

L10: Constraint Force Computation, Control

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Agenda

- Constrained Lagrangian Method
- Variational Method
- Control
- PID Control

click to jump to the section.

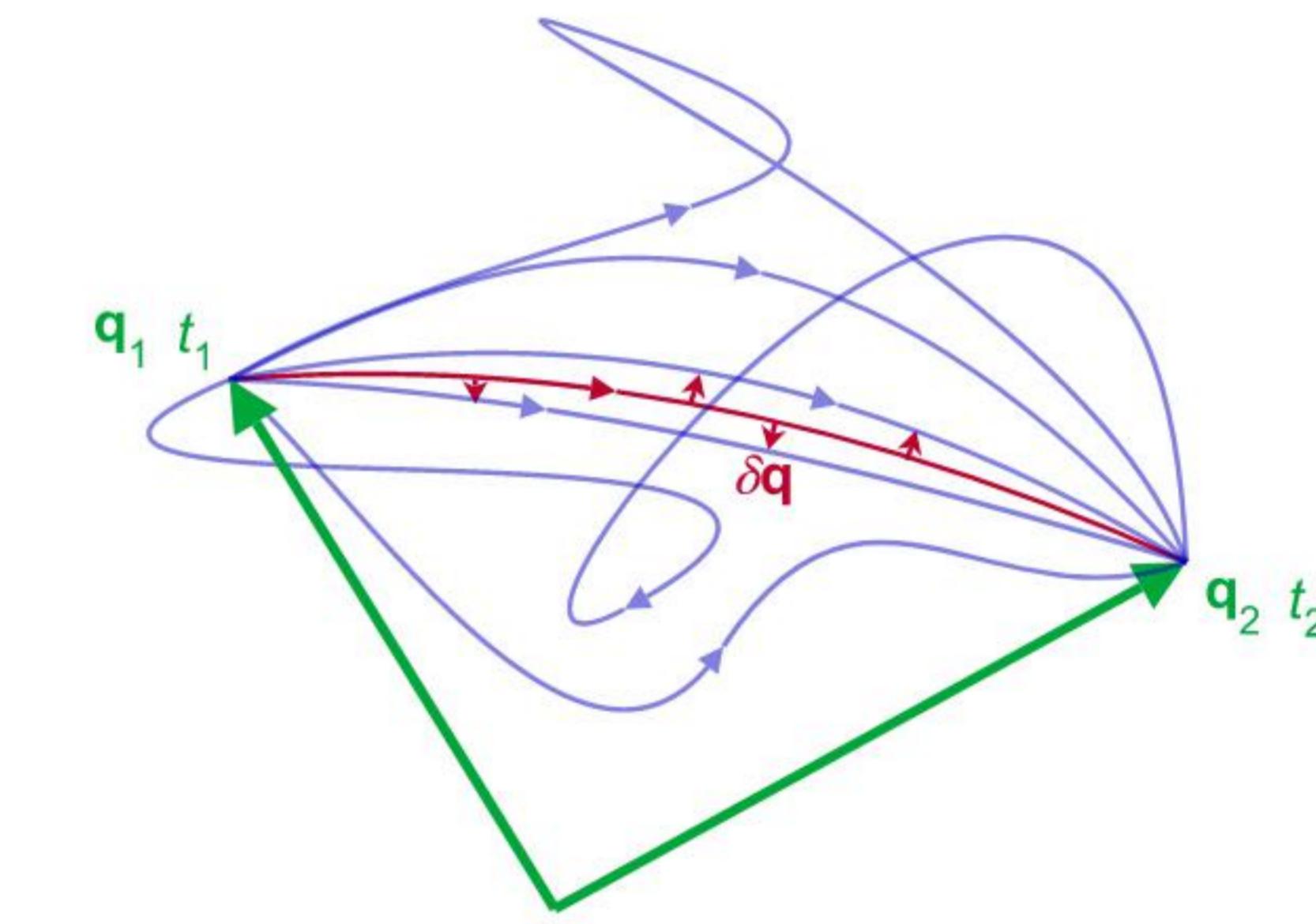
Review: Lagrangian Function

- Let $q \in \mathbb{R}^n$ be the generalized coordinates.
- **Lagrangian function:** $L(q, \dot{q}) = T(q, \dot{q}) - V(q)$
 - $T(q, \dot{q})$: kinetic energy of system
 - $V(q)$: potential energy (given by some conservative force, e.g., gravity, electric field force)

Review: The Principle of Stationary Action

- Given a pair of time instants, t_1 and t_2
- What is the curve $\mathbf{q} : [t_1, t_2] \rightarrow \mathcal{C}$ in the configuration space \mathcal{C} ?
- **Action** is defined to be a functional of $\mathbf{q}(t)$:

$$S[\mathbf{q}] = \int_{t_1}^{t_2} L(q, \dot{q}) dt = \int_{t_1}^{t_2} [T(\mathbf{q}, \dot{\mathbf{q}}) - V(\mathbf{q})] dt$$



Wikipedia: Stationary Action Principle

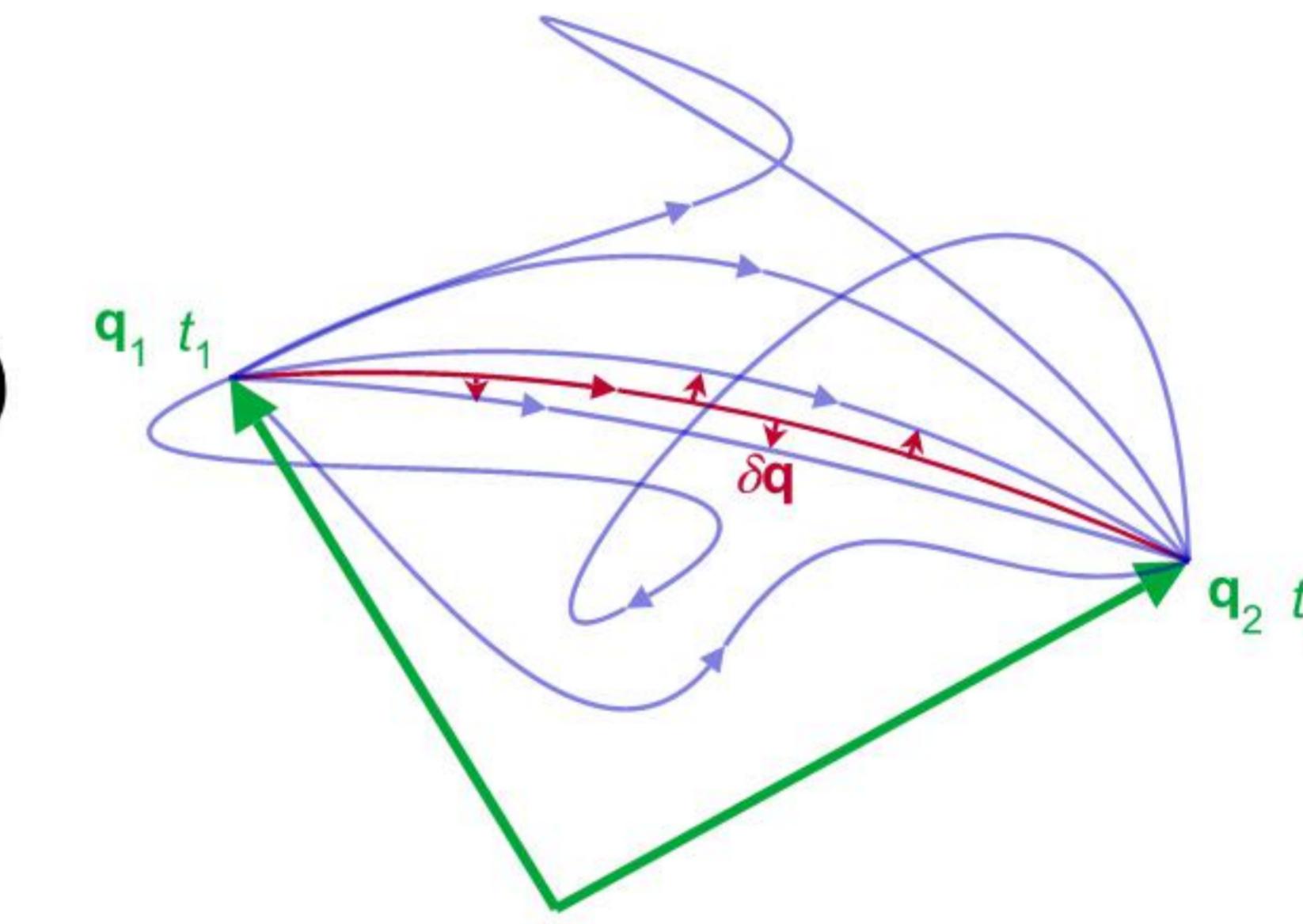
Review: The Principle of Stationary Action

- The actual curve $\mathbf{q}(t)$ is a stationary point of the $S[\mathbf{q}]$:

$$\forall \boldsymbol{\delta} : [t_1, t_2] \rightarrow \mathcal{C}, \quad \lim_{\epsilon \rightarrow 0} \frac{1}{\epsilon} (\mathcal{S}[\mathbf{q} + \epsilon \boldsymbol{\delta}] - \mathcal{S}[\mathbf{q}]) = \mathbf{0} \quad (1)$$

- Note: Treating \mathbf{q} as a variable, and (1) is an extension of the first-order optimality condition that we use in calculus:

$$\nabla_{\mathbf{q}} S[\mathbf{q}] = \mathbf{0}$$



Review: The Principle of Stationary Action

- The actual curve $\mathbf{q}(t)$ is a stationary point of the $S[\mathbf{q}]$:

