

Dynamical insights from Lagrangian-filtered structure functions

HAN WANG (UNIVERSITY OF EDINBURGH, UK)

DHRUV BALWADA (COLUMBIA UNIVERSITY, USA)

JINHAN XIE (PEKING UNIVERSITY, CHINA)

TRR 181 Eddy-Wave Meeting, .2023



Gulf of Mexico surface drifter

Similar spatial region, drogued on the surface.

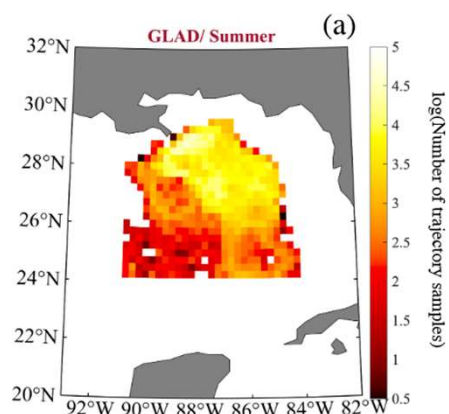
Velocities at 15 min interval. Spans 100 m – 100 km.

Lagrangian filtering 2nd order structure functions



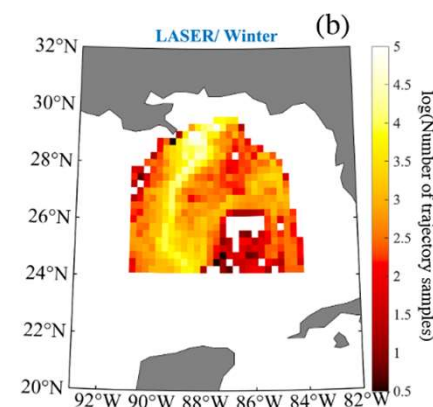
Grand LAGrangian Deployment (GLAD)

- Summer 2012



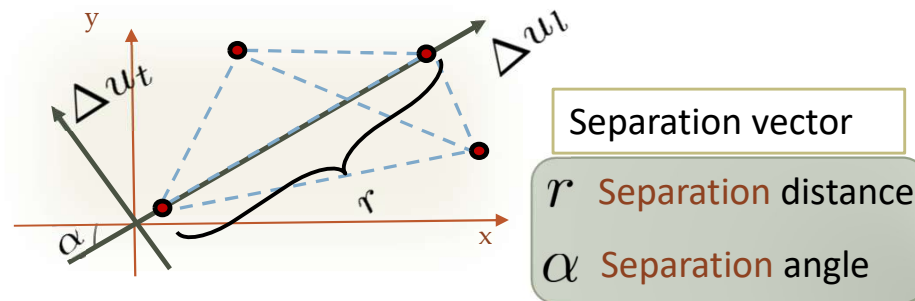
Lagrangian Submesoscale Experiment (LASER)

- Winter 2016



Balwada, Xie, Marino and Feraco 2022

2nd order structure functions



- At each snapshot, $\Delta \mathbf{u}$ projected into directions **along/across** the **separation vector** (r, α)
- Under **horizontal homogeneity and stationarity**,
 $\overline{(u_l(\mathbf{r}_0; t_0) - u_l(\mathbf{r}_0 + \mathbf{r}; t_0))^2}$ is a function of \mathbf{r} only

→ **second order structure functions** :

$$D_{ll}(r, \alpha) = \overline{\Delta u_l^2(r, \alpha)}$$

$$D_{tt} = \overline{\Delta u_t^2(r, \alpha)}, D_{lt} = \overline{\Delta u_l(r, \alpha) \Delta u_t(r, \alpha)}$$

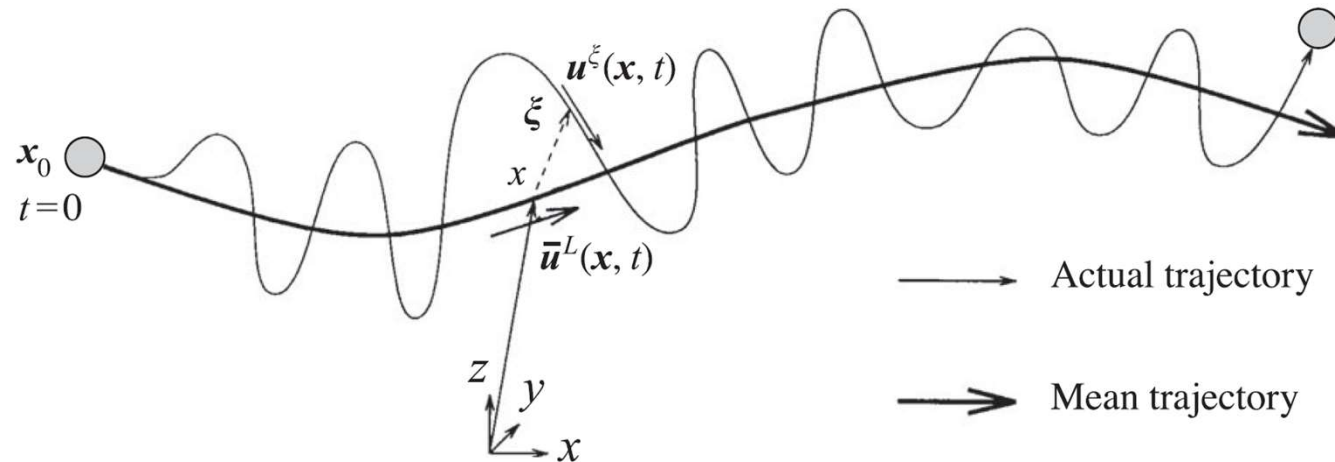
Lagrangian filtering (“LF”)

