

Exercises of Chapter 2, 3 of Analysis and Data-Based Reconstruction of Complex Nonlinear Dynamical Systems

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1 Chapter 2,3 exercises

2 Chapter 2 exercises

2.1 Exercise 2.1

2.1.1 Part a

the moment generator function is given by $Z_x(u) = \langle e^{u^T x} \rangle$ Now, let's consider the taylor's expansion of $e^{u^T x}$

$$e^{u^T x} = \sum_{i=0}^{\infty} \frac{(u^T x)^i}{i!} = 1 + u^T x + \frac{(u^T x)^2}{2!} + \dots$$

Mean operator and Σ are interchangeable. also, every $u^T x$ is:

$$u^T x = u_1 x_1 + \dots + u_n x_n$$

now, we want to calculate $\frac{\partial^k}{\partial x_1^{k_1} \dots \partial x_n^{k_n}} z_x(u)$

$$\begin{aligned} \frac{\partial^k}{\partial x_1^{k_1} \dots \partial x_n^{k_n}} Z_x(u)|_{u=0} &= \frac{\partial^k}{\partial x_1^{k_1} \dots \partial x_n^{k_n}} \sum_{i=0}^{\infty} \langle \frac{(u^T x)^i}{i!} \rangle |_{u=0} \\ &= D^{(\vec{k})} \left(\sum_{i=0}^{k-1} \langle \frac{(u^T x)^i}{i!} \rangle + \frac{1}{k!} \langle (u^T x)^k \rangle + \sum_{i=k+1}^{\infty} \frac{1}{i!} \langle (u^T x)^i \rangle \right) |_{u=0} \\ &= D^{(\vec{k})}(x) \sum_{i=0}^{k-1} \langle \frac{(u^T x)^i}{i!} \rangle + \langle \frac{1}{k!} (x_1 u_1 + \dots + x_n u_n)^k \rangle + D^{(\vec{k})}(x) \sum_{i=k+1}^{\infty} \frac{1}{i!} \langle (u^T x)^i \rangle |_{u=0} \\ &= \frac{1}{k!} (k)(k-1) \dots (1) \langle x_1^{k_1} \dots x_n^{k_n} \rangle |_{u=0} = \langle x_1^{k_1} \dots x_n^{k_n} \rangle \end{aligned}$$

2.1.2 Part b

We have:

$$\begin{aligned} Z_x(t) &= E(e^{xt}) \\ &= \int e^{xt} \frac{1}{\sqrt{(2\pi\sigma^2)}} e^{-\frac{1}{2}(x-\mu)^2/\sigma^2} dx \end{aligned}$$

Define $z = \frac{x-\mu}{\sigma}$, which implies $x = z\sigma + \mu$.

We also have:

$$\begin{aligned} Z_z(t) &= E(e^{zt}) = \int e^{zt} \frac{1}{\sqrt{2\pi}} e^{-\frac{1}{2}z^2} dz \\ &= e^{\frac{1}{2}t^2}. \end{aligned}$$

Then, using the change-of-variable technique, we get

$$\begin{aligned} Z_x(t) &= e^{\mu t} \int e^{z\sigma t} \frac{1}{\sqrt{(2\pi\sigma^2)}} e^{-\frac{1}{2}z^2} \left| \frac{dx}{dz} \right| dz \\ &= e^{\mu t} \int e^{z\sigma t} \frac{1}{\sqrt{2\pi}} e^{-\frac{1}{2}z^2} dz \\ &= e^{\mu t} e^{\frac{1}{2}\sigma^2 t^2} \end{aligned}$$

2.2 Exercise 2.2

The conditional distribution of any subset vector x , given the complement vector y , is a multivariate normal distribution

$$x | y \sim \mathcal{N}(\mu_{x|y}, \Sigma_{x|y})$$

Lets assume that by construction, the joint distribution of x and y are:

$$x, y \sim \mathcal{N}(\mu, \Sigma).$$

Where:

$$\mu = \begin{pmatrix} \mu_x \\ \mu_y \end{pmatrix}$$

$$\Sigma = \begin{pmatrix} \sigma_{xx} & \sigma_{xy} \\ \sigma_{yx} & \sigma_{yy} \end{pmatrix} = \begin{pmatrix} \text{Var}(x) & \text{Cov}(x, y) \\ \text{Cov}(y, x) & \text{Var}(y) \end{pmatrix} = \begin{pmatrix} \sigma_x^2 & \rho\sigma_x\sigma_y \\ \rho\sigma_x\sigma_y & \sigma_y^2 \end{pmatrix}$$

The marginal distribution of y is

$$y \sim \mathcal{N}(\mu_y, \sigma_{yy}).$$

According to the law of conditional probability, it holds that

$$p(x | y) = \frac{p(x, y)}{p(y)}$$

By the definition, we have:

$$p(x | y) = \frac{\mathcal{N}(X; \mu, \Sigma)}{\mathcal{N}(y; \mu_y, \sigma_{yy})}.$$

Where:

$$X = \begin{pmatrix} x \\ y \end{pmatrix}$$

Using the probability density function of the p-variate normal distribution, this becomes:

$$\begin{aligned} p(x | y) &= \frac{1/\sqrt{(2\pi)^n |\Sigma|} \cdot \exp[-\frac{1}{2}(X - \mu)^T \Sigma^{-1}(X - \mu)]}{1/\sqrt{(2\pi)^{n_y} |\sigma_{yy}|} \cdot \exp[-\frac{1}{2}(y - \mu_y)^T \sigma_{yy}^{-1}(y - \mu_y)]} \\ &= \frac{1}{\sqrt{(2\pi)^{n-n_y}}} \cdot \sqrt{\frac{|\sigma_{yy}|}{|\Sigma|}} \cdot \exp[-\frac{1}{2}(X - \mu)^T \Sigma^{-1}(X - \mu) + \frac{1}{2}(y - \mu_y)^T \sigma_{yy}^{-1}(y - \mu_y)]. \end{aligned}$$

