



# Out-of-Time-order Correlator in billiard systems

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# Outline

1. Background

2. Numerical result

- 1) Short time behavior
- 2) Long time behavior

3. Analytical analysis

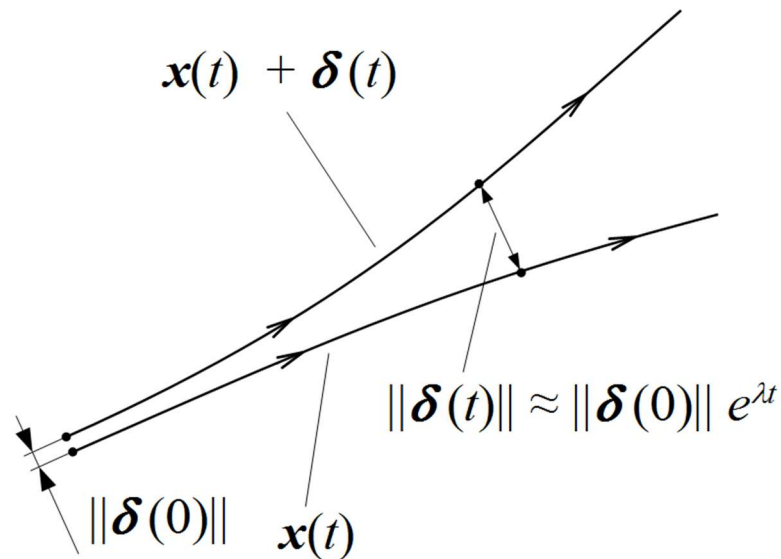
4. Conclusion

C.-D. Han, H.-Y. Xu and Y.-C. Lai, “Out-of-Time-order Correlator in billiard systems,”  
submitted to *Physical Review Letters*.



# Classical chaos

**Lyapunov exponent** characterizes the rate of separation of infinitesimally close trajectories



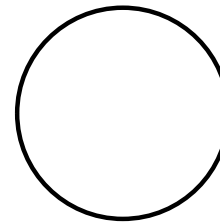
From: Wiki

# Quantum chaos

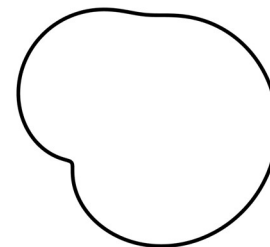
**Classical chaotic system**  $\longrightarrow$  **Quantum chaos**

Billiard system:

Integrable



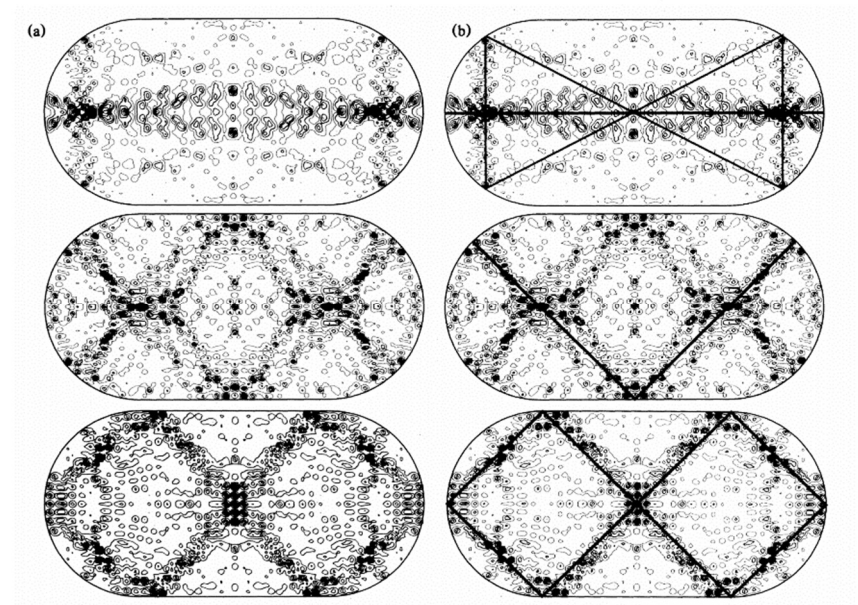
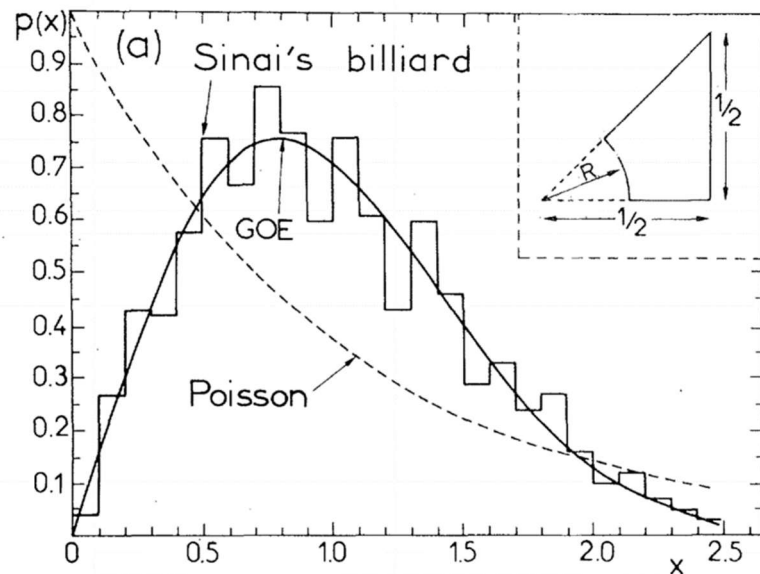
Chaotic