$$\mathbf{x} + \boldsymbol{\xi}(\mathbf{x}, t)$$

slow

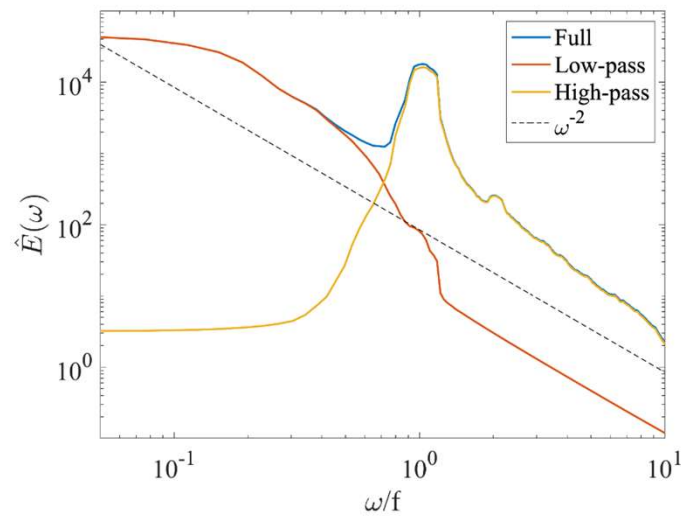
fast

$$\bar{\mathbf{u}}^L + \mathbf{u}^\xi$$

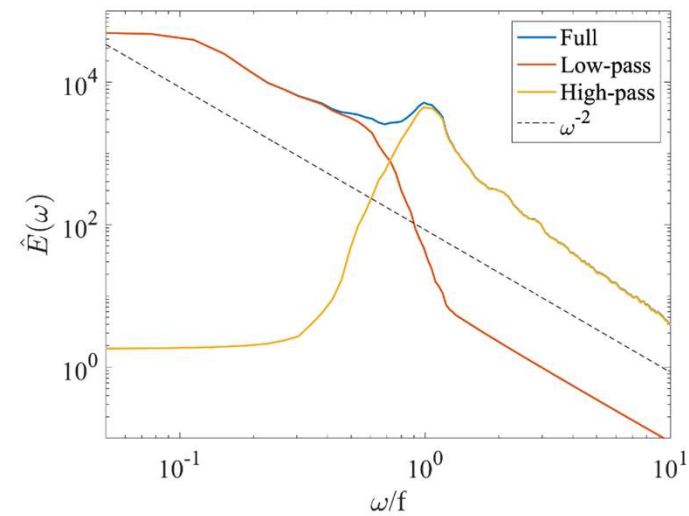


Bühler 2009

LF-ed velocities



(a) *GLAD*

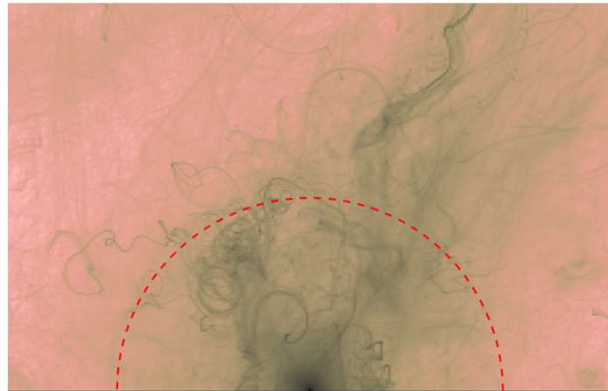
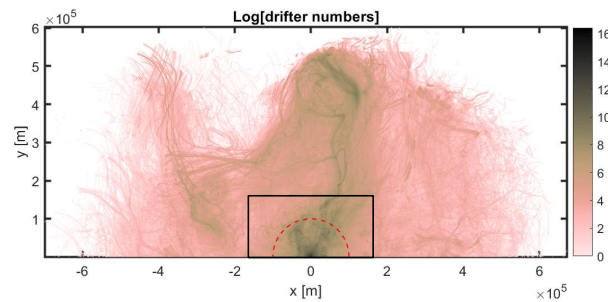
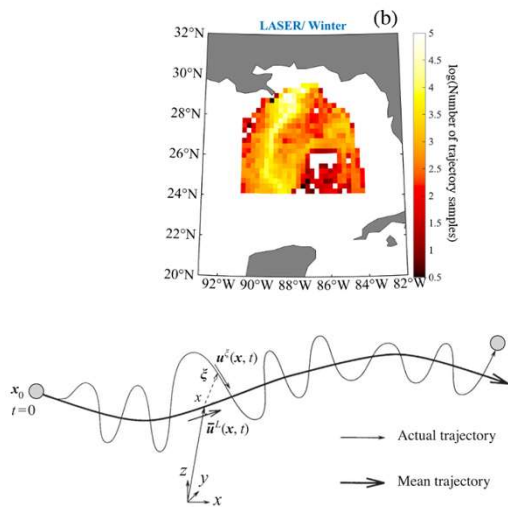


(b) *LASER*

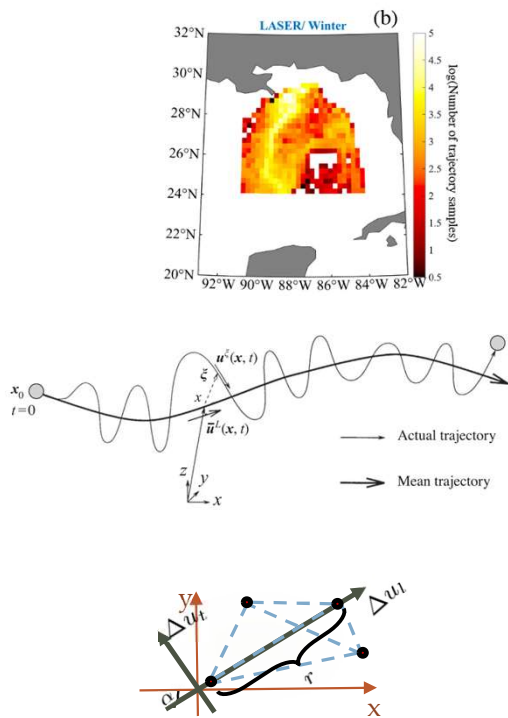
- 1.5 days
- High-pass includes f

Mean (low-pass) trajectories

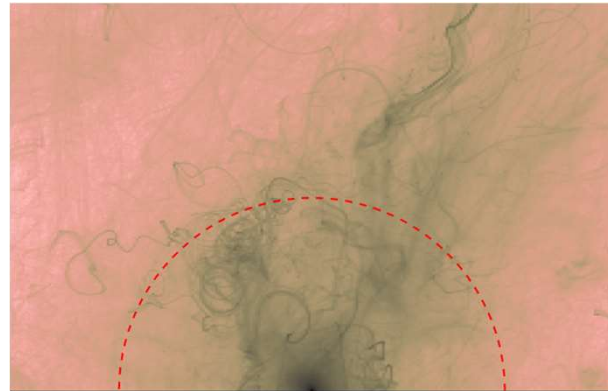
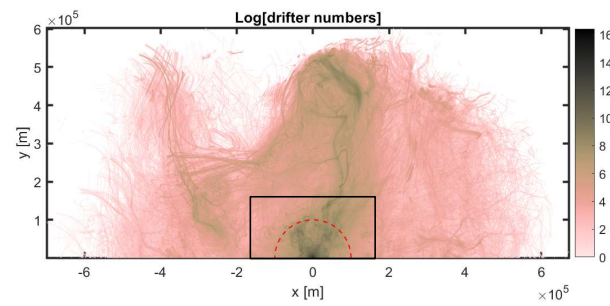
Separation vector from mean trajectories



