

# Exploring Periodic Boundary Conditions: A Study on 2D and 3D Spheres and Continuum Limit

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## Abstract

We analyze systems with periodic boundary conditions in two and three dimensions. Beginning with an  $N$ -particle system confined to a circle, we derive the governing ordinary differential equation and solve it both analytically using circulant matrices and numerically. We then take the continuum limit as  $N$  goes to infinity to obtain a wave equation on the circle with periodic boundary conditions. This partial differential equation is solved exactly, allowing us to construct visualizations of the propagating waves. Next we study an  $N$ -particle system confined to the surface of a regular polyhedron, deriving the discrete equations and solving numerically using symmetry arguments. We then discuss statistical mechanics approaches for this system. Finally, we analyze longitudinal elastic waves propagating on a sphere, deriving the continuum wave equation in spherical coordinates and discussing analytical and numerical solution methods based on spectral decomposition. Throughout we emphasize the mathematical physics, solving equations analytically where possible and numerically otherwise. Results for all systems are illustrated through figures.

**Keywords:** Particle Systems, Mathematical Physics, Fourier Analysis, Differential Equations, Continuum Limit, Periodic Symmetry, Platonic Solids

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## 1. Introduction

Periodic boundary conditions arise naturally in the study of waves and particles constrained to closed topologies. The mathematics of such systems, though conceptually simple, reveals rich physical behaviors. In one dimension, systems with periodic boundaries can be solved exactly. As the dimension increases, new challenges emerge. On a circle, the topology is uniform and easily parametrized, yielding to straightforward analysis via Fourier decomposition. On a sphere, the additional dimension introduces variability in the surface curvature, requiring more advanced methods like spherical harmonics. In three dimensions, modeling particles or waves constrained to regular convex polyhedron brings deeper obstacles, as these Platonic solids resist exact geometric parameterization.

Approximation methods must therefore complement exact mathematical techniques in exploring these systems. The interplay of analytical, numerical, and geometric ideas is what makes periodic systems fascinating to study. When symmetry allows clean solutions, the mathematics elegantly explains the physics. And when intricacies of higher dimensions obstruct analysis, numerical methods and geometric intuition guide the way.

In this Article, we embark on a mathematical tour of periodic boundaries in multiple dimensions. We begin with approachable problems, deriving and solving ordinary differential

equations through matrix techniques. Moving to continuous systems, we encounter tractable partial differential equations which yield to Fourier analysis. Finally, we contend with intricate geometries, wrestling with equations which only numerical computation can conquer. What ties these investigations together is the quest to understand systems constrained to live within themselves, endowed with a symmetry as beautiful as it is challenging.

## 2. 2D Spheres Revisited

### 2.1. Analytical Description

As established in foundational texts on wave phenomena (see [Georgi \(2015\)](#)), a prototypical model for cyclic wave propagation considers a discrete set of point particles confined to a circular domain. This framework relies on two key approximations: first, that interactions are strictly local, with each particle only directly influencing its immediate neighbors. Second, that the inter-particle forces respond linearly to changes in separation relative to an equilibrium length. While undeniably reductionist, these assumptions prove sufficient to capture the salient physics of transverse wave transmission through the medium.

Specifically, we examine a system of  $N$  identical particles of mass  $m$  regularly spaced around a ring of radius  $R$ . Identical harmonic springs, characterized by force constant  $k$ , link adjacent particles along the circumference. This geometry exhibits a continuous rotational symmetry - rotations and reflections by  $2\pi/N$  map the configuration onto itself. Our theoretical description should therefore not differentiate between the  $2N$  symmetry-equivalent orientations. The sole measurable degree of freedom lies in the particles' relative displacements, as

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absolute positions along the circle can not be experimentally determined without first arbitrarily fixing a reference point.

The simplicity of this model, alongside its inherent periodicity, motivate its study before tackling more intricate geometries in higher dimensions. Even within the confines of localized, linear interactions, analytical and computational techniques reveal rich wave behaviors arising from the cyclic topology. The mathematics provides profound and elegant insight into the underlying physics.