Writing the inverse of Σ as

$$\Sigma^{-1} = \begin{bmatrix} \sigma^{xx} & \sigma^{xy} \\ \sigma^{yx} & \sigma^{yy} \end{bmatrix} = \begin{bmatrix} (\sigma_{xx} - \sigma_{xy}\sigma_{yy}^{-1}\sigma_{yx})^{-1} & -(\sigma_{xx} - \sigma_{xy}\sigma_{yy}^{-1}\sigma_{yx})^{-1}\sigma_{xy}\sigma_{yy}^{-1} \\ -\sigma_{yy}^{-1}\sigma_{yx}(\sigma_{xx} - \sigma_{xy}\sigma_{yy}^{-1}\sigma_{yx})^{-1} & \sigma_{yy}^{-1} + \sigma_{yy}^{-1}\sigma_{yx}(\sigma_{xx} - \sigma_{xy}\sigma_{yy}^{-1}\sigma_{yx})^{-1}\sigma_{xy}\sigma_{yy}^{-1} \end{bmatrix}$$

In the matrix form, we get:

$$\begin{aligned} p(x | y) &= \frac{1}{\sqrt{(2\pi)^{n-n_y}}} \cdot \sqrt{\frac{|\sigma_{yy}|}{|\Sigma|}} \cdot \\ &\quad \exp[-\frac{1}{2}(\begin{bmatrix} x \\ y \end{bmatrix} - \begin{bmatrix} \mu_x \\ \mu_y \end{bmatrix})^T \begin{bmatrix} \sigma^{xx} & \sigma^{xy} \\ \sigma^{yx} & \sigma^{yy} \end{bmatrix} (\begin{bmatrix} x \\ y \end{bmatrix} - \begin{bmatrix} \mu_x \\ \mu_y \end{bmatrix})] \\ &\quad \cdot \exp[\frac{1}{2}(y - \mu_y)^T \sigma_{yy}^{-1}(y - \mu_y)]. \end{aligned}$$

By inserting the inverse and keeping in mind that both matrices are symmetric, we get:

$$\begin{aligned}
p(x | y) &= \frac{1}{\sqrt{(2\pi)^{n-n_y}}} \cdot \sqrt{\frac{|\sigma_{yy}|}{|\Sigma|}} \cdot \\
&\exp\left[-\frac{1}{2}((x - \mu_x)^T(\sigma_{xx} - \sigma_{xy}\sigma_{yy}^{-1}\sigma_{yx})^{-1}(x - \mu_x) - \dots \right. \\
&2(x - \mu_x)^T(\sigma_{xx} - \sigma_{xy}\sigma_{yy}^{-1}\sigma_{yx})^{-1}\sigma_{xy}\sigma_{yy}^{-1}(y - \mu_y) + \\
&\cdot (y - \mu_y)^T[\sigma_{yy}^{-1} + \sigma_{yy}^{-1}\sigma_{yx}(\sigma_{xx} - \sigma_{xy}\sigma_{yy}^{-1}\sigma_{yx})^{-1}\sigma_{xy}\sigma_{yy}^{-1}](y - \mu_y)) \\
&\left. + \frac{1}{2}((y - \mu_y)^T\sigma_{yy}^{-1}(y - \mu_y))\right].
\end{aligned}$$

Eliminating some terms, we have:

$$\begin{aligned}
p(x | y) &= \frac{1}{\sqrt{(2\pi)^{n-n_y}}} \cdot \sqrt{\frac{|\sigma_{yy}|}{|\Sigma|}} \cdot \\
&\exp\left[-\frac{1}{2}((x - \mu_x)^T(\sigma_{xx} - \sigma_{xy}\sigma_{yy}^{-1}\sigma_{yx})^{-1}(x - \mu_x) - \dots \right. \\
&2(x - \mu_x)^T(\sigma_{xx} - \sigma_{xy}\sigma_{yy}^{-1}\sigma_{yx})^{-1}\sigma_{xy}\sigma_{yy}^{-1}(y - \mu_y) + \\
&\left. \cdot (y - \mu_y)^T\sigma_{yy}^{-1}\sigma_{yx}(\sigma_{xx} - \sigma_{xy}\sigma_{yy}^{-1}\sigma_{yx})^{-1}\sigma_{xy}\sigma_{yy}^{-1}(y - \mu_y))\right].
\end{aligned}$$

Rearranging the terms, we have

$$\begin{aligned}
p(x | y) &= \frac{1}{\sqrt{(2\pi)^{n-n_y}}} \cdot \sqrt{\frac{|\sigma_{yy}|}{|\Sigma|}} \cdot \exp\left[-\frac{1}{2} \cdot \right. \\
&\cdot [(x - \mu_x) - \sigma_{xy}\sigma_{yy}^{-1}(y - \mu_y)]^T(\sigma_{xx} - \sigma_{xy}\sigma_{yy}^{-1}\sigma_{yx})^{-1}[(x - \mu_x) - \sigma_{xy}\sigma_{yy}^{-1}(y - \mu_y)] \\
&= \frac{1}{\sqrt{(2\pi)^{n-n_y}}} \cdot \sqrt{\frac{|\sigma_{yy}|}{|\Sigma|}} \cdot \exp\left[-\frac{1}{2} \cdot \right. \\
&\cdot [x - (\mu_x + \sigma_{xy}\sigma_{yy}^{-1}(y - \mu_y))]^T(\sigma_{xx} - \sigma_{xy}\sigma_{yy}^{-1}\sigma_{yx})^{-1}[x - (\mu_x + \sigma_{xy}\sigma_{yy}^{-1}(y - \mu_y))]]
\end{aligned}$$

where we have used the fact that $\sigma_{yx} = \sigma_{xy}^T$, because Σ is a covariance matrix.