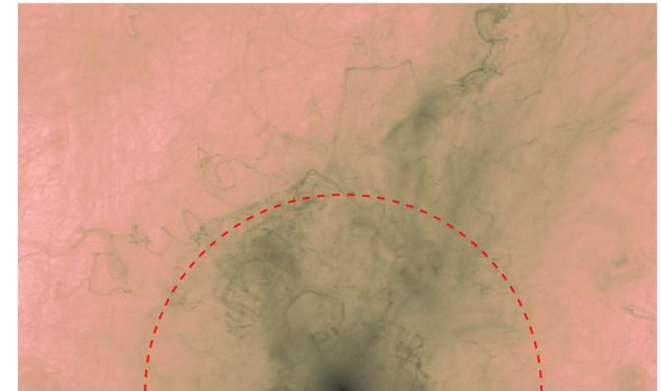
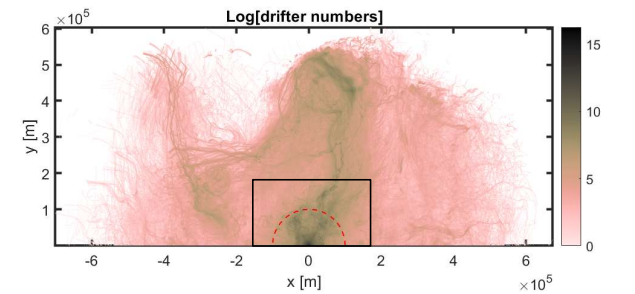
Effects of LF



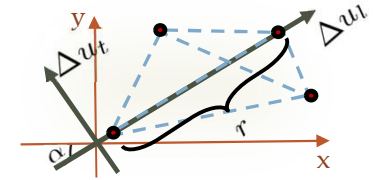
Separation vector from mean trajectories



... from raw trajectories



Helmholtz decomposition



$$\begin{aligned} u &= \phi_x - \psi_y \\ v &= \phi_y + \psi_x \end{aligned} \quad \Delta \mathbf{u} = \Delta \mathbf{u}_{\text{rot}} + \Delta \mathbf{u}_{\text{div}}$$

↓
due to rotational

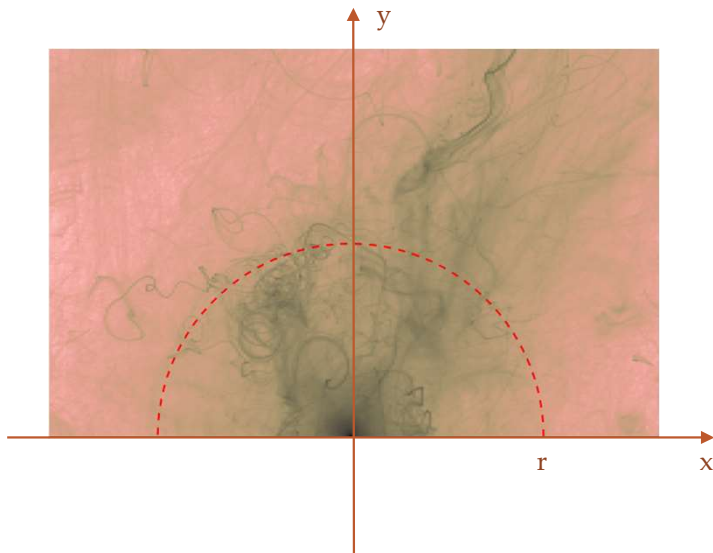
Consider **isotropy** $D_{ll}(r, \alpha) = \overline{\Delta u_l^2(r, \alpha)}$

To relax later $D_{tt} = \overline{\Delta u_t^2(r, \alpha)}, D_{lt} = \overline{\Delta u_l(r, \alpha) \Delta u_t(r, \alpha)}$

$$\begin{aligned} D_{\text{rot}}(r) &= \int_0^r \left(\frac{-D_{ll}(\tau) + D_{tt}(\tau)}{\tau} \right) d\tau + D_{tt}(r), \\ D_{\text{div}}(r) &= - \int_0^r \left(\frac{-D_{ll}(\tau) + D_{tt}(\tau)}{\tau} \right) d\tau + D_{ll}(r) \end{aligned}$$

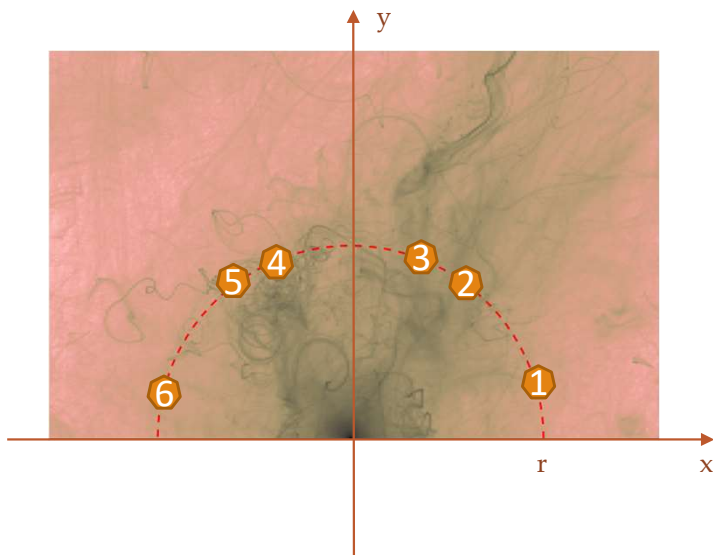
Bühler, Callies and Ferrari 2014
Lindborg 2015

Angle weighting



- Anisotropic distribution of separation vector

Angle weighting

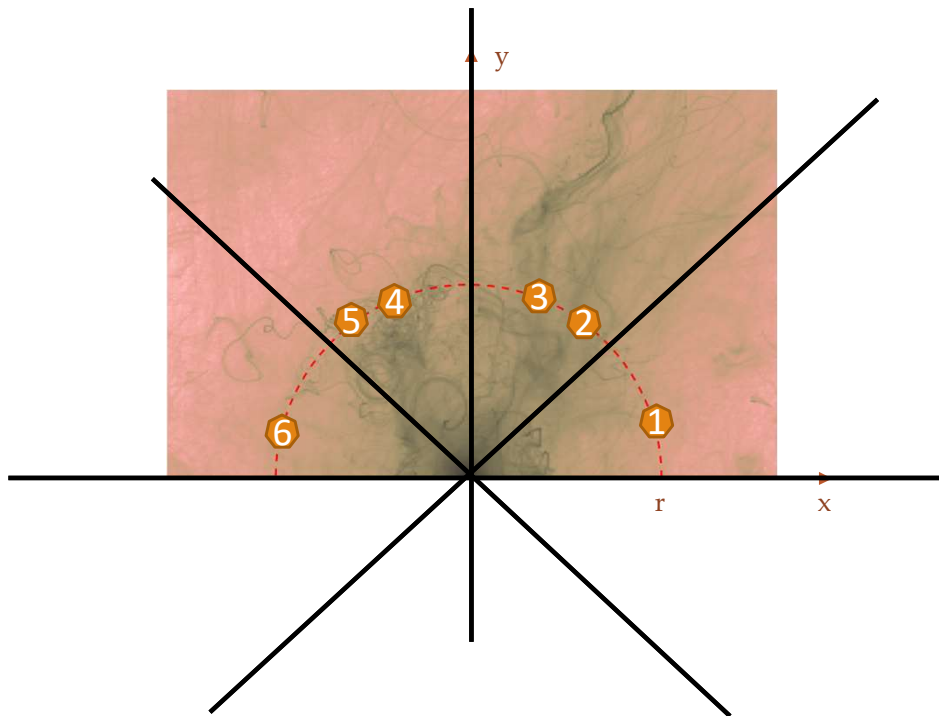


- Anisotropic distribution of separation vector

To compute $D_{ll}(r) = \overline{\Delta u_l^2(r)}$

- Traditional: unweighted average over all observations at each r
(1+2+3+4+5+6)/6
- Angle-weighted (Wang and Bühler 2021):
Weight by angles to account for nonuniformity

