

boson no spin

$$T_0^2 = 1$$

$$T_0 \psi(t, \vec{r}) = \psi^*(-t, \vec{r})$$

fermion spin $\frac{1}{2}$

$$T^2 = -1$$

$$T \begin{pmatrix} \psi_1(t, \vec{r}) \\ \psi_2(t, \vec{r}) \end{pmatrix} = e^{i\varphi} \sigma_y \begin{pmatrix} \psi_1^*(-t) \\ \psi_2^*(-t) \end{pmatrix}$$

zero spin

$$\hat{H} \psi_{nl} = E_n \psi_{nl} \quad l=1 \dots d(n)$$

general sol of Schrödinger eqn $\psi(\vec{r}, t) = \sum_n a_n \psi_{nl}(\vec{r}) e^{-\frac{i}{\hbar} E_n t}$
(\hat{H} is real, not explicitly dep. on t)

$$\hat{T}_0 \psi(\vec{r}, t) = \sum_n a_n^* \psi_{nl}^*(\vec{r}) e^{-\frac{i}{\hbar} E_n t}$$

effectively $\hat{T}_0 : \psi_{nl}(\vec{r}) \rightarrow \psi_{nl}^*(\vec{r})$

if \hat{H} is real, $\hat{H} \psi_{nl} = E_n \psi_{nl}$ ψ_{nl} is eigenstate

$\Rightarrow \hat{H} \psi_{nl}^* = E_n \psi_{nl}^*$ ψ_{nl}^* is eigenstate

ψ_{nl} basis in $\mathcal{H}^{(n)}$, ψ_{nl}^* is basis in $\mathcal{H}^{(n)*}$

whether $\mathcal{H}^{(n)}$ and $\mathcal{H}^{(n)*}$ the same?

(1) $\mathcal{H}^{(n)} \cong \mathcal{H}^{(n)*}$ equivalent rep of symm. group

(2) $\mathcal{H}^{(n)} \not\cong \mathcal{H}^{(n)*}$ not equivalent

$$\text{eigenspace} = \mathcal{H}^{(n)} \oplus \mathcal{H}^{(n)*}$$

real rep, pseudo real rep, and complex rep

Given a symm. group G , unitary irrep $D: g \mapsto D(g) \in \mathcal{H}^{(n)}$

Def $D(g)$ is unitary on $\mathcal{H}^{(n)}$, ψ_{hi} orthonormal basis in $\mathcal{H}^{(n)}$

(complex
conjugate
irrep)

$$D(g) \psi_{hi}(\vec{r}) = \sum_j \psi_{hj}(\vec{r}) D_{ji}(g)$$

$D_{ji}(g)$ unitary matrix

complex conjugate : $D(g)^* \psi_{hi}^*(\vec{r}) = \sum_j \psi_{hj}^*(\vec{r}) D_{ji}^*(g)$

$$\psi_{hi}^*(\vec{r}) = T_0 \psi_{hi}(\vec{r})$$

$D_{ij}(g)$ unitary irrep matrix
(irrep D)

$\xrightarrow{T_0} D_{ij}^*(g)$ unitary irrep matrix
irrep D^* :

Complex conjugate irrep

relation between D and D^*

(1) if \exists unitary transf. $U: \mathcal{H}^{(n)} \rightarrow \mathcal{H}^{(n)}$ s.t

$$U D(g) U^{-1} = D_0(g) \quad \forall g \in G$$

\uparrow
real matrix

$$\Rightarrow \underline{D \text{ is equivalent to } D^*}, \text{ since } U^* D(g)^* U^{*-1} = D_0(g)$$

$$\Rightarrow \underline{D^*(g) = U^{*-1} D_0(g) U^*} = \underline{U^{*-1} U D(g) U^{-1} U^*}$$

$U^{*-1} U$ is unitary

we say $D \simeq D^*$ is a real rep.

(2) D is ^{unitarily} equivalent to D^* , but they are not equivalent to real rep.

$$\nexists U \text{ s.t. } D(g) = U D_0(g) U^{-1} \quad \forall g \in G$$

\uparrow
real matrix

We say, $D \simeq D^*$ is pseudo-real

(3) D is inequivalent to D^* .

