

My dg adventures

So far everything sucks.

Update: There's light at the end.

Hydrostatic case

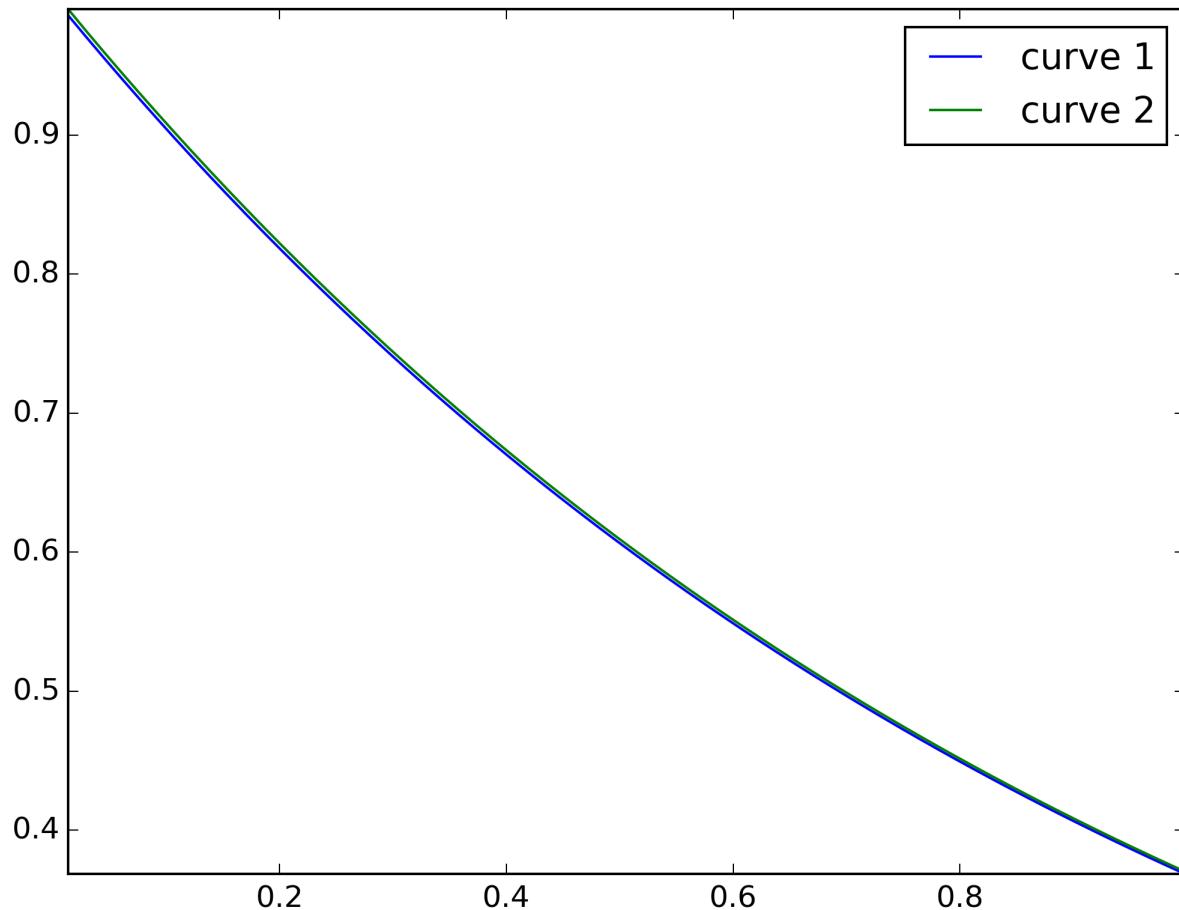
I.c.s:

$$\rho_0 = \exp(-x)$$

$$u_0 = 0.$$

$$p_0 = \exp(-x)$$

N = 64, T=100

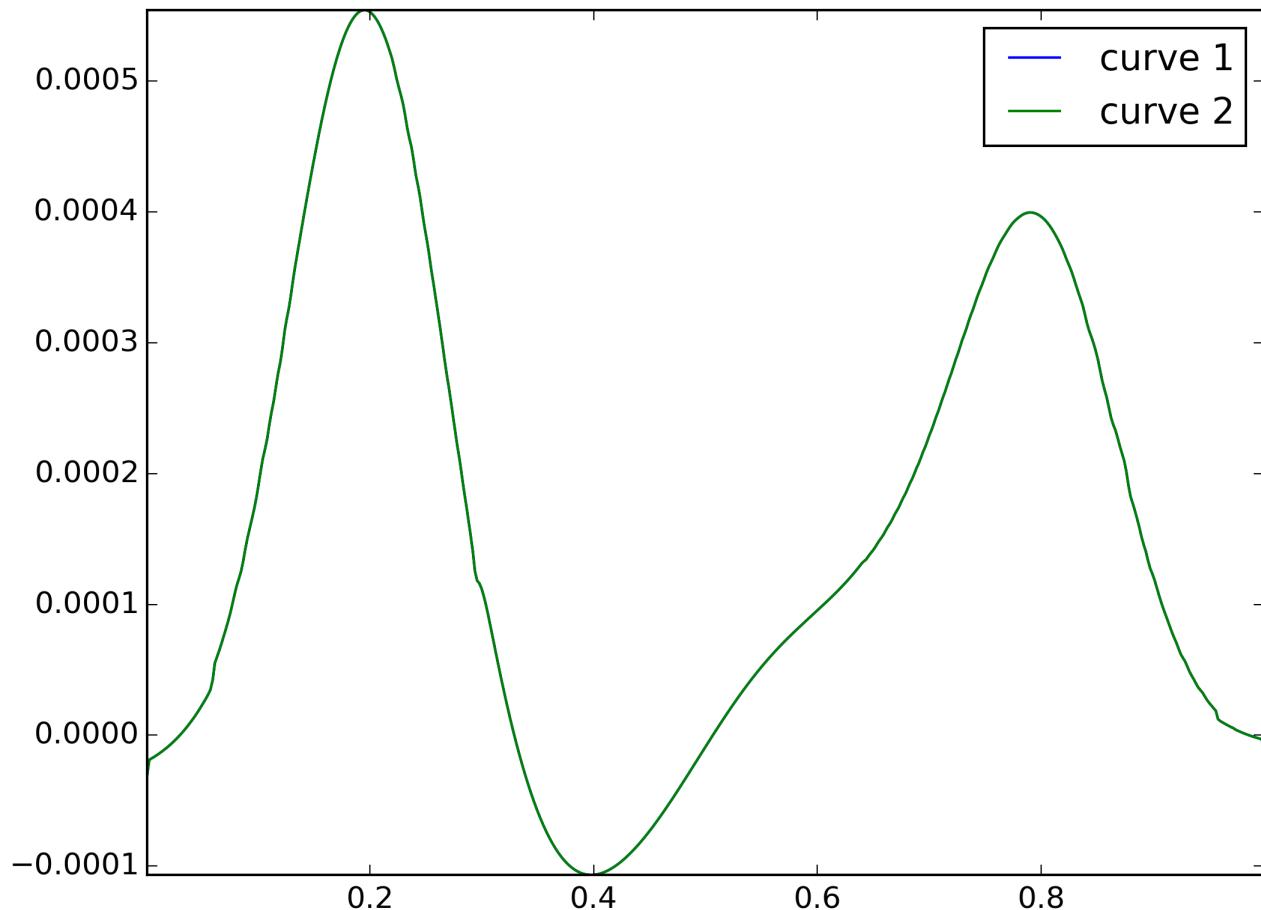


hydrostatic case + perturbation

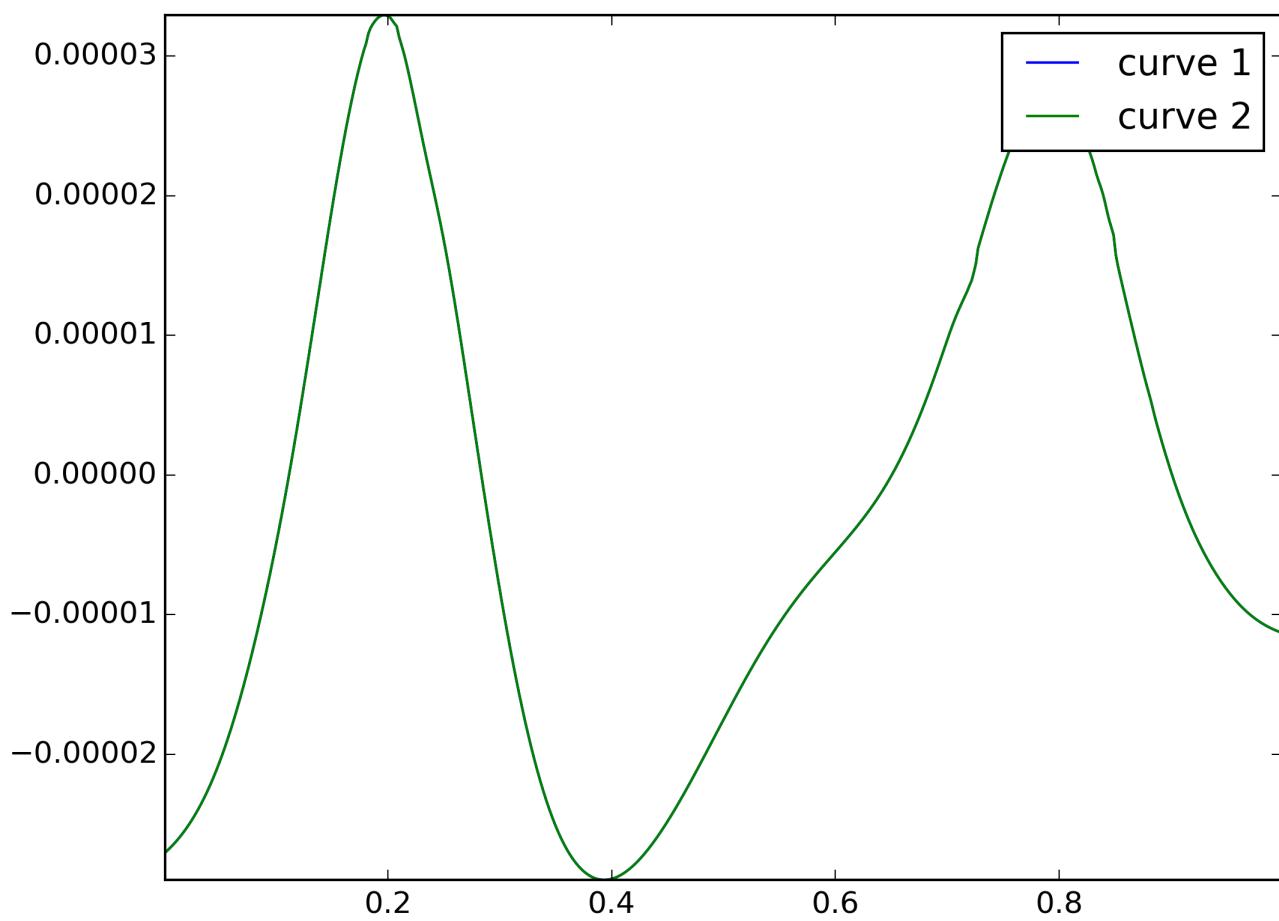
$$\begin{aligned}
\rho_0 &= \exp(-x) \\
u_0 &= 0. \\
a &= \exp(-x) + \alpha * \exp(-100 * (x - x_c)^2)
\end{aligned}$$