The determinant of a block matrix is

$$\begin{vmatrix} A & B \\ C & D \end{vmatrix} = |D| \cdot |A - BD^{-1}C|,$$

such that we have for Σ that

$$\begin{vmatrix} \sigma_{xx} & \sigma_{xy} \\ \sigma_{yx} & \sigma_{yy} \end{vmatrix} = |\sigma_{yy}| \cdot |\sigma_{xx} - \sigma_{xy}\sigma_{yy}^{-1}\sigma_{yx}|.$$

With this and $n - n_y = n_x$, we finally arrive at

$$\begin{aligned}
p(x | y) &= \frac{1}{\sqrt{(2\pi)^{n_1} |\sigma_{xx} - \sigma_{xy}\sigma_{yy}^{-1}\sigma_{yx}|}} \cdot \exp\left[-\frac{1}{2} \cdot \right. \\
&\cdot [x - (\mu_x + \sigma_{xy}\sigma_{yy}^{-1}(y - \mu_y))]^T(\sigma_{xx} - \sigma_{xy}\sigma_{yy}^{-1}\sigma_{yx})^{-1}[x - (\mu_x + \sigma_{xy}\sigma_{yy}^{-1}(y - \mu_y))]]
\end{aligned}$$

which is the probability density function of a multivariate normal distribution

$$p(x | y) = \mathcal{N}(x; \mu_{x|y}, \Sigma_{x|y})$$

with the mean $\mu_{x|y}$ and covariance $\Sigma_{x|y}$ given by:

$$\begin{aligned}
\mu_{x|y} &= \mu_x + \sigma_{xy}\sigma_{yy}^{-1}(y - \mu_y) \\
\Sigma_{x|y} &= \sigma_{xx} - \sigma_{xy}\sigma_{yy}^{-1}\sigma_{yx}
\end{aligned}$$

By plugging this equation in to the probability, we get:

$$p(x, y) = \frac{\exp\{-\frac{1}{2(1-\rho^2)}[(\frac{x-\mu_x}{\sigma_x})^2 - 2\rho(\frac{x-\mu_x}{\sigma_x})(\frac{y-\mu_y}{\sigma_y}) + (\frac{y-\mu_y}{\sigma_y})^2]\}}{2\pi\sigma_x\sigma_y\sqrt{1-\rho^2}}$$

and

$$\begin{aligned}\langle y | x \rangle &= \mu_y + \rho \frac{\sigma_y}{\sigma_x} (x - \mu_x) \\ \sigma_{y|x}^2 &= \sigma_y^2 (1 - \rho^2)\end{aligned}$$

2.3 Exercise 2.3

2.3.1 Part a

Suppose $X \sim N_k(\mu, g)$ where g is full-rank variance covariance matrix, thus it is positive definite.

By definition of p-variate normal distribution, any linear combination of X has a univariate normal distribution. That is, $t^T X \sim N(t^T \mu, t^T \Sigma t)$ for any vector $t \in R^p$.

Moment generating function of some $Z \sim N(\mu, \sigma^2)$ is

$$M_Z(s) = E[e^{sZ}] = e^{\mu s + \sigma^2 s^2 / 2}, \quad s \in R$$

By cholesky decomposition, we can decompose $g = ll^T$ for some nonsingular matrix l since g is positive definite. Transform $X \mapsto Y$ such that $Y = l^{-1}(X - \mu)$, i.e. $X = \mu + lY$. Then it follows that $Y = (Y_1, \dots, Y_p)^T \sim N_p(0, I_p)$. This is the same as a change of variables such that X is from normal distribution with zero mean and a unit variance.

In other words Y_1, \dots, Y_p are independent standard normal.

Therefore using MGF of standard normal distribution,

$$\begin{aligned}M_X(t) &= E[e^{t^T X}] \\ &= E[e^{t^T (\mu + lY)}] \\ &= e^{t^T \mu} E[e^{\ell^T Y}] \quad , \ell^T = t^T l \\ &= e^{t^T \mu} E[e^{\sum_{i=1}^p \ell_i Y_i}] \quad , \ell = (\ell_1, \dots, \ell_p) \\ &= e^{t^T \mu} \prod_{i=1}^p E[e^{\ell_i Y_i}] \\ &= e^{t^T \mu} \prod_{i=1}^p e^{\ell_i^2 / 2} \\ &= \exp(\mu^T t + \frac{1}{2} \ell^T \ell) \\ &= \exp(\mu^T t + \frac{1}{2} t^T g t)\end{aligned}$$

2.3.2 Part b

In exercise 2.1, we proved that:

$$\cdot \frac{\partial^k}{\partial u_1^{k_1} \dots \partial u_n^{k_n}} Z_{\mathbf{x}}(\mathbf{u})|_{\mathbf{u}=\mathbf{0}} = \langle x_1^{k_1} \dots x_n^{k_n} \rangle$$