# Quantum chaos

**BGS conjecture:** Random matrix theory

**Quantum scar**



Bohigas, Oriol, Marie-Joya Giannoni, and Charles Schmit. "Characterization of chaotic quantum spectra and universality of level fluctuation laws." *Physical Review Letters* 52.1 (1984): 1.

Heller, Eric J. "Bound-state eigenfunctions of classically chaotic Hamiltonian systems: scars of periodic orbits." *Physical Review Letters* 53.16 (1984): 1515.

# Quantum-Classical correspondence

Nonlinear dynamics (**Classical mechanics**)

**Quantum mechanics**

Two  
Order

$$\frac{\Delta x(t)}{\Delta x(0)} \sim e^{\gamma t} \quad \{x(t), p(0)\}_P$$

$$\frac{[\hat{x}(t), \hat{p}(0)]}{i\hbar}$$

Four  
Order

$$\langle \{x(t), p(0)\}_P^2 \rangle$$

$$- \langle [\hat{x}(t), \hat{p}(0)]^2 \rangle_\beta \quad \beta = 1/T$$

Out-of-Time-order  
Correlator

$$C_T(t) = -\langle [\hat{W}(t), \hat{V}(0)]^2 \rangle_\beta$$

Larkin, A. I., and Yu N. Ovchinnikov. "Quasiclassical method in the theory of superconductivity." *Sov Phys JETP* 28.6 (1969): 1200-1205.

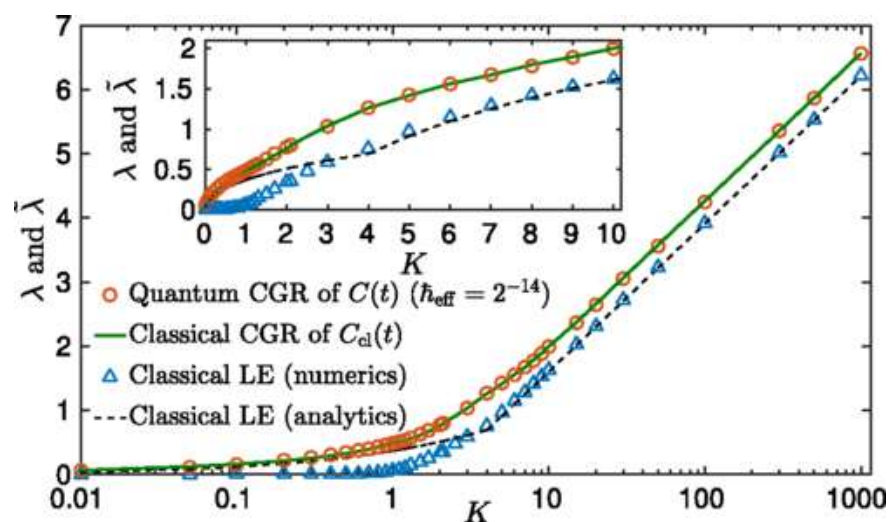
Maldacena, Juan, Stephen H. Shenker, and Douglas Stanford. "A bound on chaos." *Journal of High Energy Physics* 2016.8 (2016): 106.