$$x_c = \frac{l}{2}$$

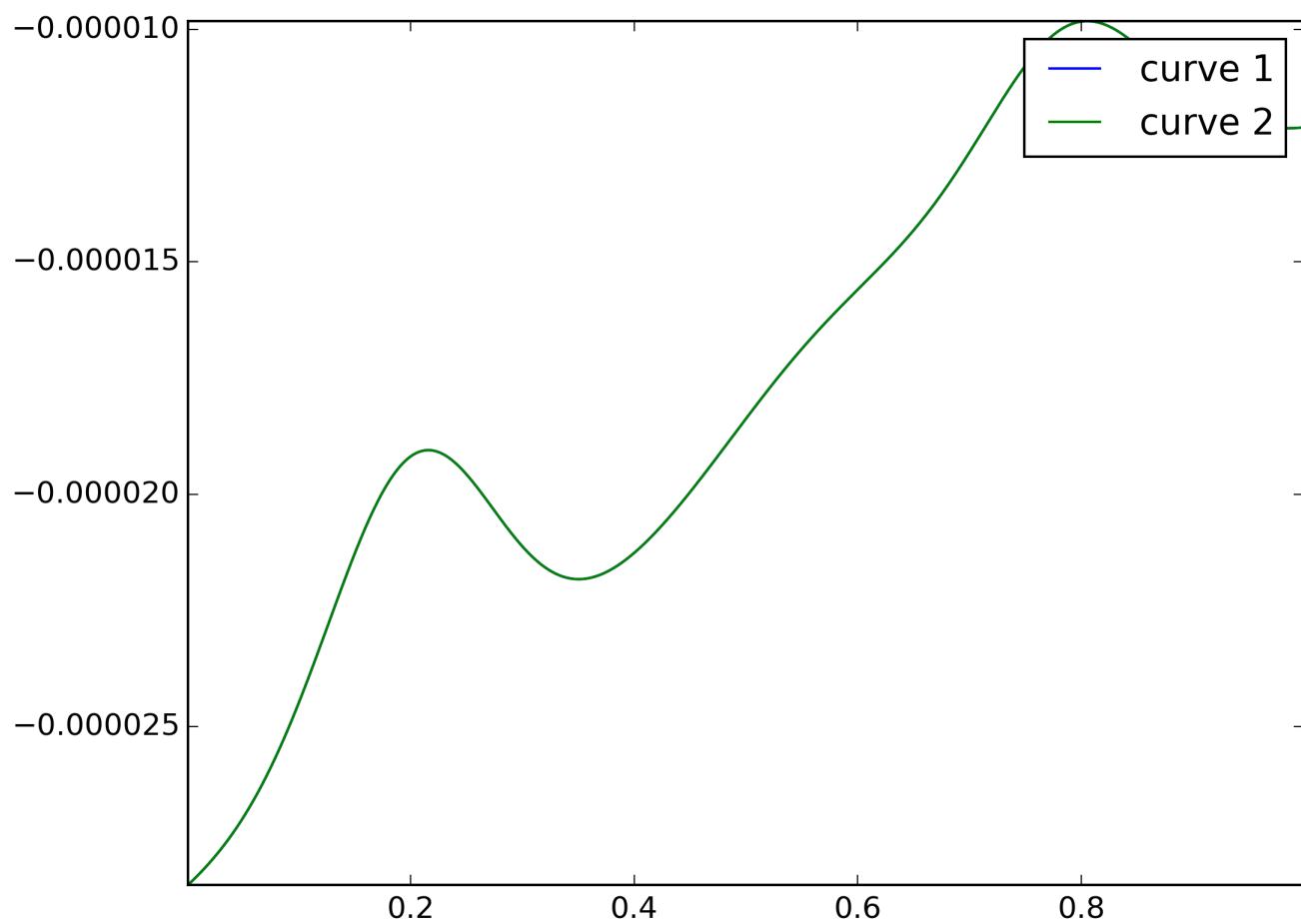
$$\alpha = 0.001$$



$$\alpha = 0.0001$$



$$\alpha = 0.00001$$



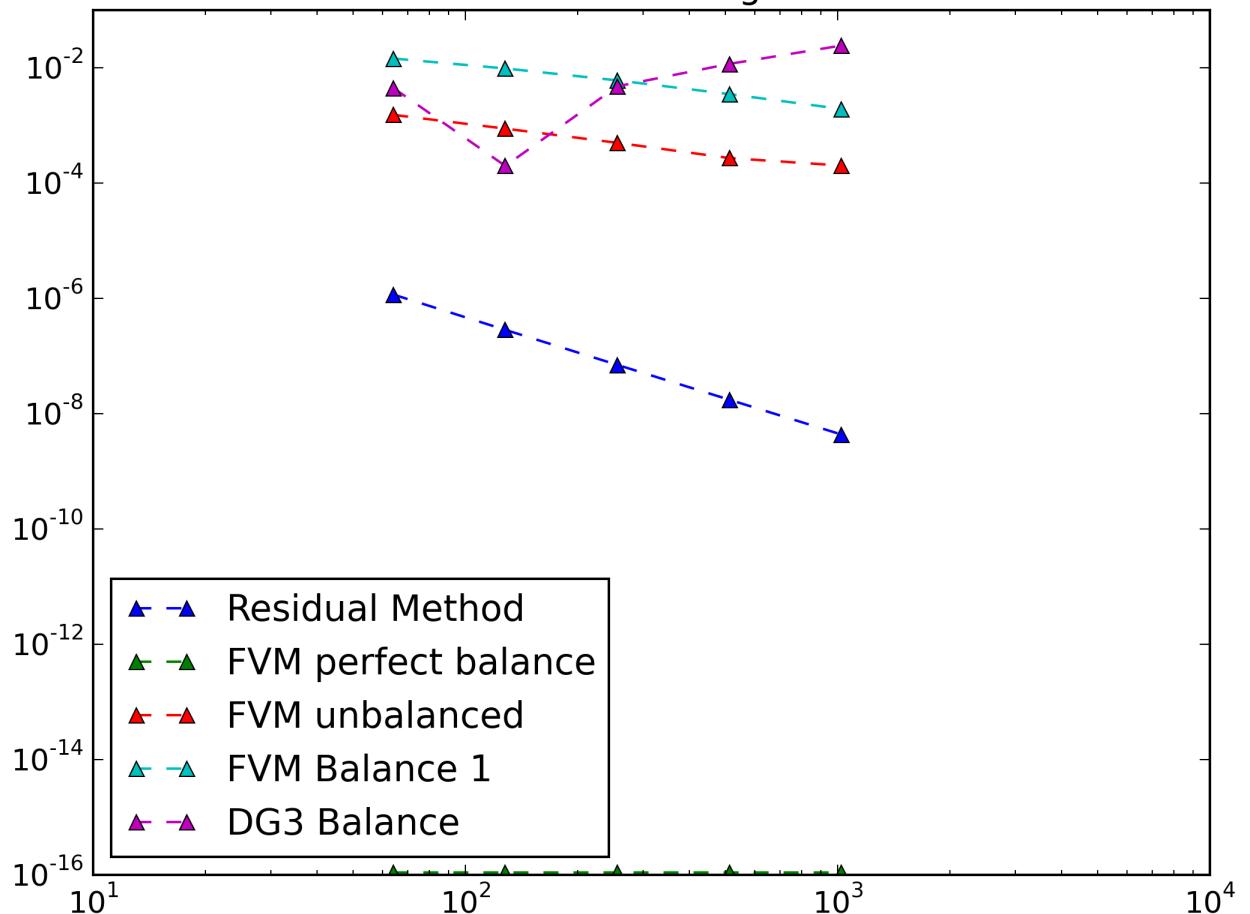
TT__TT

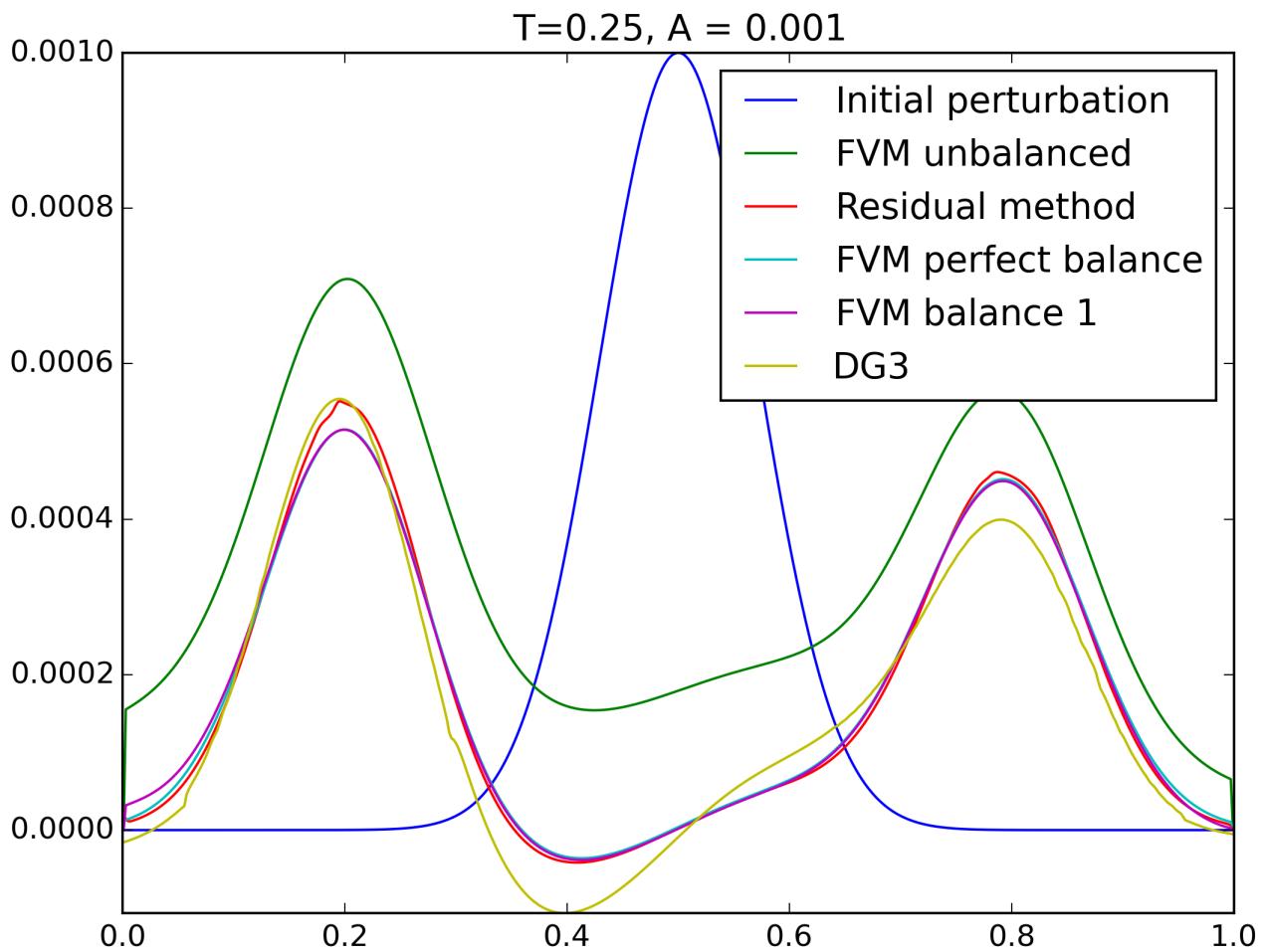
Boundary conditions don't seem to matter so much.

Convergence plot:

$$l_{\infty} = \max_i |(u(x_i, t) - \hat{u}(x_i, t)|$$

L_∞ error norm vs grid size





Dg doesn't make any sense. The growing error with refinement of the mesh is disturbing.

Debugging log

If initial condition is the equilibrium solution, delta_u should be identically 0, and all parts should cancel out.

$T = 0.0$

u_left - u_eq		
-6.2042815329732548E-011	0.000000000000000	-1.5510703832433137E-010
mode		
1		
source 1		
0.000000000000000	0.000000000000000	0.000000000000000
flux volume 1		
0.000000000000000	0.000000000000000	0.000000000000000
flux surface1		
-3.7540074831897270E-008	-3.1789056720299413E-008	-9.3850388853862138E-008
mode		
2		
source 1		
0.000000000000000	0.000000000000000	0.000000000000000
flux volume 1		
0.000000000000000	0.000000000000000	0.000000000000000
flux surface1		
-6.5021316928783755E-008	5.4980318964226171E-008	-1.6255364180498508E-007
mode		
3		
source 1		
0.000000000000000	0.000000000000000	0.000000000000000
flux volume 1		
0.000000000000000	0.000000000000000	0.000000000000000
flux surface1		
-8.3942159204551289E-008	-7.1082695285440423E-008	-2.0985584919202433E-007

T after 100 iterations

mode		
1		
source 1		
0.000000000000000	4.5757682205493921E-006	0.000000000000000
flux volume 1		
0.000000000000000	0.000000000000000	0.000000000000000
flux surface1		
-4.1529389942100458E-006	5.3033872404739668E-006	-1.5875515877461083E-005
mode		
2		
source 1		
0.000000000000000	0.000000000000000	0.000000000000000
flux volume 1		
0.000000000000000	-4.0578370776529482E-003	0.000000000000000
flux surface1		
-7.1931013387057904E-006	4.0671328542885021E-003	-2.7497199630151634E-005
mode		
3		
source 1		
0.000000000000000	5.5511151231257827E-017	0.000000000000000
flux volume 1		
0.000000000000000	2.2737367544323206E-013	0.000000000000000
flux surface1		
-9.2862538974632689E-006	1.1858734296765761E-005	-3.5498732679880204E-005

Higher modes seem to start deviating. Is this numerical error? Limiting?

<http://arxiv.org/pdf/1511.08739v1.pdf>

From Chandarashekhar and Zenk's paper, they talk about the limiting (nodal DG). " A Nonlinear TVD limiter is necessary (...). The limiter can destroy well balancedness" so they don't limit if the residual solution is close to zero.

Introduced a check which quits limiting if

$$|\delta u_i| < 1 \times 10^{-6}$$

Doesn't change much. :(

Maybe it's the quadrature? We use Newton Raphson to get the 0s of the legendre polynomials.

Increasing the iterations to 50: Error for a test case: 1.4750444904687576E-003

Increasing the iterations to 100: Error for a test case: 1.4750444904687576E-003

Nop.

Limiters?

Maybe it's the limiter which is introducing this increasing difference between u and u_{eq} , even when $u_0 = u_{eq}$.

$Nx = 128$, $N = 3$, $T = 100$, BC: $\frac{du}{dt} = 0$.

Using Lydia's limiter Vs basic TVD limiter as reported in [Cockburn and Shu](#):

Doesn't change much

$N = 1$ should be identical to FVM... what's x_i ? (gauss quadrature for $n=1$ is not defined so spit out $x_1 = 0$).