Now, consider the case of 4 variates of n-variate Gaussian MGF. first of all, we can zero-mean all the X by a n-dimensional constant shift vector. Then:

$$\cdot \frac{\partial^4}{\partial u_i \partial u_j \partial u_k \partial u_l} Z_{\mathbf{x}}(\mathbf{u})|_{\mathbf{u}=\mathbf{0}} = \langle x_i x_j x_k x_l \rangle$$

Also, we calculated MGF as below:

$$\exp(\mu^T t + \frac{1}{2} t^T g t)$$

By differentiating the same as Exercise 2.1, we get to

$$\langle x_1^{k_1} \cdots x_n^{k_n} \rangle = \frac{1}{4!} \left(\prod_{i',j',k',l' \in i,j,k,l} g_{i'j'} g_{k'l'} \right) = g_{ij} g_{kl} + g_{ik} g_{jl} + g_{il} g_{jk}$$

The last equality is true because of symmetry of Matrix g .

2.4 Exercise 2.4

2.4.1 Brownian motion

The Chapman-Kolmogorov equation for Brownian motion with setting initial condition at 0,0

$$p(x, t | 0, 0) = \int_{-\infty}^{\infty} p(x, t | x', t') p(x', t' | 0, 0) dx'$$

where $p(x, t | x', t')$ is the transitional probability of Brownian motion. Since

$$p(x, t | x', t') = P(X(t) = x | X(t') = x') = \frac{1}{\sqrt{2\pi(t-t')}} e^{-\frac{(x-x')^2}{2(t-t')}} e^{-\frac{x'^2}{2t'}}$$

In other words, We need to show

$$\frac{1}{\sqrt{2\pi t}} e^{-\frac{x^2}{2t}} = \int_{-\infty}^{\infty} \frac{1}{\sqrt{2\pi(t-t')}} e^{-\frac{(x-x')^2}{2(t-t')}} \frac{1}{\sqrt{2\pi t'}} e^{-\frac{x'^2}{2t'}} dx'$$

By simplifying the integral:

$$\begin{aligned} \int_{-\infty}^{\infty} \frac{1}{\sqrt{2\pi(t-t')}} e^{-\frac{(x-x')^2}{2(t-t')}} \frac{1}{\sqrt{2\pi t'}} e^{-\frac{x'^2}{2t'}} dx' &= \frac{1}{2\pi \sqrt{t'(t-t')}} \int_{-\infty}^{\infty} e^{-\frac{(x-x')^2}{2(t-t')}} e^{-\frac{x'^2}{2t'}} dx' \\ &= \frac{1}{2\pi \sqrt{t'(t-t')}} \int_{-\infty}^{\infty} \exp(-\frac{x^2}{2t}) \exp(\frac{t}{2t'(t-t)} (x' - \frac{xt'}{t})^2) dx' \\ &= \frac{1}{\sqrt{2\pi t}} \exp(-\frac{x^2}{2t}) \end{aligned}$$

Which shows that Brownian motion satisfies Chapman-Kolmogorov equation.

2.4.2 Cauchy Process

We need to show that:

$$p(x, t | 0, 0) = \int_{-\infty}^{\infty} p(x, t | x', t') p(x', t' | 0, 0) dx'$$

By putting the Cauchy process into the transitional probability, we are getting:

$$p(x, t | x', t') = P(X(t) = x | X(t') = x') = \frac{t' \cdot (t - t')}{\pi^2 \cdot ((x - x')^2 + (t - t')^2)(x'^2 + t'^2)}$$

Let's put them into the integral:

$$p(x, t | 0, 0) = \int_{-\infty}^{\infty} \frac{t' \cdot (t - t')}{\pi^2 \cdot ((x - x')^2 + (t - t')^2)(x'^2 + t'^2)} dx'$$

By using Mathematica to solve this integral, the result is:

$$\begin{aligned} &= \frac{1}{[\pi^2 \cdot (x^2 + (t - 2t')^2)(x^2 + t^2)|t - t'|]} ((t - t') \\ &(|t - t'| (t' x \ln((x' - x)^2 + (t - t')^2) - t' x \ln(x'^2 + t'^2) + (-x^2 - t \cdot (t - 2t')) \arctan(\frac{x'}{t'})) \\ &t' \cdot (x^2 - t \cdot (t - 2t')) \arctan(\frac{x' - x}{|t - t'|})) \end{aligned}$$

By putting the upper and lower limit of integral, this can be rewrote in this form:

$$\frac{(t-t')(2tx^2+2t^3-8t't^2+8t'^2t)}{\pi \cdot (2t-2t')(x^4+(2t^2-4t't+4t'^2)x^2+t^4-4t't^3+4t'^2t^2)}$$

And by simplifying one gets to:

$$\frac{t}{\pi \cdot (x^2+t^2)}$$

Which is the desired solution to Chapman-Kolmogorov equation.