We say, D and D^* are complex reps

distinguish (3) from (1) and (2), check character

if $D \simeq D^*$, then $\chi(g) = \chi^*(g) \quad \forall g \in G$

if $\chi(g) \neq \chi^*(g)$, then D & D^* are complex reps.

distinguish between (1) & (2)

$$D^* \simeq D \rightarrow D^*(g) = Z D(g) Z^{-1} \quad \forall g \in G$$

\uparrow
unitary matrix

Complex conjugate $D(g) = Z^* D^*(g) Z^{*-1}$

$$\Rightarrow D^*(g) = Z Z^* D^*(g) (Z Z^*)^{-1}$$

$$\therefore [Z Z^*, D^*(g)] = 0 \quad \forall g \in G$$

by Schur's lemma, when D^* is irrep, $Z Z^* = c \mathbb{1}$
 $c \in \mathbb{C}$

Z is unitary, $Z^* = (Z^T)^{-1}$

$$\Rightarrow Z (Z^T)^{-1} = c \mathbb{1} \Rightarrow \underline{Z = c Z^T}$$

$$\xRightarrow{\text{transpose}} Z^T = c Z$$

$$\Rightarrow Z = c^2 Z \Rightarrow c^2 = 1, c = \pm 1$$

$$D^* \simeq D \Rightarrow Z Z^* = \pm \mathbb{1}$$

Thm: $Z Z^* = \mathbb{1}$ iff D is real

$Z Z^* = -\mathbb{1}$ iff D is pseudo-real

pf if D is real $\rightarrow D^*(g) = (U^*)^{-1} D_0(g) U^*$
 $= (U^*)^{-1} U D(g) U^{-1} U^*$

on the other hand $D^*(g) = Z D(g) Z^{-1}$

$$\Rightarrow (U^*)^{-1} U = b Z \quad b \in \mathbb{C}$$

Complex conjugate $U^{-1} U^* = b^* Z^*$

$$U^{-1} U^* (U^*)^{-1} U = b^* b Z^* Z = |b|^2 Z^* Z$$

$$\parallel$$
$$1$$

$$\underline{Z^* Z = |b|^{-2} \mathbb{1}}$$

$$\underline{Z Z^* = c \mathbb{1}} \quad \Rightarrow \quad |b|^2 c = 1$$

$$c = 1 \text{ or } -1 \quad c = 1 \quad |b|^2 > 0$$

conversely if $c = 1$ ($Z Z^* = \mathbb{1}$)

Z unitary, $Z = e^{iA}$ A hermitian

$$Z^* = Z^{-1} \Leftrightarrow Z^T = Z, \quad Z \text{ symmetric}$$

$$Z^* = (Z^T)^{-1}$$

$$\text{def. } U = Z^{\frac{1}{2}} = e^{iA/2} \quad \text{unitary}$$

$$U^2 = Z$$

$$\underline{U^* Z = e^{-iA/2} e^{iA} = e^{\frac{1}{2}iA} = U}$$

$$\begin{aligned} \forall g \in G, [U D(g) U^{-1}]^* &= U^* D(g)^* U^{*-1} \\ &= U^* Z D(g) Z^{-1} U^{*-1} \\ &= U D(g) U^{-1} = D_o(g) \text{ real} \end{aligned}$$

$\Rightarrow D$ is real rep.

$$Z Z^* = \mathbb{1} \Leftrightarrow D \text{ is real rep}$$

$$\Rightarrow Z Z^* = -\mathbb{1} \Leftrightarrow D \text{ is pseudo-real rep.} \quad \square$$

Distinguish real & pseudo-real reps by characters

$$\underline{D^*(g) = Z D(g) Z^{-1}}$$

orthogonality $\sum_g D_{\alpha\delta}^*(g) D_{\beta\gamma}(g) = \delta_{\alpha\beta} \delta_{\gamma\delta} \frac{h}{\dim(D)=d}$

$$\sum_{\alpha, \delta} \underline{Z_{\tau\alpha}^{-1}} \sum_g \sum_{\sigma, \rho} \underline{Z_{\alpha\sigma}} D_{\sigma\rho}(g) \underline{Z_{\rho\delta}^{-1}} D_{\beta\gamma}(g) \times \underline{Z_{\delta\chi}}$$

$$\Rightarrow \sum_g D_{\tau\chi}(g) D_{\beta\gamma}(g) = Z_{\gamma\chi} Z_{\tau\beta}^{-1} \frac{h}{d}$$

let $\chi = \rho$. \sum_{β} , $\sum_g \sum_{\beta} D_{\tau\beta}(g) D_{\beta\gamma}(g)$

$$\begin{aligned} \sum_g D_{\tau\gamma}(g^2) &= (Z Z^*)_{\tau\gamma} \frac{h}{d} \\ &= \pm \delta_{\tau\gamma} \frac{h}{d} \end{aligned}$$

let $\tau = \gamma$, \sum_{τ} ,

$$\boxed{\sum_g \chi(g^2) = \pm h}$$

for complex rep. $D \neq D^* \Rightarrow \sum_g [D_{\alpha\beta}^*(g)]^* D_{\gamma\delta}(g) = 0$

$$\sum_j \text{Dop}(j) \text{Drs}(j)$$

$$\Rightarrow \sum_j \chi(j^2) = 0$$

Summary :

$$\frac{1}{h} \sum_{j \in G} \chi(j^2) = \begin{cases} 1 & \text{real rep} \\ -1 & \text{pseudo-real rep} \\ 0 & \text{complex.} \end{cases}$$

Extra-degeneracy of \hat{H} due to time-reversal inv.