Reformulating the scheme

Represent:

$$u(x, t) = u_{eq}(x, t) + \sum_{i=0}^n \delta \hat{u}_i(t) \phi_i(x)$$

Compute the modes of the perturbation using DG framework. Let $\phi(x) \in V$ where V is the chosen function space.

$$\delta \hat{u}_i(t) = \int_{\Omega_i} (u(x) - u_{eq}(x)) \phi_i(x)$$

$$\frac{\partial \hat{u}_i}{\partial t} = \int_{\Omega_i} F[u_{eq}(x) + \sum_{i=0}^n \delta \hat{u}_i(t)] \frac{\partial \phi_i(x)}{\partial x} dx - \int F[u_{eq}(x) + \sum_{i=0}^n \delta \hat{u}_i(t)] \phi_i(x) \cdot \hat{n} dS + \int_{\Omega_i} S[u_{eq}(x) + \sum_{i=0}^n \delta \hat{u}_i(t)] \phi_i(x) dx$$

$$\frac{\partial u}{\partial t} = \frac{\partial \delta u}{\partial t} + \frac{\partial u_{eq}}{\partial t}$$

Assuming equilibrium solution, the following should hold: $\frac{\partial u_{eq}}{\partial t} = 0$

Pseudo code:

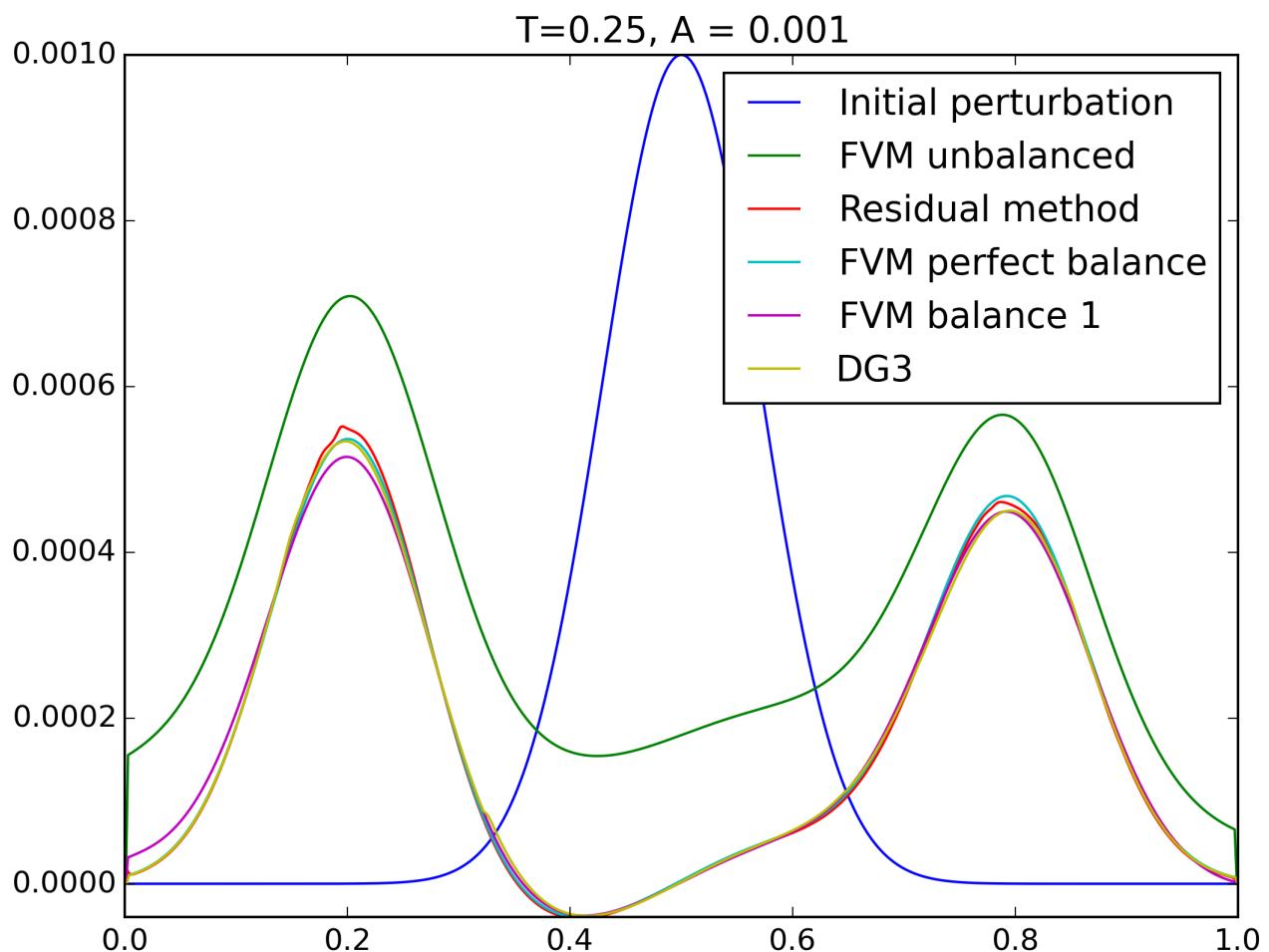
```
I
while t<T:

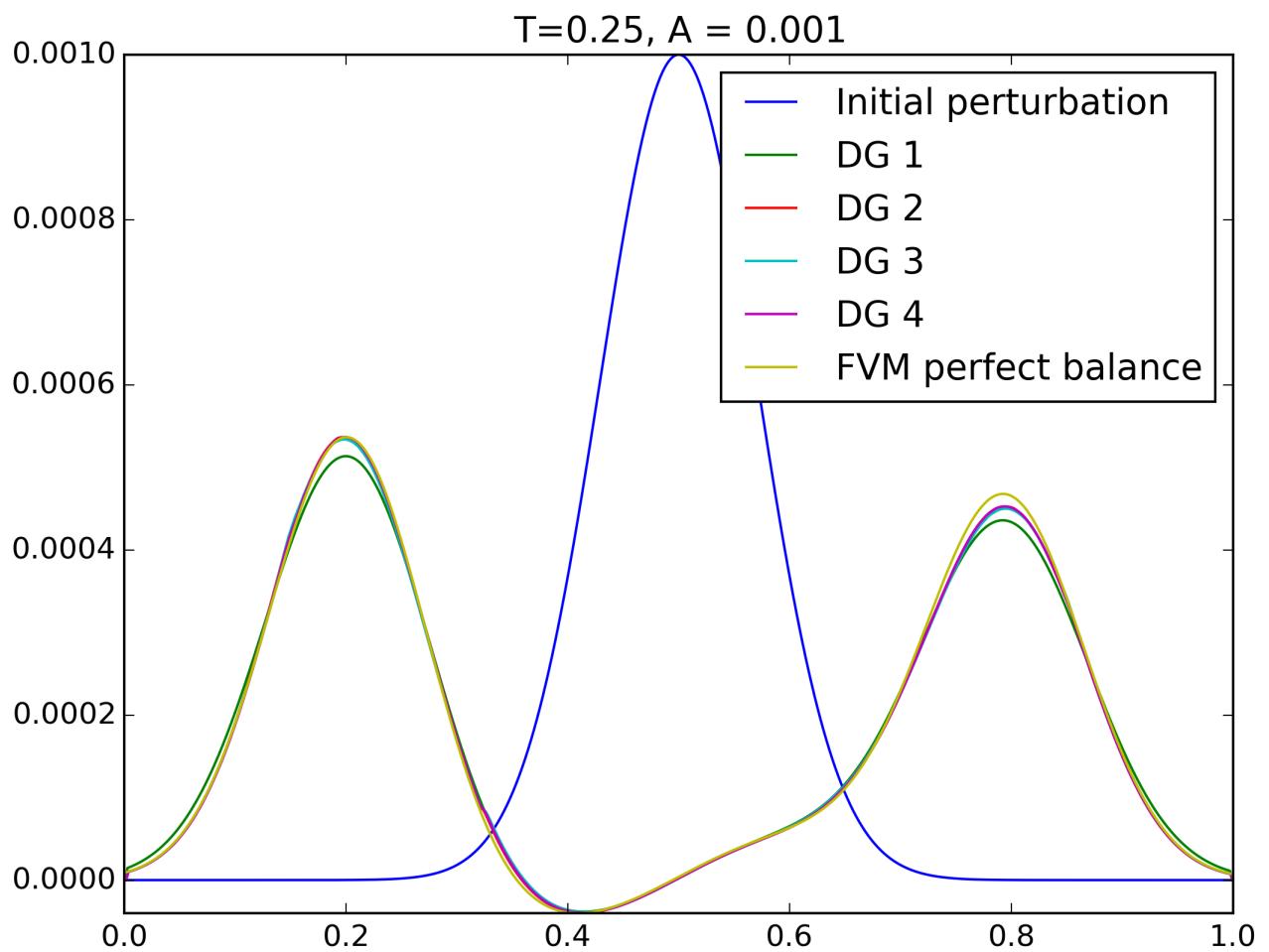
    w1 = dudt(delta_u)
    delta_u_new = rgkt(w1)

    delta_u_new_modes = get_modes(u_new)
    u_new = u + delta_u_new_modes # we are updating the deltas

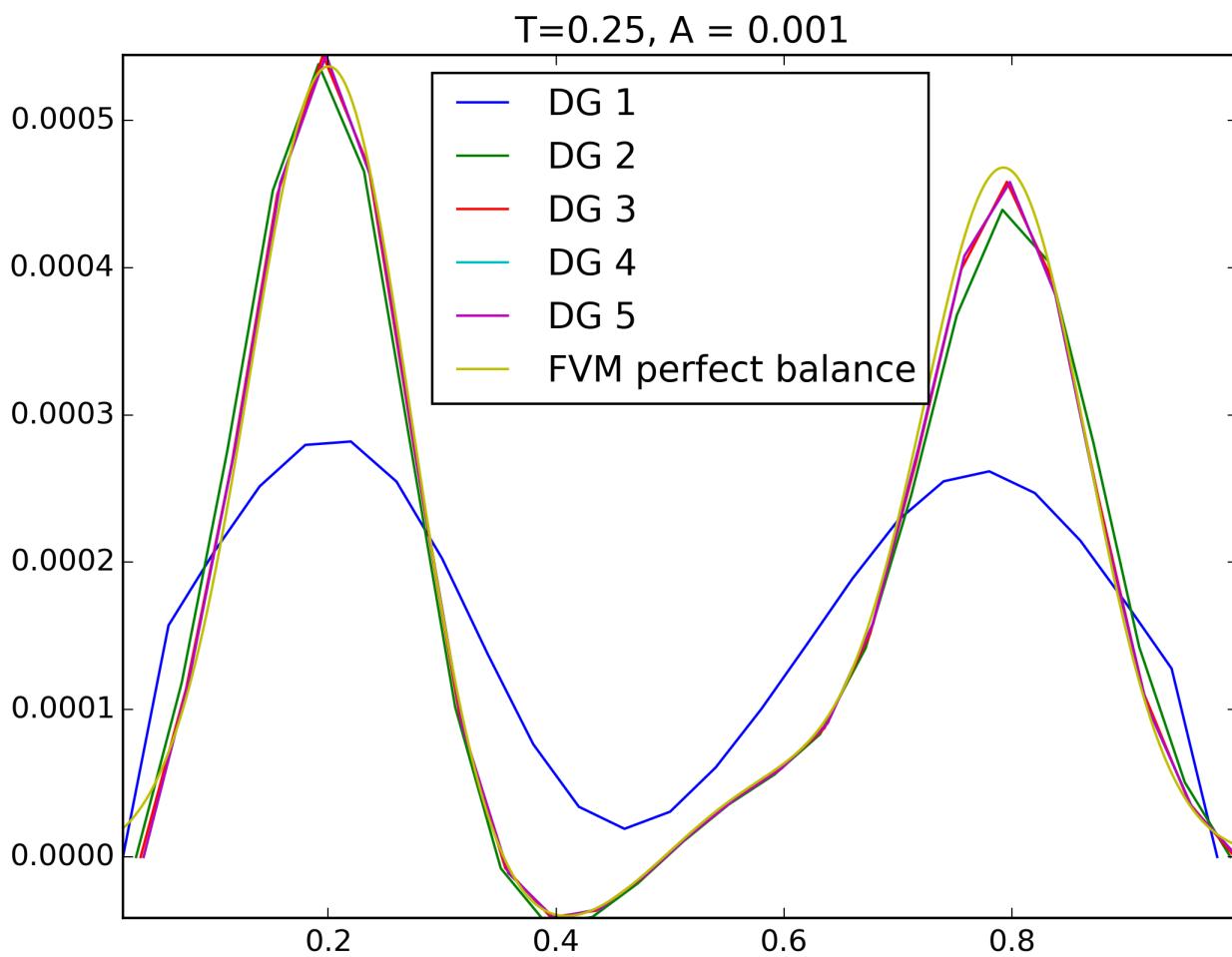
    u = u_new
    delta_u = w1
```