$$C_T(t) = \frac{\sum_n c_n \exp(-|E_n|/T)}{\exp(-|E_n|/T)}$$

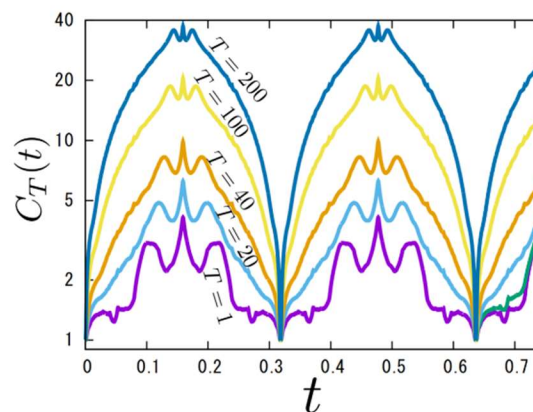
$$c_n(t) = -\langle n | [\hat{W}(t), \hat{V}]^2 | n \rangle$$

# Quantum-Classical discordance

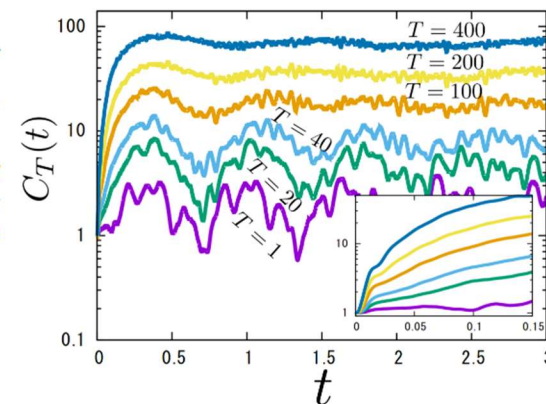
## Kicked Rotor



## Schrodinger Billiard



(a) Particle in a box



(b) Stadium billiard

Rozenbaum, Efim B., Sriram Ganeshan, and Victor Galitski. "Lyapunov exponent and out-of-time-ordered correlator's growth rate in a chaotic system." *Physical review letters* 118.8 (2017): 086801.

Hashimoto, Koji, Keiju Murata, and Ryosuke Yoshii. "Out-of-time-order correlators in quantum mechanics." *Journal of High Energy Physics* 2017.10 (2017): 138.

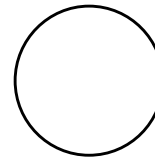


# Model

Schrodinger system

$$\hat{H} = \frac{\hat{\mathbf{p}}^2}{2m} + V(\mathbf{r})$$

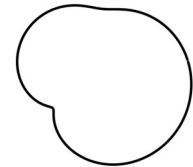
$$V(\mathbf{r}) = \begin{cases} 0 & \mathbf{r} \in D, \\ \infty & \mathbf{r} \notin D \end{cases} \quad \text{Circle billiard}$$



Dirac system

$$\hat{H} = \hat{\boldsymbol{\sigma}} \cdot \hat{\mathbf{p}} + V(\mathbf{r})\hat{\sigma}_z$$

African billiard



$$C_T(t) = -\langle [\hat{W}(t), \hat{V}(0)]^2 \rangle_\beta$$

$$\hat{W} = \hat{V} = \frac{\partial \hat{H}}{\partial \hat{p}}$$

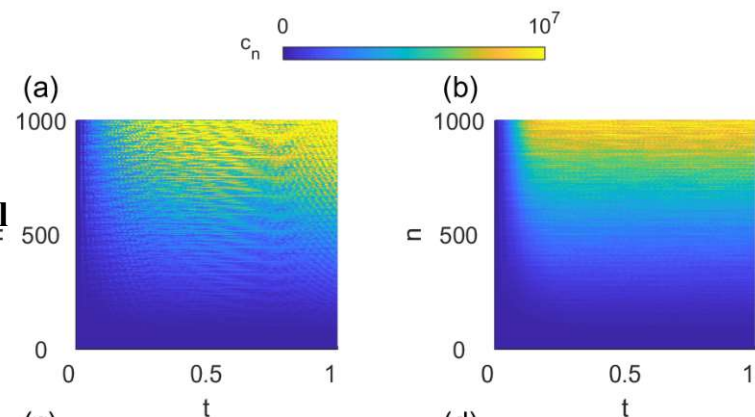
$$C_T(t) = -\langle [\hat{p}_x(t), \hat{p}_x]^2 \rangle_\beta \sim \left( \frac{\Delta p_x(t)}{\Delta x(0)} \right)^2$$