As an elementary example, Figure 1 illustrates the case for  $N = 4$  particles. The fourfold rotational symmetry is evident, as is the periodic nature of the circular boundary. With the model established, we now embark on an in-depth analysis.

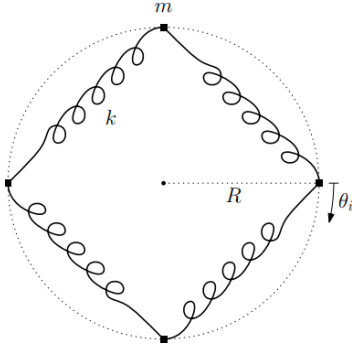


Figure 1: 4 identically distributed point masses on a circle

The configuration of the system can be fully specified by the central angle displacement  $\theta_i$  of each particle  $i$  relative to its initial angular position. As per the rotational symmetry, the labeling of individual particles carries no physical significance. We therefore characterize the state through the  $N$ -dimensional vector  $\boldsymbol{\theta} = (\theta_1, \theta_2, \dots, \theta_N)$  encompassing each particle's angular shift. For small oscillations, we expect the time evolution of  $\boldsymbol{\theta}$  to follow a linear equation of the form:

$$M\ddot{\boldsymbol{\theta}} + K\boldsymbol{\theta} = 0 \Rightarrow \ddot{\boldsymbol{\theta}} + W\boldsymbol{\theta} = 0 \quad (1)$$

The matrix  $W$  belongs to the special class of circulant matrices which respect circular permutation symmetry. To derive its structure, consider the  $i$ th particle's interactions: since it only couples to its nearest neighbors, its motion is governed by:

$$\ddot{\theta}_i = a\theta_{i-1} + b\theta_i + c\theta_{i+1}$$

By symmetry, the coefficient  $a$  equals  $c$ . For a system of  $N$  particles, this yields an interaction matrix  $W$  of the form:

$$W = \begin{pmatrix} A & B & 0 & \dots & B \\ B & A & B & \dots & 0 \\ 0 & B & A & \dots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ B & 0 & 0 & \dots & A \end{pmatrix}_{N \times N} \quad (2)$$

While the precise values of  $A$  and  $B$  do not qualitatively impact the dynamics, specific models may impose further constraints such as  $A = -2B$  (see Georgi (2015)). Analytical solu-

tions rely on the circulant structure, as discussed in Gray et al. (2006).

As discussed, the normal modes of this system should respect the circular symmetry. This mandates mode shapes of the form:

$$\boldsymbol{\beta}^{(k)} = (1, \omega^k, \omega^{2k}, \dots, \omega^{(N-1)k}) \quad \omega = e^{2\pi i/N}$$

where  $\omega$  denotes the complex  $N$ th root of unity. The angular frequency  $\omega_k$  of the  $k$ th mode takes the form:

$$\omega_k^2 = A + 2B \cos\left(\frac{2k\pi}{N}\right) \quad k = 0, 1, \dots, N-1$$

This reveals intrinsic degeneracies arising from the system's symmetry. For even  $N$ , the  $k = 0$  and  $k = N/2$  modes have real-valued eigenvectors, while the remaining  $N - 2$  modes come in complex-conjugate pairs which can be combined into degenerate real modes. The connection to Fourier analysis will become transparent in examining the continuum limit. The normal modes indeed correspond to a Discrete Fourier Transform, as elaborated in Gray et al. (2006). The circulant structure enables analytical decomposition into these symmetry-dictated eigenmodes.

In the Case where  $A = -2B$  the relation simplifies to the simpler case below:

$$\omega_k^2 = \omega_0^2 \sin^2\left(\frac{k\pi}{N}\right)$$

## 2.2. Continuum Limit

In taking the continuum limit, the number of particles  $N$  approaches infinity, transforming the discrete mechanical system into a continuous cyclic domain. Our coordinate  $\varphi$  becomes a continuous angular variable, with the wave displacement  $u(\varphi, t)$  encapsulating the discrete particle shifts  $\theta_j(t)$  through:

$$u(\varphi_j, t) = \theta_j(t) \quad \text{where:} \quad \varphi_j = \frac{2\pi j}{N} \quad j = 0, \dots, N-1$$