Angle weighting

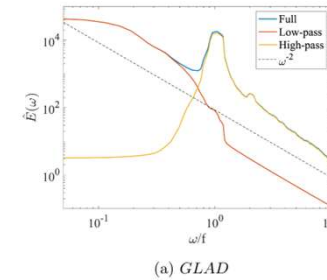


- Anisotropic distribution of separation vector

To compute $D_{ll}(r) = \overline{\Delta u_l^2(r)}$

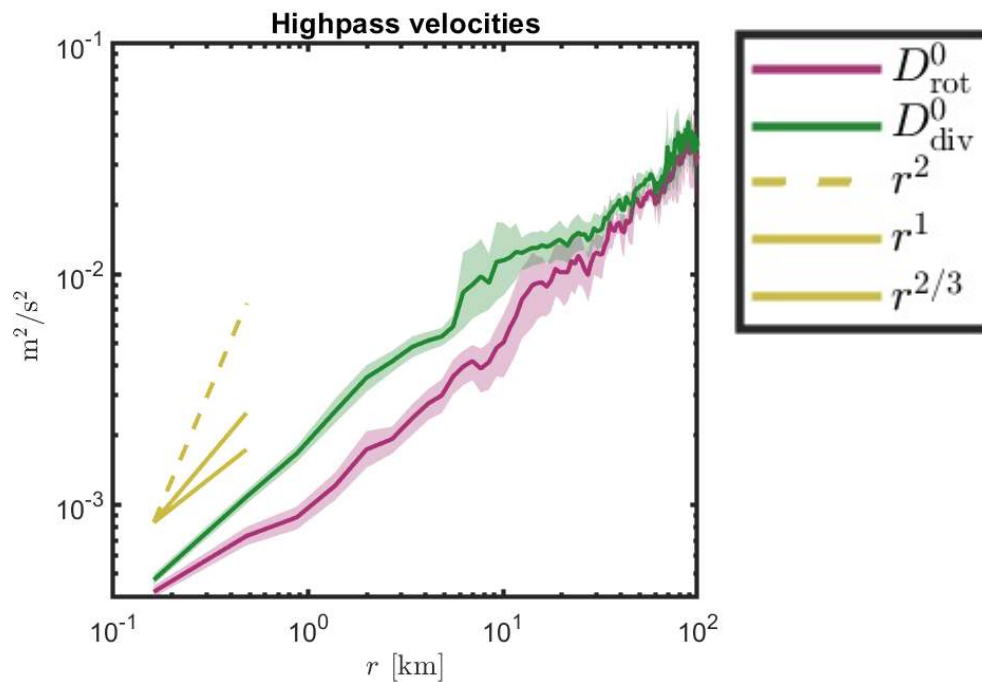
- Traditional: unweighted average over all observations at each r
($\textcircled{1} + \textcircled{2} + \textcircled{3} + \textcircled{4} + \textcircled{5} + \textcircled{6}$)/6
- Angle-weighted (Wang and Bühler 2021):
Weight by angles to account for nonuniformity
[$\textcircled{1} + (\textcircled{2} + \textcircled{3})/2 + (\textcircled{4} + \textcircled{5})/2 + \textcircled{6}$]/4

Helmholtz: GLAD, highpass



$$u = \phi_x - \psi_y$$

$$v = \phi_y + \psi_x$$



- div \geq rot, consistent with linear theory

$$\mathbf{u}_t + \mathbf{u} \cdot \nabla \mathbf{u} + \mathbf{f} \times \mathbf{u} + g \nabla \eta = 0$$

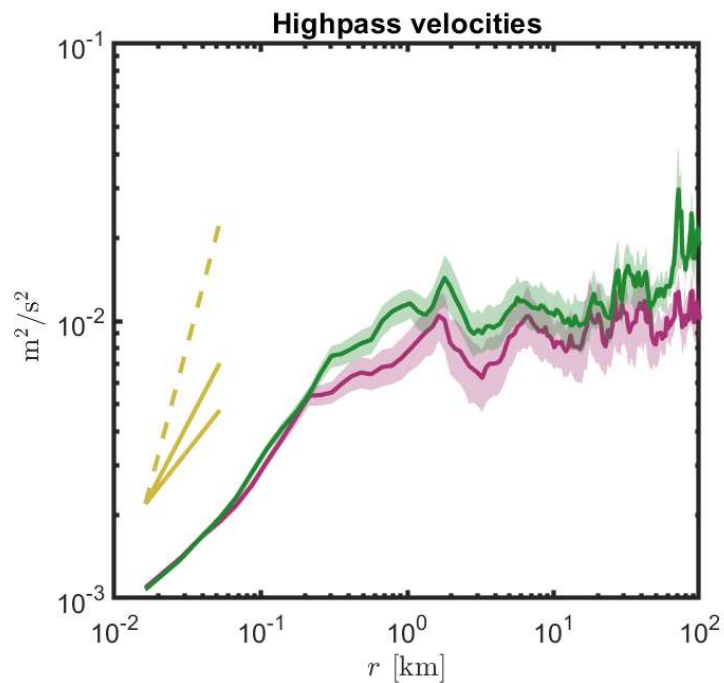
$$\Rightarrow \frac{\partial \psi}{\partial t} + f \phi = 0$$

$$\Rightarrow \frac{D_{\text{div}}(r)}{D_{\text{rot}}(r)} \approx \left| \frac{\omega(k)}{f} \right|^2, \quad k = 2\pi/r$$

$$|\omega| \geq f \Rightarrow D_{\text{div}}(r) \geq D_{\text{rot}}(r)$$

- div \sim rot at larger scales: NIWs! $|\omega| \approx f, \quad k \ll m$
- Slope