$$C_T(t) = -\langle [\hat{\sigma}_x(t), \hat{\sigma}_x]^2 \rangle_\beta$$

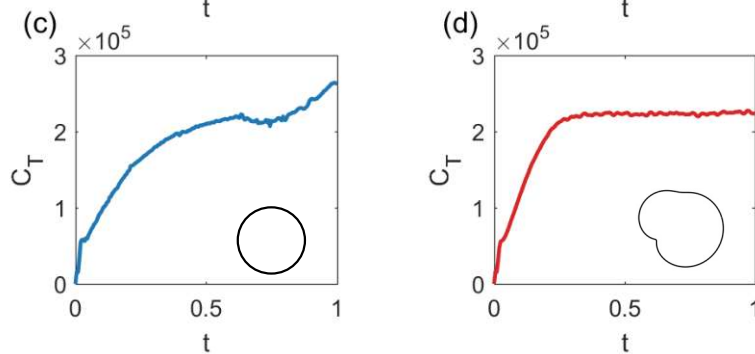
Classical view: Wave-packet

**Result: Small time**  $t < t_R$

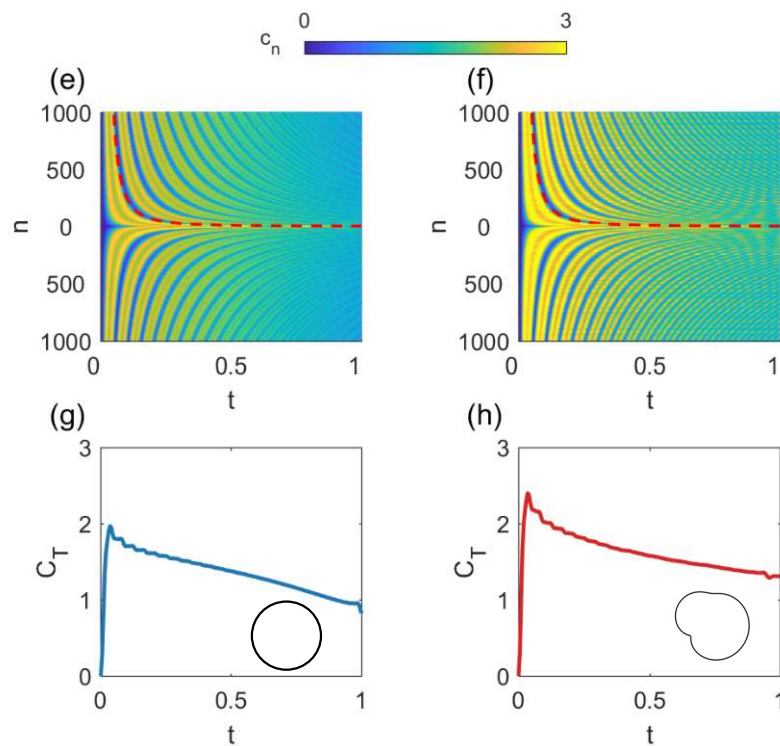
**Micro-canonical  
OTOC**



**Canonical  
OTOC**



Schrodinger system (a) – (d)

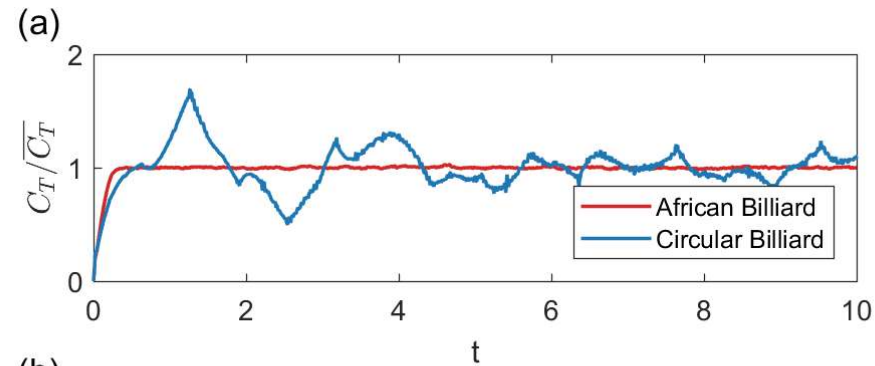


Dirac system (e) – (h)