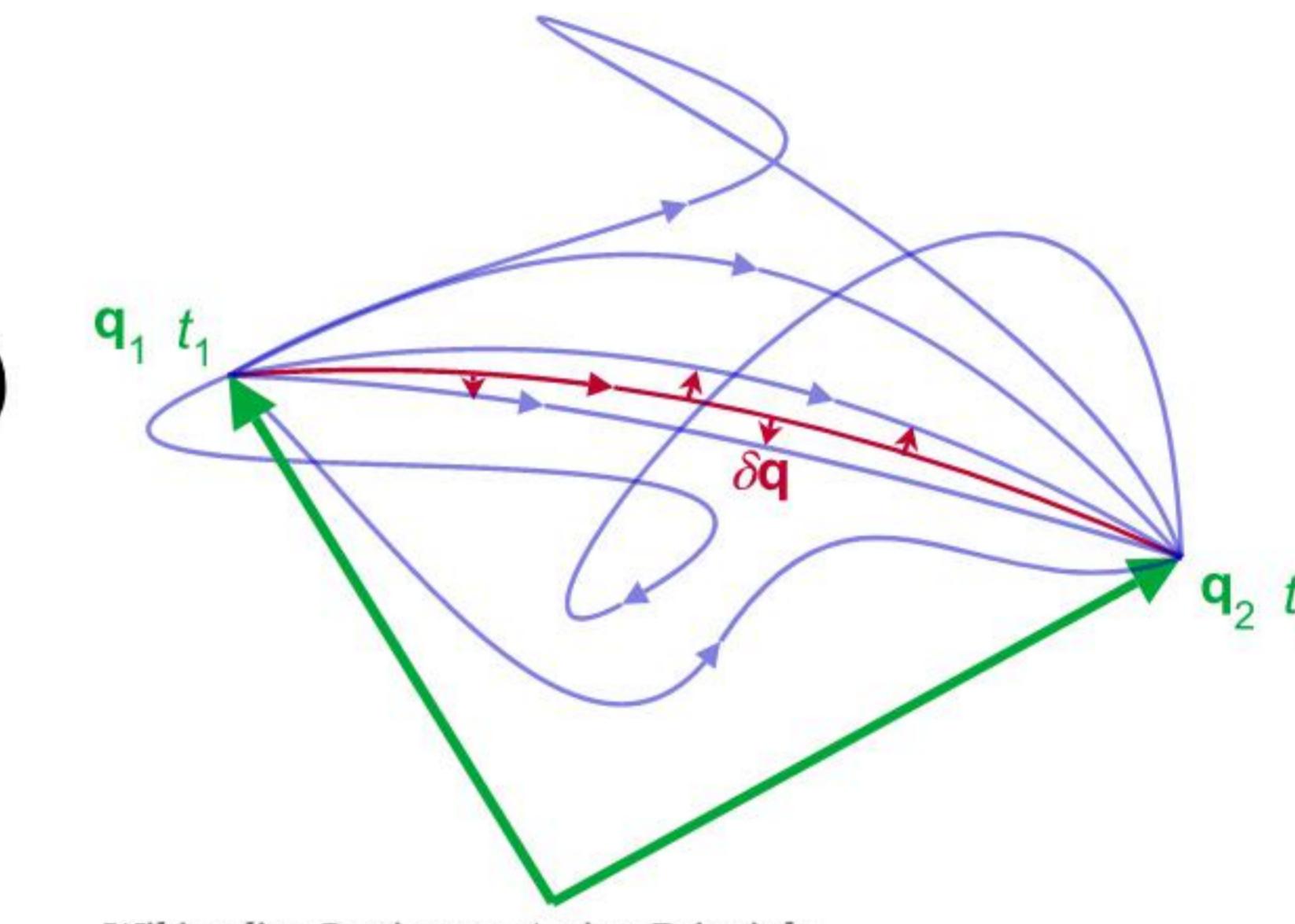
$$\forall \boldsymbol{\delta} : [t_1, t_2] \rightarrow \mathcal{C}, \quad \lim_{\epsilon \rightarrow 0} \frac{1}{\epsilon} (\mathcal{S}[\mathbf{q} + \epsilon \boldsymbol{\delta}] - \mathcal{S}[\mathbf{q}]) = \mathbf{0} \quad (1)$$

- Note: Treating \mathbf{q} as a variable, and (1) is an extension of the first-order optimality condition that we use in calculus:

$$\nabla_{\mathbf{q}} S[\mathbf{q}] = \mathbf{0}$$

- Using *variational method*, condition (1) becomes

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q} = 0$$



Review: Euler-Lagrange Equation

- When there are external non-conservative generalized force $\mathbf{F} \in \mathbb{R}^n$ added to the system (e.g., torque at robot arm joints), we have the following Euler-Lagrange equation:

$$\mathbf{F} = \frac{d}{dt} \frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q} \quad (\text{Euler-Lagrange Equation})$$

Example: Robot Arm

Robot Arm

- For kinematic chains with n joints, it is convenient and always possible to choose the joint angles $\theta = (\theta_1, \dots, \theta_n)$ and the joint torques $\tau = (\tau_1, \dots, \tau_n)$ as the generalized coordinates and generalized forces, respectively.
 - If joint i is revolute: θ_i joint angle and τ_i is joint torque
 - If joint i is prismatic: θ_i joint position and τ_i is joint force
 - Then, the joint force/torque τ and joint angular/linear velocity $\dot{\theta}$ form dual pairs.
- Lagrangian Equations:

$$\tau_i = \frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}_i} - \frac{\partial L}{\partial \theta_i}$$

Review of Notations for Robot Arms

For each link $i \in \{1, \dots, n\}$, the corresponding body-frame \mathcal{F}_i is attached to the center of mass of link i . All the following quantities are expressed in \mathcal{F}_i

- $\boldsymbol{\xi}_i^b$: body twist of the link frame \mathcal{F}_i
- $\mathfrak{M}_i^b = \begin{bmatrix} m_i \text{Id}_{3 \times 3} & 0 \\ 0 & \mathbf{I}_i^b \end{bmatrix}$: body inertia matrix
- Kinetic energy of link i : $T_i = \frac{1}{2}(\boldsymbol{\xi}_i^b)^T \mathfrak{M}_i^b \boldsymbol{\xi}_i^b$
- $J_i^b \in \mathbb{R}^{6 \times n}$: body Jacobian of link i (so that we have the forward dynamics $\boldsymbol{\xi}_i^b = J_i^b \dot{\theta}$)

Kinetic and Potential Energies

- Total kinetic energy:

$$T(\theta, \dot{\theta}) = \frac{1}{2} \sum_{i=1}^n (\boldsymbol{\xi}_i^b)^T \mathcal{M}_i^b \boldsymbol{\xi}_i^b = \frac{1}{2} \dot{\theta}^T \underbrace{\left(\sum_{i=1}^n (J_i^b(\theta) \mathcal{M}_i^b J_i^b(\theta)) \right)}_{\mathbf{M}^b(\theta)} \dot{\theta} := \frac{1}{2} \dot{\theta}^T \mathbf{M}^b(\theta) \dot{\theta}$$

- Potential energy:

$$V(\theta) = \sum_{i=1}^n m_i g h_i(\theta)$$

- $h_i(\theta)$: height of center of mass of link i

Lagrangian Equation

- Plug $L = T - V$ into $F = \frac{d}{dt} \frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q}$, and we have
- $\tau_i = \frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}_i} - \frac{\partial L}{\partial \theta_i}$

$$\tau_i = \sum_{j=1}^n M_{ij}^b(\theta) \ddot{\theta}_j + \sum_{j=1}^n \sum_{k=1}^n \Gamma_{ijk}^b(\theta) \dot{\theta}_j \dot{\theta}_k + \frac{\partial V}{\partial \theta_i}$$

M_{ij}^b is the (i, j) -th entry of matrix \mathbf{M}^b

- $\Gamma_{ijk}^b(\theta)$ is called the **Christoffel symbols of the first kind**

$$\Gamma_{ijk}^b(\theta) = \frac{1}{2} \left(\frac{\partial M_{ij}^b}{\partial \theta_k} + \frac{\partial M_{ik}^b}{\partial \theta_j} - \frac{\partial M_{jk}^b}{\partial \theta_i} \right)$$

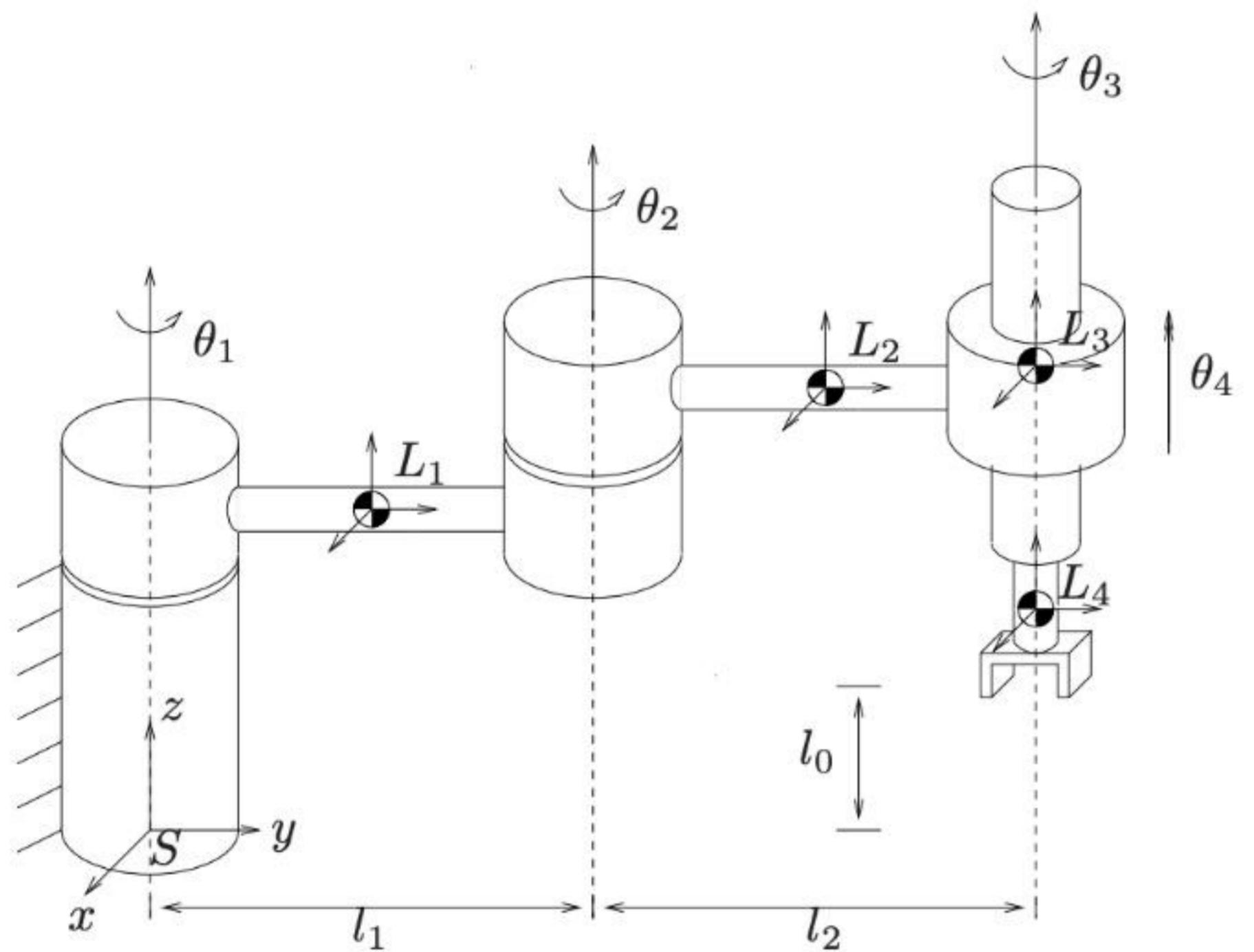
Lagrangian Equation

- Lagrangian equation in vector form:

$$\tau = \mathbf{M}^b(\theta)\ddot{\theta} + C^b(\theta, \dot{\theta})\dot{\theta} + g^b(\theta)$$

- $C_{ij}^b(\theta, \dot{\theta}) := \sum_{k=1}^n \Gamma_{ijk}^b \dot{\theta}_k$ is called the **Coriolis matrix**
 - Recall that in the body-frame Newton Euler equation, we also have a Coriolis term that comes from the derivative of the rotational inertia ($\boldsymbol{\tau}^b = \mathbf{I}^b \dot{\boldsymbol{\omega}}^b + \boldsymbol{\omega}^b \times \mathbf{I}^b \boldsymbol{\omega}^b$). It was used to compensate for the rotational acceleration of the body frame.
 - This $C_{ij}^b(\theta, \dot{\theta})$ also comes from taking the derivative of \mathbf{M}^b w.r.t. θ . Because \mathbf{M}^b and $\boldsymbol{\xi}^b$ are described in the body frame in our derivation, we also need this Coriolis term to compensate for the movement of the body frame.
- $g^b(\theta)$ is due to gravity in our derivation. If there are other external forces (e.g., friction), it would also show up here.