New plots:

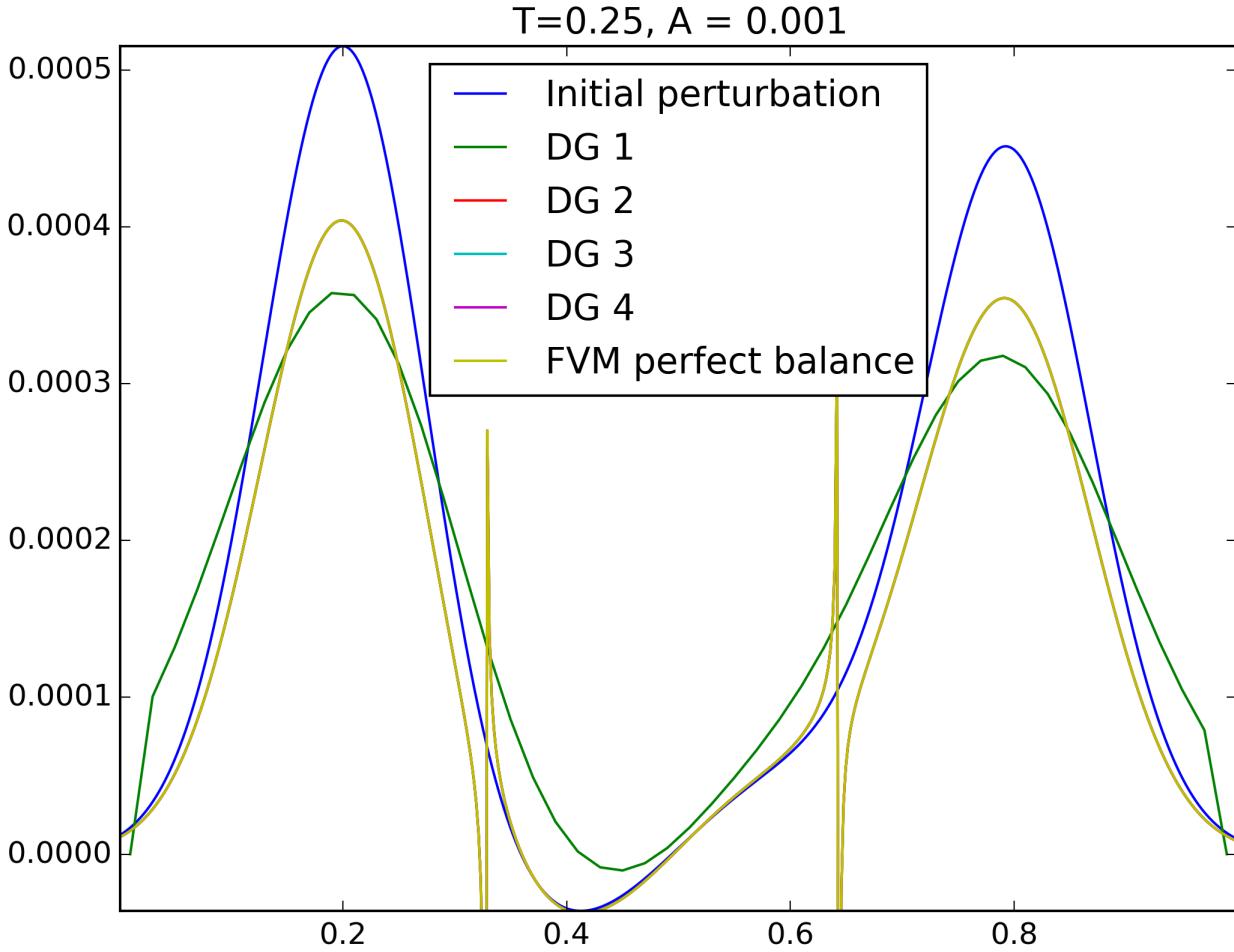




$N = 25$ versus 8126 as perfect balance.



If I plot $\text{const2prim}(\delta\tau)$ instead of $\text{const2prim}(u_{\text{eq}} + \delta\tau - \delta\tau)$



!!! There is something fishy with the normalization.

DG 2d Implementation

We look at how to solve the 2D version of a hyperbolic PDE, given as:

$$\frac{\partial u}{\partial t} + \nabla \cdot F(u) = 0$$

For $u(x, y, t) : \mathbb{R} \times \mathbb{R} \times \mathbb{R}^+ \rightarrow \mathbb{R}$.

