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Report Title	Fourier Series Analysis using Matlab

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Introduction

The purpose of this lab is to use MATLAB to analyze periodic signals through Fourier series. The experiment involves calculating Fourier coefficients, plotting magnitude and phase spectra, and reconstructing signals from a limited number of harmonics. It also examines how a signal's period and the number of coefficients affect the accuracy of reconstruction.

Part A)

A.1

$$A_1: x_1(t) = \cos \frac{3\pi}{10} t + \frac{1}{2} \cos \frac{\pi}{10} t \quad \omega_0 = \frac{\pi}{10} \quad T_0 = \frac{2\pi}{\omega_0} = 20$$

$$\cos(k\omega_0 t) = \frac{1}{2} (e^{jk\omega_0 t} + e^{-jk\omega_0 t})$$

$$x_1(t) = \frac{1}{2} (e^{j3\omega_0 t} + e^{-j3\omega_0 t}) + \frac{1}{4} (e^{j\omega_0 t} + e^{-j\omega_0 t}) = \sum_{n=-\infty}^{\infty} D_n e^{jn\omega_0 t}$$

$$\therefore D_{\pm 3} = \frac{1}{2}, \quad D_{\pm 1} = \frac{1}{4}, \quad D_n = 0 \text{ otherwise}$$

A.2

$$A_2: D_n = \frac{1}{T_0} \int_{t_0}^{t_0 + T_0} x(t) e^{-jn\omega_0 t} dt \quad \omega_0 = \frac{2\pi}{T_0}$$

$$D_n = \frac{A}{T_0} \int_{t_s}^{t_s + \tau} e^{-jn\omega_0 t} dt = \frac{A}{T_0} \frac{1 - e^{-jn\omega_0 \tau}}{jn\omega_0} e^{-jn\omega_0 t_s}$$

$$D_0 = \frac{A}{T_0} \int_{t_s}^{t_s + \tau} 1 dt = \frac{A\tau}{T_0} \quad D_n = \frac{A\tau}{T_0} \operatorname{sinc}\left(\frac{n\omega_0 \tau}{2}\right) e^{-jn\omega_0 (t_s + \tau/2)} \left\{ \operatorname{sinc}(x) = \frac{\sin x}{x} \right.$$

$$x_2(t): A=1, \tau=10, T_0=20 \quad \omega_0 = \pi/10$$

$$D_n = \frac{\tau}{T_0} \operatorname{sinc}\left(\frac{n\omega_0 \tau}{2}\right) = \frac{10}{20} \operatorname{sinc}\left(\frac{n(\pi/10) \cdot 10}{2}\right)$$

$$D_n = \frac{1}{2} \operatorname{sinc}\left(\frac{n\pi}{2}\right)$$

$$x_3(t): A=1, \tau=10, T_0=40 \quad \omega_0 = \pi/20$$

$$D_n = \frac{\tau}{T_0} \operatorname{sinc}\left(\frac{n\omega_0 \tau}{2}\right) = \frac{10}{40} \operatorname{sinc}\left(\frac{n(\pi/20) \cdot 10}{2}\right)$$

$$D_n = \frac{1}{4} \operatorname{sinc}\left(\frac{n\pi}{4}\right)$$

A.3

```

Dn_x1.m × Dn_rect.m × A4.m × +
/Users/hamzah/Documents/MATLAB/Dn_x1.m
1 function D = Dn_x1(nvec)
2     D = zeros(size(nvec)); % preallocate
3     for k = 1:numel(nvec)
4         n = nvec(k);
5         if abs(n) == 3
6             D(k) = 1/2; % from  $\cos(3\omega_0 t) = 0.5(e^{j3\omega_0 t} + e^{-j3\omega_0 t})$ 
7         elseif abs(n) == 1
8             D(k) = 1/4; % from  $0.5 \cos(\omega_0 t) = 0.25(e^{j\omega_0 t} + e^{-j\omega_0 t})$ 
9         else
10            D(k) = 0;
11        end
12    end
13 end

```

Figure 1: Dn X1 code

```

/Users/hamzah/Documents/MATLAB/Dn_rect.m
1 function D = Dn_rect(nvec, A, T0, tau, ts)
2     w0 = 2*pi/T0;
3     D = zeros(size(nvec));
4     for k = 1:numel(nvec)
5         n = nvec(k);
6         if n == 0
7             D(k) = A * (tau / T0);
8         else
9             D(k) = (A/T0) * (1 - exp(-1j*n*w0*tau)) / (1j*n*w0 * exp(-1j*n*w0*ts));
10        end
11    end
12 end

