

# PHY3110 SP23 HW11

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**Problem 1.** Verify the Jacobi identity for the Poisson brackets.

*Solution.* Let  $A_q = \partial_q A$ , then

$$\begin{aligned}\{A, \{B, C\}\} &= \{A, B_{q_i} C_{p_i} - B_{p_i} C_{q_i}\} \\ &= A_{q_j} (B_{q_i p_j} C_{p_i} + B_{q_i} C_{p_i p_j} - B_{p_i p_j} C_{q_i} - B_{p_i} C_{q_i p_j}) - A_{p_j} (B_{q_i q_j} C_{p_i} + B_{q_i} C_{p_i q_j} - B_{p_i q_j} C_{q_i} - B_{p_i} C_{q_i q_j})\end{aligned}$$

Hence  $J = \{A, \{B, C\}\} + \{B, \{C, A\}\} + \{C, \{A, B\}\}$  is

$$\begin{aligned}J &= A_{q_j} (B_{q_i p_j} C_{p_i} + B_{q_i} C_{p_i p_j} - B_{p_i p_j} C_{q_i} - B_{p_i} C_{q_i p_j}) - A_{p_j} (B_{q_i q_j} C_{p_i} + B_{q_i} C_{p_i q_j} - B_{p_i q_j} C_{q_i} - B_{p_i} C_{q_i q_j}) \\ &\quad + B_{q_j} (C_{q_i p_j} A_{p_i} + C_{q_i} A_{p_i p_j} - C_{p_i p_j} A_{q_i} - C_{p_i} A_{q_i p_j}) - B_{p_j} (C_{q_i q_j} A_{p_i} + C_{q_i} A_{p_i q_j} - C_{p_i q_j} A_{q_i} - C_{p_i} A_{q_i q_j}) \\ &\quad + C_{q_j} (A_{q_i p_j} B_{p_i} + A_{q_i} B_{p_i p_j} - A_{p_i p_j} B_{q_i} - A_{p_i} B_{q_i p_j}) - C_{p_j} (A_{q_i q_j} B_{p_i} + A_{q_i} B_{p_i q_j} - A_{p_i q_j} B_{q_i} - A_{p_i} B_{q_i q_j}) \\ &= 0\end{aligned}$$

**Problem 2.** Show by the use of Poisson brackets that for a one-dimensional harmonic oscillator there is a constant of motion  $u$  defined as

$$u(q, p, t) = \ln(p + im\omega q) - i\omega t, \quad \omega = \sqrt{\frac{k}{m}} \quad (1)$$

*Solution.* The Hamiltonian of a one-dimensional harmonic oscillator is

$$H(q, p, t) = \frac{1}{2m}(p^2 + m^2\omega^2 q^2)$$

Hence

$$\begin{aligned}\frac{du}{dt} &= \{u, H\} + \frac{\partial u}{\partial t} \\ &= \frac{im\omega}{p + im\omega q} \frac{p}{q} - \frac{m\omega^2 q}{p + im\omega q} - i\omega \\ &= \frac{1}{p + im\omega q} (i\omega p - m\omega^2 q + m\omega^2 q - i\omega p) \\ &= 0\end{aligned}$$

**Problem 3.** Show that the following transformation is canonical ( $\alpha$  is a fixed parameter):

$$\begin{aligned}x &= \frac{1}{\alpha}(\sqrt{2P_1} \sin Q_1 + P_2), & p_x &= \frac{\alpha}{2}(\sqrt{2P_1} \cos Q_1 - Q_2) \\ y &= \frac{1}{\alpha}(\sqrt{2P_1} \cos Q_1 + Q_2), & p_y &= -\frac{\alpha}{2}(\sqrt{2P_1} \sin Q_1 - P_2)\end{aligned} \quad (2)$$

Apply this transformation to the problem of a particle of charge  $q$  moving in a plane that is perpendicular to a constant magnetic field  $\mathbf{B}$ . Express the Hamiltonian for this problem in the  $(Q_i, P_i)$  coordinates letting the parameter  $\alpha$  take the form

$$\alpha^2 = \frac{qB}{c} \quad (3)$$

From this Hamiltonian, obtain the motion of the particle as a function of time.