We have 2 terms to approximate:

$$\frac{\partial u_{lk}}{\partial t} = - \int_{\partial l} (F v_{lk}) \cdot \hat{n} dl + \int_{\Omega_c} F \cdot \nabla v_{lk} dx dy$$

Let us treat the boundary terms first:

$$\int_{\partial l} F(x_i, y_i) v_{lk}(x_i, y_i) \cdot \hat{n} dl$$

We are in a cartesian grid, so dealing with the positive y-direction first, we have $\hat{n} = (0, 1)$:

$$\int_{y_i}^{y_{i+1}} F_2(x_i, y_i) v_{lk}(x_i, y_i) dy$$

Using a gaussian quadrature and noting that $v_{lk}(x, y) = \phi_l(x) \times \phi_k(y)$ which is given by the tensor product of 1D basis:

$$\int_{y_i}^{y_{i+1}} F_2(x, y) \phi_l(x) \times \phi_k(y) dy \approx \sum_{j=1}^{N_y} F_2(x_i, y_j) \phi_l(x_i) \phi_k(y_j) w(j)$$

Timestepping:

$$\Delta t \leq |a|^{-1} \frac{CFL \sqrt{\Delta x \Delta y}}{(2m_x + 1)(2m_y + 1)}$$

CFL = 0.5

This is a pretty restricted timestep...

Transverse velocity problem - Comes from having a velocity transverse to the shock front - "cross flow instability". Is it because riemann solver is not truly solving a 2d problem. This comes from solution containing discontinuities and the corner points from the squares?

Is it true that for a smooth solution it shouldn't matter.

For the initial conditions:

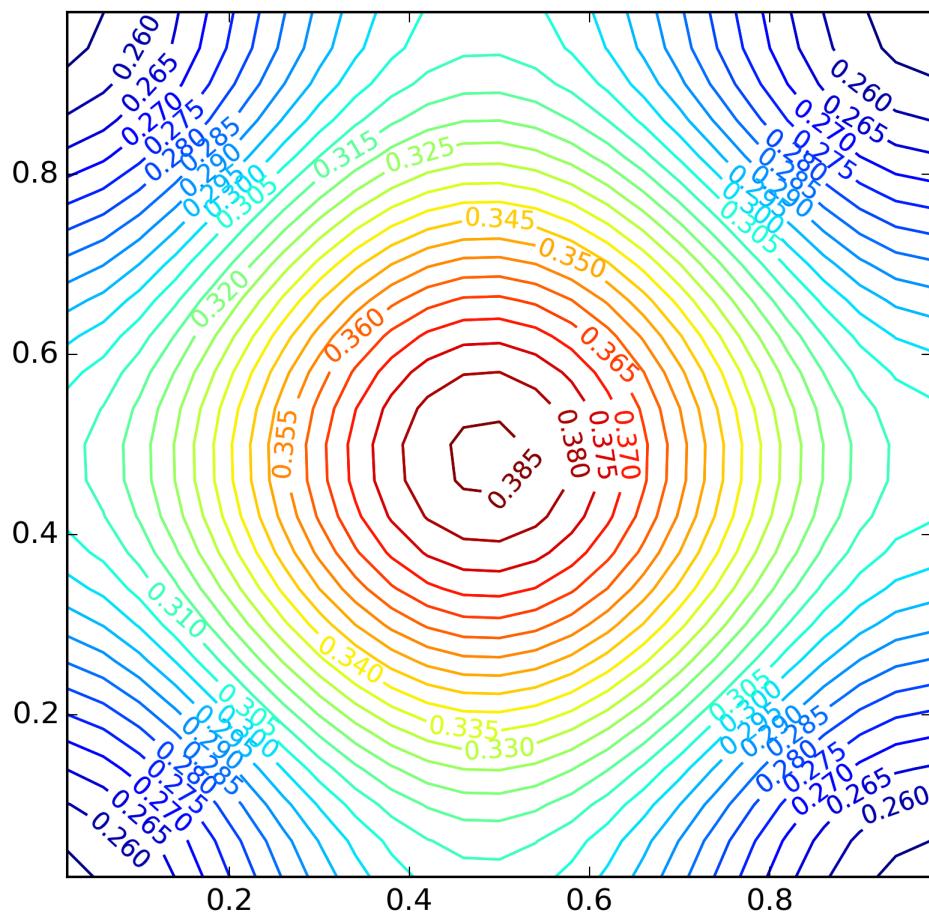
$$\rho = \exp(-40 * (x - 0.5)^2 + (y - 0.5)^2)$$

$$u = 1$$

$$v = 1$$

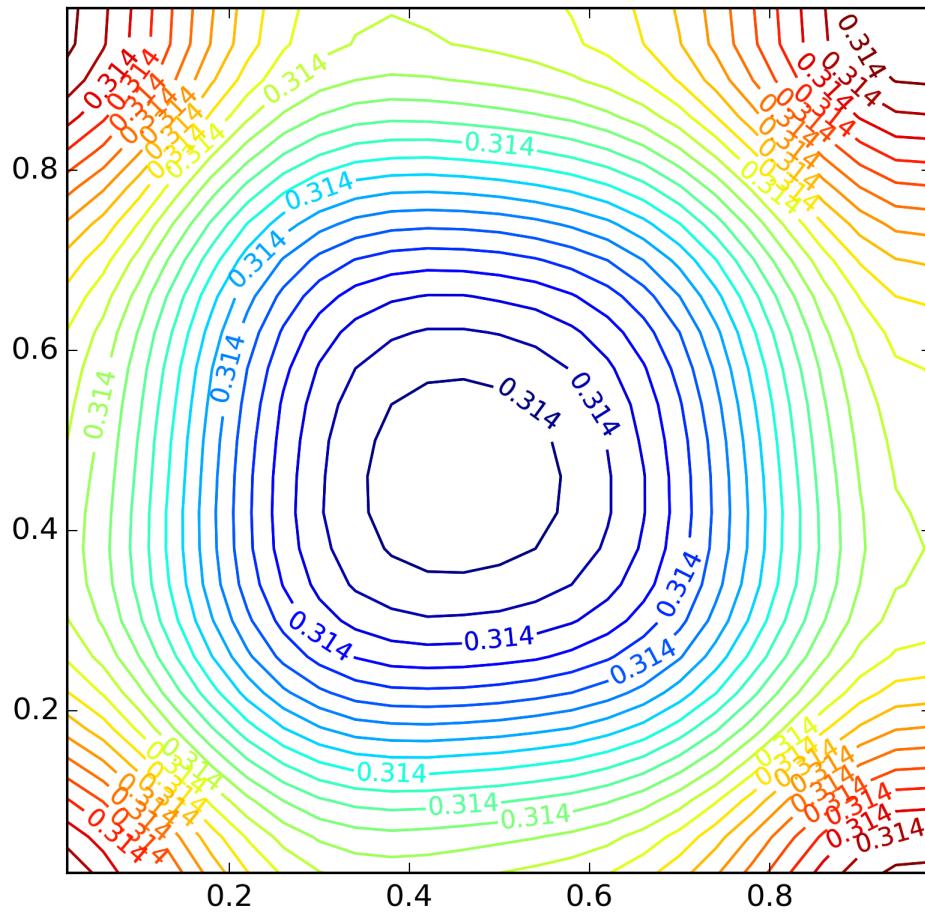
$$p = \exp(-40 * (x - 0.5)^2 + (y - 0.5)^2)$$

t = 1.



- First order - Diffusion and deforms the pulse....

$t = 2$



Smooth IC and pure advection

Well, turns out the initial conditions were stupid, and this wouldn't lead to a linear advection (look at the pressure). So in reality, for linear advection with the Euler system, we want the following IC:

$$\rho = \exp(-40 * (x - 0.5)^2 + (y - 0.5)^2)$$

$$u = 1$$

$$v = 1$$

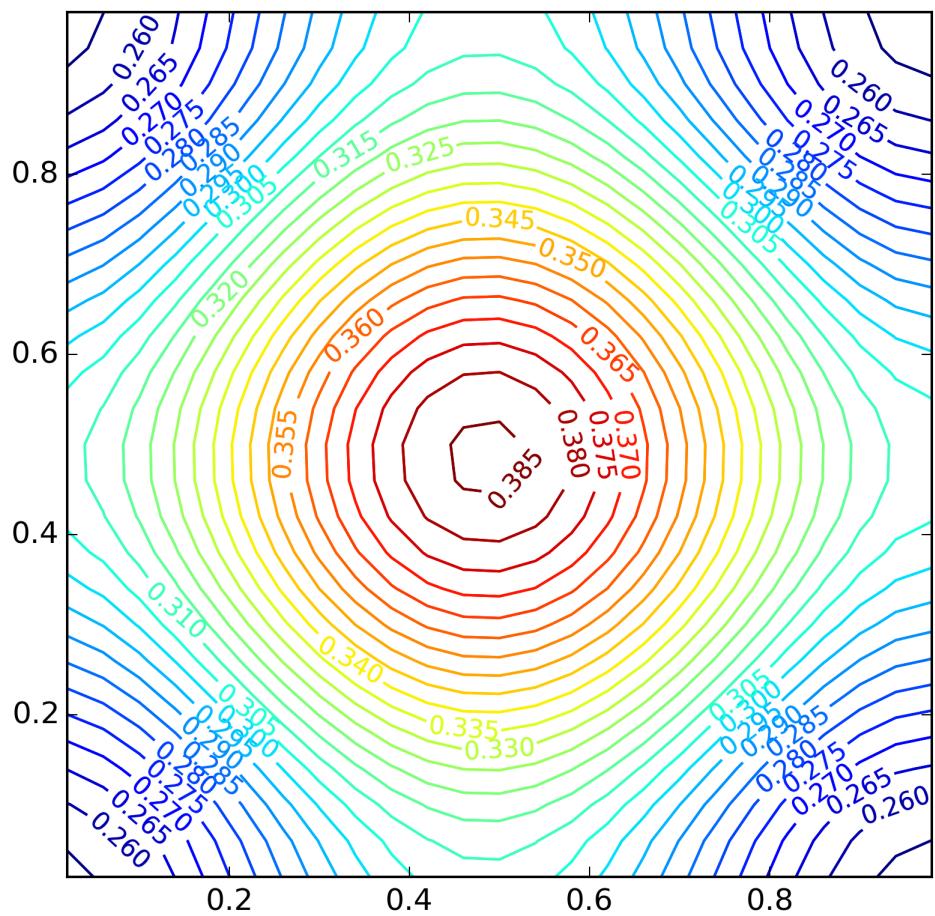
$$p = 1$$

(write the PDE and you can see it!!)