```

Figure 2: Dn X2/X3 code

A.4

```
1 % Parameters
2 A_x2 = 1; T0_x2 = 20; tau_x2 = 10; ts_x2 = -5;
3 A_x3 = 1; T0_x3 = 40; tau_x3 = 10; ts_x3 = -5;
4
5 ranges = { -5:5, -20:20, -50:50, -500:500 };
6 labels = { 'n=-5..5', 'n=-20..20', 'n=-50..50', 'n=-500..500' };
7
8 % x1
9 for i = 1:numel(ranges)
10     nvec = ranges{i};
11     D = Dn_x1(nvec);
12     figure;
13     subplot(2,1,1); stem(nvec,abs(D),'filled'); grid on
14     ylabel('|D_n|'); title(['x_1 ', labels{i}])
15     subplot(2,1,2); stem(nvec,angle(D),'filled'); grid on
16     xlabel('n'); ylabel('\angle D_n (rad)')
17 end
18
19 % x2
20 for i = 1:numel(ranges)
21     nvec = ranges{i};
22     D = Dn_rect(nvec, A_x2, T0_x2, tau_x2, ts_x2);
23     figure;
24     subplot(2,1,1); stem(nvec,abs(D),'filled'); grid on
25     ylabel('|D_n|'); title(['x_2 ', labels{i}])
26     subplot(2,1,2); stem(nvec,angle(D),'filled'); grid on
27     xlabel('n'); ylabel('\angle D_n (rad)')
28 end
29
30 % x3
31 for i = 1:numel(ranges)
32     nvec = ranges{i};
33     D = Dn_rect(nvec, A_x3, T0_x3, tau_x3, ts_x3);
34     figure;
35     subplot(2,1,1); stem(nvec,abs(D),'filled'); grid on
36     ylabel('|D_n|'); title(['x_3 ', labels{i}])
37     subplot(2,1,2); stem(nvec,angle(D),'filled'); grid on
38     xlabel('n'); ylabel('\angle D_n (rad)')
39 end
```