$$\hat{H} \psi_{nl} = E_n \psi_{nl} \quad l = 1 \dots d$$

$$\psi_{nl} \in \mathcal{H}^{(n)} \quad \text{unitary irrep. of symmetry group } G$$

$$l = 1 \dots d \quad \dim \mathcal{H}^{(n)} = d$$

find degeneracy at E_n , deg. at E_n is either d or $2d$ (spin-zero)

if \hat{H} is real and not explicitly dep on t

$$\hat{H} \psi_{nl} = E_n \psi_{nl}$$

$$\hat{H} \psi_{nl}^* = E_n \psi_{nl}^*$$

$$\psi_{nl}^*(\vec{r}) = \hat{T}_0 \psi_{nl}(\vec{r})$$

• if $\forall \psi_{nl}^*$, $\psi_{nl}^* \in \mathcal{H}^{(n)}$ spanned by ψ_{nl}

then $\mathcal{H}^{(n)}$ is the final eigenspace i.e. $\mathcal{H}^{(n)} = \mathcal{H}^{(n)*}$
 degeneracy at E_n is d

• otherwise, eigenspace $= \mathcal{H}^{(n)} \oplus \mathcal{H}^{(n)*}$

degeneracy at $E_n = 2d$ (extra degeneracy)

Spin-zero

$$\hat{H} \psi_i = E \psi_i \quad \psi_i \in \underline{\mathcal{H}}^{(n)}$$

$$D(g) \psi_i = \sum_j \psi_j D_{ji}(g) \quad \forall g \in G$$

\hat{H} is real $\hat{H} \psi_i^* = E \psi_i^* \quad \psi_i^* \in \mathcal{H}^{(n)*}$

$$D^*(g) \psi_i^* = \sum_j \psi_j^* D_{ji}^*(g) \quad \forall g \in G$$

$\mathcal{H}^{(n)}$ carries irrep D of G

$\mathcal{H}^{(n)*}$ carries complex conjugate irrep D^* of G

$\mathcal{H}^{(n)} \simeq \mathcal{H}^{(n)*}$ or not relates to $D \simeq D^*$ or not

firstly if $D \not\simeq D^*$ complex irrep, $\mathcal{H}^{(n)} \perp \mathcal{H}^{(n)*}$ by
 orthogonality theorem.

\Rightarrow extra degeneracy $2d$

let's look at case (1) real rep and (2) pseudo-real rep.

$$\exists \text{ unitary } Z, D^*(g) = Z D(g) Z^{-1}, \quad ZZ^* = \begin{cases} 1 & \text{real} \\ -1 & \text{pseudo-real} \end{cases}$$

Lemma if $\mathcal{H}^{(n)} = \mathcal{H}^{(n)*}$, then D is of case (1) i.e. real rep.

pf. $\mathcal{H}^{(n)} = \mathcal{H}^{(n)*}$, $\hat{H} \psi_i = E \psi_i$
 $\hat{H} \psi_i^* = \bar{E} \psi_i^*$

both $\{\psi_i\}$, $\{\psi_i^*\}$ are both orthonormal basis

then \exists unitary U s.t. $\psi_i = \sum_k \psi_k^* U_{ki}$

$$\psi_i^* = \sum_k \psi_k U_{ki}^*$$

$$\Rightarrow \psi_i = \sum_{kl} \psi_k U_{lk}^* U_{ki} \quad \text{i.e. } UU^* = \mathbb{1}$$

$$D(g) \psi_i = \sum_j \psi_j D_{ji}(g)$$

$$\parallel \quad \sum_j \sum_k \psi_k^* U_{kj} D_{ji}(g)$$

$$\times U_{il}^{-1} \sum_i$$

$$D(g) \psi_i^* = \sum_k \psi_k^* (U D(g) U^{-1})_{ki}$$

\uparrow
same as $D^*(g) \psi_i^*$

compare to $D^*(g) \psi_i^* = \sum_j \psi_j^* D_{ji}^*(g)$

$$D^*(g) = U D(g) U^{-1} \quad \} = D \text{ is real}$$

$$U U^* = I$$

J

□

$$U = Z$$

Lemma if D is real, then $\mathcal{H}^{(n)} = \mathcal{H}^{(n)*}$

pf D is real, $Z D Z^{-1} = D^*$ & $Z Z^* = I$

$$\begin{aligned} D^*(g) \psi_i^* &= \sum_j \psi_j D_{ji}^*(g) \\ &= \sum_j \psi_j \sum_{k,l} Z_{jk} D_{kl}(g) Z_{li}^{-1} \end{aligned}$$

$$\Rightarrow D^*(g) \left(\sum_i \psi_i^* Z_{im} \right) = \sum_k \left(\sum_j \psi_j^* Z_{jk} \right) \underline{D_{km}(g)}$$