As  $N$  tends to infinity,  $\varphi_j$  transitions to a smooth continuous coordinate  $\varphi$ . The partial differential wave equation governing disturbances naturally emerges from the initial discrete equations of motion (Equation 1):

$$\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial \varphi^2}, \quad u(\varphi, t) = u(\varphi + 2\pi, t) \quad (3)$$

The wave velocity  $c$  descending directly from the interaction matrix elements  $W$  reflects intrinsic timescales rather than typical speeds. More profoundly, the periodic continuity condition connecting values of  $u$  at  $\varphi$  and  $\varphi + 2\pi$  represents the imprint of the cyclic group symmetry. This maintenance of key features accompanies the facilitation of an analytical treatment of waves on the smooth circle.

We can solve this equation (3) by separation of variables. if we assume the form below for the disturbance the wave equation becomes:

$$u(\varphi, t) = \phi(\varphi)\tau(t) \Rightarrow \frac{1}{c^2\tau} \frac{d^2\tau}{dt^2} = \frac{1}{\phi} \frac{d^2\phi}{d\varphi^2}, \quad \phi(\varphi) = \phi(\varphi + 2\pi)$$

The SLODEs above have the periodic boundary condition imposed on the part involving  $\varphi$ . So each part can only be a negative constant which respects periodic conditions:

$$\frac{d^2\phi}{d\varphi^2} + m^2\phi = 0 \Rightarrow m = 0, 1, 2, \dots \quad (4)$$

$$\frac{d^2\tau}{dt^2} + m^2c^2\tau = 0 \Rightarrow \tau(t) = C_1 \cos(mct) + C_2 \sin(mct) \quad (5)$$

Also, the time translation invariance implies that there is no preferred phase for the time factor that cannot be considered in the angle factor. So, we can represent Eigenstates of the PDE by the following functions:

$$\begin{aligned} u_0(\varphi, t) &= a_0 + c_0t \\ u_m(\varphi, t) &= a_m \cos(m\varphi) \cos(mct) + b_m \sin(m\varphi) \cos(mct) \\ &+ c_m \cos(m\varphi) \sin(mct) + d_m \sin(m\varphi) \sin(mct) \end{aligned} \quad (6)$$

The set above creates a complete orthogonal set of answers for the wave equation 3 respecting the periodic boundary conditions. What is fascinating about the continuum limit is that the completely alternating case does not happen anymore and it is not an eigenmode. As you may recall we had the first normal mode in the discrete case with zero frequency. this is exactly  $u_0$ . As it turns out, the equation below is completely solved (see Evans (2022)) and has some beautiful remarks.

$$\begin{aligned} \frac{\partial^2 u}{\partial t^2} &= c^2 \frac{\partial^2 u}{\partial \varphi^2} \\ u(\varphi, t) &= u(\varphi + 2\pi, t) \\ u(\varphi, 0) &= f(\varphi) = f(\varphi + 2\pi) \\ \dot{u}(\varphi, 0) &= g(\varphi) = g(\varphi + 2\pi) \end{aligned} \quad (7)$$

We assume here that  $f$  and  $g$  are well-behaved functions that have a Fourier Series. Every answer of the equation 7 is of the form  $\sum_{m=0}^{\infty} u_m(\varphi, t)$  where  $u_m$ s are given by the equation 6. Imposing the initial values in time we would have the following formulas for determining the constants:

$$\begin{aligned} u_0 &= \frac{1}{2\pi} \int_0^{2\pi} f(\varphi) d\varphi \\ c_0 &= \frac{1}{2\pi} \int_0^{2\pi} g(\varphi) d\varphi \\ a_m &= \frac{1}{\pi} \int_0^{2\pi} f(\varphi) \cos(m\varphi) d\varphi \\ b_m &= \frac{1}{\pi} \int_0^{2\pi} f(\varphi) \sin(m\varphi) d\varphi \\ c_m &= \frac{1}{\pi mc} \int_0^{2\pi} g(\varphi) \cos(m\varphi) d\varphi \\ d_m &= \frac{1}{\pi mc} \int_0^{2\pi} g(\varphi) \sin(m\varphi) d\varphi \end{aligned} \quad (8)$$