Figure 3: A4 Code

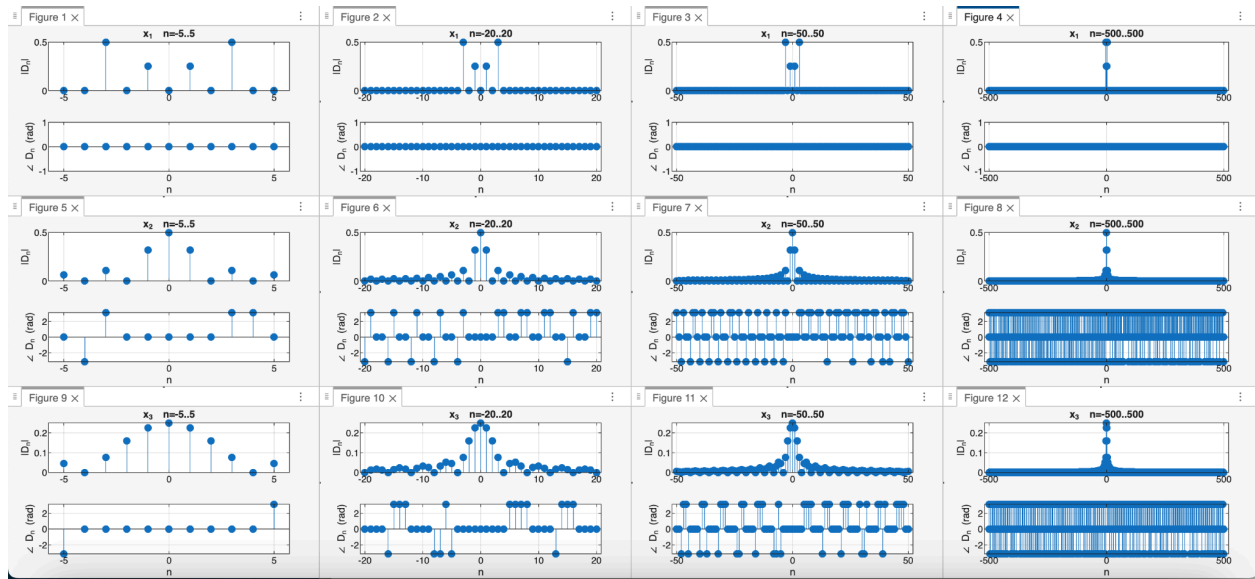


Figure 4: $X_1(t)$, $X_2(t)$, $X_3(t)$ magnitude and phase spectra plots

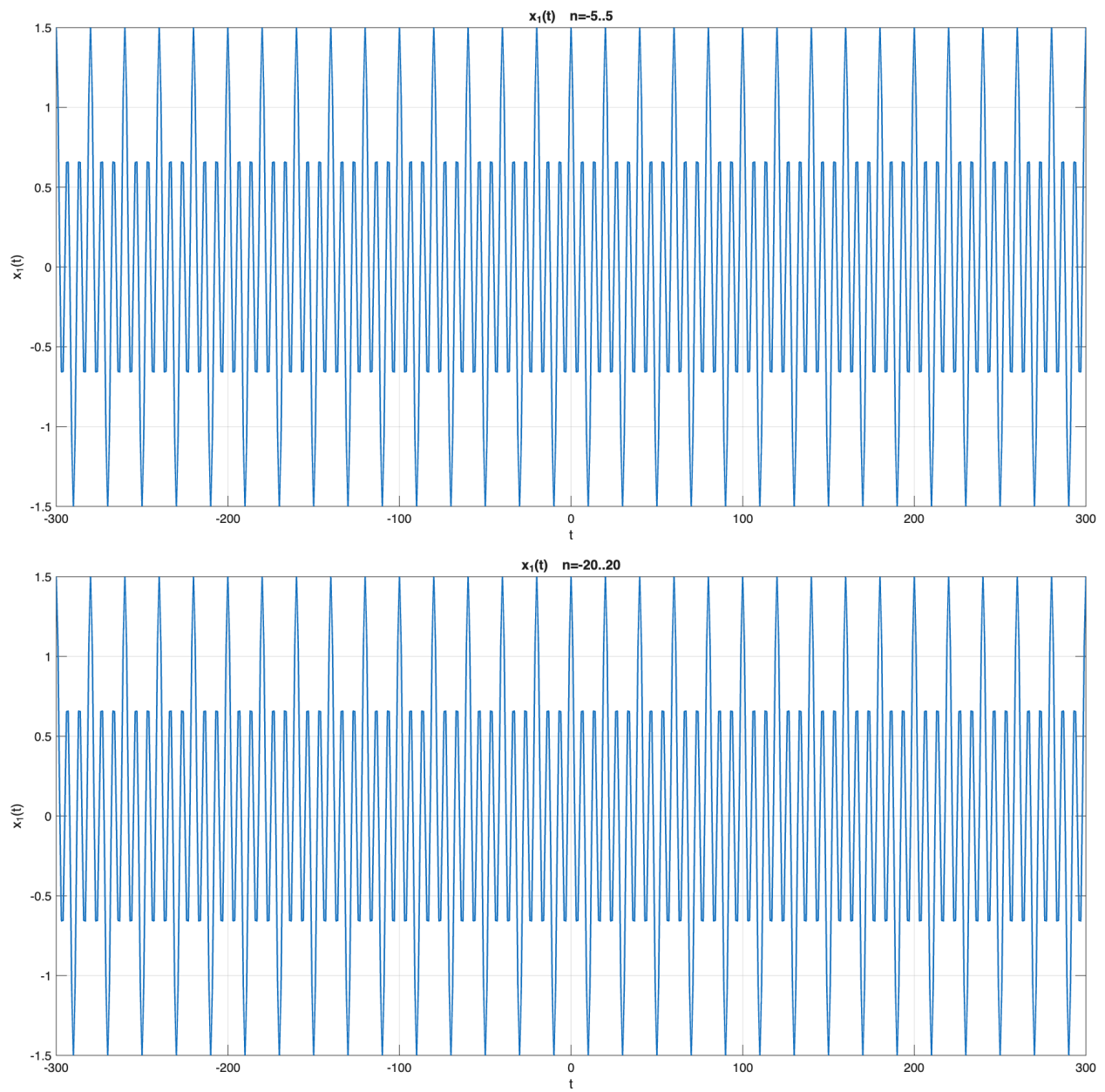
A.5/A.6

```

1      % A5/A6
2
3      % Parameters
4      T0_x1 = 20;
5      A_x2 = 1; T0_x2 = 20; tau_x2 = 10; ts_x2 = -5;
6      A_x3 = 1; T0_x3 = 40; tau_x3 = 10; ts_x3 = -5;
7
8      ranges = { -5:5, -20:20, -50:50, -500:500 };
9      labels = { 'n=-5..5', 'n=-20..20', 'n=-50..50', 'n=-500..500' };
10
11     % x1 reconstructions
12     for i = 1:numel(ranges)
13         nvec = ranges{i};
14         D = Dn_x1(nvec);
15         [t,xrec] = reconstruct_from_Dn(nvec, D, T0_x1);
16         figure; plot(t, xrec, 'LineWidth', 1.2); grid on
17         title(['x_1(t) ', labels{i}]); xlabel('t'); ylabel('x_1(t)')
18     end
19
20     % x2 reconstructions
21     for i = 1:numel(ranges)
22         nvec = ranges{i};
23         D = Dn_rect(nvec, A_x2, T0_x2, tau_x2, ts_x2);
24         [t,xrec] = reconstruct_from_Dn(nvec, D, T0_x2);
25         figure; plot(t, xrec, 'LineWidth', 1.2); grid on
26         title(['x_2(t) ', labels{i}]); xlabel('t'); ylabel('x_2(t)')
27     end
28
29     % x3 reconstructions
30     for i = 1:numel(ranges)
31         nvec = ranges{i};
32         D = Dn_rect(nvec, A_x3, T0_x3, tau_x3, ts_x3);
33         [t,xrec] = reconstruct_from_Dn(nvec, D, T0_x3);
34         figure; plot(t, xrec, 'LineWidth', 1.2); grid on
35         title(['x_3(t) ', labels{i}]); xlabel('t'); ylabel('x_3(t)')
36     end
37
38     % local functions
39
40     function [t, xrec] = reconstruct_from_Dn(nvec, D, T0)
41         w0 = 2*pi/T0;
42         t = -300:300;
43         xrec = zeros(size(t));
44         for k = 1:numel(nvec)
45             n = nvec(k);
46             xrec = xrec + D(k).*exp(1j*n*w0*t);
47         end
48         xrec = real(xrec);
49     end
50
51     function D = Dn_x1(nvec)
52         D = zeros(size(nvec));
53         for k = 1:numel(nvec)
54             n = nvec(k);
55             if abs(n)==3, D(k)=1/2;
56             elseif abs(n)==1, D(k)=1/4;
57             else, D(k)=0;
58             end
59         end
60     end
61
62     function D = Dn_rect(nvec, A, T0, tau, ts)
63         w0 = 2*pi/T0; D=zeros(size(nvec));
64         for k=1:numel(nvec)
65             n=nvec(k);
66             if n==0
67                 D(k)=A*(tau/T0);
68             else
69                 D(k)=(A/T0)*(1-exp(-1j*n*w0*tau))/(1j*n*w0)*exp(-1j*n*w0*ts);
70             end
71         end
72     end