× Z_{im}
and \sum_i

$$D(g) \psi_m = \sum_k \psi_k D_{km}(g)$$

$$\Rightarrow \sum_i \psi_i^* Z_{im} = \psi_m$$

$$\psi_i^* \in \mathcal{H}^{(n)*} \quad \psi_i \in \mathcal{H}^{(n)}$$

they are linear dep. by $\sum_i \psi_i^* Z_{im} = \psi_m$

$$\Rightarrow \mathcal{H}^{(n)} = \mathcal{H}^{(n)*}$$

above 2 lemmas $\Rightarrow \mathcal{H}^{(h)} = \mathcal{H}^{(h)*}$ iff D is real
(no extra degeneracy) d

then if D is pseudo-real then $\mathcal{H}^{(h)} \neq \mathcal{H}^{(h)*}$
 \Rightarrow extra degeneracy $2d$

if D is $\begin{cases} \text{real} \\ \text{pseudo-real} \\ \text{complex} \end{cases}$ degeneracy $= d = \dim(D)$
degeneracy $= 2d$
degeneracy $= 2d$

Examples (1) 1d free particle $\hat{H} = \frac{1}{2m} \hat{p}^2$

$$\hat{p} = -i\hbar \frac{\partial}{\partial x}$$

symmetry: transl. inv. $Q(\lambda)x = x + \lambda \quad \lambda \in \mathbb{R}$

$$D(\lambda) = e^{-\frac{i}{\hbar} \lambda \hat{p}}$$

$$D(\lambda)\psi(x) = \psi(x + \lambda)$$

$G = \mathbb{R} = \{\lambda\}$ group multiplication: $+$; $\lambda_1 + \lambda_2$

irrep of \mathbb{R} : $D^{(k)}(\lambda) = e^{ik\lambda}$

$$\mathcal{H}^{(k)} = \mathbb{C} \quad \dim(D^{(k)}) = 1$$

all irrep of G are 1-dim

$$\underline{D^{(k)}(\lambda_1) D^{(k)}(\lambda_2) = e^{ik(\lambda_1 + \lambda_2)} = D^{(k)}(\lambda_1 + \lambda_2)}$$

Complex conjugate of $D^{(k)}$: $D^{(k)}(\lambda)^* = e^{-ik\lambda}$
 $\neq e^{ik\lambda}$ ↗

$$D^{(k)} \neq D^{(k)*}$$

cannot transform
between them
by unitary on \mathbb{C}

Complex irrep.

$$\text{energy level degeneracy} = 2 \dim(D^{(k)}) \\ = 2$$

eigenstates of H : $\hat{H} \psi_k = \frac{\hbar^2 k^2}{2m} \psi_k$

$$\psi_k = \underbrace{e^{ikx}}_{\psi^{(k)}}, \quad \psi_{-k} = \underbrace{e^{-ikx}}_{\psi^{(k)*}}$$

indeed degeneracy = 2

$$U \underbrace{e^{ikx}}_{\in \mathbb{C}} U^{-1} = e^{-ikx}$$

$$e^{ikx} \neq e^{-ikx} \Rightarrow D^{(k)} \neq D^{(k)*}$$

(2) central potential : $\hat{H} \psi_{nlm}(\vec{r}) = E_{nl} \psi_{nlm}$
 $l = 0, 1, 2, \dots$

$$m = -l, -l+1, \dots, l$$

symmetry group $G = SO(3)$

eigenspace $\mathcal{H}^{(n,l)}$ relates to irrep of $SO(3)$

\forall Rotation $Q(\alpha, \beta, \gamma) \in SO(3)$

$\mathcal{H}^{(l)}$ is spanned by $Y_{lm}(\theta, \varphi)$
 is labelled by l , $m = -l, \dots, l$
 Euler angles

$$(l = 0, 1, 2, 3, \dots)$$

$$\dim(\mathcal{H}^{(l)}) = 2l+1$$

$$\langle l m | D^{(l)}(\alpha, \beta, \gamma) | l m' \rangle = \int_{S^2} d\theta d\varphi \sin\theta Y_{lm}^* \hat{D}^{(l)} Y_{lm'}$$

$$= D_{mm'}^{(l)}(\alpha, \beta, \gamma) = e^{-im'\alpha} d_{mm'}^{(l)}(\beta) e^{-im\gamma}$$

