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- (a)  $\int_a^b \rho(x) dx$
- (b)  $\int_a^b \rho(x) u_{tt}(x, t) dx$
- (c)  $\int_a^b -\rho(x) g dx$
- (d) The tension is pulling in opposite directions, making the tension at  $x = a$  negative and the tension at  $x = b$  positive. Thus, the left vertical force is  $-T \sin(\theta_a)$  and the right vertical force  $T \sin(\theta_b)$ . We are ignoring horizontal force because we are assuming there is a small displacement in the string from  $x = a$  to  $x = b$ , making the horizontal force negligible. We see this by noting that when  $\theta$  is small,  $\cos(\theta) \approx 1$ , so  $T_x \approx 0$
- (e) When the angle between the tangent line and the string is small at a given point, it means that  $\cos \theta \approx \sin \theta$  (since the hypotenuse and adjacent side are similar in length, while the opposite side is very small).
- (f) We know that  $\tan(\theta)$  is equal to the length of the opposite side divided by the adjacent side of a right triangle. In this scenario,  $\theta$  is the angle between the line formed by the string and the horizontal force. Therefore, there is a right triangle between the horizontal and the vertical forces, so  $\tan(\theta)$  is equal to the magnitude of the vertical force divided by the horizontal force. We can note, therefore, that  $\tan(\theta)$  is now also equivalent to the slope of the line segment of our string (rise over run!). We also know that  $u_x(x, t)$  is equivalent to the slope of the line at a certain point  $x$ . Therefore,  $u_x(a, t)$  is the slope of the line at  $x = a$ , which we explained is also the same as  $\tan(\theta_a)$ . Logically,  $u_x(b, t)$  is the slope of the line at  $x = b$ , which is also the same as  $\tan(\theta_b)$ . So,  $u_x(a, t) = \tan(\theta_a)$  and  $u_x(b, t) = \tan(\theta_b)$
- (g)  $-Tu_x(a, t) + Tu_x(b, t)$
- (h)  $\int_a^b -Tu_{xx}(x, t) dx$
- (i)  $\int_a^b \rho(x) u_{tt}(x, t) dx = \int_a^b -\rho(x) g dx + \int_a^b Tu_{xx}(x, t) dx$   
 $\int_a^b \rho(x) u_{tt}(x, t) dx = \int_a^b Tu_{xx}(x, t) - \rho(x) g dx$
- (j) Since we know  $a$  and  $b$  are arbitrary, the only way (i) can be true is if the elements inside the integrals on both sides are equal, so  $\rho(x) u_{tt} = Tu_{xx} - \rho(x) g$   
 Now, we are assuming density is constant along the string, so therefore the string tension can be treated as a constant independent of  $x$ . Thus, we can divide everything by  $\rho(x) = \rho$  and end up with  $u_{tt} = \frac{T}{\rho} u_{xx} - g$   
 Now, we can let a constant  $c = \sqrt{\frac{T}{\rho}}$  so that  $c^2 = \frac{T}{\rho}$  (since we know that both  $T$  and  $\rho$

are unchanging constants)

Finally, this gives us the desired expression:

$$u_{tt} = c^2 u_{xx} - g$$

(k) Since  $c = \sqrt{\frac{T}{\rho}}$ ,  $c$  is of the units  $\sqrt{(kgm/s^2) \cdot (m/kg)} = \sqrt{(m^2/s^2)} = m/s$

The units of  $c$  are thus  $m/s$ . In plain english, this is the velocity of the string!