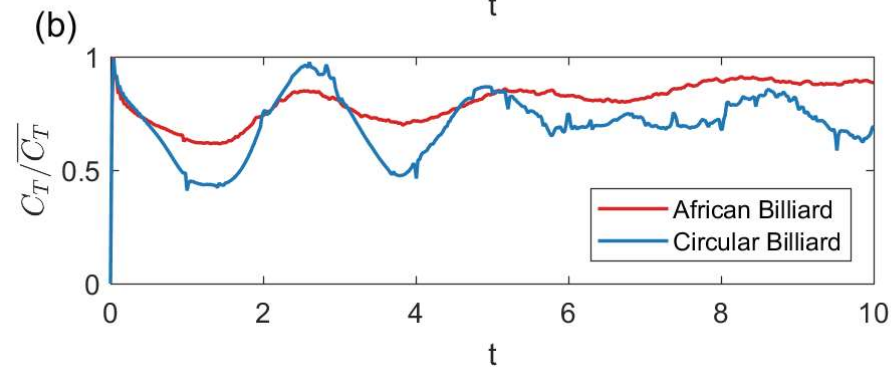
$$t = \frac{2\pi}{2|E_n|}$$

# Result: Long time $t > t_R$

Schrodinger system



Dirac system



Long time behavior could diagnosed quantum chaos

$$\overline{C}_T(t) = \langle \hat{W}(t) \hat{V} \hat{V} \hat{W}(t) + \hat{V} \hat{W}(t) \hat{W}(t) \hat{V} \rangle_{\beta}$$

# Analytical analysis

$$C_T(t) = -\langle [\hat{W}(t), \hat{V}(0)]^2 \rangle_\beta$$

1. Find the maximum value
2. Truncate and make approximation

**1D system:** Harmonic oscillator

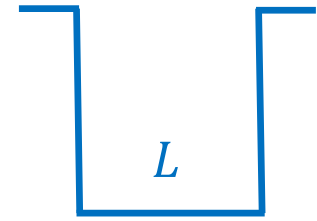
$$C_T = -\langle [\hat{x}(t), \hat{p}_x]^2 \rangle_\beta = \left\langle \left( \frac{\Delta x(t)}{\Delta x(0)} \right)^2 \right\rangle = (\cos t)^2$$

$$C_T = -\langle [\hat{p}(t), \hat{p}_x]^2 \rangle_\beta = \left\langle \left( \frac{\Delta p_x(t)}{\Delta x(0)} \right)^2 \right\rangle = (\omega \sin t)^2$$

Hashimoto, Koji, Keiju Murata, and Kentaroh Yoshida.  
 "Chaos in chiral condensates in gauge theories." *Physical review letters* 117.23 (2016): 231602.