\uparrow
Wigner D-function

\uparrow
Wigner d-function

$$d(0) = 1$$

$$\chi^l(\alpha) = \text{tr } D^{(l)}(\alpha, \beta, \gamma) = \text{tr } D^{(l)}(\alpha, 0, 0)$$

$$= \sum_{m=-l}^l e^{-im\alpha} = \frac{\sin((l+\frac{1}{2})\alpha)}{\sin \alpha/2} \quad \text{real}$$

$$\chi^{l*} = \chi^l \quad D^l \text{ is not complex}$$

basis in $\mathcal{H}^{(l)}$: $Y_{lm}(\theta, \varphi)$, $Y_{lm}^*(\theta, \varphi) = Y_{l, -m}(\theta, \varphi)$

$$Y_{lm}^* \in \mathcal{H}^{(l)}$$

$$\Rightarrow \mathcal{H}^{(l)*} = \mathcal{H}^{(l)}$$

$\Rightarrow D^{(l)}$, $l=0, 1, \dots$ are all real rep., no extra degeneracy

$\mathcal{H}^{(h, l)}$ is the eigenspace, $\deg. = 2l+1$

single spin- $\frac{1}{2}$ particle: $T^2 = -1$

lemma $\langle \psi | T\psi \rangle = 0$

pf. Let $\varphi = T\psi \quad \forall \psi \in \mathcal{H}^{(0)} \otimes \mathbb{C}^2$

$$\langle \psi | \varphi \rangle = \langle \psi | T\psi \rangle = \langle T\psi | T^2\psi \rangle^*$$

↑
T is antiunitary

$$= -\langle T\psi | \psi \rangle^* = -\langle \psi | \varphi \rangle$$

$$\Rightarrow \langle \psi | \varphi \rangle = 0$$

$\psi, T\psi$ are orthogonal. □

$$\hat{H} \psi_i = E \psi_i \quad i = 1 \dots d \quad \psi_i \in \mathcal{H}^{(d)} \text{ carry irrep } D^{(d)} \text{ of } G$$

$$\hat{H}(T\psi_i) = E(T\psi_i) \quad \varphi_i \equiv T\psi_i$$

$$(1) \quad \langle \psi_i | \psi_i \rangle = 0$$

$$(2) \quad \langle \psi_i | \psi_j \rangle = \delta_{ij}$$

$$(3) \quad \langle \varphi_i | \varphi_j \rangle = \delta_{ij}$$

Thm (Kramer's Thm) single spin- $\frac{1}{2}$ particle, \forall energy level degeneracy d' is always even and

$$d \leq d' \leq 2d$$

pf we have (1), (2), & (3) whether $\{\psi_i, \varphi_i\}$ form a complete basis

but it is possible $\langle \psi_i | \varphi_j \rangle \neq 0 \quad i \neq j$

if $|\varphi_k\rangle = \sum_i |\psi_i\rangle c_i$ the φ_k should be removed from the set of basis

but the $|\psi_k\rangle$ should be removed as well

$$\text{since } T^2 = -1$$

$$T|\varphi_k\rangle = \sum_i T(|\psi_i\rangle c_i)$$

$$\parallel$$

$$T^2|\varphi_k\rangle$$

$$\parallel$$

$$-|\varphi_k\rangle$$

$$\parallel$$

$$\sum_i c_i^* |\varphi_i\rangle$$

$$\Rightarrow |\varphi_k\rangle = -\sum_i c_i^* |\varphi_i\rangle$$

$$\psi_1 \dots \widehat{\psi_k} \dots \psi_d \quad \psi_1 \dots \widehat{\psi_k} \dots \psi_d$$

\uparrow removed in pair.

find degeneracy $d' = 2d - 2n$ # of removed pairs

d' is even

$$d \leq d' \leq 2d$$

□

Theory of angular momentum

spatial rotation & SO(3) group

finite rotation: Q 3×3 matrix s.t. $Q\vec{r}_1 \cdot Q\vec{r}_2 = \vec{r}_1 \cdot \vec{r}_2$

$$\Rightarrow Q^T Q = I \quad Q \in O(3)$$

$$\det(Q^T Q) = 1$$

$$\det(Q)^2 = 1$$

$$\det Q = \pm 1$$

$$SO(3) : Q^T Q = I \text{ \& \& } \det Q = 1 \quad P = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix}$$

HW Tutorial

problem 1.

$$\hat{H} = \hat{H}_0 + V_{\text{lattice}}(\vec{r})$$

↑
central force, $SO(3)$ symmetry

eigenspace of H_0 : $\mathcal{H}_{L=0,1,\dots}$
carries irrep of $SO(3)$

V_{lattice} breaks $SO(3)$ to O (cubic lattice group)

irrep of $SO(3)$ $\left\{ \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \right\}$ irreps of O
 $l=0,1,\dots$ $T_i=1,\dots,5$

$$D^L = \bigoplus_T D_T$$

↑ irrep of $SO(3)$ ↑ irreps of O

$$\forall g \in O, \chi_L(g) = \sum_T \chi_T(g)$$

character of O

$\chi_T(g)$	E	$8C_3$	$6C_2$	$2C_2$	$6C_4$
T_1	1	1	1	1	1
T_2	1	1	-1	1	-1
T_3	2	-1	0	2	0
T_4	3	0	-1	-1	1
T_5	3	0	1	-1	-1

character of $SO(3)$ evaluated at elements in O

$\chi_l(g)$	E	$8C_3$	$6C_2$	$2C_2$	$6C_4$
$l=0$	1	1	1	1	1
$l=1$	3	0	-1	-1	1
$l=2$	5	-1	1	1	-1
$l=3$	7	1	-1	-1	-1
$l=4$	9	0	1	1	1