```

Figure 5: A5/A6 Code



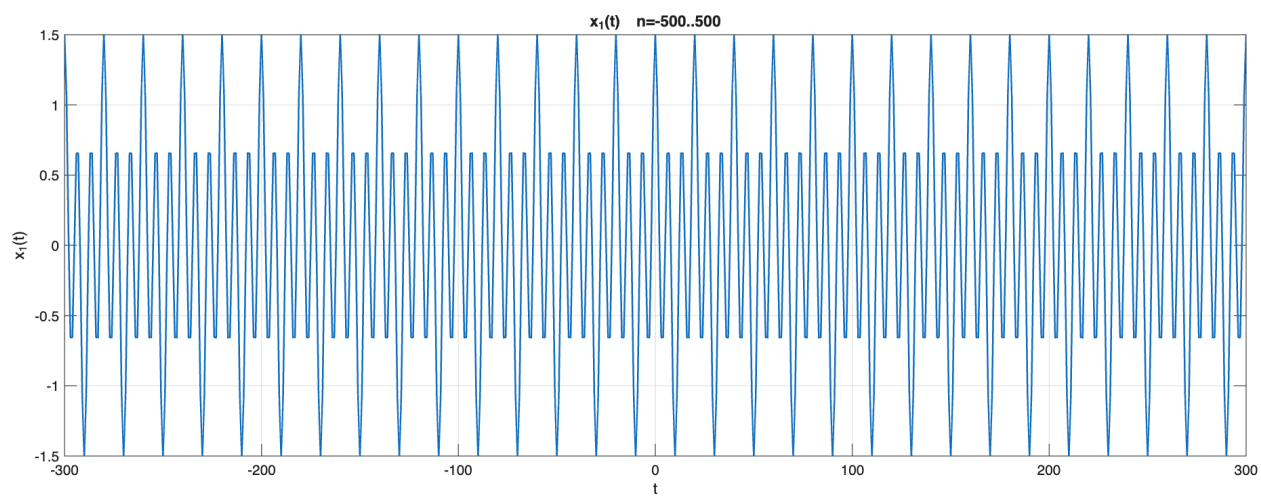
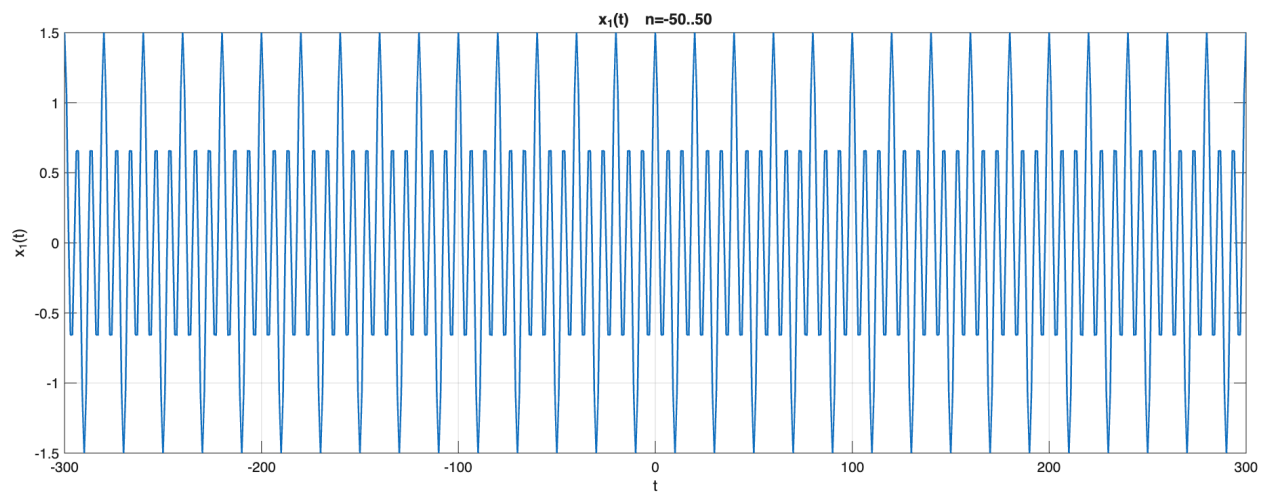
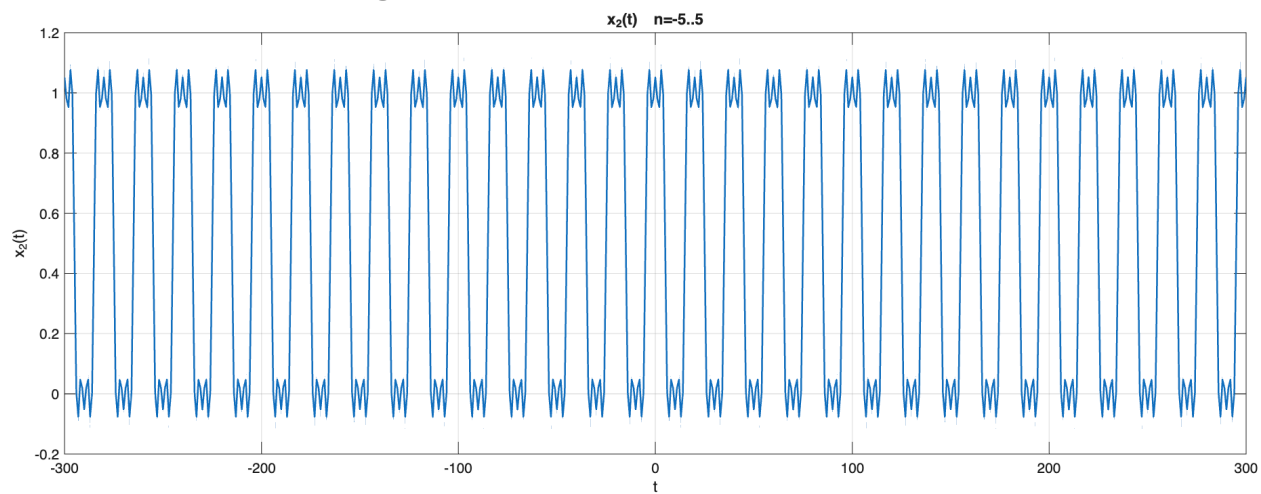