3 Chapter 3 exercises

3.1 Exercise 3.1

Let's start with Kramers-Moyal expansion.

$$\begin{aligned} & \frac{\partial P(x, t)}{\partial t} \sum_{n=1}^{\infty} \left(-\frac{\partial}{\partial x}\right)^n [D^{(n)}(x, t) P(x, t)] \\ & \int x^n \frac{\partial P'(x, t)}{\partial t} dx = \frac{\partial}{\partial t} \int x^n P(x, t) dx = \frac{\partial}{\partial t} \langle x^n \rangle \\ & \int x^n \sum_{k=1}^{\infty} \left(-\frac{\partial}{\partial x}\right)^k [D^{(k)}(x, t) P(x, t)] dx \\ & = \sum_{k=1}^{\infty} \int x^n \left(-\frac{\partial}{\partial x}\right)^k [D^{(k)}(x, t) P(x, t)] dx \\ & = \sum_{k=1}^{\infty} \int \left(-\frac{\partial}{\partial x}\right) [x^n \left(-\frac{\partial}{\partial x}\right)^{k-1} [D^{(k)}(x, t) P(x, t)]] dx \\ & + \sum_{k=1}^{\infty} \int n x^{n-1} \left(-\frac{\partial}{\partial x}\right)^{k-1} [D^{(k)}(x, t) P(x, t)] dx \end{aligned}$$

Now, with recursion:

$$\begin{aligned} & = - \int \frac{\partial}{\partial x} \left\{ \sum_{k=1}^{\infty} [x^n \left(-\frac{\partial}{\partial x}\right)^{k-1} + n x^{n-1} \left(-\frac{\partial}{\partial x}\right)^{k-2}] [D^{(k)}(x, t) P(x, t)] \right\} dx \\ & + \sum_{k=1}^{\infty} \int n(n-1) x^{n-2} \left(-\frac{\partial}{\partial x}\right)^{k-2} [D^{(k)}(x, t) P(x, t)] dx \end{aligned}$$

Now, this can be rewritten in the form:

$$\begin{aligned} & - \int \frac{\partial}{\partial x} \left\{ \sum_{k=1}^{\infty} [x^n \left(-\frac{\partial}{\partial x}\right)^{k-1} + n x^{n-1} \left(-\frac{\partial}{\partial x}\right)^{k-2} + \dots \right. \\ & \quad \left. \dots + n(n-1) \dots (n-k+2) x^{n-k+1}] [D^{(k)}(x, t) p(x, t)] \right\} dx \\ & + \sum_{k=1}^{\infty} n(n-1) \dots (n-k+1) \int x^{n-k} D^{(k)}(x, t) P(x, t) dx \end{aligned}$$

The first term, is always zero. And the second term, is equivalent to:

$$\sum_{k=1}^n \frac{n!}{(n-k)!} \langle x^{n-k} D^{(k)}(x, t) \rangle$$

Which proves:

$$\frac{\partial}{\partial t} \langle x^n \rangle = \sum_{k=1}^n \frac{n!}{(n-k)!} \langle x^{n-k} D^{(k)}(x, t) \rangle$$

3.2 Exercise 3.2

3.2.1 Part a

We are going to demonstrate one step, then by induction, it will be true for all steps.

$$P(x, t) = \int P(x, t; x', t') dx'$$

$$P(x, t; x', t') = P(x, t | x', t') P(x', t')$$

By chapman-kolmogorov equation we can write:

$$P(x, t | x', t') = \int P(x, t | x'', t'') P(x'', t'' | x', t') dx'' \quad (t > t'' > t')$$

$$\rightarrow P(x, t) = \int P(x, t | x', t') P(x', t') dx'$$

$$= \int dx' \int dx'' P(x, t | x'', t'') P(x'', t'' | x', t') \quad (t > t'' > t')$$

now by defining the $\tau = (t - t_0)/N$ and $t_n = t_0 + n\tau$, we can do this for any given N .

$$p(x, t) = \lim_{N \rightarrow \infty} \int dx_{N-1} \cdots \int dx_0 p(x, t | x_{N-1}, t_{N-1}) \cdots p(x_1, t_1 | x_0, t_0) p(x_0, t_0).$$

3.2.2 Part b

Let's begin by The Fokker-Plank Equation and calculate for small τ :

$$\frac{\partial}{\partial t} P(x, t | x', t') = [-\frac{\partial}{\partial x} D^{(1)}(x, t) + \frac{\partial^2}{\partial x^2} D^{(2)}(x, t)] P(x, t | x', t')$$

Use the definition of \mathcal{L}_{FP} and write for small τ :

$$P(x, t + \tau | x', t) = e^{\tau \mathcal{L}_{FP}} P(x, t | x', t) = e^{\tau \mathcal{L}_{FP}} \delta(x - x') \delta(x - x') = \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{iu(x-x')} du$$

Now using Kramers-Moyal Coefficients:

$$P(x, t + \tau | x', t) = \{1 - \tau \frac{\partial}{\partial x} D^{(1)}(x, t) + \tau \frac{\partial^2}{\partial x^2} D^{(2)}(x, t) + O(\tau^2)\} \delta(x - x')$$

In the integral form:

$$\begin{aligned} P(x, t + \tau | x', t) &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{iu(x-x')} [1 - iu\tau D^{(1)}(x', t) - u^2\tau D^{(2)}(x', t) + O(\tau^2)] du \\ &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{iu(x-x') - iu\tau D^{(1)}(x', t) - u^2\tau D^{(2)}(x', t) + O(\tau^2)} du \\ &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} \exp\{\tau D^{(2)}(x', t) [-(u - \frac{i(x-x' - \tau D^{(1)}(x', t))}{2\tau D^{(2)}(x', t)})^2 - (\frac{x-x' - \tau D^{(1)}(x', t)}{2\tau D^{(2)}(x', t)})^2]\} du \end{aligned}$$