* $l=0$: $D^{l=0} = T_1$ * $l=1$: $D^{l=1} = T_4$

* $l=2$: $D^{l=2} = T_3 \oplus T_5$ $\chi_{l=2}(g) = \chi_{T_3}(g) + \chi_{T_5}(g)$

check: E : $\chi_{l=2}(E) = 5$ $\forall g \in O$

$\chi_{T_3}(E) + \chi_{T_5}(E) = 2 + 3 = 5$

$8C_3$: $\chi_2(8C_3) = -1$

$\chi_{T_3}(8C_3) + \chi_{T_5}(8C_3) = -1 + 0 = -1$

$D_l = \bigoplus_T D_T$ $\forall g \in O$, $D_l(g) = \bigoplus_T D_T(g)$

$$\left(\begin{array}{ccc} \boxed{D_{T_1}(g)} & & \\ & \boxed{D_{T_2}(g)} & \\ & & \boxed{D_{T_3}(g)} \\ & & & \ddots \end{array} \right) \begin{array}{c} \circ \\ \circ \end{array} \left(\begin{array}{c} - \\ - \\ - \\ - \end{array} \right) \quad \text{---} \begin{array}{c} \vdots \\ \vdots \\ \vdots \end{array}$$

$\chi_l(g) = \text{tr } D_l(g) = \sum_T \text{tr } D_T(g) = \sum_T \chi_T(g)$

$$6C_2: \chi_{l=2}(6C_2) = 1$$

$$\chi_{T_3}(6C_2) + \chi_{T_5}(6C_2) = 0 + 1 = 1$$

$$2C_2: \chi_{l=2}(2C_2) = 1$$

$$\chi_{T_3}(2C_2) + \chi_{T_5}(2C_2) = 2 + (-1) = 1$$

$$6C_4: \chi_{l=2}(6C_4) = -1$$

$$\chi_{T_3}(6C_4) + \chi_{T_5}(6C_4) = 0 + (-1) = -1$$

$$\bullet l=3 \quad D_3 = T_2 \oplus T_4 \oplus T_5$$

$$E: \chi_3 = 7$$

$$\chi_{T_2} + \chi_{T_4} + \chi_{T_5} = 1 + 3 + 3 = 7$$

$$8C_3: \chi_3 = 1$$

$$\chi_{T_2} + \chi_{T_4} + \chi_{T_5} = 1 + 0 + 0 = 1$$

$$6C_2: \chi_3 = -1$$

$$\chi_{T_2} + \chi_{T_4} + \chi_{T_5} = -1 + (-1) + 1 = -1$$

$$3C_2: \chi_3 = -1$$

$$\chi_{T_2} + \chi_{T_4} + \chi_{T_5} = 1 + (-1) + (-1) = -1$$

$$6C_4: \chi_3 = -1$$

$$\chi_{T_2} + \chi_{T_4} + \chi_{T_5} = -1 + 1 + (-1) = -1$$

$$\bullet l = 4 \quad D_4 = T_1 \oplus T_3 \oplus T_4 \oplus T_5$$

$$E: \quad \chi_4 = 9$$

$$\chi_{T_1} + \chi_{T_3} + \chi_{T_4} + \chi_{T_5} = 1 + 2 + 3 + 3 = 9$$

$$8C_3: \quad \chi_4 = 0$$

$$\chi_{T_1} + \chi_{T_3} + \chi_{T_4} + \chi_{T_5} = 1 + (-1) + 0 + 0 = 0$$

$$6C_2: \quad \chi_4 = 1$$

$$\chi_{T_1} + \chi_{T_3} + \chi_{T_4} + \chi_{T_5} = 1 + 0 + (-1) + 1 = 1$$

$$3C_2: \quad \chi_4 = 1$$

$$\chi_{T_1} + \chi_{T_3} + \chi_{T_4} + \chi_{T_5} = 1 + 2 + (-1) + (-1) = 1$$

$$6C_4: \quad \chi_4 = 1$$

$$\chi_{T_1} + \chi_{T_3} + \chi_{T_4} + \chi_{T_5} = 1 + 0 + 1 + (-1) = 1$$

$$\vec{r} = \sum_{i=1}^3 r_i \vec{e}_i \quad \text{choose basis } \vec{e}_{i=1,2,3}$$

$$\vec{r}' = Q \vec{r} = \sum_{i=1}^3 r'_i \vec{e}_i$$

Rotation of basis

$$\vec{r}' = Q \vec{r} = \sum_{i=1}^3 r_i Q \vec{e}_i = \sum_{i=1}^3 r_i \vec{e}'_i$$

$$\vec{e}_i' = Q \vec{e}_i = \sum_j \vec{e}_j (\underbrace{\vec{e}_j \cdot Q \vec{e}_i}_{Q_{ji}}) = \sum_{j=1}^3 \vec{e}_j Q_{ji}$$

