t^i = \frac{\sigma^i}{2}$$

$SU(N)$ t^i : basis of traceless Hermitian

$SU(3)$: t^i Gell-Mann matrices

$N \times N$ matrices.

α^i are small

$$u = 1 + i \alpha^i t^i$$

$$u u^\dagger = 1 \rightarrow (1 + i \alpha^i t^i)(1 - i \alpha^i t^{i\dagger}) = 1 + i \alpha^i (t^i - t^{i\dagger})$$

$$\Rightarrow t^i = t^{i\dagger} \quad \text{Hermitian}$$

$$1 = \det(u) = 1 + i \alpha^i \text{tr}(t^i)$$

$$\Rightarrow \text{tr}(t^i) = 0 \quad \text{traceless.}$$

$$\psi^a = 1 \dots N$$

N flavors of quark, $\psi^a \rightarrow u^a_\alpha \psi^a$ global sym

of $\mathcal{L}_0 = \bar{\psi} (i \gamma^\mu \partial_\mu - m) \psi$ if u indep. of x

upgrade $u \rightarrow u(x)$

$$\partial_\mu \rightarrow D_\mu = \partial_\mu - ig A_\mu$$

$$\begin{matrix} \uparrow \\ N \times N \text{ matrix} \end{matrix} \sum_{i=1}^{\dim(SU(N))} A_\mu^i t^i \equiv A_\mu$$

$$(A_\mu)^a_b = A_\mu^i (t^i)^a_b$$

$$A_\mu \rightarrow u A_\mu u^{-1} + \frac{i}{g} u \partial_\mu u^{-1}$$

$$\psi \rightarrow u \psi \quad \psi_\alpha^a(x) \rightarrow u^a_b(x) \psi_\alpha^b(x)$$

$$\bar{\psi} \rightarrow \bar{\psi} u^{-1} \quad \bar{\psi}_\alpha^a(x) \rightarrow \bar{\psi}_\alpha^b(x) (u^{-1})^a_b(x)$$

$$\mathcal{L}_{SU(N)} = -\frac{1}{2} \text{tr}_N (F_{\mu\nu} F^{\mu\nu}) + \bar{\psi} (i \gamma^\mu D_\mu - m) \psi$$

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu - ig [A_\mu, A_\nu]$$

$$D_\mu \psi \rightarrow D_\mu (u \psi) = u D_\mu \psi$$

HW check this ↗

$$F_{\mu\nu}^i = \partial_\mu A_\nu^i - \partial_\nu A_\mu^i + g f^{ijk} A_\mu^j A_\nu^k, \quad i, j, k = 1 \dots \dim(SU(N))$$

\uparrow
structure const.
 $N^2 - 1$

$$[t^i, t^j] = i f^{ijk} t^k, \quad f^{ijk} \text{ is totally skew}$$

$$[D_\mu, D_\nu] \psi = -ig F_{\mu\nu} \psi \quad \rightarrow \quad F_{\mu\nu} \rightarrow u F_{\mu\nu} u^{-1}$$

HW check this ↗

QCD: $N=3$

Path integral of gauge theory

gauge index: $i = 1 \dots \dim(SU(N)) \rightarrow a = 1 \dots \dim(SU(N))$

rep. index: $a = 1 \dots N \rightarrow i = 1 \dots N$

Generators : $t^a = 1 \dots \dim(SU(N)) \rightarrow T^a = 1 \dots \dim(SU(N))$

Gauge transf : $u(x) \in SU(N) \rightarrow U(x) \in SU(N)$

$$S = \int d^d x \left[-\frac{1}{4} F_{\mu\nu}^a F^{\mu\nu a} + \bar{\psi}^i (\gamma^\mu D_\mu^{ij} - m) \psi^j \right]$$

$$A_\mu^a T^a \equiv A_\mu(x) \rightarrow U(x) A_\mu(x) U(x)^{-1} + \frac{i}{g} U(x) \partial_\mu U(x)^{-1} \equiv A'_\mu$$

$$\psi^i(x) \rightarrow U^{ij}(x) \psi^j(x) \equiv \psi'^i$$

Naively

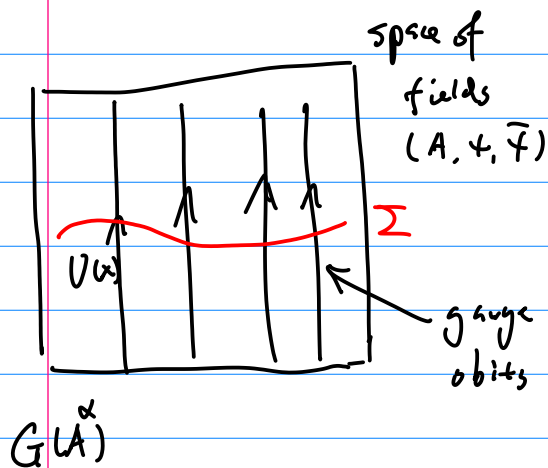
$$Z = \int DA_\mu D\psi D\bar{\psi} e^{iS}$$

assume $DA D\psi D\bar{\psi}$ gauge inv.