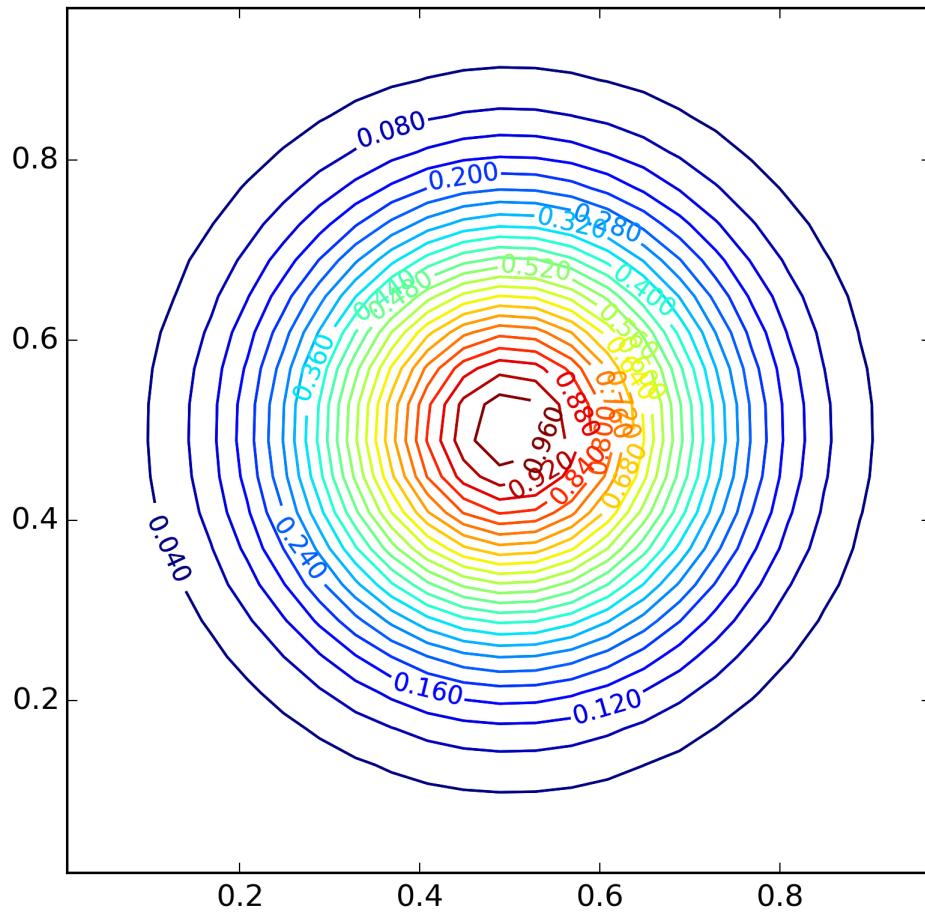
So, testing for higher order:

Today, the 26th of April, we have a working 2D DG solver for any order*. Oh my god! (that people are willing to implement)

$\circ = 1$



$\alpha = 2$



Limiter

Moving onto non smooth initial conditions, we need to include a limiter.

2d limiting

The idea is the same as 1d, but there are some small considerations.

Shu limiter

Details are given here: <http://lsec.cc.ac.cn/lcf/LCFD/DEWENO/paper/cs5.pdf>

To compute u_h , we rely on the assumption that spurious oscillations are present in u_h , only if they are present in its u_h^1 part, which is its L_2 -projection into the space of piecewise linear functions V_h^1 . A theoretical justification of this assumption is still an open problem.

What is implemented in the code is the following. We approximate the solution as

$$u = \bar{u} + u_x * \Phi_1(x, y) + u_y * \Phi_2(x, y) + h.o$$

The intuition is that if the "derivatives", essentially the 2nd and 3rd coefficients (which correspond to the linear

approximation of the solution), are *wiggly*, then the h.o. stuff is also wiggly. This argument has no real theoretical backing but somehow seems to work.

So we replace u_x and u_y by

$$\minmod(u_x, \bar{u}_{i+1,j} - \bar{u}_{i,j}, \bar{u}_{i,j} - \bar{u}_{i-1,j})$$

$$\minmod(u_y, \bar{u}_{i,j+1} - \bar{u}_{i,j}, \bar{u}_{i,j} - \bar{u}_{i,j-1})$$

respectively.

And if $u_x \neq \minmod(u_x, \bar{u}_{i+1,j} - \bar{u}_{i,j}, \bar{u}_{i,j} - \bar{u}_{i-1,j})$, we set the h.o. terms to zero. Effectively, worst case scenario, we revert to 1st order approximation.

This clips extrema.

2-D riemann problem (case 3)

Configuration 1:

The initial data are

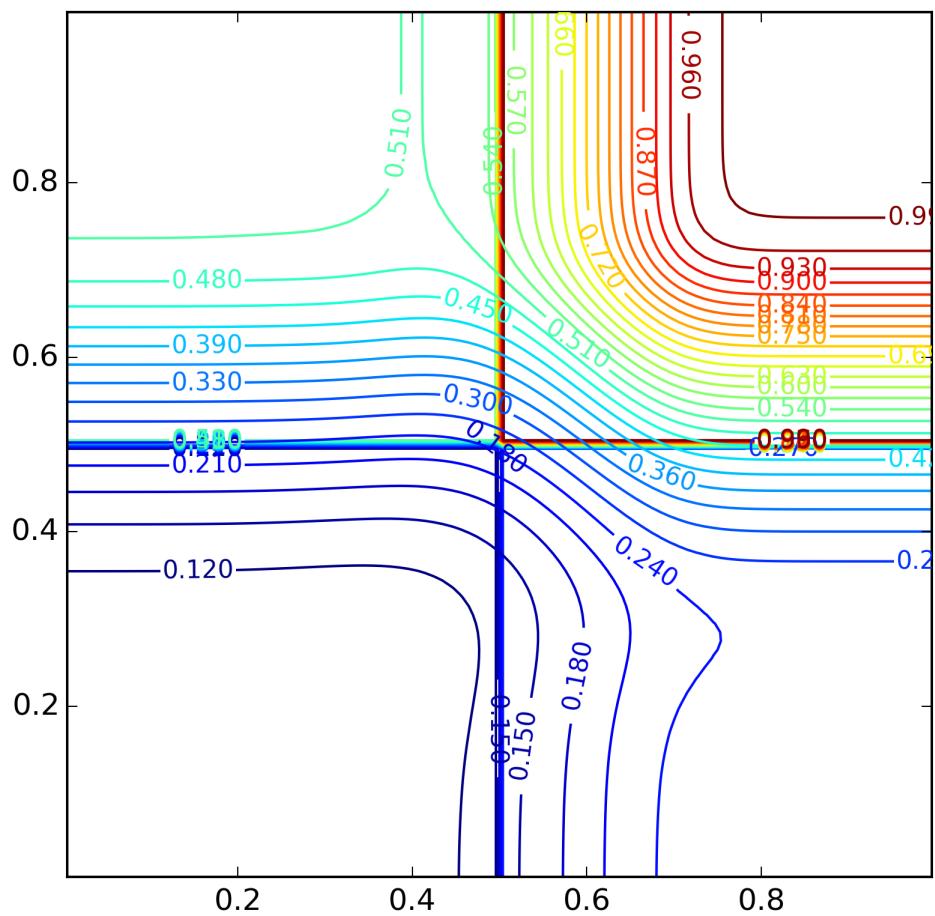
$$\begin{array}{llll} p_2 = 0.4 & \rho_2 = 0.5197 & p_1 = 1 & \rho_1 = 1 \\ u_2 = -0.7259 & v_2 = 0 & u_1 = 0 & v_1 = 0 \end{array}$$

$$\begin{array}{llll} p_3 = 0.0439 & \rho_3 = 0.1072 & p_4 = 0.15 & \rho_4 = 0.2579 \\ u_3 = -0.7259 & v_3 = -1.4045 & u_4 = 0 & v_4 = -1.4045 \end{array}$$