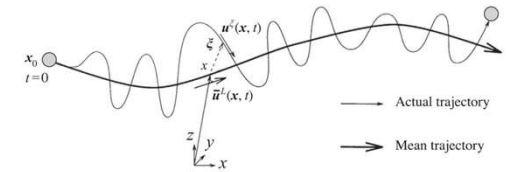
Helmholtz: LASER, highpass



- $\text{div} \geq \text{rot}$, consistent with linear theory

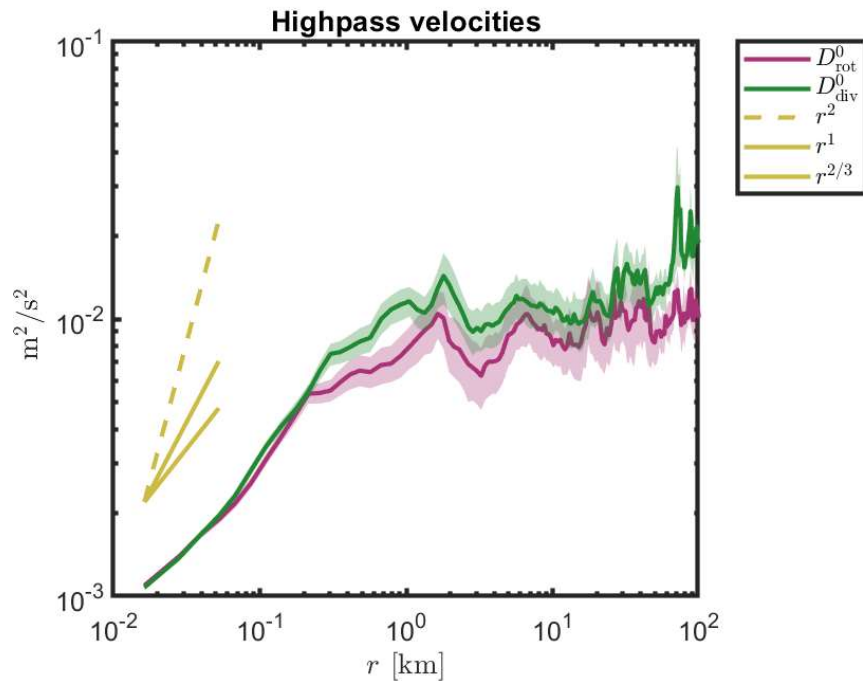
$$\frac{D_{\text{div}}(r)}{D_{\text{rot}}(r)} \approx \left| \frac{\omega(k)}{f} \right|^2, \quad k = 2\pi/r$$

- Decorrelation; may be impacts from balanced flows?

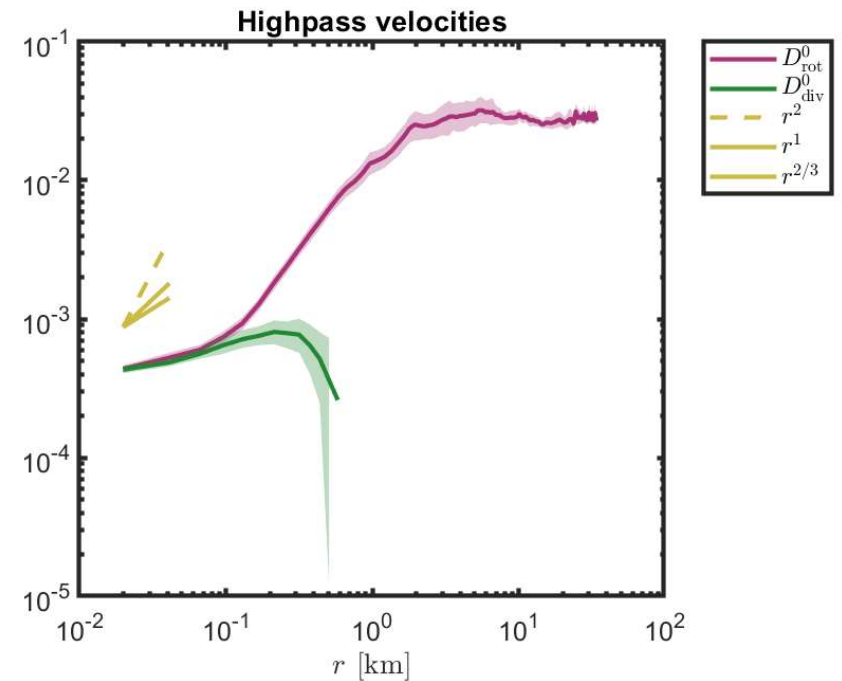


Mean trajectories make more sense!