*Solution.* To show the transformation is a canonical transformation, consider the Jacobian  $\mathbf{J}$

$$\begin{aligned}\mathbf{J} &= \frac{\partial[\mathbf{Q} \ \mathbf{P}]^T}{\partial[\mathbf{q} \ \mathbf{p}]^T} \\ &= \begin{bmatrix} (2P_1)^{0.5} \cos Q_1/\alpha & 0 & (2P_1)^{-0.5} \sin Q_1/\alpha & 1/\alpha \\ -(2P_1)^{0.5} \sin Q_1/\alpha & 1/\alpha & (2P_1)^{-0.5} \cos Q_1/\alpha & 0 \\ -\alpha(2P_1)^{0.5} \sin Q_1 & -\alpha/2 & \alpha(2P_1)^{-0.5} \cos Q_1/2 & 0 \\ -\alpha(2P_1)^{0.5} \cos Q_1 & 0 & \alpha(2P_1)^{-0.5} \sin Q_1/2 & \alpha/2 \end{bmatrix}\end{aligned}$$

We can show that

$$\mathbf{J} \begin{bmatrix} & -\mathbb{1} \\ \mathbb{1} & \end{bmatrix} \mathbf{J}^T = \begin{bmatrix} & -\mathbb{1} \\ \mathbb{1} & \end{bmatrix}$$

Hence, it is a canonical transformation.

To simplify calculation, let  $c = 1$ , then  $\alpha^2 = qB$ . Hence

$$L = \frac{1}{2}m\mathbf{v}^2 - q\phi + q\mathbf{A} \cdot \mathbf{v}, \quad H = \mathbf{v} \cdot \mathbf{p} - L = \frac{1}{2}m\mathbf{v}^2 = \frac{1}{2m}(\mathbf{p} - q\mathbf{A})^2$$

Since  $\mathbf{B} = [0 \ 0 \ B]^T$ , there are multiple choices of  $\mathbf{A}$ , for example

$$\begin{bmatrix} 0 \\ Bx \\ 0 \end{bmatrix}, \quad \begin{bmatrix} -By \\ 0 \\ 0 \end{bmatrix}, \quad \begin{bmatrix} -By/2 \\ Bx/2 \\ 0 \end{bmatrix}$$

I will adopt the last one to simplify the calculation. Suppose we have  $A_x = -By/2$ ,  $A_y = Bx/2$ ,  $A_z = 0$ , then

$$\begin{aligned}H &= \frac{1}{2m}[(p_x - qA_x)^2 + (p_y - qA_y)^2] \\ &= \frac{1}{2m}(\alpha^2 2P_1 \sin^2 Q_1 + \alpha^2 2P_1 \cos^2 Q_1) \\ &= \frac{qB}{m}P_1\end{aligned}$$

In this case, the solution of EOM is

$$\begin{aligned}Q_1 &= \frac{qB}{m}t + C_1 \\ Q_2 &= C_2 \\ P_1 &= C_3 \\ P_2 &= C_4\end{aligned}$$

**Problem 4.** Use the method of infinitesimal canonical transformation to solve the motion of a one-dimensional harmonic oscillator as a function of time.

*Solution.* Let  $G = H$ ,  $U = q$ , then

$$\begin{aligned}U' &= [q, H] = \frac{p}{m} \\ U'' &= [p/m, H] = -\omega^2 q \\ U''' &= -\omega^2 \frac{p}{m}\end{aligned}$$

$$U^{(4)} = \omega^4 q$$

Therefore

$$\begin{aligned} q(t) &= q_0 + \frac{p_0}{m}t - \frac{1}{2!}\omega^2 q_0 t^2 - \frac{1}{3!}\omega^3 \frac{p_0}{m}t^3 + \frac{1}{4!}\omega^4 q_0 t^4 + \dots \\ &= q_0 \cos \omega t + \frac{p_0}{m\omega} \sin \omega t \end{aligned}$$

Follow the same way, we can also get  $p(t)$

$$p(t) = -m\omega q_0 \sin \omega t + p_0 \cos \omega t$$

where  $q_0 = q(0)$  and  $p_0 = p(0)$ .