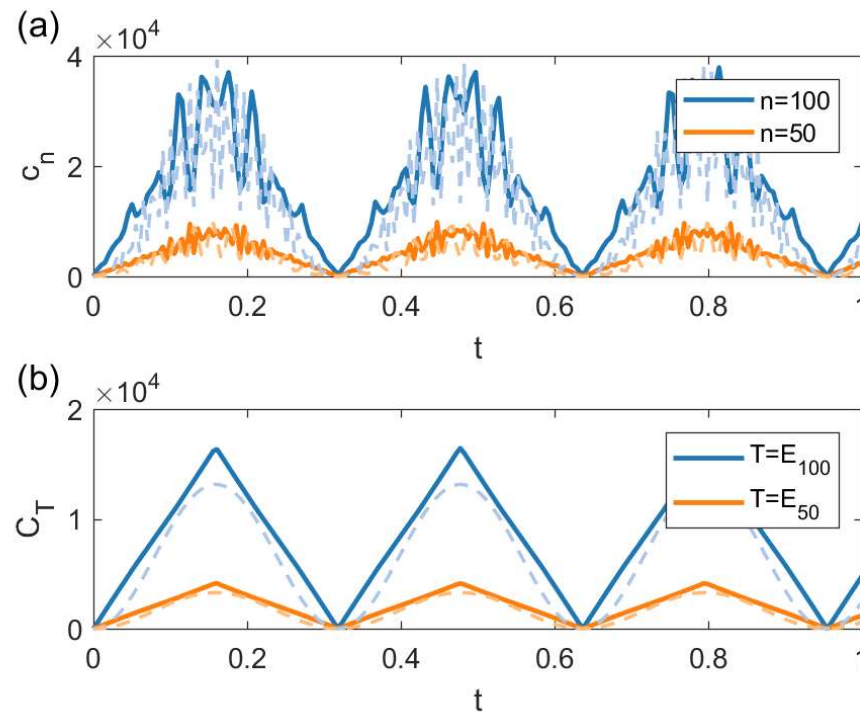
# Analytical analysis

**1D system:** Infinite potential well



$$c_n(t) = \frac{64n^2L^2}{\pi^4} \left[ 4 \sin^4 \left( \frac{\pi^2}{L^2} t \right) \sin^2 \left( \frac{\pi^2 n}{L^2} t \right) + 1 - \cos \left( \frac{2\pi^2}{L^2} t \right) \right]$$

$$C_T(t) = \frac{64L^4}{\pi^6} T \left[ 1 - \cos \left( \frac{2\pi^2}{L^2} t \right) \right]$$



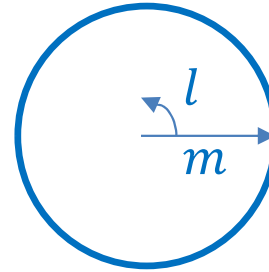
— Numerical

- - - Analytical

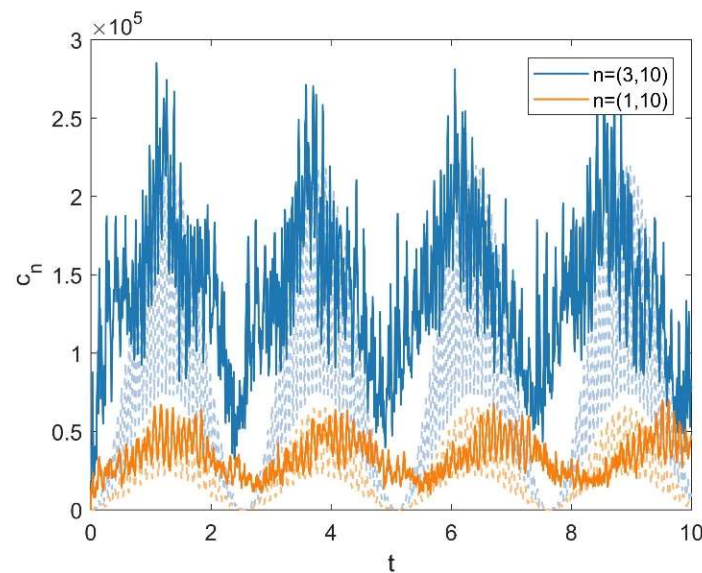
# Analytical analysis

**2D system:**

Circular billiard  
(Schrodinger system)



$$c_{ml} \approx 16x_m^4 k_{ml}^4 d_m^4 \left[ \sin^2 \left( \frac{d_m^2}{\hbar} \right) \cos^2 \left( \frac{2k_{ml}d_m}{\hbar} \right) \right] + 4x_m^4 k_{ml}^4 d_m^4 \left( 1 - \cos \frac{2d_m^2}{\hbar} \right) \quad l \gg m$$



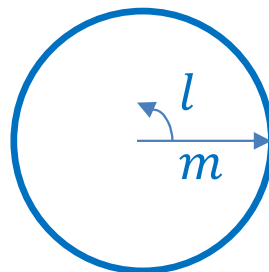
— Numerical

- - - - Analytical

# Analytical analysis

**2D system:**

Circular billiard  
(Dirac system)



$$C_T(t) = -\langle [\hat{W}(t), \hat{V}(0)]^2 \rangle_\beta$$

Particle-Particle: **Canceled**

Particle-Hole: **Particle-Hole  
symmetry**

$$\hat{A} = \hat{U} \hat{K}$$

$$\hat{H}' = \hat{A} \hat{H} \hat{A}^{-1} = -\hat{H}$$

**Zitterbewegung**

Mondragon, R. J. "Neutrino billiards: time-reversal symmetry-breaking without magnetic fields." *Proc. R. Soc. Lond. A*. Vol. 412. No. 1842. The Royal Society, 1987.  
Bjorken, James D., and Sidney David Drell. "Relativistic quantum mechanics." (1964).