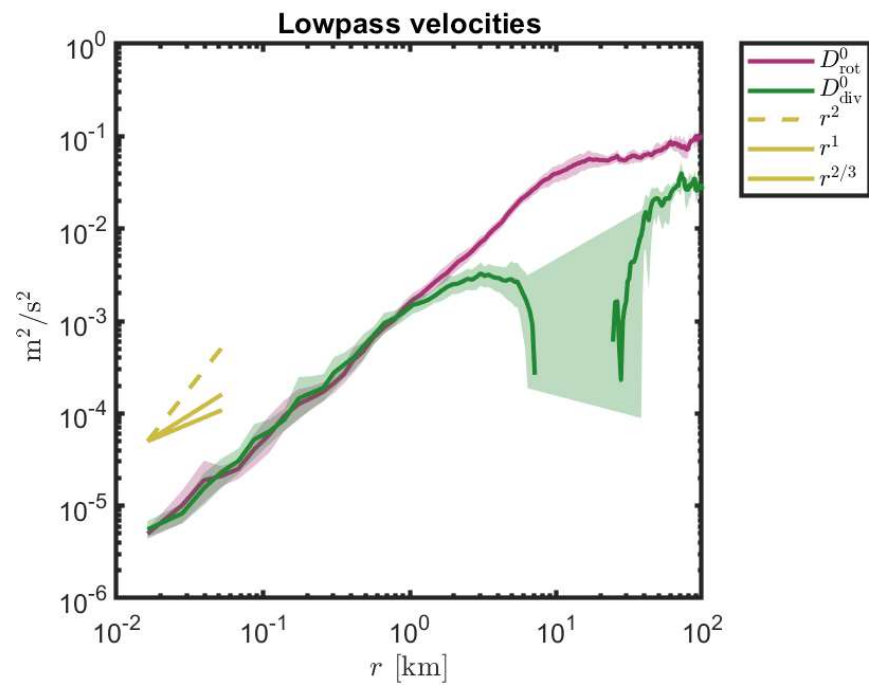
On mean trajectories, LASER



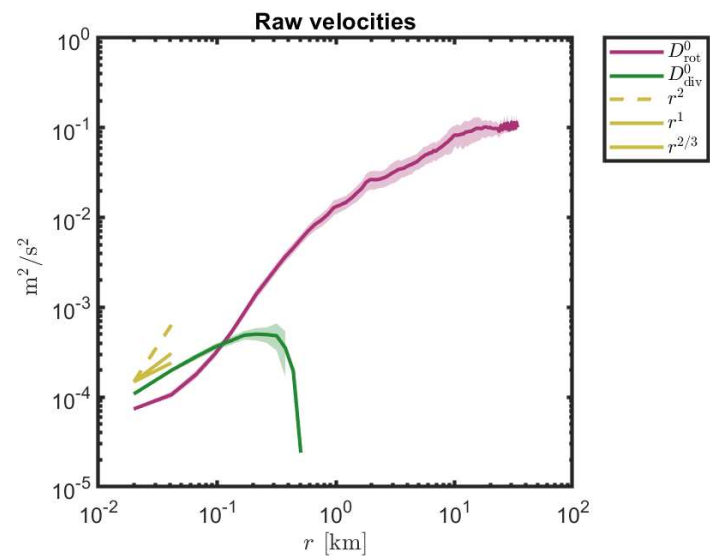
On raw trajectories, LASER



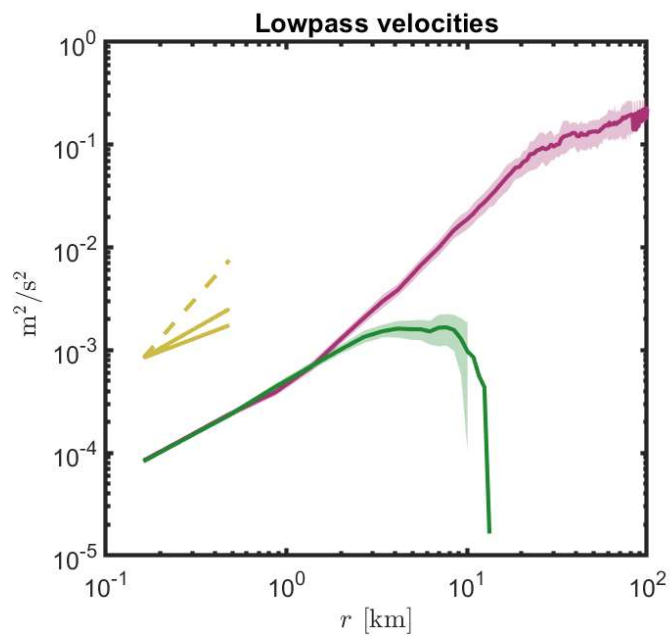
Helmholtz: LASER, lowpass



- Dominance of rot at > 10 km
- At < 1 km, equipartition
- Compare with raw:

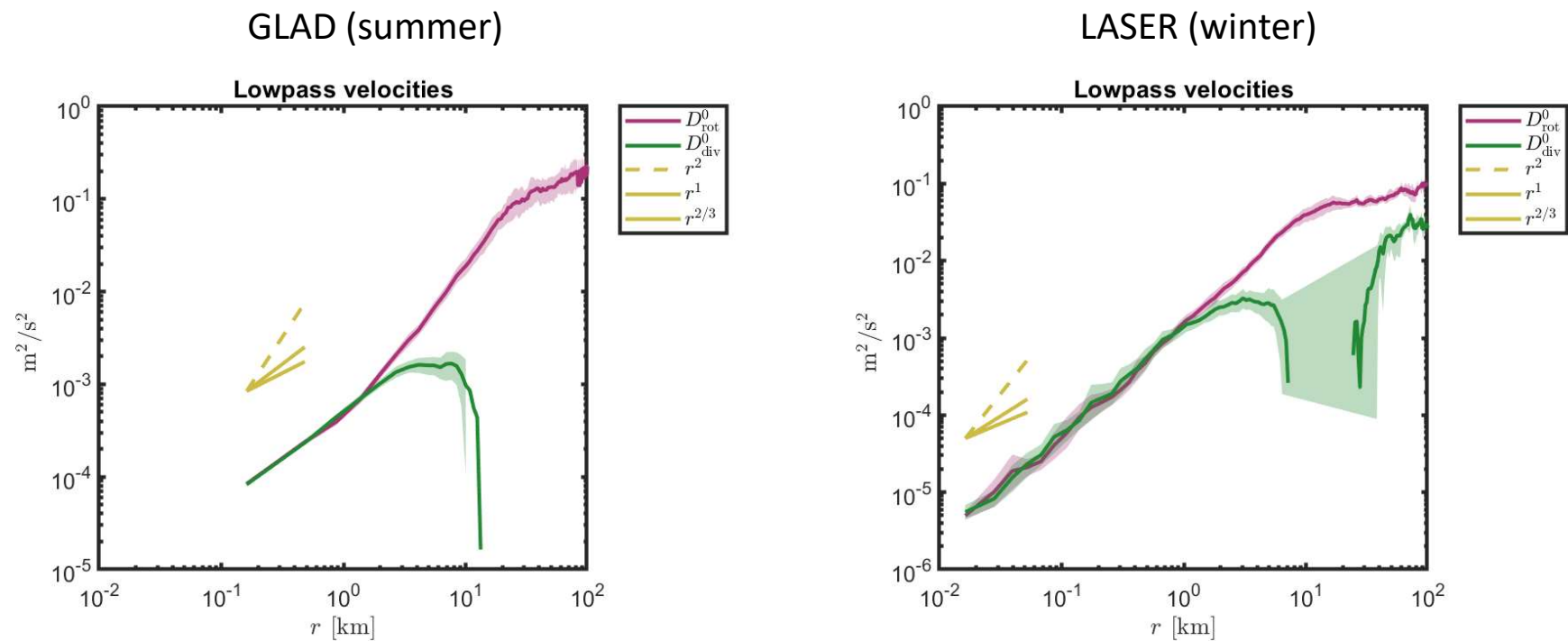


Helmholtz: GLAD, lowpass

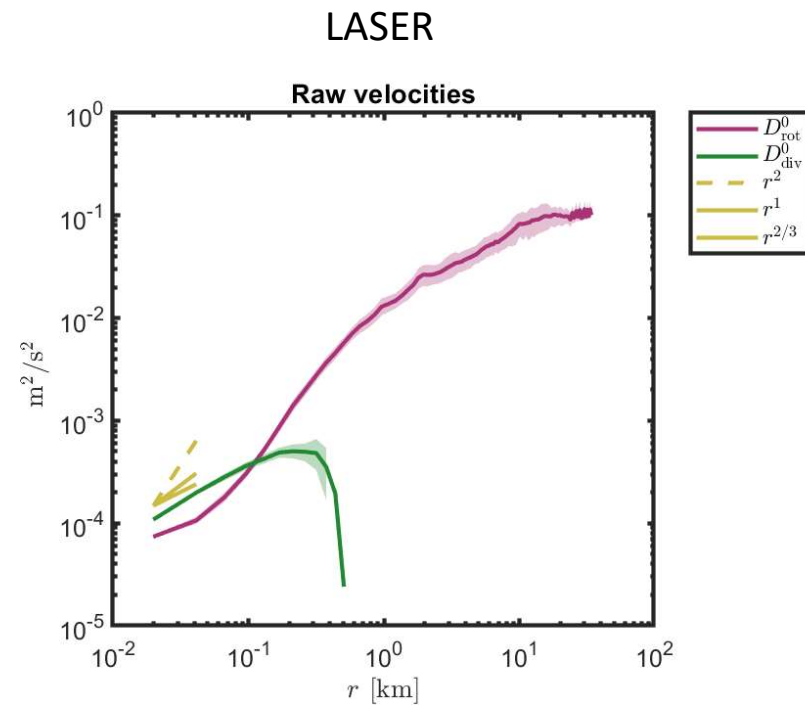
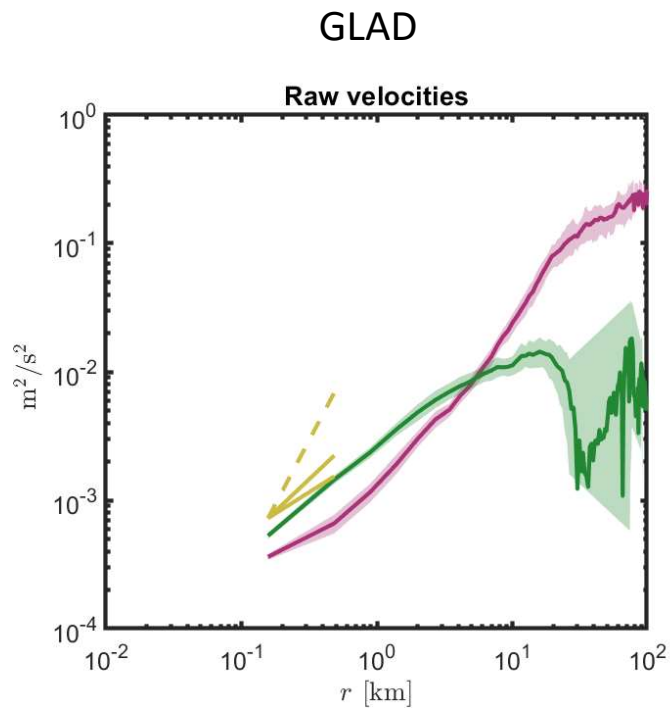


- Dominance of rot at > 10 km
- At < 1 km, equipartition

Transitional scales insensitive to mixed layer thickness



... which can't be seen in unfiltered data



Anisotropic Helmholtz decomp.

$$f(r, \alpha) = \sum_{n=0}^{\infty} f^{cn}(r) \cos(n\alpha) + f^{sn}(r) \sin(n\alpha),$$

→ Evaluate $D_{\text{rot}}^{c0}(r), D_{\text{rot}}^{c2}(r), D_{\text{rot}}^{s2}(r), \dots$

→ Compare

$$D_{\text{rot}}^0 = D_{\text{rot}}^{c0}$$

with

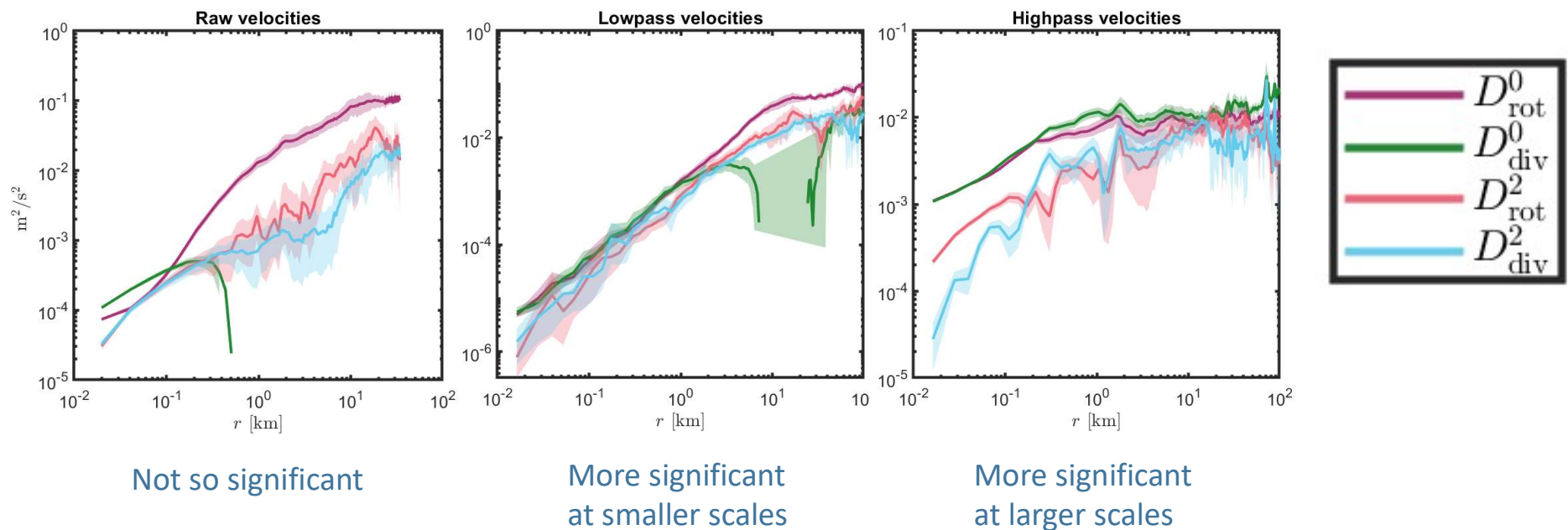
$$D_{\text{rot}}^2 = \sqrt{(D_{\text{rot}}^{c2})^2 + (D_{\text{rot}}^{s2})^2}$$

2: “2nd Fourier mode”

$$\begin{aligned} D_{\text{rot}}^{cn}(r) &= \int_0^r \left(\frac{-D_{ll}^{cn}(\tau) + D_{tt}^{cn}(\tau)}{\tau} - n \frac{D_{lt}^{sn}(\tau)}{\tau} \right) d\tau + D_{tt}^{cn}(r), \\ D_{\text{rot}}^{sn}(r) &= \int_0^r \left(\frac{-D_{ll}^{sn}(\tau) + D_{tt}^{sn}(\tau)}{\tau} + n \frac{D_{lt}^{cn}(\tau)}{\tau} \right) d\tau + D_{tt}^{sn}(r), \\ D_{\text{div}}^{cn}(r) &= - \int_0^r \left(\frac{-D_{ll}^{cn}(\tau) + D_{tt}^{cn}(\tau)}{\tau} - n \frac{D_{lt}^{sn}(\tau)}{\tau} \right) d\tau + D_{ll}^{cn}(r), \\ D_{\text{div}}^{sn}(r) &= - \int_0^r \left(\frac{-D_{ll}^{sn}(\tau) + D_{tt}^{sn}(\tau)}{\tau} + n \frac{D_{lt}^{cn}(\tau)}{\tau} \right) d\tau + D_{ll}^{sn}(r). \end{aligned}$$