$$\hookrightarrow DA' D\psi' D\bar{\psi}' = DA D\psi D\bar{\psi} \cdot J$$

\uparrow
field indep

Why Z is wrong?



$$Z \sim \int D\mu \int D\psi(x) e^{iS}$$

\uparrow
 \perp to gauge orbits
 \uparrow
along gauge orbits

$$= \underbrace{\text{Vol}[\text{gauge orbits}]}_{\infty} \underbrace{\int_{\Sigma} D\mu e^{iS}}$$

Z is divergent.

We want to define " $\frac{Z}{\text{Vol}[\text{gauge orbits}]} = \int_{\Sigma} D\mu e^{iS}$ "

secondly Z violate unitarity

(P&S sec 16.1)

solution: Faddeev-Popov trick: a gauge fixed version of path integral

gauge fixing condition: $G(A) = 0$

e.g. $G(A) = \partial^\mu A_\mu^a(x) - \omega^a(x)$

$\partial^\mu A_\mu^a = 0$ covariant gauge

any function $u = e^{i\alpha^a T^a}$

gauge transf: $A_\mu \rightarrow u \left[A_\mu + \frac{i}{g} \partial_\mu \right] u^{-1} = u A_\mu u^{-1} + \frac{i}{g} u \partial_\mu u^{-1}$ finite gauge transf

$\partial^\mu A_\mu \rightarrow \partial^\mu A_\mu + \dots \partial_\mu u \text{ or } \partial_\mu u^{-1}$

infinitesimal gauge transf.

when $\alpha^a(x)$ is small: $u(x) = 1 + i\alpha^a(x) T^a + O(\alpha^2)$

$$A_\mu^a \rightarrow (A^\alpha)_\mu^a = A_\mu^a + \frac{1}{g} \partial_\mu \alpha^a + \underbrace{f^{abc} A_\mu^b \alpha^c}_{= D_\mu^{ac} \alpha^c} + O(\alpha^2)$$
$$= A_\mu^a + \frac{1}{g} D_\mu^{ac} \alpha^c + O(\alpha^2)$$

HW check this

$$D_\mu^{ac} = \partial_\mu \delta^{ac} + g f^{abc} A_\mu^b$$

covariant derivative w.r.t. adj. rep.

$$G = \partial^\mu A_\mu^a - \omega^a$$

$$\partial^\mu A_\mu^a \rightarrow \partial^\mu (A^\alpha)_\mu^a = \partial^\mu A_\mu^a + \frac{1}{g} \partial^\mu D_\mu^{ac} \alpha^c$$

FP trick: insert in path integral Z an identity

$$1 = \int DG \delta(G)$$

$$1 = \int \delta(G(A^\alpha)) D\alpha \det \left(\frac{\delta G(A^\alpha)}{\delta \alpha} \right) \quad \checkmark \quad \begin{matrix} \text{FP} \\ \text{determinant} \\ \Delta_{\text{FP}} \end{matrix}$$

along gauge orbit

$$G^a(A^a(y)) = \partial^\mu A_\mu^a(y) + \frac{1}{g} \partial^\mu D_\mu^{ac} a^c(y) - \omega^a(y)$$

$$\frac{\delta G^a(A^a(y))}{\delta a^c(x)} = \frac{1}{g} \partial^\mu D_\mu^{ac} \delta(y-x)$$

$$Z = \int D A D \psi D \bar{\psi} e^{iS} = \int D \alpha \int \underline{D A D \psi D \bar{\psi}} e^{\underline{iS[A, \psi, \bar{\psi}]}} \delta(G(A^\alpha)) \det\left(\frac{\delta G}{\delta \alpha}\right)(A^\alpha)$$

gauge inv. of S & $D A D \psi D \bar{\psi}$

$$= \int D \alpha \int \underline{D A^\alpha D \psi^\alpha D \bar{\psi}^\alpha} e^{iS[A^\alpha, \psi^\alpha, \bar{\psi}^\alpha]} \delta(G(A^\alpha)) \det\left(\frac{\delta G}{\delta \alpha}\right)(A^\alpha)$$

change of variable $A^\alpha \rightarrow A$ or $A \rightarrow A^{\alpha^{-1}}$

$$= \underbrace{\int D \alpha}_{\text{Vol [gauge orbits]}} \underbrace{\int D A D \psi D \bar{\psi} e^{iS[A, \psi, \bar{\psi}]}}_{\int D A e^{iS}} \delta(G(A)) \det\left(\frac{\delta G}{\delta \alpha}\right)(A)$$

$$\int f(x) dx = \int f(y) dy$$