Figure 6: $X_1(t)$ Reconstructed



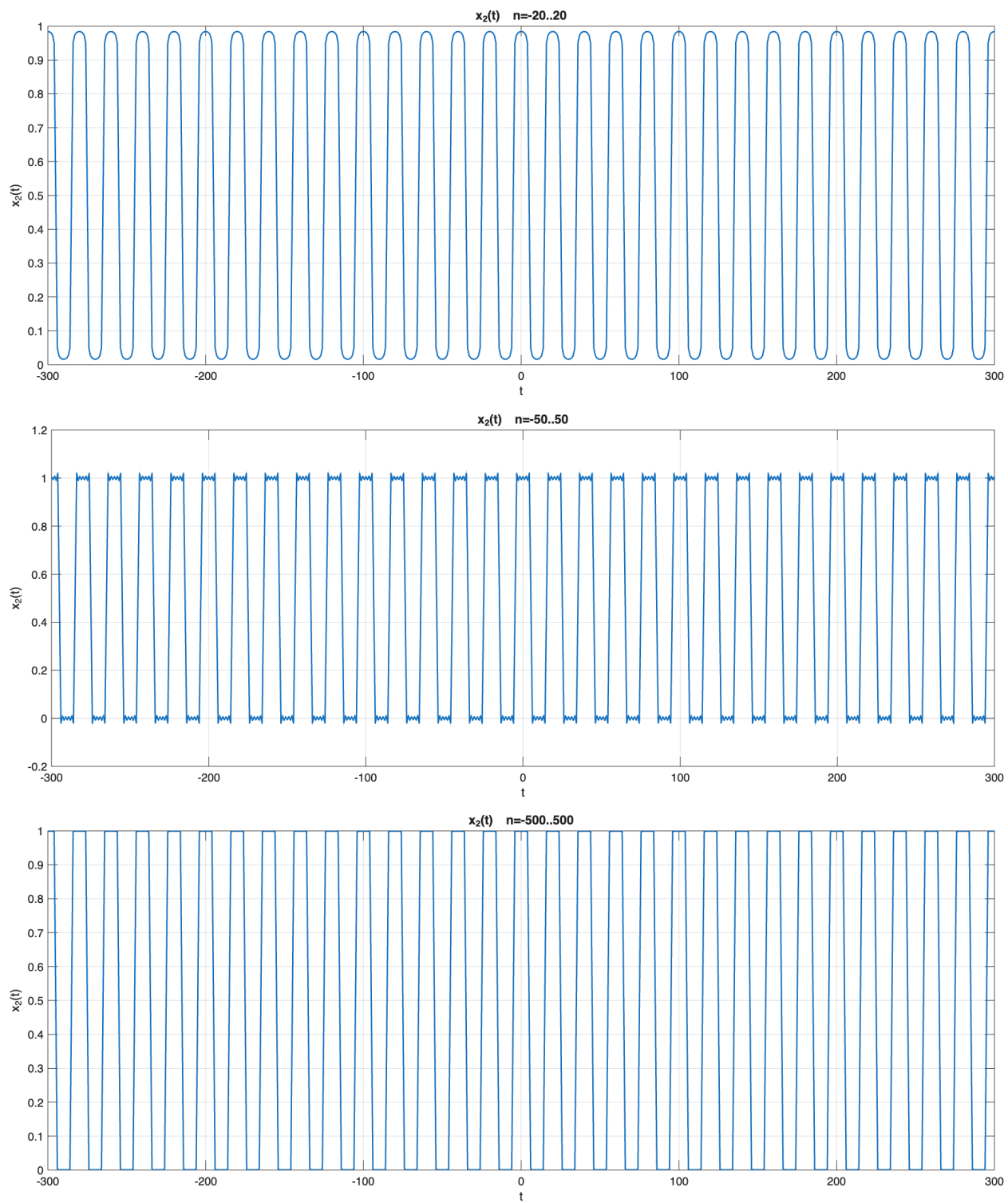
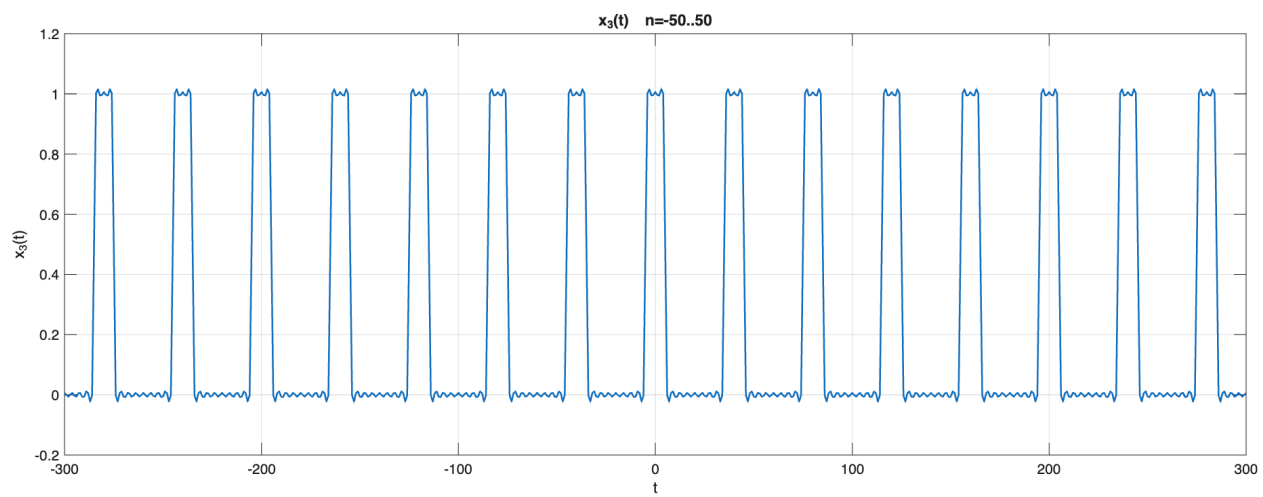