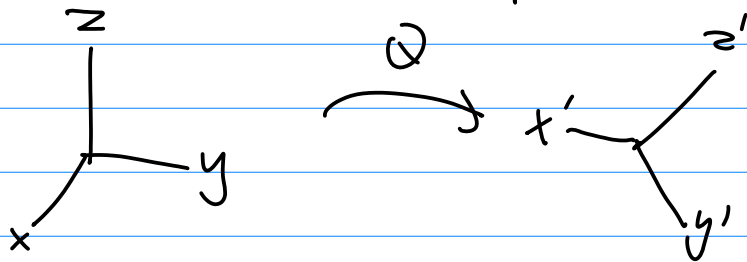
$$\vec{r}' = \sum_{i=1}^3 r_i' \vec{e}_i$$

$$\sum_{i=1}^3 r_i \sum_{j=1}^3 \vec{e}_j Q_{ji} = \sum_{j=1}^3 \left(\sum_{i=1}^3 Q_{ji} r_i \right) \vec{e}_j$$

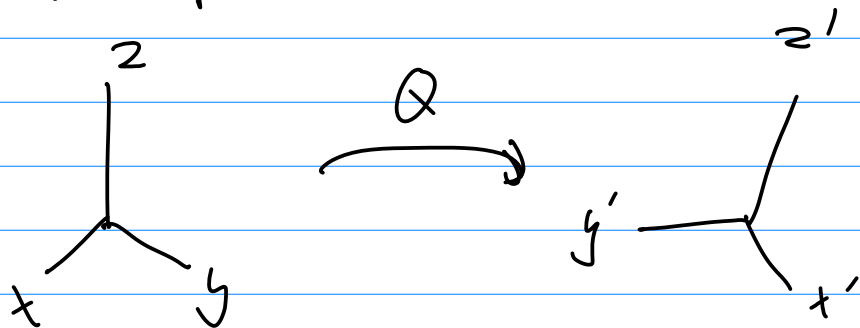
$$r_j' = \sum_{i=1}^3 Q_{ji} r_i$$

Parity $P: \vec{r} \rightarrow -\vec{r}$ $P_{ij} = -\delta_{ij}$ $\det P = -1$
 $P \notin SO(3)$ $P \in O(3)$

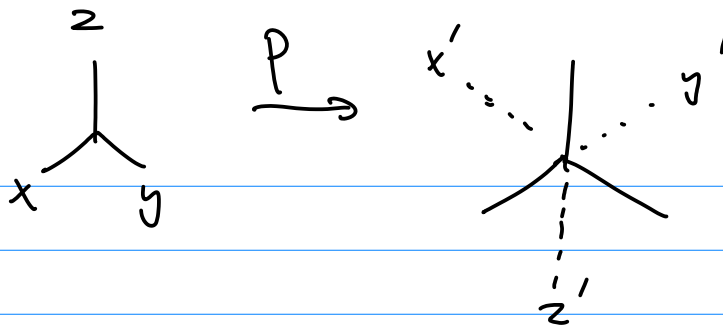
$Q \in SO(3)$ transf. right-hand frame to right-hand frame



$Q \notin SO(3)$, $Q \in O(3)$ transf. right-hand frame to left-hand frame



P:



Proof:

Let $\vec{z} = \vec{x} \times \vec{y}$ we should show that

$$Q\vec{x} \times Q\vec{y} = Q\vec{z} \quad \forall Q \in SO(3)$$

$$Q\vec{x} \times Q\vec{y} = -Q\vec{z} \quad \forall Q \notin SO(3)$$

$$Q \in O(3)$$

\forall vector $\vec{u} \in \mathbb{R}^3$

$$\underline{\underline{Q\vec{u}}} \cdot (Q\vec{x} \times Q\vec{y}) = \sum_{ijk} \varepsilon_{ijk} (\underline{\underline{Q\vec{u}}})_i (Q\vec{x})_j (Q\vec{y})_k$$

$$= \sum_{ijk} \varepsilon_{ijk} Q_{i\alpha} u_{\alpha} Q_{j\beta} x_{\beta} Q_{k\gamma} y_{\gamma}$$

$$= \sum_{\alpha\beta\gamma} \left(\underbrace{\sum_{ijk} \varepsilon_{ijk} Q_{i\alpha} Q_{j\beta} Q_{k\gamma}}_{\varepsilon_{\alpha\beta\gamma} \det Q} \right) u_{\alpha} x_{\beta} y_{\gamma}$$

$$= \det Q \sum_{\alpha\beta\gamma} \varepsilon_{\alpha\beta\gamma} u_{\alpha} x_{\beta} y_{\gamma}$$