Then we get to:

$$p(x, t) = \lim_{N \rightarrow \infty} \int \underbrace{\cdots}_{N \text{ times}} \int \{\prod_{i=0}^{N-1} (4\pi D^{(2)}(x_i, t_i)\tau)^{-1/2} dx_i\} \exp(-\sum_{i=0}^{N-1} \frac{[x_{i+1} - x_i - D^{(1)}(x_i, t_i)]^2}{4D^{(2)}(x_i, t_i)\tau}) p(x_0, t_0)$$

Now, for the derivative:

$$\begin{aligned} \lim_{N \rightarrow \infty} \sum_{i=1}^{N-1} \frac{(x_{i+1} - x_i - \tau D^{(1)}(x_i, t_i))}{4\tau D^{(2)}(x_i, t_i)} &= \lim_{N \rightarrow \infty} \sum_{i=1}^{N-1} \frac{1}{4D^{(2)}(x_i, t_i)} [\frac{x_{i+1} - x_i}{\tau} - D^{(1)}(x_i, t_i)] \\ &= \int_{t_0}^t \frac{\dot{x}(t') - D^{(1)}(x(t'), t')}{4D^{(2)}(x(t'), t')} dt' \\ \lim_{N \rightarrow \infty} \sum_{i=0}^{N-1} \frac{[x_{i+1} - x_i - D^{(1)}(x_i, t_i)\tau]^2}{4D^{(2)}(x_i, t_i)\tau} &= \int_{t_0}^t \frac{[\dot{x}(t') - D^{(1)}(x(t'), t')]^2}{4D^{(2)}(x(t'), t')} dt' \end{aligned}$$

3.2.3 Part c

Both exponential and KM-coefficients part are positive, The integral must remain positive for all t.

3.3 Exercise 3.3

3.3.1 Part a

From CK equation:

$$P(x, t | x', t') = \int P(x, t | x'', t' + \tau) P(x'', t' + \tau | x', t') dx'' = \int dx'' P(x, t | x'', t' + \tau) \int dy \delta(y - x'') P(y, t' + \tau | x', t')$$

Now, by using $P(x, t + \tau | x', t) = \int \delta(y - x) P(y, t + \tau | x', t) dy$ and Taylor expansion of delta distribution:

$$\delta(y - x'') = \delta(x' - x'' + y - x') = \sum_{n=0}^{\infty} \frac{(y - x')^n}{n!} \left(\frac{\partial}{\partial x'} \right)^n \delta(x' - x'')$$

We get to:

$$\begin{aligned} P(x, t | x', t') &= \int dx'' P(x, t | x'', t' + \tau) \int dy \delta(y - x'') P(y, t' + \tau | x', t') \\ &= \int dx'' P(x, t | x'', t' + \tau) \left[\sum_{n=1}^{\infty} \frac{1}{n!} \int dy (y - x')^n P(y, t' + \tau | x', t') \left(\frac{\partial}{\partial x'} \right)^n \delta(x' - x'') \right] \\ &= \int dx'' P(x, t | x'', t' + \tau) \left[\int dy P(y, t' + \tau | x', t') + \sum_{n=1}^{\infty} \frac{1}{n!} \left(\int dy (y - x')^n P(y, t' + \tau | x', t') \right) \left(\frac{\partial}{\partial x'} \right)^n \delta(x' - x'') \right] \end{aligned}$$

Also, $K^{(n)}(x', t', \tau) = \int dy (y - x')^n P(y, t' + \tau | x', t')$. Then we can rewrite:

$$\begin{aligned} &= \left[1 + \sum_{n=1}^{\infty} \frac{K^{(n)}(x', t', \tau)}{n!} \left(\frac{\partial}{\partial x'} \right)^n \right] \int dx'' P(x, t | x'', t' + \tau) \delta(x' - x'') \\ &= P(x, t | x', t' + \tau) + \left(\sum_{n=1}^{\infty} \frac{K^{(n)}(x', t', \tau)}{n!} \left(\frac{\partial}{\partial x'} \right)^n \right) P(x, t | x', t' + \tau) \end{aligned}$$

For difference of probability:

$$P(x, t | x', t') - P(x, t | x', t' + \tau) = \left(\sum_{n=1}^{\infty} \frac{K^{(n)}(x', t', \tau)}{n!} \left(\frac{\partial}{\partial x'} \right)^n \right) P(x, t | x', t' + \tau)$$

As for the limit:

$$\lim_{\tau \rightarrow 0} \frac{P(x, t | x', t') - P(x, t | x', t' + \tau)}{\tau} = - \frac{\partial}{\partial t} P(x, t | x', t') = \sum_{n=1}^{\infty} \frac{1}{n!} \lim_{\tau \rightarrow 0} \frac{K^{(n)}(x', t', \tau)}{\tau} \left(\frac{\partial}{\partial x'} \right)^n P(x, t | x', t' + \tau)$$

By putting $D^{(n)}(x', t') = \frac{1}{n!} \lim_{\tau \rightarrow 0} \frac{K^{(n)}(x', t', \tau)}{\tau}$ We get to the desired relation:

$$\frac{\partial}{\partial t} P(x, t | x', t') = - \sum_{n=1}^{\infty} D^{(n)}(x', t') \left(\frac{\partial}{\partial x'} \right)^n P(x, t | x', t')$$

3.3.2 Part b

Let's construct The adjoint operator, directly.