- Equations for a simple arm



$$M_{11} = I_{y2}s_2^2 + I_{y3}s_{23}^2 + I_{z1} + I_{z2}c_2^2 + I_{z3}c_{23}^2 + m_2r_1^2c_2^2 + m_3(l_1c_2 + r_2c_{23})^2$$

$$M_{12} = 0$$

$$M_{13} = 0$$

$$M_{21} = 0$$

$$M_{22} = I_{x2} + I_{x3} + m_3l_1^2 + m_2r_1^2 + m_3r_2^2 + 2m_3l_1r_2c_3$$

$$M_{23} = I_{x3} + m_3r_2^2 + m_3l_1r_2c_3$$

$$M_{31} = 0$$

$$M_{32} = I_{x3} + m_3r_2^2 + m_3l_1r_2c_3$$

$$M_{33} = I_{x3} + m_3r_2^2.$$

$$\Gamma_{112} = (I_{y2} - I_{z2} - m_2r_1^2)c_2s_2 + (I_{y3} - I_{z3})c_{23}s_{23} - m_3(l_1c_2 + r_2c_{23})(l_1s_2 + r_2s_{23})$$

$$\Gamma_{113} = (I_{y3} - I_{z3})c_{23}s_{23} - m_3r_2s_{23}(l_1c_2 + r_2c_{23})$$

$$\Gamma_{121} = (I_{y2} - I_{z2} - m_2r_1^2)c_2s_2 + (I_{y3} - I_{z3})c_{23}s_{23} - m_3(l_1c_2 + r_2c_{23})(l_1s_2 + r_2s_{23})$$

$$\Gamma_{131} = (I_{y3} - I_{z3})c_{23}s_{23} - m_3r_2s_{23}(l_1c_2 + r_2c_{23})$$

$$\Gamma_{211} = (I_{z2} - I_{y2} + m_2r_1^2)c_2s_2 + (I_{z3} - I_{y3})c_{23}s_{23} + m_3(l_1c_2 + r_2c_{23})(l_1s_2 + r_2s_{23})$$

$$\Gamma_{223} = -l_1m_3r_2s_3$$

$$\Gamma_{232} = -l_1m_3r_2s_3$$

$$\Gamma_{233} = -l_1m_3r_2s_3$$

$$\Gamma_{311} = (I_{z3} - I_{y3})c_{23}s_{23} + m_3r_2s_{23}(l_1c_2 + r_2c_{23})$$

$$\Gamma_{322} = l_1m_3r_2s_3$$

$$\begin{bmatrix} 0 \\ -(m_2gr_1 + m_3gl_1)\cos\theta_2 - m_3r_2\cos(\theta_2 + \theta_3) \\ -m_3gr_2\cos(\theta_2 + \theta_3) \end{bmatrix}$$

• Equations for PUMA 560 Arm

$$\begin{aligned}
I_2 &= I_{xx2} + m_2 * (r_{x2}^2 + r_{y2}^2) + (m_3 + m_4 + m_5 + m_6) * a_2^2 ; \\
I_3 &= -I_{xz2} + I_{yz2} + (m_3 + m_4 + m_5 + m_6) * a_2^2 \\
&\quad m_2 * r_{x2}^2 - m_2 * r_{y2}^2 ; \\
I_4 &= m_2 * r_{x2} * (d_2 + r_{z2}) + m_3 * a_2 * r_{z2} \\
&\quad + (m_3 + m_4 + m_5 + m_6) * a_2 * (d_2 + d_3) ; \\
I_5 &= -m_5 * a_2 * r_{x3} + (m_4 + m_5 + m_6) * a_2 * d_4 + m_4 * a_2 * r_{z4} ; \\
I_6 &= I_{zz5} + m_5 * r_{z5}^2 + m_4 * a_2^2 + m_4 * (d_4 + r_{z4})^2 + I_{yy4} \\
&\quad + m_5 * a_3^2 + m_5 * d_4^2 + I_{zz5} + m_6 * a_2^2 + m_6 * d_4 \\
&\quad + m_6 * r_{z6}^2 + I_{zz6} ; \\
I_7 &= m_5 * r_{x2}^2 + I_{xz5} - I_{yy5} + m_4 * r_{z4}^2 + 2 * m_4 * d_4 * r_{z4} \\
&\quad + (m_4 + m_5 + m_6) * (d_4^2 - a_3^2) + I_{yy5} - I_{xz4} + I_{zz5} ; \\
I_8 &= -m_4 * (d_2 + d_3) * (d_4 + r_{z4}) - (m_5 + m_6) * (d_2 + d_3) * d_4 \\
&\quad m_5 * r_{x2} * r_{z5} + m_5 * (d_2 + d_3) * r_{y5} ; \\
I_9 &= m_2 * r_{y2} * (d_2 + r_{z2}) ; \\
I_{10} &= 2 * m_4 * a_2 * r_{z4} + 2 * (m_4 + m_5 + m_6) * a_3 * d_4 ; \\
I_{11} &= -2 * m_5 * r_{x2} * r_{y2} ; \\
I_{12} &= (m_4 + m_5 + m_6) * a_2 * a_3 ; \\
I_{13} &= (m_4 + m_5 + m_6) * a_3 * (d_2 + d_3) ; \\
I_{14} &= I_{zz4} + I_{yy5} + I_{zz6} ; \\
I_{15} &= m_6 * d_4 * r_{z6} ; \\
I_{16} &= m_6 * a_2 * r_{z6} ; \\
I_{17} &= I_{zz5} + I_{zz6} + m_6 * r_{z6}^2 ; \\
I_{18} &= m_6 * (d_2 + d_3) * r_{z6} ; \\
I_{19} &= I_{yy5} - I_{xz4} + I_{zz5} - I_{yy5} + m_6 * r_{x2}^2 + I_{zz6} - I_{zz6} ; \\
I_{20} &= I_{yy5} - I_{zz5} - m_6 * r_{z6}^2 + I_{zz6} - I_{zz6} ; \\
I_{21} &= I_{zz4} - I_{yy4} + I_{zz5} - I_{zz5} ; \\
I_{22} &= m_6 * a_3 * r_{z6} ; \\
I_{23} &= I_{zz6} ;
\end{aligned}$$

Part II. Gravitational Constants

$$\begin{aligned}
g_1 &= -g * ((m_3 + m_4 + m_5 + m_6) * a_2 + m_2 * r_{z2}) ; \\
g_2 &= g * (m_3 * r_{y5} - (m_4 + m_5 + m_6) * d_4 - m_4 * r_{z4}) ; \\
g_3 &= g * m_2 * r_{y2} ; \\
g_4 &= -g * (m_4 + m_5 + m_6) * a_3 ; \\
g_5 &= -g * m_6 * r_{z6} ;
\end{aligned}$$

Table A3. Computed Values for the Constants Appearing in the Equations of Forces of Motion.
(Inertial constants have units of kilogram meters-squared)

$$\begin{aligned}
I_1 &= 1.43 \pm 0.05 & I_2 &= 1.75 \pm 0.07 \\
I_3 &= 1.38 \pm 0.05 & I_4 &= 6.90 \times 10^{-1} \pm 0.20 \times 10^{-1} \\
I_5 &= 3.72 \times 10^{-1} \pm 0.31 \times 10^{-1} & I_6 &= 3.33 \times 10^{-1} \pm 0.16 \times 10^{-1} \\
I_7 &= 2.98 \times 10^{-1} \pm 0.29 \times 10^{-1} & I_8 &= -1.34 \times 10^{-1} \pm 0.14 \times 10^{-1} \\
I_9 &= 2.38 \times 10^{-2} \pm 1.20 \times 10^{-2} & I_{10} &= -2.13 \times 10^{-2} \pm 0.22 \times 10^{-2} \\
I_{11} &= -1.42 \times 10^{-2} \pm 0.70 \times 10^{-2} & I_{12} &= -1.10 \times 10^{-2} \pm 0.11 \times 10^{-2} \\
I_{13} &= -3.79 \times 10^{-3} \pm 0.90 \times 10^{-3} & I_{14} &= 1.64 \times 10^{-3} \pm 0.07 \times 10^{-3} \\
I_{15} &= 1.25 \times 10^{-3} \pm 0.30 \times 10^{-3} & I_{16} &= 1.24 \times 10^{-3} \pm 0.30 \times 10^{-3} \\
I_{17} &= 6.42 \times 10^{-4} \pm 3.00 \times 10^{-4} & I_{18} &= 4.31 \times 10^{-4} \pm 1.30 \times 10^{-4} \\
I_{19} &= 3.00 \times 10^{-4} \pm 1.40 \times 10^{-4} & I_{20} &= -2.02 \times 10^{-4} \pm 8.00 \times 10^{-4} \\
I_{21} &= -1.00 \times 10^{-4} \pm 6.00 \times 10^{-4} & I_{22} &= -5.80 \times 10^{-5} \pm 1.50 \times 10^{-5} \\
I_{23} &= 4.00 \times 10^{-5} \pm 2.00 \times 10^{-5} & & \\
I_{m1} &= 1.14 \pm 0.27 & I_{m2} &= 4.71 \pm 0.54 \\
I_{m3} &= 8.27 \times 10^{-1} \pm 0.93 \times 10^{-1} & I_{m4} &= 2.00 \times 10^{-1} \pm 0.16 \times 10^{-1} \\
I_{m5} &= 1.79 \times 10^{-1} \pm 0.14 \times 10^{-1} & I_{m6} &= 1.93 \times 10^{-1} \pm 0.16 \times 10^{-1}
\end{aligned}$$

(Gravitational constants have units of newton meters)

$$\begin{aligned}
g_1 &= -37.2 \pm 0.05 & g_2 &= -8.44 \pm 0.20 \\
g_3 &= 1.02 \pm 0.50 & g_4 &= 2.49 \times 10^{-1} \pm 0.25 \times 10^{-1} \\
g_5 &= -2.82 \times 10^{-2} \pm 0.56 \times 10^{-2} & &
\end{aligned}$$

Table A4. The expressions giving the elements of the kinetic energy matrix.
(The Abbreviated Expressions have units of kg-m².)

$$\begin{aligned}
a_{11} &= I_{m1} + I_1 + I_3 * CC2 + I_7 * SC23 + I_{10} * SC23 + I_{11} * SC2 \\
&\quad + I_{20} * (SS5 * (SS23 * (1 + CC4) - 1) - 2 * SC23 * C4 * SC5) \\
&\quad + I_{22} * ((1 - 2 * SS23) * C5 - 2 * SC23 * C4 * S5) \\
&\quad + I_{10} * (1 - 2 * SS23) + I_{11} * (1 - 2 * SS2) ; \\
a_{12} &= I_{m1} + I_8 * C23 + I_9 * C2 + I_{13} * S23 + I_{12} * C2 * C23 \\
&\quad + I_{15} * (SS23 * C5 + SC23 * C4 * S5) \\
&\quad + I_{16} * C2 * (S23 * C5 + C23 * C4 * S5) \\
&\quad + I_{18} * S4 * S5 + I_{22} * (SC23 * C5 + C23 * C4 * S5) ; \\
a_{13} &= I_4 * S2 + I_8 * C23 + I_9 * C2 + I_{13} * S23 - I_{15} * C23 * S4 * S5 \\
&\quad + I_{16} * S2 * S4 * S5 + I_{21} * (S23 * C4 * S5 - C23 * C5) \\
&\quad + I_{19} * S23 * SC4 + I_{20} * S4 * (S23 * C4 * CC5 + C23 * SC5) \\
&\quad + I_{22} * S23 * S4 * S5 ; \\
a_{14} &= 2 * (-I_{15} * S23 * S4 * S5 - I_{16} * C2 * C23 * S4 * S5) \\
&\quad + I_{18} * C4 * S5 - I_{20} * (S23 * SS5 * SC4 - SC23 * S4 * SC5) \\
&\quad - I_{22} * C23 * S4 * S5 - I_{21} * SS23 * SC4) ; \\
a_{15} &= 2 * (I_{20} * (SC5 * (CC4 * (1 - CC23) - CC23) \\
&\quad - SC23 * C4 * (1 - 2 * SS5)) - I_{15} * (SS23 * S5 - SC23 * C4 * C5) \\
&\quad - I_{16} * C2 * (S23 * S5 - C23 * C4 * C5) + I_{18} * S4 * C5 \\
&\quad + I_{22} * (CC23 * C4 * C5 - SC23 * S5)) ; \\
a_{16} &= 2 * (I_{20} * (C5 * (CC4 * (1 - CC23) - CC23) \\
&\quad - SC23 * C4 * (1 - 2 * SS5)) - I_{15} * (SS23 * S5 - SC23 * C4 * C5) \\
&\quad - I_{16} * C2 * (S23 * S5 - C23 * C4 * C5) + I_{20} * (S23 * C4 * CC5 + C23 * SS5)) \\
&\quad + I_{22} * (CC23 * C4 * C5 - SC23 * S5)) ; \\
a_{17} &= I_{15} * S23 * S4 * C5 + I_{16} * C2 * S4 * C5 + I_{17} * S23 * S4 * S5 \\
&\quad + I_{18} * (S23 * S5 - C23 * C4 * C5) + I_{22} * C23 * S4 * C5 ; \\
a_{18} &= I_{15} * C23 * C5 - S23 * C4 * S5 ; \\
a_{19} &= I_{yy5} - I_{xz4} + I_{zz5} - I_{yy5} + m_6 * r_{x2}^2 + I_{zz6} - I_{zz6} ; \\
I_{20} &= I_{yy5} - I_{zz5} - m_6 * r_{z6}^2 + I_{zz6} - I_{zz6} ; \\
I_{21} &= I_{zz4} - I_{yy4} + I_{zz5} - I_{zz5} ; \\
I_{22} &= m_6 * a_3 * r_{z6} ; \\
I_{23} &= I_{zz6} ;
\end{aligned}$$