$$= \det Q (\vec{u} \cdot (\vec{x} \times \vec{y}))$$

$$= \det Q (\underline{\underline{\vec{u}}} \cdot \underline{\underline{\vec{z}}}) = \det Q (\underline{\underline{Q\vec{u}}} \cdot Q\vec{z})$$

$$Q\vec{x} \times Q\vec{y} = \underbrace{\det Q}_{\pm 1} (Q\vec{z})$$

all rotations of rigid body are $SO(3)$

$\forall Q \in SO(3)$ can be composed by 2 types of simple rotations

$$Q(\hat{k}, \alpha), \quad Q(\hat{j}, \beta)$$

rotation around
z-axis

rotation around
y-axis

$\hat{i}, \hat{j}, \hat{k}$ basis
of x, y, z axis

$$Q(\hat{k}, \alpha) = \begin{pmatrix} \cos \alpha & -\sin \alpha & 0 \\ \sin \alpha & \cos \alpha & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

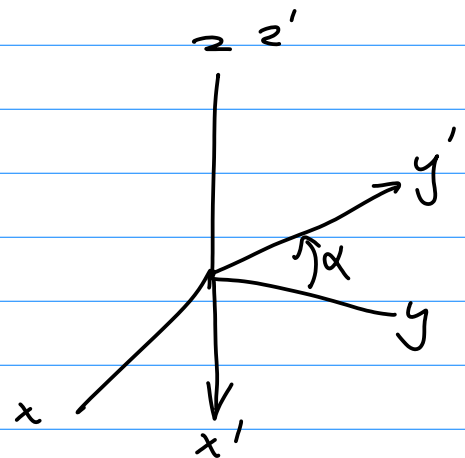
$$\vec{e}_i' = \sum_j \vec{e}_j Q_{ji}$$

$$\hat{i} \rightarrow \hat{i}' = \hat{i} \cos \alpha + \hat{j} \sin \alpha$$

$$\hat{j} \rightarrow \hat{j}' = \hat{i} (-\sin \alpha) + \hat{j} \cos \alpha$$

$$\hat{k} \rightarrow \hat{k}' = \hat{k}$$

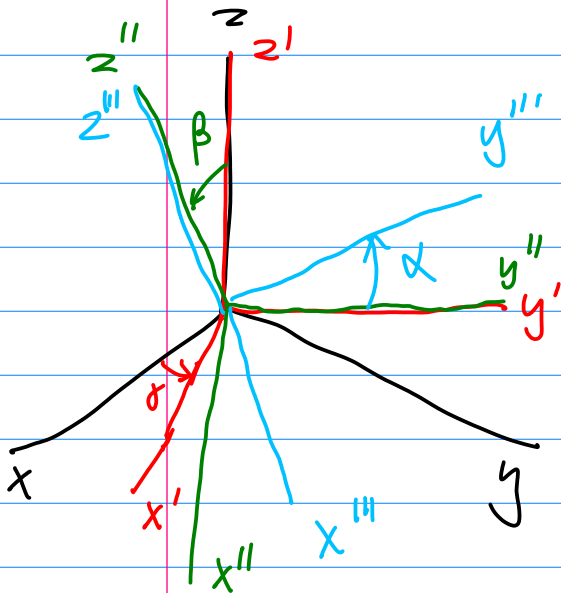
$$Q(\hat{j}, \beta) = \begin{pmatrix} \cos \beta & 0 & \sin \beta \\ 0 & 1 & 0 \\ -\sin \beta & 0 & \cos \beta \end{pmatrix}$$



Euler angles:
give a basis
 $\hat{i}, \hat{j}, \hat{k}$

- (1) rotation around z : γ
 - (2) rotation around y : β
 - (3) rotation around x : α
- } compose any rotations

$$\alpha \in [0, 2\pi), \quad \beta \in [0, \pi], \quad \gamma \in [0, 2\pi)$$



$$Q(\alpha, \beta, \gamma)$$

$$= Q(\hat{k}, \alpha) Q(\hat{j}, \beta) Q(\hat{i}, \gamma)$$

$$\vec{e}_i' = \sum_j \vec{e}_j Q_{ji}(\alpha, \beta, \gamma)$$

$$Q(\alpha, \beta, \gamma) = \begin{pmatrix} \cos\alpha \cos\beta \cos\gamma - \sin\alpha \sin\gamma & -\cos\alpha \cos\beta \sin\gamma - \sin\alpha \cos\gamma & \sin\alpha \cos\beta \\ \sin\alpha \cos\beta \cos\gamma + \cos\alpha \sin\gamma & \sin\alpha \cos\beta \sin\gamma + \cos\alpha \cos\gamma & -\sin\alpha \sin\beta \\ -\sin\alpha \cos\beta & \cos\alpha \cos\beta & \sin\beta \end{pmatrix}$$

$$\forall Q(\alpha, \beta, \gamma) \in SO(3)$$

all group elements in $SO(3)$ can be parametrized by

3 Euler angles

$SO(3)$ is a 3-dim Lie group