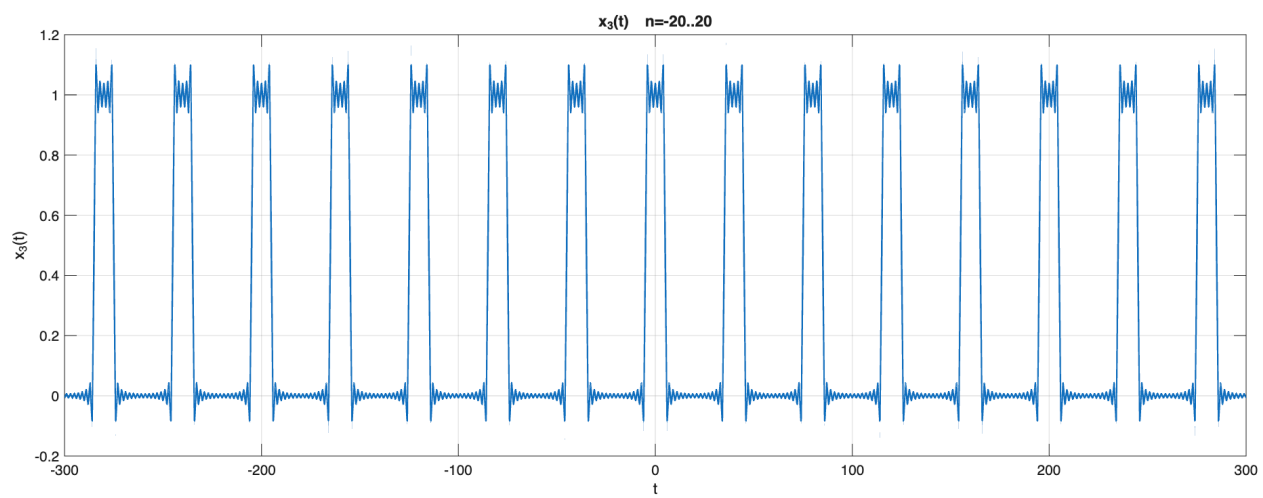
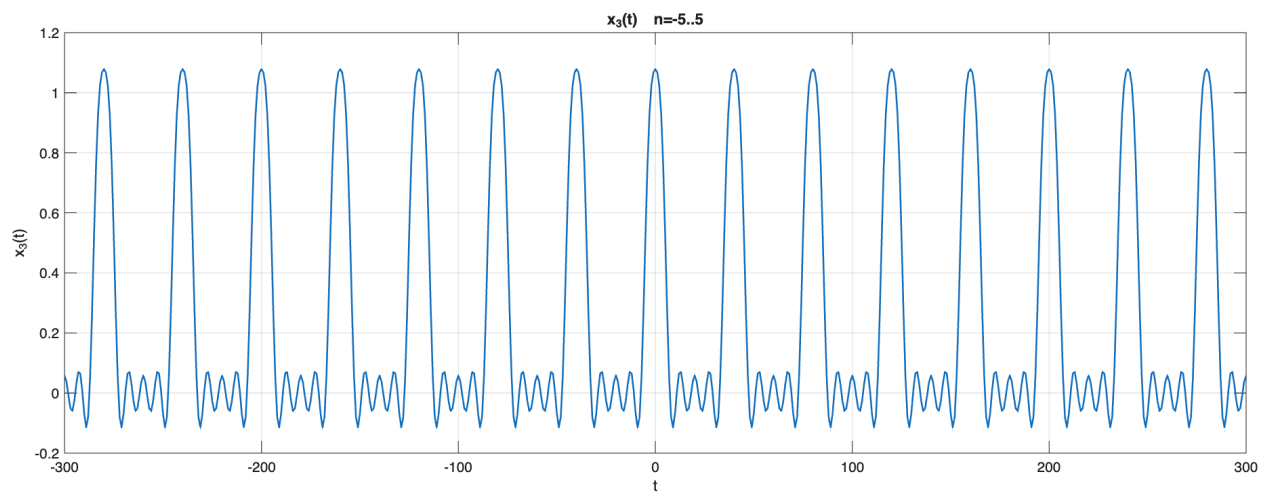