Table A5. The expressions giving the elements of the Coriolis matrix.
(The Abbreviated Expressions have units of kg-m².)

$$\begin{aligned}
b_{112} &= 2 * \{-I_3 * SC2 + I_5 * C223 + I_7 * SC23 - I_{12} * S223 \\
&\quad + I_{13} * (2 * SC23 * C5 + (1 - 2 * SS23) * C4 * S5)\} ; \\
b_{113} &= 2 * (I_5 * C23 + I_7 * SC23 - I_{12} * C2 * S23 \\
&\quad + I_{13} * (2 * SC23 * C5 + (1 - 2 * SS23) * C4 * S5) \\
&\quad + I_{16} * C2 * (C23 * C5 - S23 * C4 * S5) + I_{21} * SC23 * CC4 \\
&\quad + I_{20} * ((1 + CC4) * SC23 * SS5 - (1 - 2 * SS23) * C4 * SC5) \\
&\quad + I_{22} * ((1 - 2 * SS23) * C5 - 2 * SC23 * C4 * S5) ; \\
b_{114} &= -2.76 * SC2 + 7.44 \times 10^{-1} * C223 + 0.60 * SC23 \\
&\quad - 2.13 \times 10^{-2} * (1 - 2 * SS23) . \\
b_{115} &= 2 * (I_5 * C23 + I_7 * SC23 - I_{12} * C2 * S23 \\
&\quad + I_{13} * (2 * SC23 * C5 + (1 - 2 * SS23) * C4 * S5) \\
&\quad + I_{16} * C2 * (C23 * C5 - S23 * C4 * S5) + I_{21} * SC23 * CC4 \\
&\quad + I_{20} * ((1 + CC4) * SC23 * SS5 - (1 - 2 * SS23) * C4 * SC5) \\
&\quad + I_{22} * ((1 - 2 * SS23) * C5 - 2 * SC23 * C4 * S5) ; \\
b_{116} &= 2 * (-I_{15} * SC23 * S4 * S5 - I_{16} * C2 * C23 * S4 * S5) \\
&\quad + I_{18} * C4 * S5 - I_{20} * (S23 * SS5 * SC4 - SC23 * S4 * SC5) \\
&\quad - I_{22} * C23 * S4 * S5 - I_{21} * SS23 * SC4) ; \\
b_{117} &= 2 * (I_{20} * (SC5 * (CC4 * (1 - CC23) - CC23) \\
&\quad - SC23 * C4 * (1 - 2 * SS5)) - I_{15} * (SS23 * S5 - SC23 * C4 * C5) \\
&\quad - I_{16} * C2 * (S23 * S5 - C23 * C4 * C5) + I_{18} * S4 * C5 \\
&\quad + I_{22} * (CC23 * C4 * C5 - SC23 * S5)) ; \\
b_{118} &= 2 * (I_{20} * (C5 * (CC4 * (1 - CC23) - CC23) \\
&\quad - SC23 * C4 * (1 - 2 * SS5)) - I_{15} * (SS23 * S5 - SC23 * C4 * C5) \\
&\quad - I_{16} * C2 * (S23 * S5 - C23 * C4 * C5) + I_{20} * (S23 * C4 * CC5 + C23 * SS5)) \\
&\quad + I_{22} * (CC23 * C4 * C5 - SC23 * S5)) ; \\
b_{119} &= I_{15} * S23 * S4 * S5 ; \\
b_{120} &= 2 * \{-I_8 * S23 + I_{13} * C23 + I_{15} * S23 * S4 * S5 \\
&\quad + I_{18} * (C23 * C4 * S5 + S23 * C5) + I_{21} * C23 * SC4 \\
&\quad + I_{20} * S4 * (C23 * C4 * CC5 + C23 * SC5) \\
&\quad + I_{22} * C23 * S4 * S5\} ; \\
b_{121} &= 2.67 \times 10^{-1} * S23 - 7.58 \times 10^{-2} * C23 . \\
b_{122} &= -I_{18} * 2 * S23 * S4 * S5 + I_{19} * S23 * (1 - (2 * SS4)) \\
&\quad + I_{20} * S23 * (1 - 2 * SS4 * CC5) - I_{14} * S23 ; \\
b_{123} &= I_{17} * C23 * S4 + I_{18} * 2 * (S23 * C4 * C5 + C23 * S5) \\
&\quad + I_{20} * S4 * (C23 * (1 - 2 * SS5) - S23 * C4 * 2 * SC5) ; \\
b_{124} &= -I_{18} * 2 * S23 * S4 * S5 + I_{19} * S23 * (1 - (2 * SS4)) \\
&\quad + I_{20} * S23 * (1 - 2 * SS4 * CC5) - I_{14} * S23 ; \\
b_{125} &= I_{17} * C23 * S4 + I_{18} * 2 * (S23 * C4 * C5 + C23 * S5) \\
&\quad + I_{20} * S4 * (C23 * (1 - 2 * SS5) - S23 * C4 * 2 * SC5) ; \\
b_{126} &= -I_{23} * (S23 * C5 + C23 * C4 * S5) ; \\
b_{127} &= I_{15} * S3 + I_6 + I_{12} * C3 + I_{16} * (S3 * C5 + C3 * C4 * S5) \\
&\quad + I_{20} * SS4 * SS5 + I_{21} * SS4 + 2 * (I_{15} * C5 + I_{22} * C4 * S5) ; \\
b_{128} &= 0.33 + 3.72 \times 10^{-1} * S3 - 1.10 \times 10^{-2} * C3 . \\
b_{129} &= I_{15} * S4 * S5 - I_{16} * S3 * S4 * S5 + I_{20} * S4 * SC5 ; \\
b_{130} &= 0. \\
b_{131} &= I_{15} * C4 * C5 + I_{16} * (C3 * S5 + S3 * C4 * C5) \\
&\quad + I_{17} * C4 * S5 + I_{22} * S5 ; \\
b_{132} &= I_{23} * S4 * S5 ; \\
b_{133} &= I_{m5} + I_6 + I_{20} * SS4 * SS5 + I_{21} * SS4 \\
&\quad + 2 * (I_{15} * C5 + I_{22} * C4 * S5) ; \\
b_{134} &= -I_{13} * S4 * S5 + I_{20} * S4 * S5 + I_{21} * SS4 \\
&\quad \approx 1.16 . \\
b_{135} &= I_{15} * C4 * C5 + I_{16} * (C3 * S5 + S3 * C4 * C5) \\
&\quad + I_{17} * C4 * S5 + I_{22} * S5 ; \\
b_{136} &= -I_{13} * S4 * S5 + I_{20} * S4 * SC5 ; \\
b_{137} &= I_{15} * C4 * C5 + I_{17} * C4 + I_{22} * S5 ; \\
b_{138} &= 1.25 \times 10^{-3} * C4 * C5 . \\
b_{139} &= I_{15} * C4 * C5 + I_{17} * C4 + I_{22} * S5 ; \\
b_{140} &= 0. \\
b_{141} &= I_{15} * C4 * C5 + I_{16} * (C3 * S5 + S3 * C4 * C5) \\
&\quad + I_{17} * C4 * S5 + I_{22} * S5 ; \\
b_{142} &= I_{15} * C4 * C5 + I_{16} * (C3 * S5 + S3 * C4 * C5) \\
&\quad + I_{17} * C4 * S5 + I_{22} * S5 ; \\
b_{143} &= 0. \\
b_{144} &= I_{m4} + I_{14} - I_{20} * S55 ; \\
b_{145} &= 0. \\
b_{146} &= I_{23} * S4 * S5 ; \\
b_{147} &= I_{15} * C23 * S4 * C5 + I_{17} * C23 * S4 * C5 ; \\
b_{148} &= 0. \\
b_{149} &= I_{14} * S23 + I_{19} * S23 + I_{20} * S23 \\
&\quad + I_{21} * S23 * C4 * S5 ; \\
b_{150} &= 1.64 \times 10^{-3} * S23 - 2.50 \times 10^{-3} * C23 * C4 * S5 + \\
&\quad 2.48 \times 10^{-3} * S2 * C4 * S5 + 0.30 \times 10^{-3} * S23 * (1 - 2 * SS4) . \\
b_{151} &= 2 * \{-I_{15} * C23 * S4 * C5 + I_{17} * C23 * S4 * C5 \\
&\quad + I_{16} * S2 * S4 * C5\} - I_{17} * C23 * S4 \\
&\quad + I_{20} * (C23 * S4 * (1 - 2 * SS5) - 2 * S23 * SC4 * SC5) ; \\
b_{152} &= \approx -2.50 \times 10^{-3} * C23 * S4 * C5 + 2.48 \times 10^{-3} * S2 * S4 * C5 \\
&\quad - 6.42 \times 10^{-4} * C23 * S4 . \\
b_{153} &= -b_{126} . \\
b_{154} &= 2 * \{-I_{12} * S3 + I_5 * C3 + I_{16} * (C3 * C5 - S3 * C4 * S5)\} ; \\
b_{155} &= \approx 2.20 \times 10^{-2} * S3 + 7.44 \times 10^{-1} * C3 .
\end{aligned}$$

Table A6. The expressions for the terms of the centrifugal matrix.
(The Abbreviated Expressions have units of kg-m².)

$$\begin{aligned}
b_{224} &= 2 * \{-I_{16} * C3 * S4 * S5 + I_{20} * SC4 * SS5 \\
&\quad + I_{21} * SC4 - I_{22} * S4 * S5\} ; \\
b_{225} &= \approx -2.48 \times 10^{-2} * C3 * S4 * S5 . \\
b_{226} &= 0 . \\
b_{227} &= b_{224} . \\
b_{228} &= b_{225} . \\
b_{229} &= 0 . \\
b_{230} &= \{-I_{15} * S4 + I_{16} * (C3 * C4 * C5 - S3 * S5) \\
&\quad + I_{20} * SS4 * SC5 + I_{22} * C4 * C5\} ; \\
b_{231} &= \approx -2.50 \times 10^{-3} * S5 + 2.48 \times 10^{-3} * (C3 * C4 * C5 - S3 * S5) . \\
b_{232}$$

Are We Done Yet?