$$\begin{aligned}
\int h(x') \mathcal{L}_{KM} f(x') dx' &= \sum_{n=1}^{\infty} \int h(x') \left(-\frac{\partial}{\partial x}\right)^n D^{(n)}(x', t') f(x') dx' \\
&= \sum_{n=1}^{\infty} \left[-h(x') \left(-\frac{\partial}{\partial x}\right)^{n-1} D^{(n)}(x', t') f(x') + \int \left(\frac{\partial}{\partial x'} h(x')\right) \left(-\frac{\partial}{\partial x'}\right) D^{(n)}(x', t') f(x') dx' \right] \\
&= \sum_{n=1}^{\infty} \int D^{(n)}(x', t') \left(\left(\frac{\partial}{\partial x'}\right)^n h(x')\right) f(x') dx' = \\
&\int f(x') \left[\sum_{n=1}^{\infty} \left(\frac{\partial}{\partial x'}\right)^n D^{(n)}(x', t') \right] h(x') dx' = \int f(x') \mathcal{L}_{KM}^+ h(x') dx'
\end{aligned}$$

By defining $\mathcal{L}_{KM}^+ = \sum_{n=1}^{\infty} \left(\frac{\partial}{\partial x'}\right)^n D^{(n)}(x', t')$ we get to the desired equation.

3.4 Exercise 3.4

3.4.1 Part a

Let's start by Master equation:

$$\frac{\partial}{\partial t} P(n, t) = \sum_{n'} (\omega_{n, n'} P_{n'}(t) - \omega_{n', n} P_n(t))$$

With $\omega_{n, n'} = \frac{1}{2\Delta t}$, we get to:

$$\left(\frac{1}{\Delta t}\right)(P(n, N) - P(n, N-1)) = \omega_{n, n+1} P(n+1, N-1) + \omega_{n, n-1} P(n-1, N-1) - \omega_{n+1, n} P(n, N-1) - \omega_{n-1, n} P(n, N-1)$$

Transitions of random walk are from neighbours, so that:

$$P(i, N) - P(i, N-1) = \frac{1}{2}(P(i+1, N-1) + P(i-1, N-1)) - P(i, N-1)$$

And we get to the desired answer:

$$P(i, N) = \frac{1}{2}[P(i+1, N-1) + P(i-1, N-1)]$$

As the question asked for.

3.4.2 Part b

Above equation in terms of continuity is written as:

$$P(x, t) = \frac{1}{2}(P(x+a, t-\tau) + P(x-a, t-\tau))$$

Now, let's calculate the derivatives:

$$\tau \lim_{\tau \rightarrow 0} \frac{P(x, t) - P(x, t-\tau)}{\tau} = \frac{a^2}{2} \lim_{\tau \rightarrow 0} \frac{1}{a^2} [P(x+a, t) + P(x-a, t) - 2P(x, t)]$$

By making parameter $D = \frac{a^2}{2\tau}$ we get to the partial deviates with respect to units of space and time, we get to the equation.

$$\frac{\partial}{\partial t} P = D \frac{\partial^2 P}{\partial x^2}$$

3.4.3 Part c

We got to the diffusion equation:

$$\frac{\partial}{\partial t}P(x, t) = D \frac{\partial^2}{\partial x^2}P(x, t)$$

The diffusion equation is a partial differential equation (Also known as heat equation). The unknown quantity is a function $P(x, t)$. To complete the problem statement we need to specify an initial condition (at $t = 0$) and boundary conditions. boundary conditions are at infinity, so we take

$$P(x, t) \rightarrow 0, x \rightarrow \pm\infty.$$

We take a delta function initial condition:

$$P(x, 0) = \delta(x).$$

The equation can be solved by using the Fourier transform:

$$P(x, t) = \int_{-\infty}^{\infty} \frac{1}{2\pi} e^{ikx} P_k(t) dk$$

The inverse transform is

$$P_k(t) = \int_{-\infty}^{\infty} e^{-ikx} P(x, t) dx$$

So the transform of the initial condition is

$$P_k(0) = 1$$

Substituting $P(x, t)$ in the diffusion equation gives

$$\int_{-\infty}^{\infty} \frac{1}{2\pi} e^{ikx} (\dot{P}_k(t) + Dk^2 P_k(t)) dk = 0$$

This simplifies to

$$\dot{P}_k(t) + Dk^2 P_k(t) = 0$$

With the solution

$$P_k(t) = P_k(0) e^{-Dk^2 t} = e^{-Dk^2 t}$$

Putting it all together

$$P(x, t) = \int_{-\infty}^{\infty} \frac{1}{2\pi} e^{ikx} e^{-Dk^2 t} dk$$

And all that's left is to do the k integral. Note that the k integral is a Gaussian:

$$\int_{-\infty}^{\infty} e^{-y^2} dy = \sqrt{\pi}.$$

We should get

$$P(x, t) = \frac{1}{\sqrt{4\pi Dt}} e^{-\frac{x^2}{4Dt}}$$

3.4.4 Part d

We should do the same calculations as above.

$$P(x', t + \tau | x, t) = \delta(x' - x)$$

Transformation of initial condition is still same. now we should substitute it in the diffusion equation:

$$\int_{-\infty}^{\infty} \frac{1}{2\pi} e^{ik(x'-x)} (\dot{P}_k(\tau) + Dk^2 P_k(\tau)) dk = 0$$

Which simplifies to:

$$\dot{P}_k(\tau) + Dk^2 P_k(\tau) = 0$$

This has the same solution as above, but keep in mind that this time, $t \rightarrow \tau$ and k is adjoint Fourier variable of $x' - x$ instead of x .