Figure 7: $X_2(t)$ Reconstructed



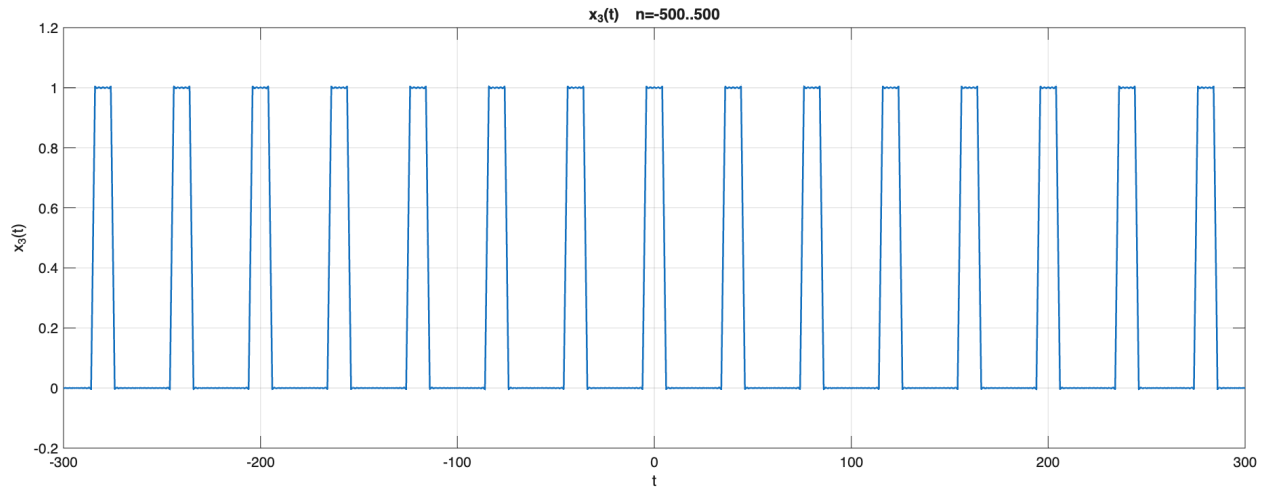


Figure 8: X3(t) Reconstructed

Part B)