- Let us use the tools to solve our grasping example

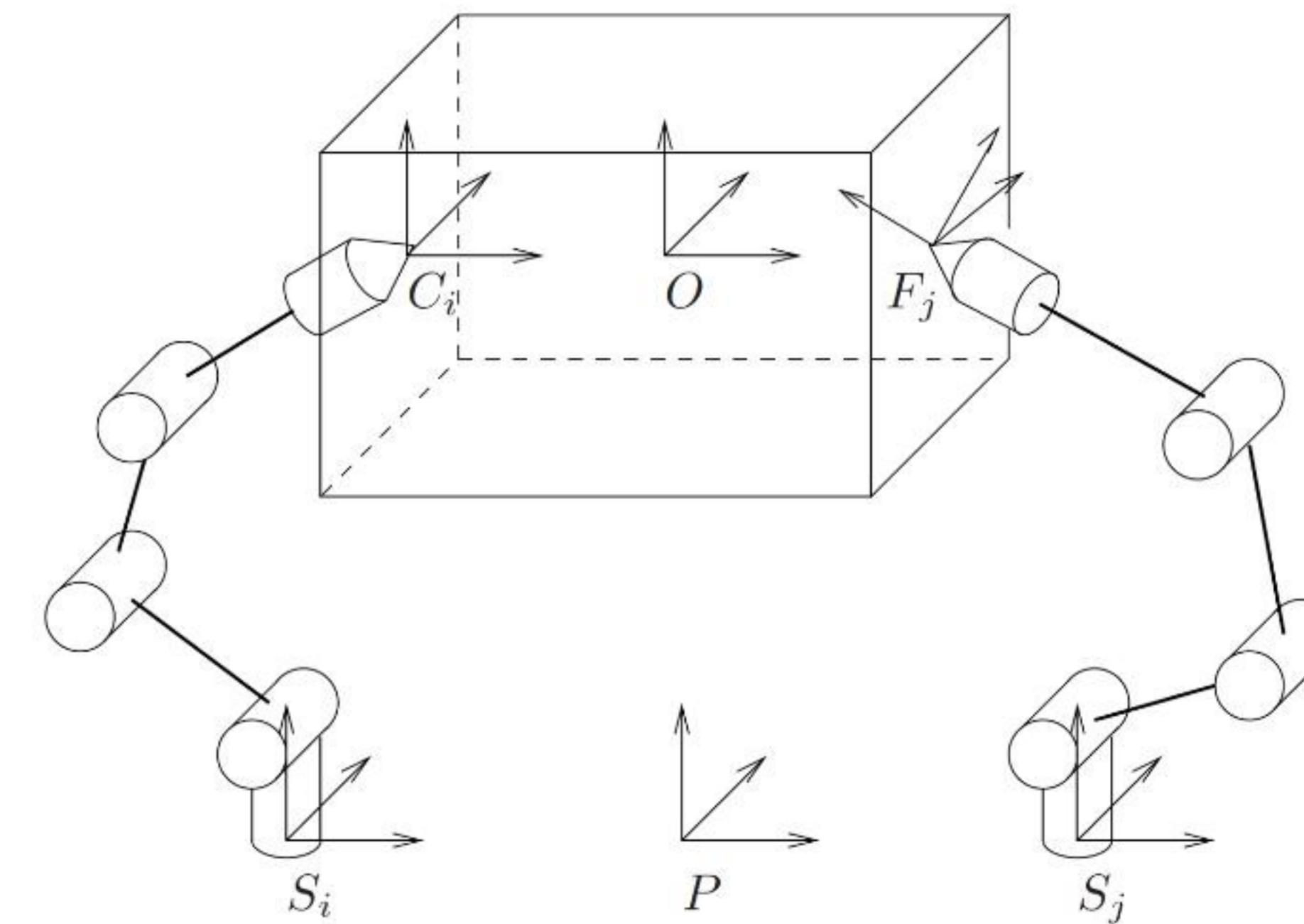


Figure 5.14: Grasp coordinate frames.

Are We Done Yet?

- Let us use the tools to solve our grasping example
 - Suppose that each arm has 3 joints with motors installed.

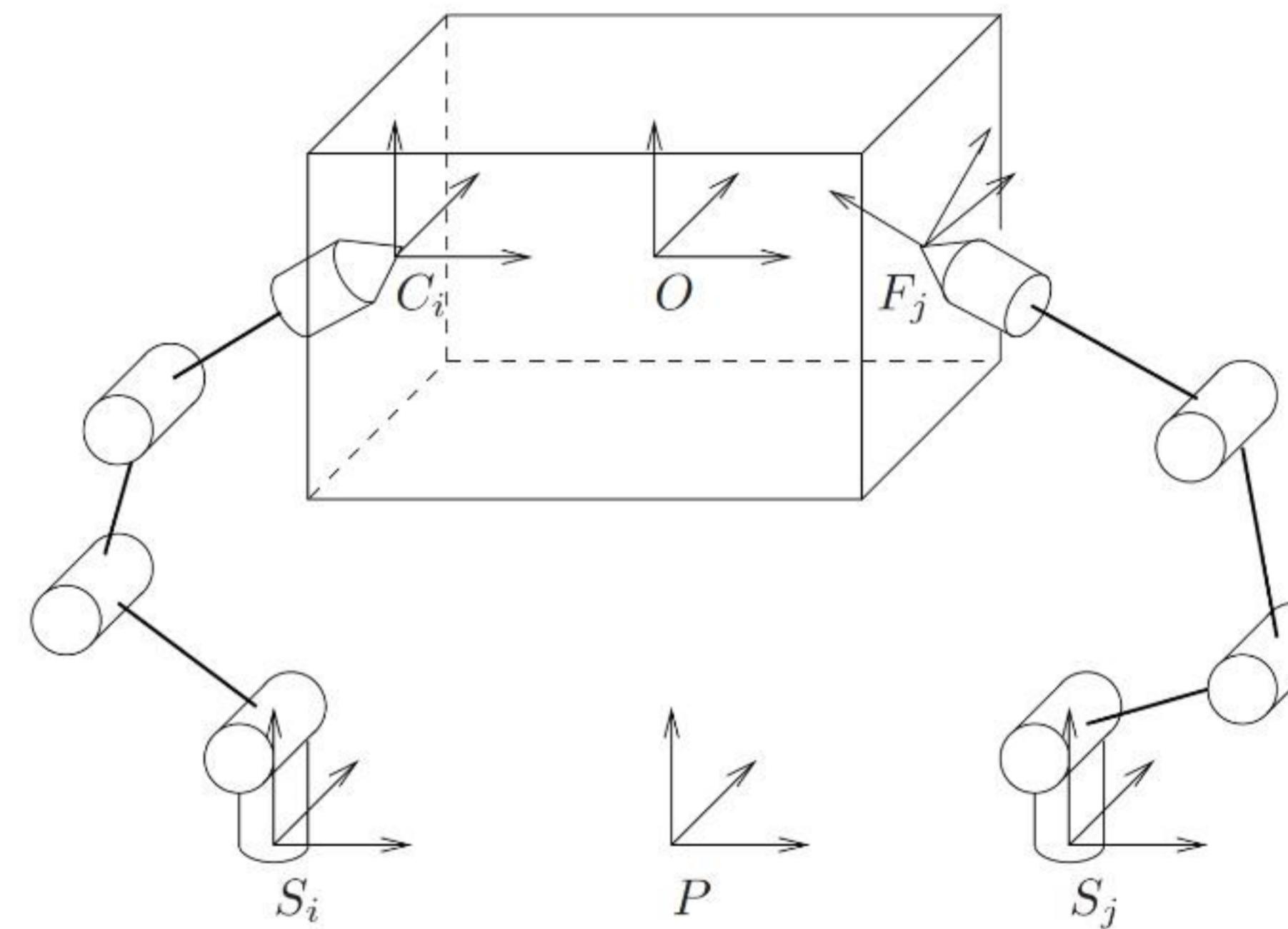


Figure 5.14: Grasp coordinate frames.

Are We Done Yet?

- Let us use the tools to solve our grasping example
 - Suppose that each arm has 3 joints with motors installed.
 - The **degree-of-freedom** (independent coordinates) of the system is: $2 \times 3 + 6 = 12$ including both two arms and the rigid object.

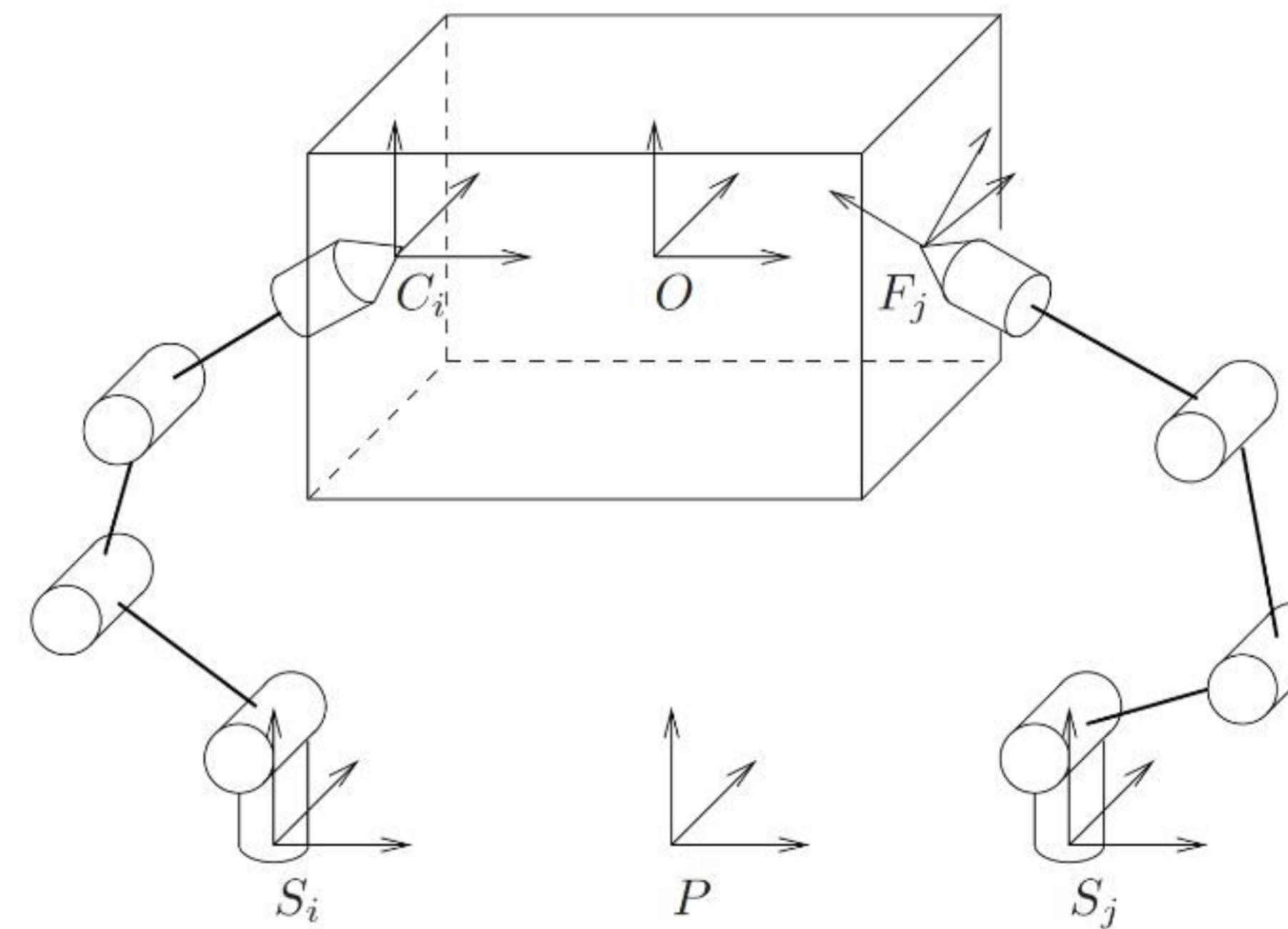


Figure 5.14: Grasp coordinate frames.