So we get to the answer as:

$$P(x', t + \tau | x, t) = \frac{1}{\sqrt{4\pi D\tau}} e^{-\frac{(x'-x)^2}{4D\tau}}$$

3.4.5 Part e

Starting by:

$$P(x, t) = \frac{1}{\sqrt{4\pi Dt}} e^{-\frac{x^2}{4Dt}}$$

Notice that this is similar to a Gaussian distribution. So let's prove a more general case of this distribution that initial condition is not as $x_0 = 0$, instead, it is starting at μ which we know that mean of a random walk, is the position of start. so this is a solution for the conditional part aswell.

Change of variables $\sigma = \sqrt{2D\tau}$

$$\begin{aligned} \langle x^2 \rangle &= \frac{1}{\sigma\sqrt{2\pi}} \int_{-\infty}^{\infty} x^2 \exp\left(-\frac{(x-\mu)^2}{2\sigma^2}\right) dx - \mu^2 \\ &= \frac{\sqrt{2}\sigma}{\sigma\sqrt{2\pi}} \int_{-\infty}^{\infty} (\sqrt{2}\sigma t + \mu)^2 \exp(-t^2) dt - \mu^2 \\ &= \frac{1}{\sqrt{\pi}} (2\sigma^2 \int_{-\infty}^{\infty} t^2 \exp(-t^2) dt + 2\sqrt{2}\sigma\mu \int_{-\infty}^{\infty} t \exp(-t^2) dt + \mu^2 \int_{-\infty}^{\infty} \exp(-t^2) dt) - \mu^2 \\ &= \frac{1}{\sqrt{\pi}} (2\sigma^2 \int_{-\infty}^{\infty} t^2 \exp(-t^2) dt + 2\sqrt{2}\sigma\mu [-\frac{1}{2} \exp(-t^2)]_{-\infty}^{\infty} + \mu^2 \sqrt{\pi}) - \mu^2 \\ &= \frac{1}{\sqrt{\pi}} (2\sigma^2 \int_{-\infty}^{\infty} t^2 \exp(-t^2) dt + 2\sqrt{2}\sigma\mu \cdot 0) + \mu^2 - \mu^2 \\ &= \frac{2\sigma^2}{\sqrt{\pi}} \int_{-\infty}^{\infty} t^2 \exp(-t^2) dt \\ &= \frac{2\sigma^2}{\sqrt{\pi}} \left(\left[-\frac{t}{2} \exp(-t^2) \right]_{-\infty}^{\infty} + \frac{1}{2} \int_{-\infty}^{\infty} \exp(-t^2) dt \right) \\ &= \frac{2\sigma^2}{\sqrt{\pi}} \cdot \frac{1}{2} \int_{-\infty}^{\infty} \exp(-t^2) dt \\ &= \frac{2\sigma^2 \sqrt{\pi}}{2\sqrt{\pi}} \\ &= \sigma^2 = 2D\tau \end{aligned}$$

3.5 Exercise 3.5

3.5.1 Part a

The master equation:

$$\frac{\partial}{\partial t} P_n(t) = \sum_{n'} (\omega_{n,n'} P_{n'}(t) - \omega_{n',n} P_n(t))$$

Question's hypothesis:

$$\omega_{n,n'} = \lambda \delta_{n,n'+1}$$

Which λ is > 0 .

Putting $\omega_{n,n'}$ in master equation:

$$\frac{\partial}{\partial t} P_n(t) = \sum_{n'} (\lambda \delta_{n,n'+1} P_{n'}(t) - \lambda \delta_{n'+1,n} P_n(t))$$

We get to:

$$\dot{P}_n(t) = \lambda P_{n-1}(t) - \lambda P_n(t).$$

As the question needed.

3.5.2 Part b

We put the given equation in question in the equation we got in last part:

$$\dot{P}_n = \frac{d}{dt} \left[\frac{(\lambda t)^n}{n!} \exp(-\lambda t) \right] = \frac{n\lambda(\lambda t)^{n-1}}{n!} \exp(-\lambda t) - \frac{\lambda(\lambda t)^n}{n!} \exp(-\lambda t) = \lambda P_{n-1} - \lambda P_n$$

3.5.3 Part c

for $\langle n \rangle$ we can write:

$$\begin{aligned}\langle n \rangle &= \sum n \frac{(\lambda t)^n}{n!} e^{-\lambda t} \\ &= \lambda \sum \frac{(\lambda t)^{n-1}}{(n-1)!} e^{-\lambda t} \\ &= \lambda e^{-\lambda t} \sum \frac{(\lambda t)^{n-1}}{(n-1)!} \\ &= \lambda e^{-\lambda t} \sum \frac{(\lambda t)^m}{m!} \\ &= \lambda e^{-\lambda t} (\lambda t)^n e^{-\lambda t} \\ &= \lambda e^{-\lambda t} e^{\lambda t} \\ &= \lambda\end{aligned}$$

And also for $\langle n^2 \rangle$ we can write:

$$\begin{aligned}\langle n^2 \rangle &= \sum n^2 \frac{(\lambda t)^n}{n!} e^{-\lambda t} \\ &= \sum \frac{n}{(n-1)!} (\lambda t)^n e^{-\lambda t} \\ &= \sum \frac{n-1+1}{n-1} \frac{1}{(n-2)!} (\lambda t)^n e^{-\lambda t} \\ &= \lambda^2 \sum \frac{(\lambda t)^{n-2}}{(n-2)!} e^{-\lambda t} + \lambda \sum \frac{(\lambda t)^{n-1}}{(n-1)!} e^{-\lambda t} \\ &= \lambda^2 + \lambda\end{aligned}$$

So variance is the same as mean $\sigma^2 = \langle n^2 \rangle - (\langle n \rangle)^2 = \lambda$.