## Analytical analysis

**Wave-packet revival:** periodic motion for wave-packet

$$E(n) \approx E(n_0) + E'(n_0)(n - n_0) + \frac{E''(n_0)}{2}(n - n_0)^2 + \frac{E'''(n_0)}{6}(n - n_0)^3 + \dots ,$$

$$T_{\text{cl}} = \frac{2\pi\hbar}{|E'(n_0)|}, \quad T_{\text{rev}} = \frac{2\pi\hbar}{|E''(n_0)|/2}, \quad \text{and} \quad T_{\text{super}} = \frac{2\pi\hbar}{|E'''(n_0)|/6} .$$

Infinite potential well(Schrodinger system)	$T_0 = T_{\text{rev}} = \frac{4mL^2}{\pi\hbar}$
Circular billiard (Schrodinger system)	$T_0 = 4T_{\text{rev}} = \frac{8mR^2}{\pi\hbar}$
Circular billiard (Dirac system)	$T_0 = T_{\text{cl}} = 4R$

Robinett, R. W., and S. Heppelmann. "Quantum wave-packet revivals in circular billiards." *Physical Review A* 65.6 (2002): 062103.

Robinett, Richard W. "Quantum wave packet revivals." *Physics Reports* 392.1-2 (2004): 1-119.

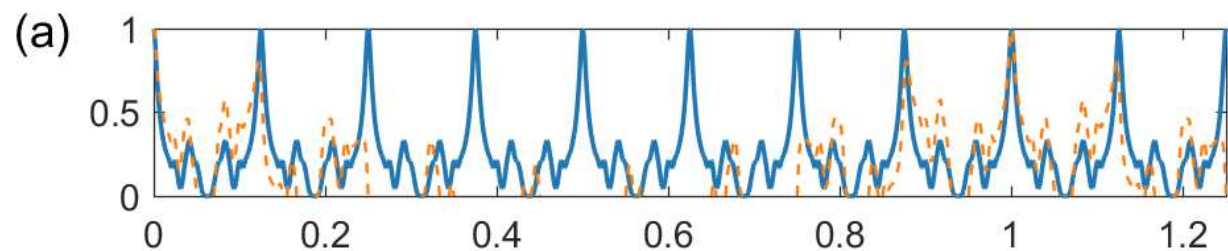


# Analytical analysis

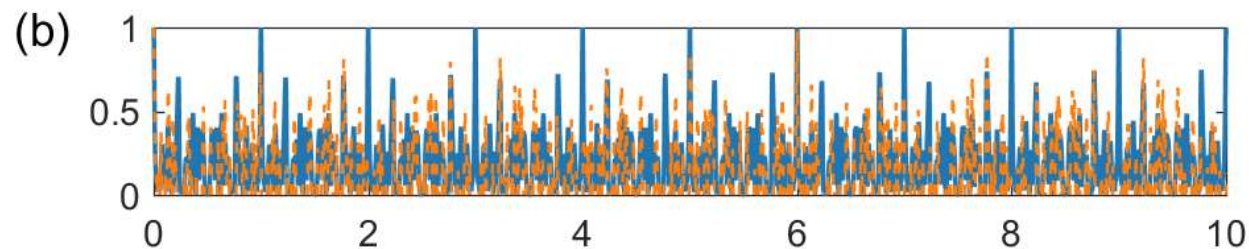
**Wave-packet revival:** periodic motion for wave-packet

—  $A = |\langle \psi(t) | \psi(0) \rangle|^2$ 
- - -  $\text{Re} \langle \psi(t) | \psi(0) \rangle$

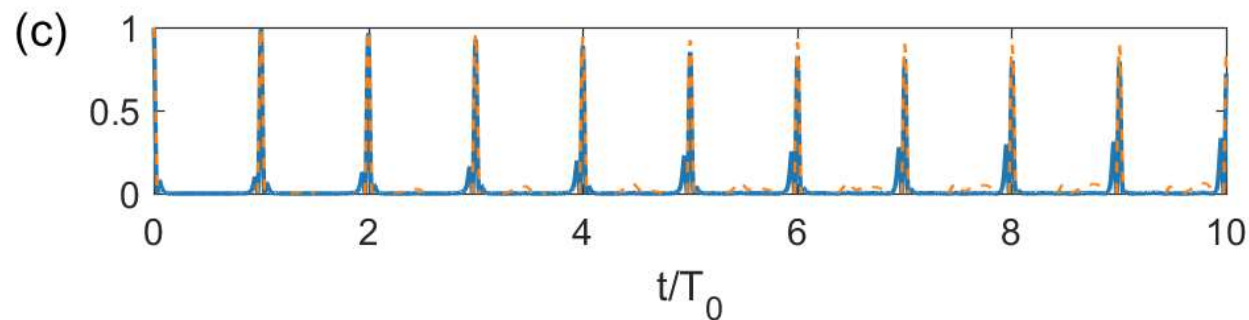
Infinite potential well



Circular billiard  
(Schrodinger system)