Are We Done Yet?

- Let us use the tools to solve our grasping example
 - Suppose that each arm has 3 joints with motors installed.
 - The **degree-of-freedom** (independent coordinates) of the system is: $2 \times 3 + 6 = 12$ including both two arms and the rigid object.
 - If we would like control the robot to hold the box firmly, the degree-of-freedom will reduce due to constraints.

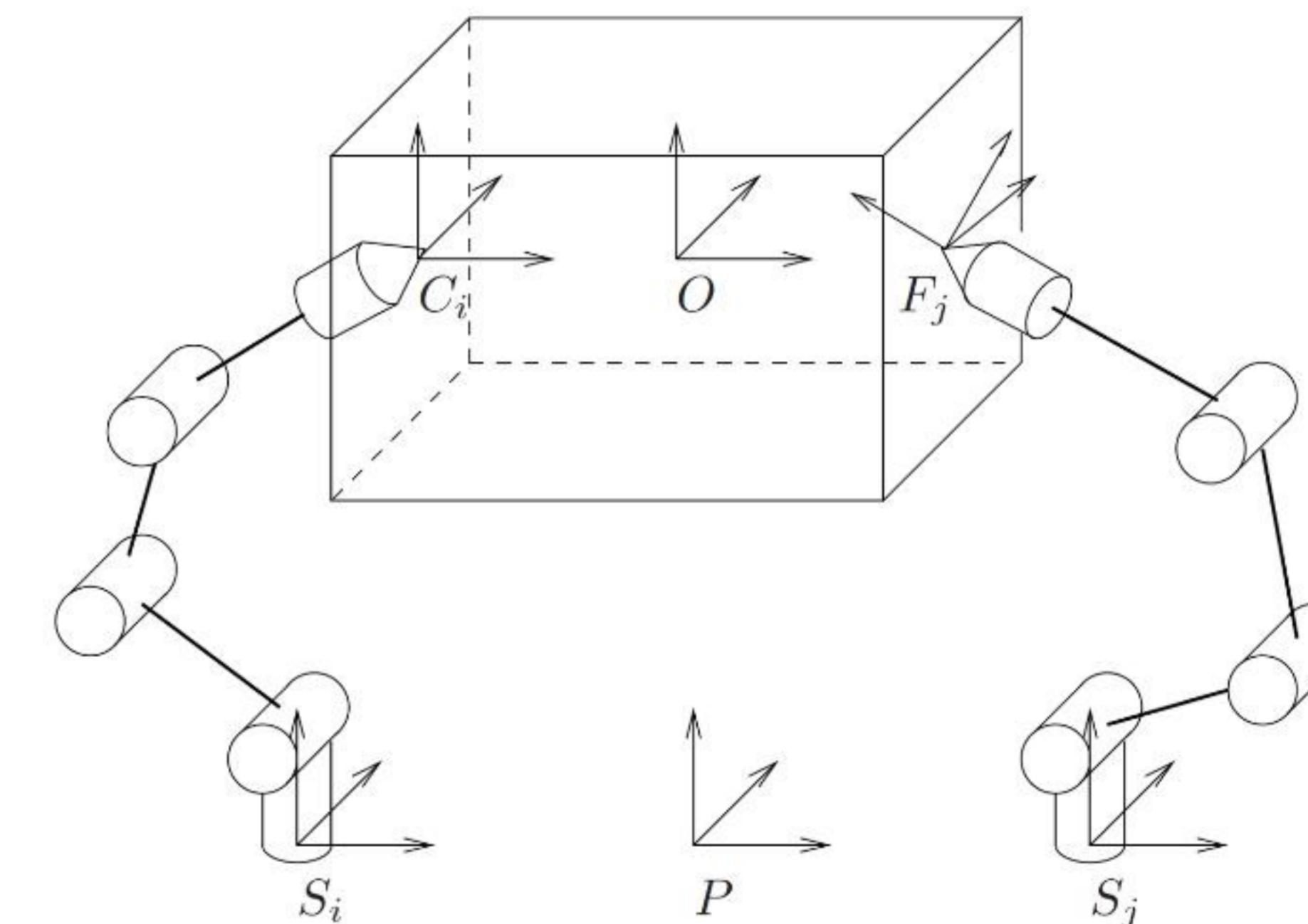


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 - Suppose that each arm has 3 joints with motors installed.
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 - If we would like control the robot to hold the box firmly, the degree-of-freedom will reduce due to constraints.
 - How to solve FD/ID for this system?

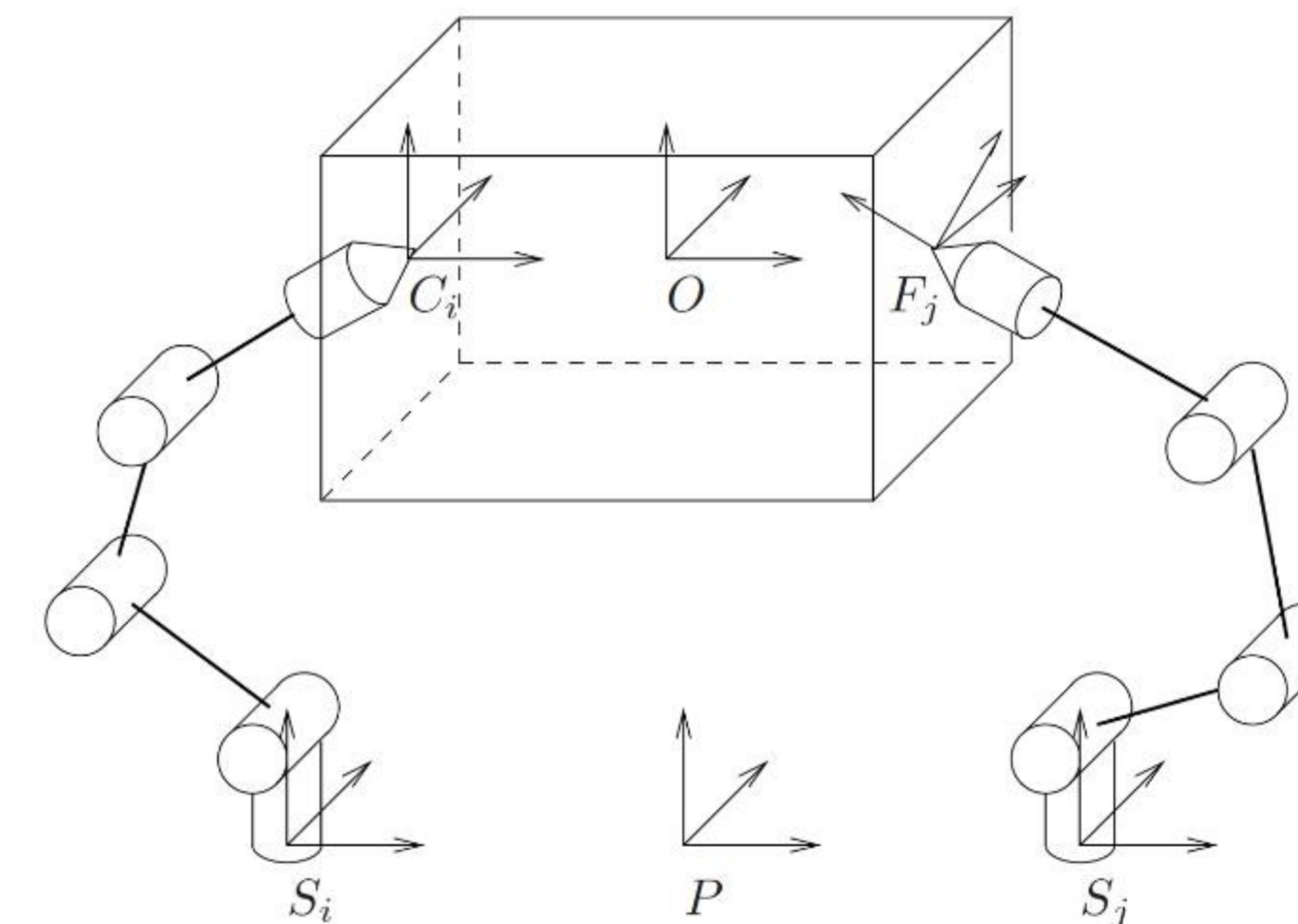


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Are We Done Yet?

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 - Suppose that each arm has 3 joints with motors installed.
 - The **degree-of-freedom** (independent coordinates) of the system is: $2 \times 3 + 6 = 12$ including both two arms and the rigid object.
 - If we would like control the robot to hold the box firmly, the degree-of-freedom will reduce due to constraints.
 - How to solve FD/ID for this system?
 - Additionally, if we are to keep the box static, among many τ candidates, we want to choose one by the contact force and torque. However, our previously developed tool does not tell us about these quantities!