We can now appreciate the power of Fourier Series in this part. it is astonishing that even the couple degenerate Eigenvalues that we encountered previously are present here. This is

another reason why the alternating case does not exist in continuum limit. just because infinity is neither even nor odd and the alternating mode only happened in the even case. Another notable fact is that we had a relation between the number of mode and its angular velocity. Here we will replace some terms in it to make it compliant with continuous case. we would call such relation the *Dispersion Relation* of system.

$$\omega_k^2 = \omega_0^2 \sin^2\left(\frac{k\pi}{N}\right), \quad N \gg 1 \Rightarrow \omega_k \approx \omega_0 \frac{k\pi}{N} = mc$$

Evidently the angular frequency of the mode  $m$  is now given by the much simpler relation below:

$$\omega_m = mc, \quad m = 0, 1, 2, \dots \quad (9)$$

### 2.3. Simulation of Longitudinal Waves

In this part we will use the power of programming to Simulate different Initial Values for the wave equation driven and we will create beautiful image and plot from it. You can see them in [this link](#).

## 3. 3D Spheres: A Thorough Analysis

### 3.1. Tetrahedron and Cube

The five Platonic solids - tetrahedron, cube, octahedron, dodecahedron, and icosahedron - represent the only possible regular, convex polyhedra. Their high degree of symmetry makes them fascinating geometries within which to explore periodic boundaries and particle interactions. However, fully parametrizing their intricate shapes for mathematical analysis poses challenges.

We focus first on the tetrahedron, the simplest Platonic solid with 4 equilateral triangular faces. Modeling an N-particle system on its surface requires using the polar and azimuthal angles  $\theta$  and  $\varphi$  to specify positions. The symmetry arguments we leveraged successfully on simpler topologies unfortunately do not apply directly here. Nevertheless, the uniform connectivity of the vertices through edges and 2D faces manifests itself in the form of the equations governing the angular particle displacements.

Specifically, the interaction matrix attains a sparse block circulant form, with each subdivision corresponding to coupling between particles on a given triangular face. While deriving analytical solutions remains intractable, this structure lends itself naturally to computational techniques. By numerically evolving the coupled oscillator equations in time under different initial displacements, we can simulate waves propagating through the tetrahedral network.

Beyond the tetrahedron lies the remaining Platonic solids, with more intricate shapes described by a greater multitude of faces and vertices. The cube, for example, requires additionally modeling motion perpendicular to the faces alongside angular displacements. The octahedron and icosahedron feature more complex networks of faces meeting at vertices. And the dodecahedron has twelve regular pentagonal faces arranged in a ring-like topology around the polyhedron.

Studying particles and waves constrained to these geometries will push the boundaries of our mathematical methods. While solutions always formally exist, deriving analytical expressions proves prohibitive, and numerical techniques must bear the brunt of revealing the dynamics. The elegance of decomposition into normal modes gives way to directly integrating the equations of motion through time. And the orderly continuum limit subsides into intricacies of granularity and connectivity between finite nodes.

Yet what rewards persistence is insight - the physics of disturbance transmission through these topological networks holds surprises mathematical symmetry tends to preclude. Pursuing the complexity of Platonic geometries through a numerical investigation promises revelations and even aesthetically pleasing simulations as waves wrap around these symmetric shapes.

### 3.2. Spherical Symmetry

In transitioning from discrete particles on Platonic solids to continuous waves on a sphere, a crucial concept emerges - that of spherical symmetry. Unlike the straightforward cyclic symmetry of a circle, the surface of a sphere treats all orientations equivalently rather than exhibiting discrete rotational periodicity. Mathematically harnessing this continuous invariance proves vital for solving physical systems.

We characterize positions on the sphere through angular coordinates  $\theta$  and  $\varphi$ . Spherical symmetry manifests that the physics must remain unchanged under rotations of these angles by any amount.

Additionally, the symmetry assures interactions must occur through scalar combinations of the angles  $\theta$  and  $\varphi$  rather than depending on their specific values. This enables expanding disturbing potentials or displacement fields in basis functions of the angular coordinates that transform appropriately under rotations. Practically, the vector spherical harmonics provide one such basis set that satisfies the covariance requirement.