Wang and Bühler, 2021

Anisotropic modes in LASER



Take home message

Methods

- Combine **spatial** and **temporal** info

Helmholtz

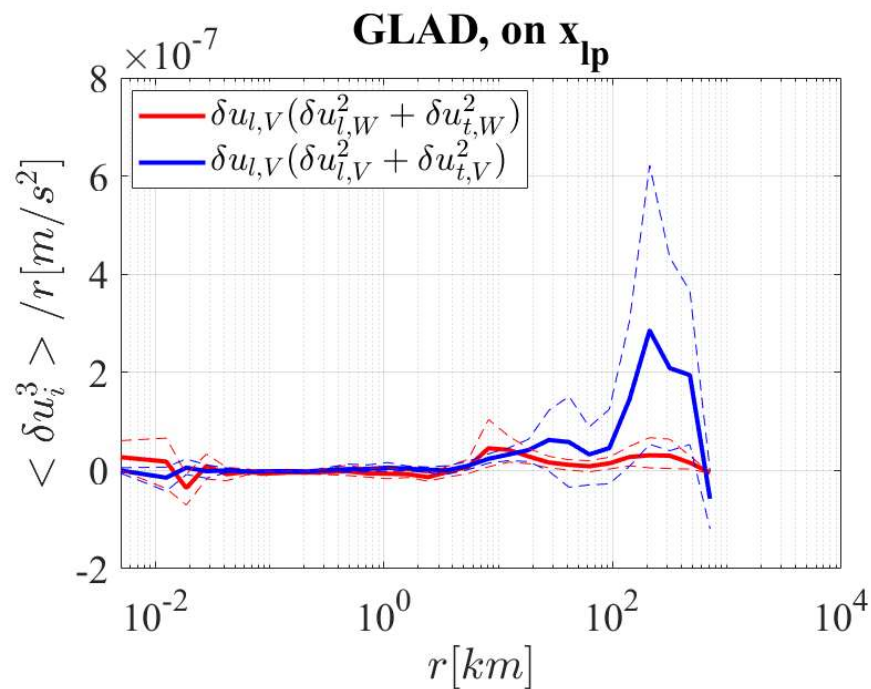
LF

- Angle weighting: simple idea to avoid bias from nonuniform angle distribution. Useful for **isotropic** modes too
- LF: High/low-pass velocities on **mean trajectories**

Observations (LASER and GLAD)

- High-pass component consistent with **linear wave theory**, esp. near-inertial waves
- Low-pass component: transitional scale insensitive to mixed layer thickness
- See anisotropy more in frequency-filtered fields

Third-order structure functions



- Reflects energy transfer rate
- Recent work on unfiltered data:
Balwada, Xie, Marino and Feraco 2022

Evidence of a dual kinetic energy cascade by surface drifter observation in the Walvis Ridge Region

Julia Draeger-Dietel¹, Alexa Griesel¹ and Jochen Horstmann²

¹Universität Hamburg, Faculty of Mathematics, Informatics and Natural Sciences, Theoretical Oceanography, Germany; ²Institute for Coastal Research, GKSS Research Center, Geesthacht, Germany

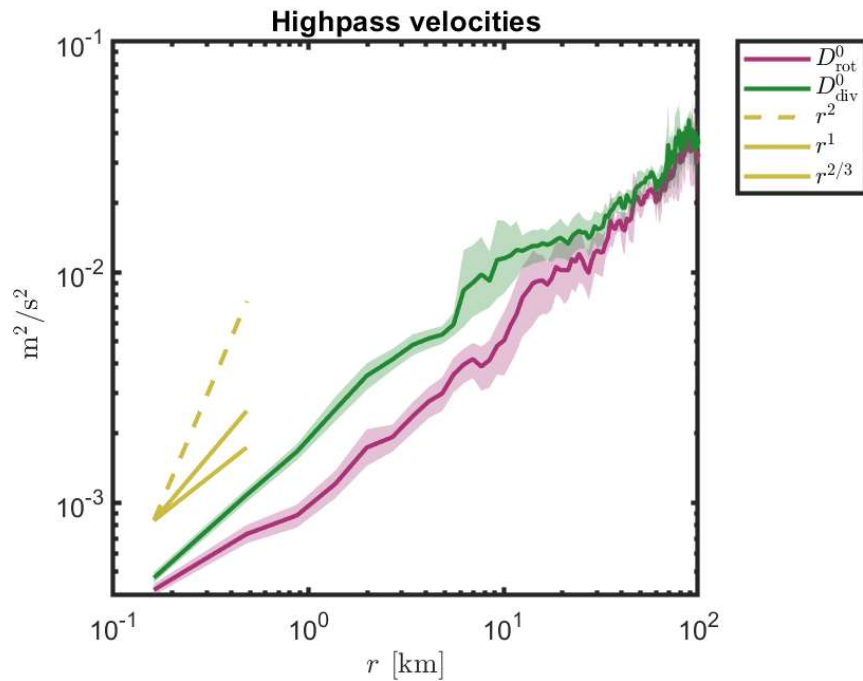
16 Feb
15:10h

- Under some theoretical constraints, we can decompose slow-fast and slow-slow transfer rates
- Collaboration **welcome**

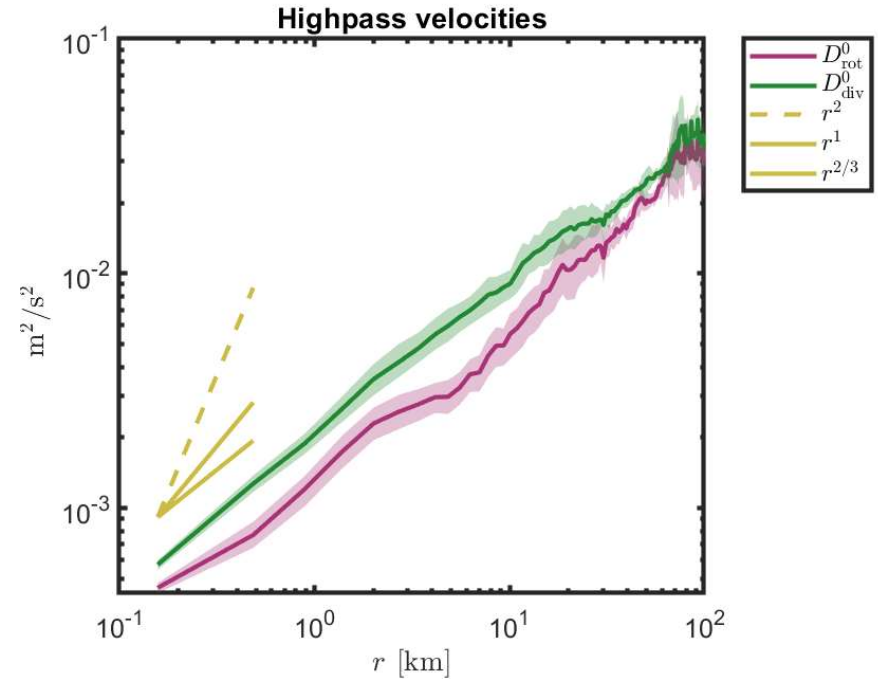
Backup slides

...and it's not coding error

On LF-ed locations, GLAD



On raw locations, GLAD



GLAD