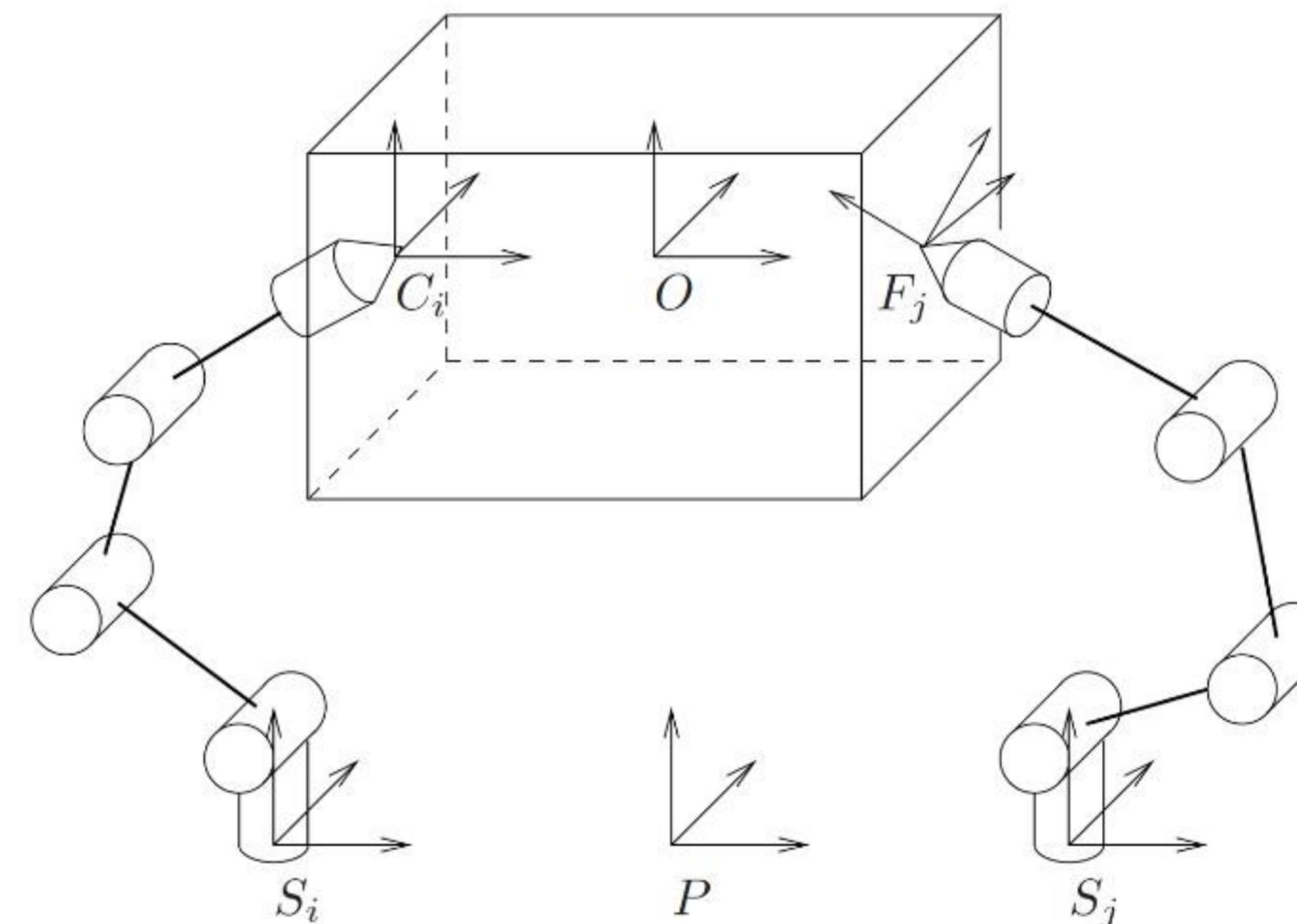


Figure 5.14: Grasp coordinate frames.

Constrained Lagrangian Method

Constrained Lagrangian Method

- Our previous Lagrangian method uses independent coordinates to describe the state of the system

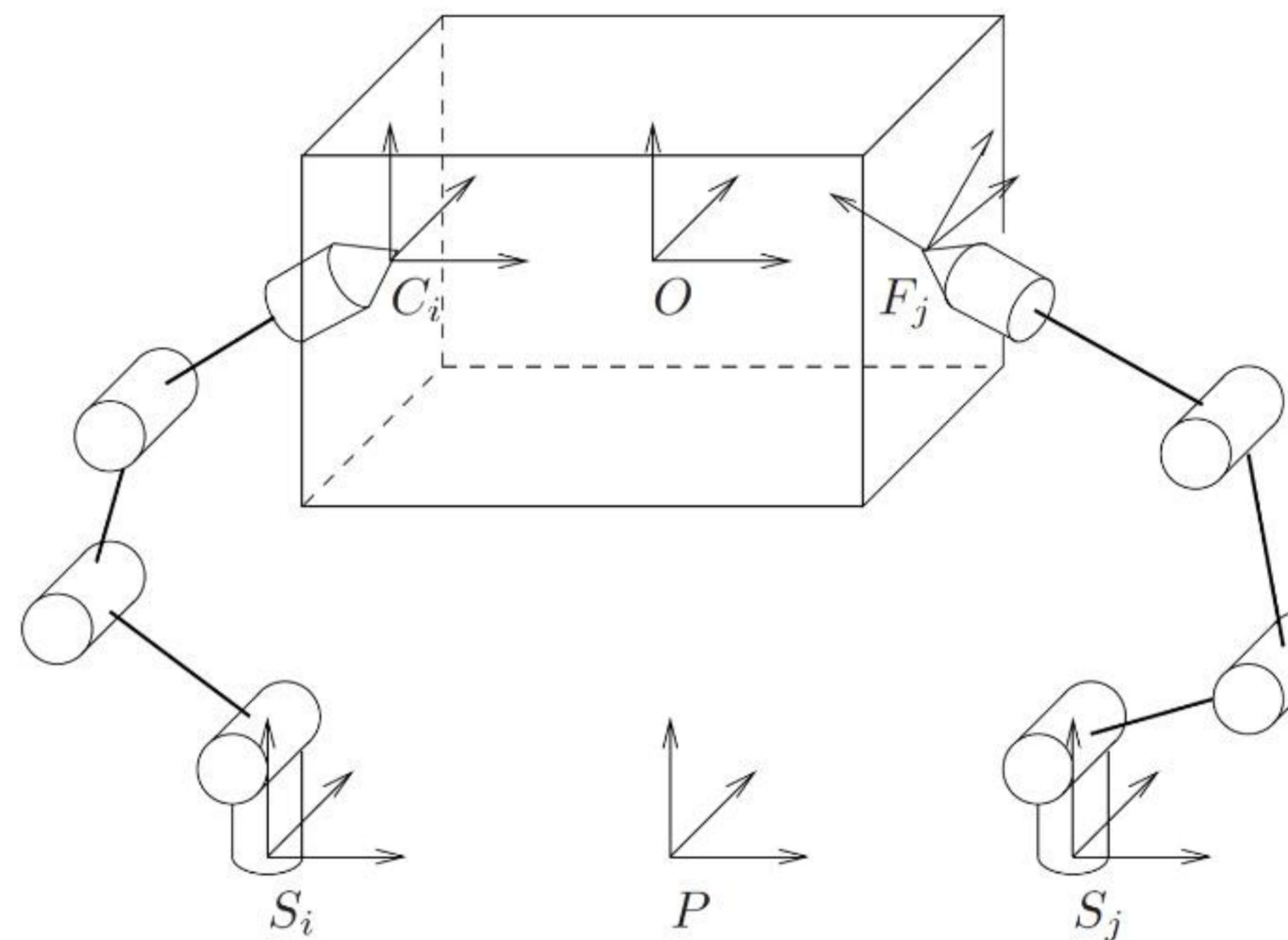


Figure 5.14: Grasp coordinate frames.

Constrained Lagrangian Method

- Our previous Lagrangian method uses independent coordinates to describe the state of the system
- Sometimes, we are interested in the situation happening not at these independent coordinates (e.g., in the grasping example, we are interested in the internal forces at the contact points).

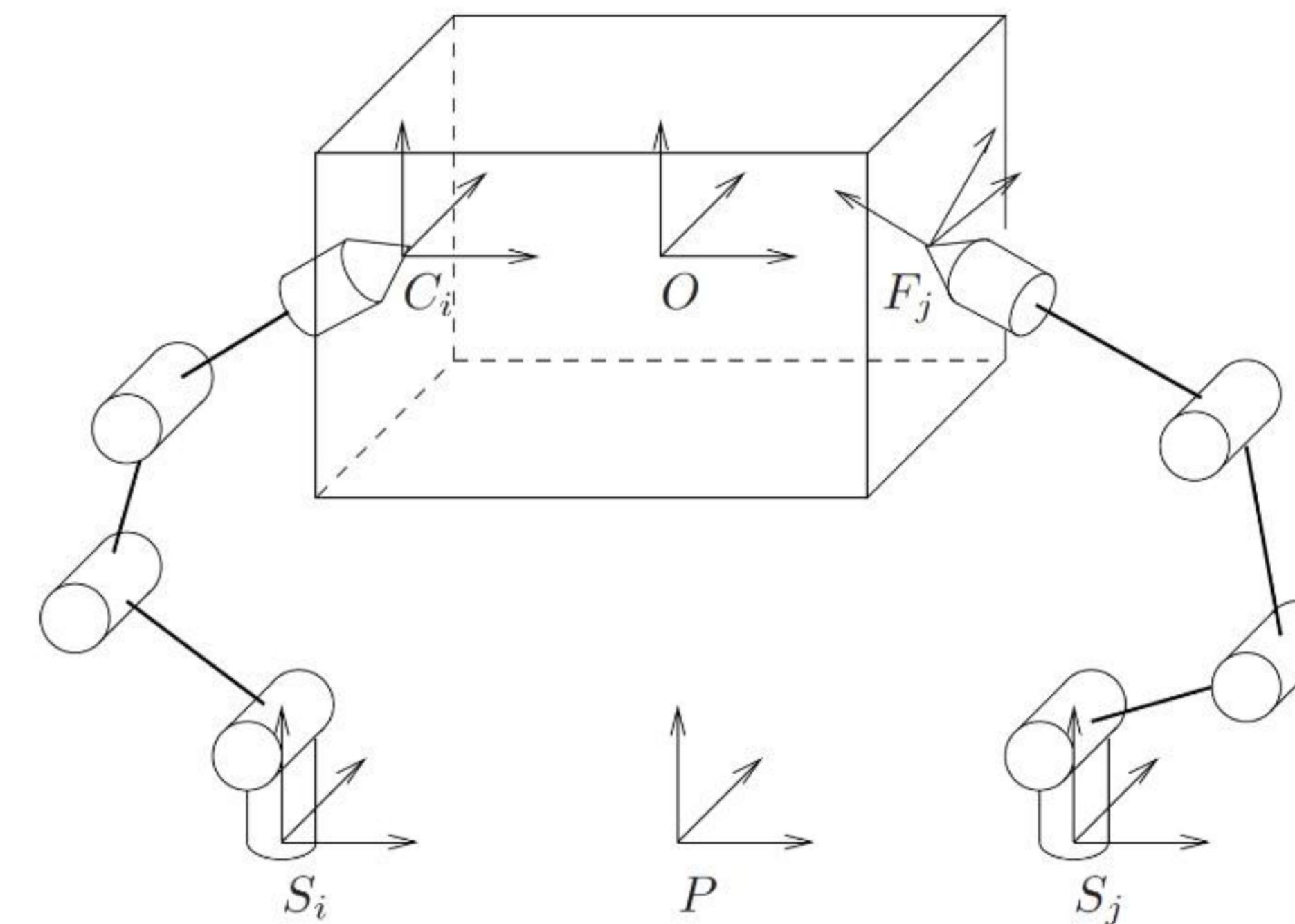


Figure 5.14: Grasp coordinate frames.

Constrained Lagrangian Method

- Our previous Lagrangian method uses independent coordinates to describe the state of the system
- Sometimes, we are interested in the situation happening not at these independent coordinates (e.g., in the grasping example, we are interested in the internal forces at the contact points).
- We augment our previous Lagrangian methods to answer such query.