Through these techniques, the complexity of dynamics on the sphere give way to a simpler and more insightful description. The details become encapsulated in the expansion coefficients, while the governing physics reduces to algebraic relationships between conserved quantities. And visualizations of disturbances in terms of harmonic modes reveal Propagation patterns difficult to anticipate intuitively. The power of spherical symmetry emerges in converting intricacy into elegance - waves on a sphere become approachable when symmetry guides the mathematics.

### 3.3. 2D Wave equation and its boundary condition

In taking the continuum limit of particles interacting on a sphere, a few key simplifications emerge due to the geometry and assumed physics. As the number of particles approaches infinity, the discrete granularity disappears, leaving only a smooth surface. The short-range linear interactions then yield only second-order spatial derivatives in the wave equation, producing the familiar Laplace operator. Additionally, the inherent spherical symmetry requires that the wave speed remain constant, rather than depending on angular orientation. And confinement to the exact spherical surface means there can be no

dynamics along the radial direction - motion is restricted to tangent planes.

Together, these conditions give rise to an elegant wave equation with suitable boundary conditions:

$$\nabla^2 u = \frac{1}{c^2} \frac{\partial^2 u}{\partial t^2}$$

Here  $u(\theta, \varphi, t)$  denotes the wave displacement at time  $t$  on the sphere parametrized by angular coordinates. The associated boundary condition is one of periodicity - wave amplitudes, frequencies, and phases must match when compared at opposing boundaries corresponding to the same physical location. We can expand the Laplacian to more expanded form below:

$$\begin{aligned} \frac{1}{\sin \theta} \left[ \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial u}{\partial \theta} \right) + \frac{1}{\sin \theta} \frac{\partial^2 u}{\partial \varphi^2} \right] &= \frac{1}{v^2} \frac{\partial^2 u}{\partial t^2} \\ u(\theta + \pi, \varphi, t) &= u(\theta, \varphi, t) \\ u(\theta, \varphi + 2\pi, t) &= u(\theta, \varphi, t) \end{aligned} \quad (10)$$

As discussed earlier, this equation is completely solvable and the PDE is consistent too. In order to derive the general formula we would need to use separation of variables and periodic boundary condition together. which follows:

$$u_{\ell,m}(\theta, \varphi, t) = T_{\ell}(t) \Theta_{\ell,m}(\theta) \Phi_m(\varphi)$$

The answer to above Equation is the Famous "Spherical Harmonics" and once again, they create a complete orthogonal set of answers which we use with Fourier's trick to expand initial values based on the harmonics.

The important point here is that we only have full integer  $\ell$ s and there is no half-integer ones. this is due to periodic behavior of the wave function.

## 4. Discussion

Throughout this paper, we have tackled the problem of particles and waves confined to spherical topologies through a mathematical physics approach, seeking exact or numerical solutions to the equations of motion and associated boundary conditions. While enabling elegant analysis in some regimes, as particle numbers grow large, fundamental limitations arise in tracking individual coordinates. This motivates adopting a statistical description in terms of collective dynamics and emerging macrostates.

Specifically, statistical mechanics provides a rigorous framework connecting the microscopic interactions studied here to macroscopic thermodynamic observables. Rather than following each trajectory, we partition the high-dimensional phase space into "cells" and ask how these populations evolve. Conservation laws derived from symmetries constrain allowable transitions between cells. And key quantities like free energy, entropy, temperature, and chemical potential emerge from counting arrangements accessible to the system under its constraints. Connecting these back to appropriate ensembles of the solutions derived earlier remains an open avenue.

Pursuing such a program promises new insights into periodic systems where the topology itself plays the role of a constraint on the dynamics. Spherical symmetry and continuity present conserved quantities akin to energy or particle number in simpler systems. And computational modeling already provides phase space evolution akin to molecular dynamics. Bringing information-theoretic ideas into this geometric setting may reveal deep and unexpected connections between curvature, causality, and statistical emergence in continuous media. The languages of dynamics, topology, and thermodynamics stand poised for synthesis.

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