Circular billiard  
(Dirac system)



Center wave-packet

$$2\sigma \sim \frac{1}{10} \text{ system size} \quad k_0 = 0$$

# Analytical analysis

Fourier Analysis for  $C_T(t)$   $T = E_{100}$

Infinite potential well

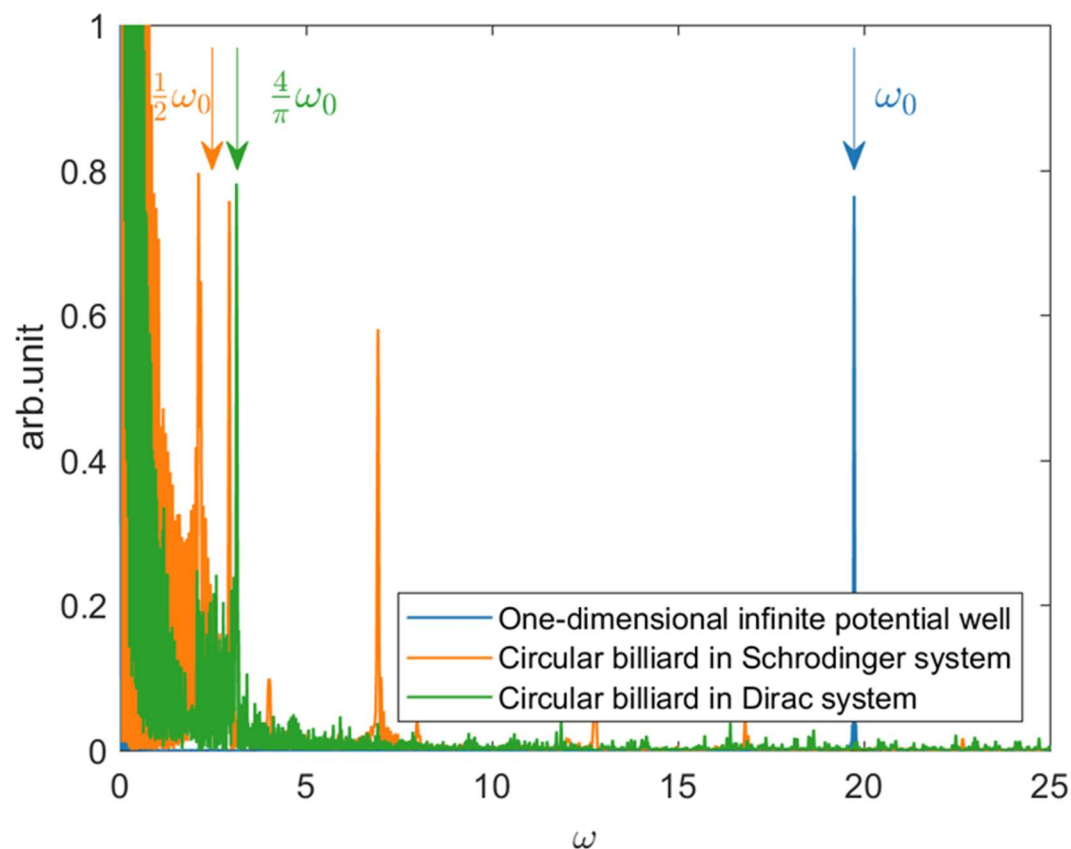
Circular billiard  
(Schrodinger system)

Circular billiard  
(Dirac system)

$T_0$

$2T_0$

$\frac{\pi}{4}T_0$



# Conclusion

1. Dynamical independent Zitterbewegung motion in Dirac billiard

**Particle-Hole symmetry**

2. Long time behavior of OTOC can distinguish quantum chaos

**Wave-packet revival**



**Thank you**