B.1

B.1) $\omega_0 = \frac{2\pi}{T_0}$ $x_1(t) = \cos\left(\frac{3\pi}{10}t\right) + \frac{1}{2}\cos\left(\frac{\pi}{10}t\right)$

Component ω are $\frac{3\pi}{10}$, $\frac{\pi}{10}$ rad/s

$\therefore \omega_{0,1} = \frac{\pi}{10}$ rad/s (gcd)

$T_{0,1} = \frac{2\pi}{\omega_{0,1}} = \frac{2\pi}{\pi/10} = 20s$

$x_2(t)$: $T_{0,2} = 20s$ (shown in graph)

$\omega_{0,2} = \frac{2\pi}{20} = \frac{\pi}{10}$ rad/s

$x_3(t)$: $T_{0,3} = 40s$ (graph)

$\omega_{0,3} = \frac{2\pi}{40} = \frac{\pi}{20}$ rad/s

B.2

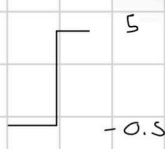
The main difference between the Fourier coefficients of $x_1(t)$ and $x_2(t)$ is that $x_1(t)$ contains only a few discrete frequencies because it is a sum of two cosines, while $x_2(t)$ has infinitely many harmonics since it is a rectangular pulse train. The coefficients for $x_1(t)$ exist only at certain n values, whereas $x_2(t)$ has coefficients that decay gradually following a sinc shaped envelope and include a DC component.

B.3

Despite having the same pulse shape, $x_2(t)$ and $x_3(t)$ differ in their periods, which have an impact on how far apart their Fourier coefficients are. The fundamental frequency separation between harmonics decreases with increasing period and increases with decreasing period. Consequently, the harmonics of $x_3(t)$ are closer together than those of $x_2(t)$.

B.4

$x_4(t)$



$T = 20$

Odd signal

$$D_n = \frac{2}{T} \int_{-T/2}^{T/2} x(t) dt$$

$$= \frac{2}{20} \left[\int_0^5 \frac{1}{2} dt - \int_5^{15} \frac{1}{2} dt + \int_{15}^{20} \frac{1}{2} dt \right]$$

$$= \frac{2}{20} \left[\frac{1}{2} t \Big|_0^5 - \frac{1}{2} t \Big|_5^{15} + \frac{1}{2} t \Big|_{15}^{20} \right]$$

$$= \frac{2}{20} (2.5 - 5 + 2.5)$$

$$= 0$$

$x_2(t)$: $T = 20$

$$D_n = \frac{1}{T} \int_{-T/2}^{T/2} x_2(t) dt$$

$$= \frac{1}{20} \left[\int_{-10}^{-5} 0 dt + \int_{-5}^5 1 dt + \int_5^{10} 0 dt \right]$$

$$= \frac{1}{20} \left[t \Big|_{-5}^5 \right]$$

$$D_n^{(2)} = \frac{1}{20} (10)$$

$$= 0.5$$

$D_n^{(4)} = 0$

$D_n(x_4) = D_n(x_2) - 0.5 = 0$

B.5

With an increase in the number of Fourier coefficients, the reconstructed signal approaches the original waveform more closely. Once all of the current harmonics are taken into account, flawless reconstruction for $x_1(t)$ happens. The approximation error for $x_2(t)$ decreases with increasing coefficients however, the Gibbs phenomenon causes oscillations to persist near the edges. With the addition of more terms, these oscillations become narrower but remain visible.

B.6

A finite number of coefficients is enough for $x_1(t)$ because it contains only a few harmonic components. However, $x_2(t)$ and $x_3(t)$ are discontinuous and therefore require infinitely many coefficients for perfect reconstruction.

B.7

Since a periodic signal generally has an infinite number of Fourier coefficients D , storing all of them is not practical. However, if the signal is finite, such as $x_1(t)$, then the corresponding D values can be stored. This approach is not recommended for signals with a large number of significant coefficients, as it would be inefficient and waste memory.

Conclusion

In this lab, MATLAB was used to generate and reconstruct signals using their Fourier series coefficients. The results showed that smooth signals can be represented exactly with a few harmonics, while discontinuous signals require many terms and exhibit Gibbs oscillations. Overall, the lab demonstrated how Fourier series link time-domain periodic signals to their frequency-domain representations.