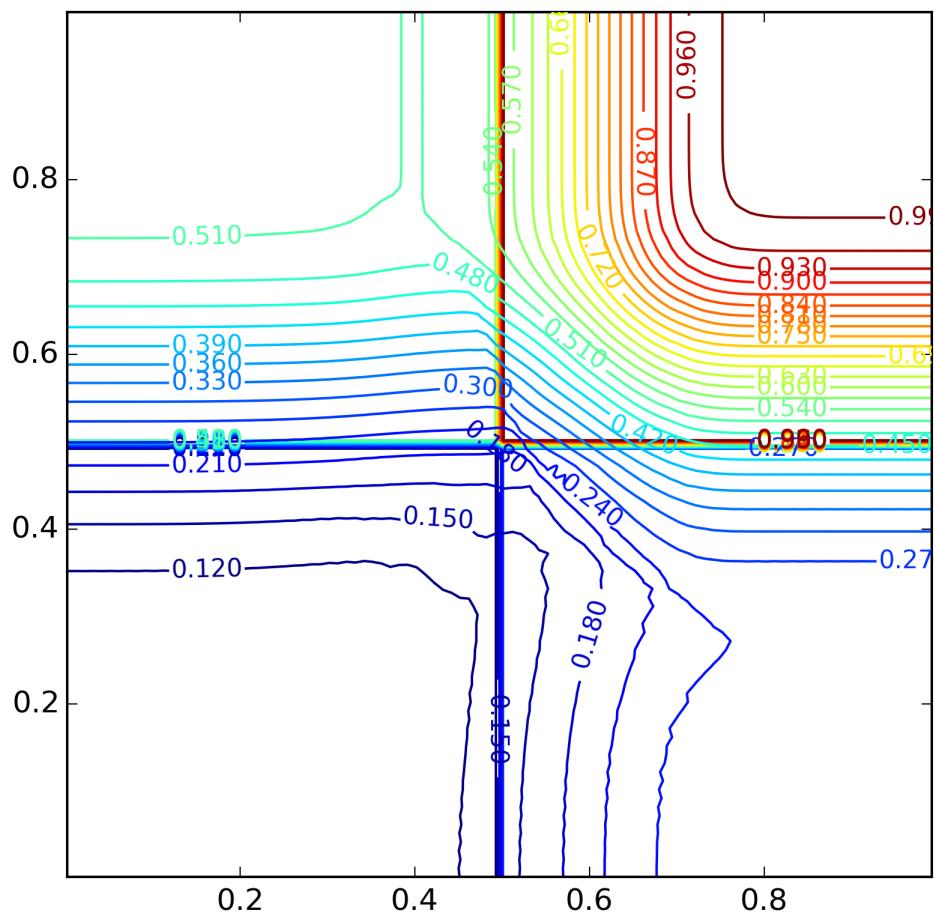
(Rarefaction waves)

100x100, t=0.3

order = 1

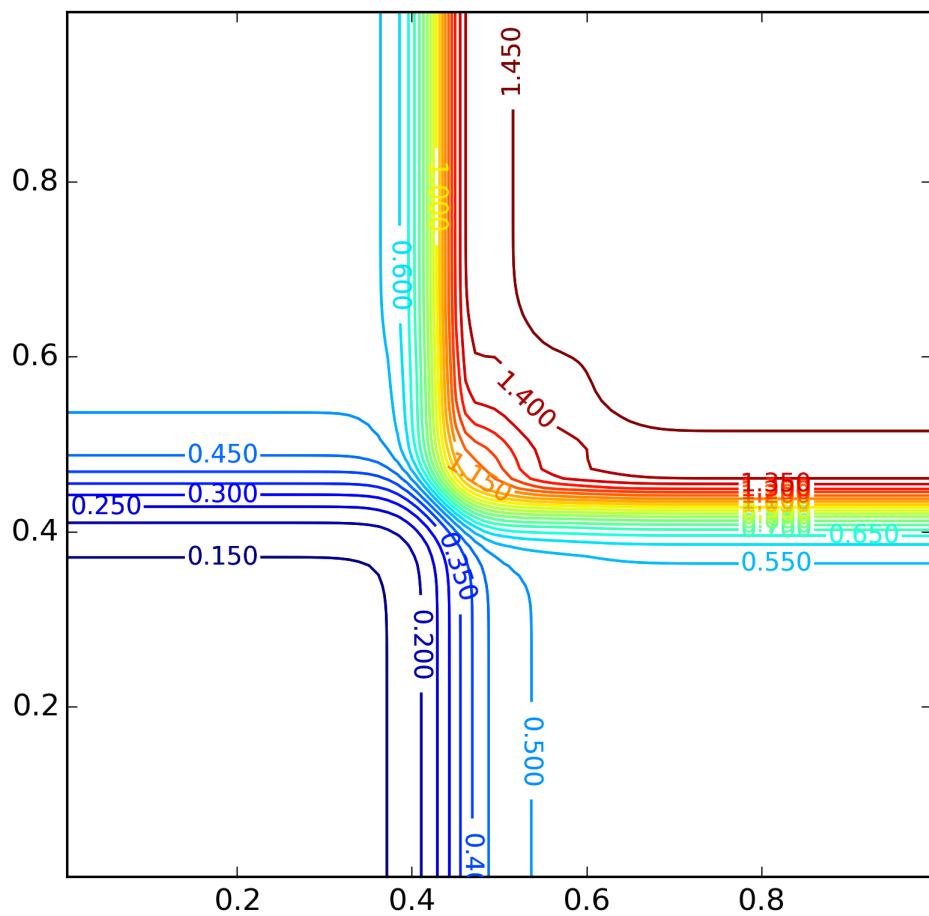


order = 2

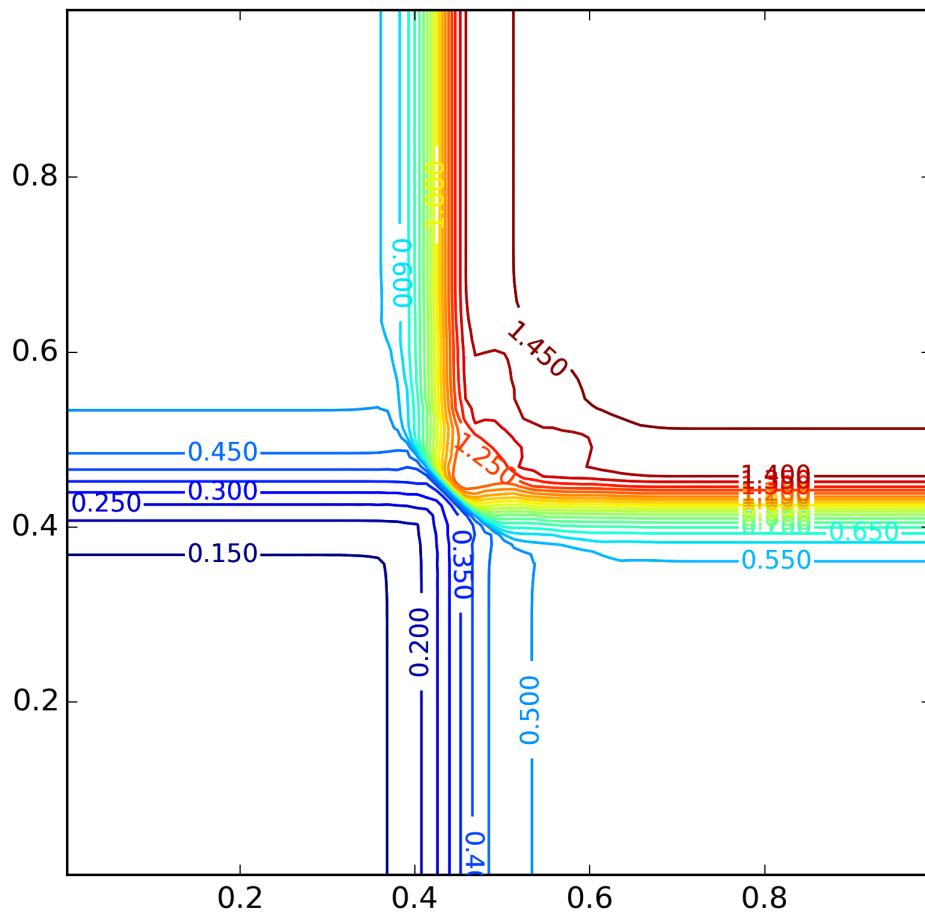


Configuration 2 (Shocks)

100x100, $t=0.2$ order = 1



order = 2



to do

- limit in characteristics - bit of oscillations
- Krivodonova limiter
- hydrodynamical implosion

Issues

- cfl
- interesting ics
- residual methods uses triangular grid :(