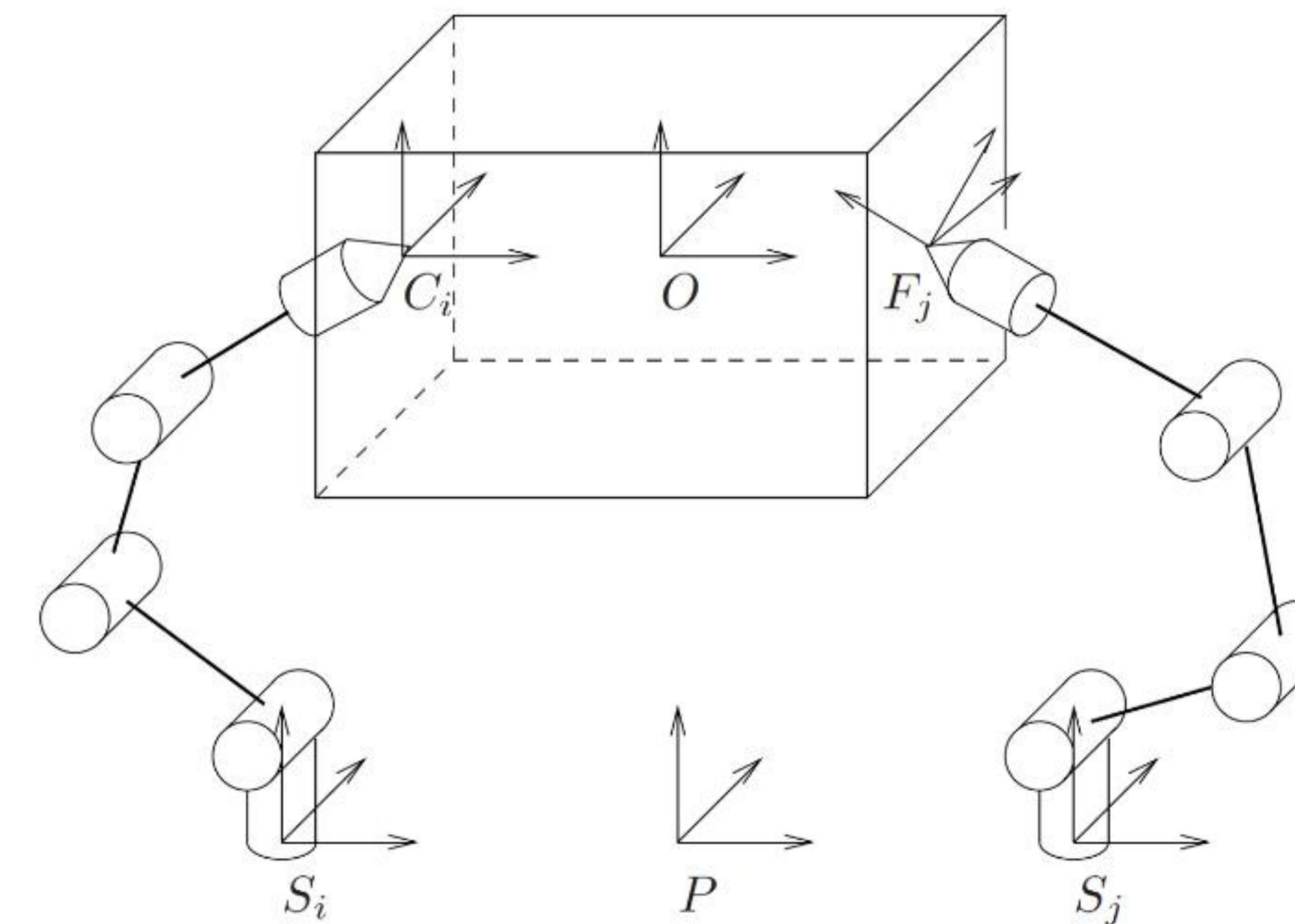


Figure 5.14: Grasp coordinate frames.

Review: Contact Coordinate Frame

- We build a **contact frame** C_i at each contact point
- The z -axis of the frame points inward along *surface normal*
- When recording force and torque at the contact point, it is natural to set C_i as the *observer's frame*, i.e.,

$$\mathbf{F}^{C_i} = \begin{bmatrix} \mathbf{f}^{C_i} \\ \boldsymbol{\tau}^{C_i} \end{bmatrix}$$

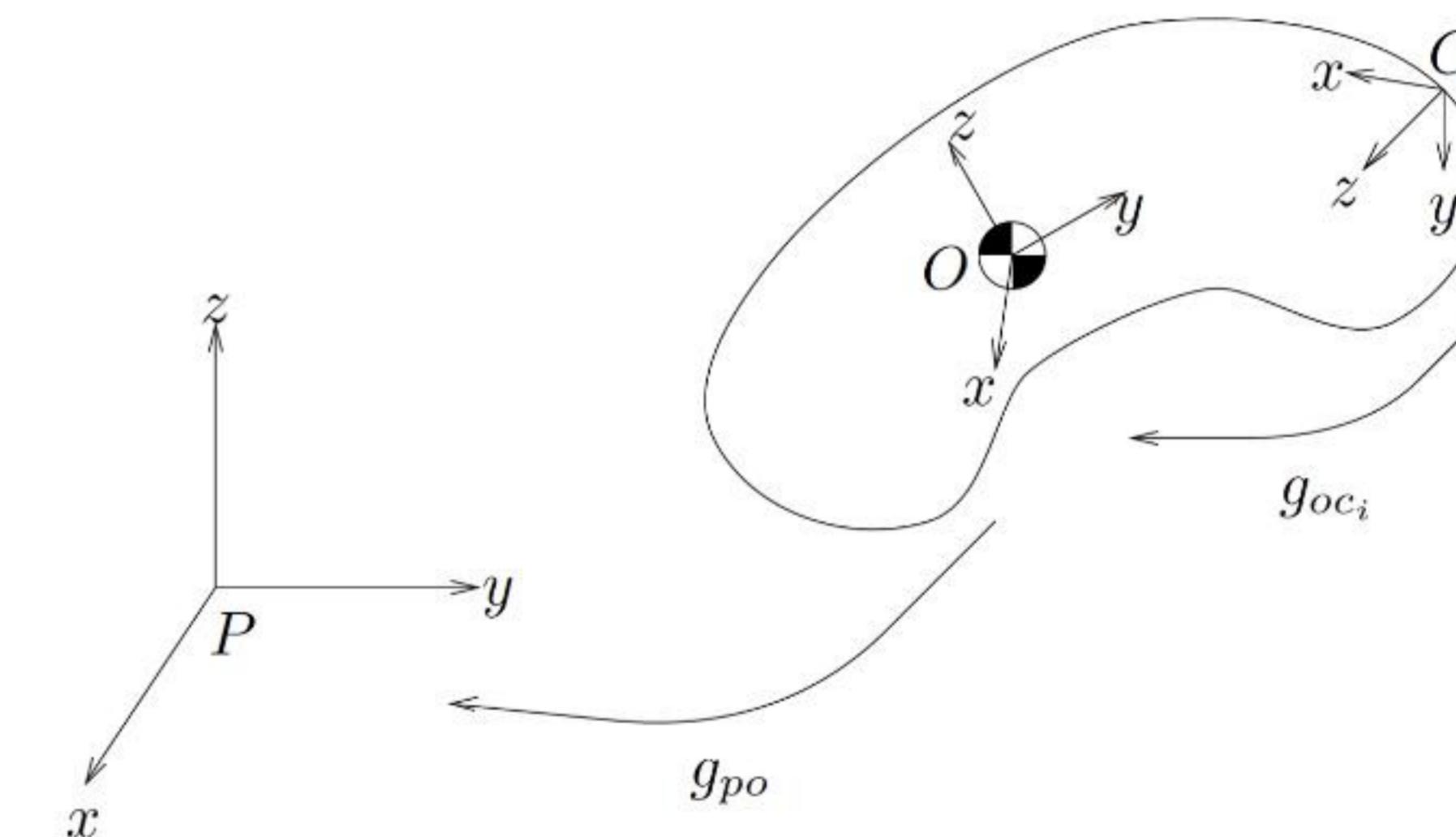


Figure 5.2: Coordinate frames for contact and object forces.

Step 1: Adding Generalized Coordinates for Constraint Force

- We assign additional generalized coordinates at the constraints.
 - E.g., we add the 6D pose of the contact frames at the end of the two finger tips as the generalized coordinates.
 - So the generalized velocities are the twists at the finger tips ξ^{C_i} .

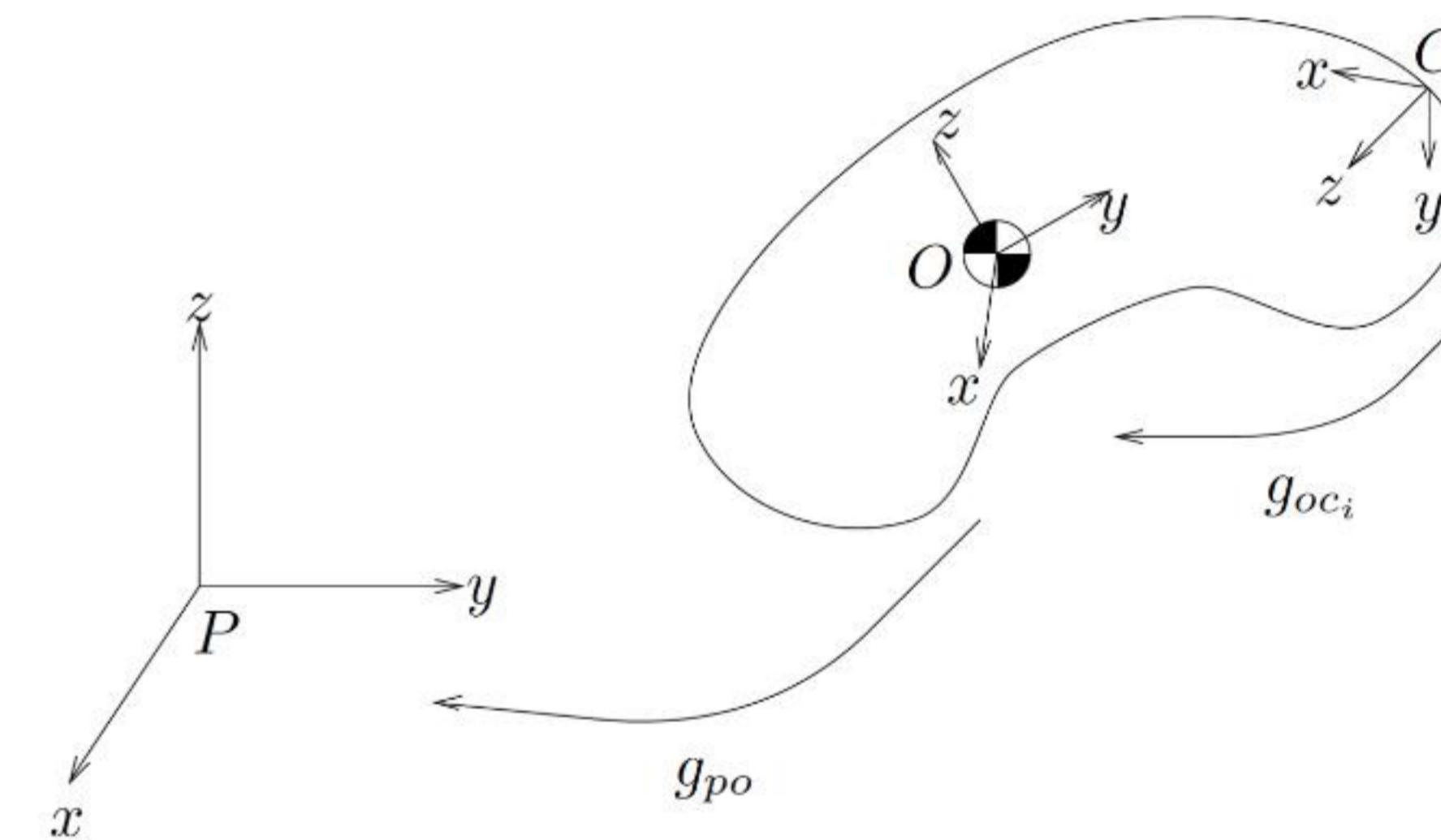


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Step 1: Adding Generalized Coordinates for Constraint Force

- We assign additional generalized coordinates at the constraints.
 - E.g., we add the 6D pose of the contact frames at the end of the two finger tips as the generalized coordinates.
 - So the generalized velocities are the twists at the finger tips ξ^{C_i} .
- These additional generalized coordinates have their corresponding generalized forces (our query).
 - In our example, the additional generalized forces are the internal forces (\mathbf{f}^{C_i} and $\boldsymbol{\tau}^{C_i}$) at the contact points \mathbf{F}^{C_i} .

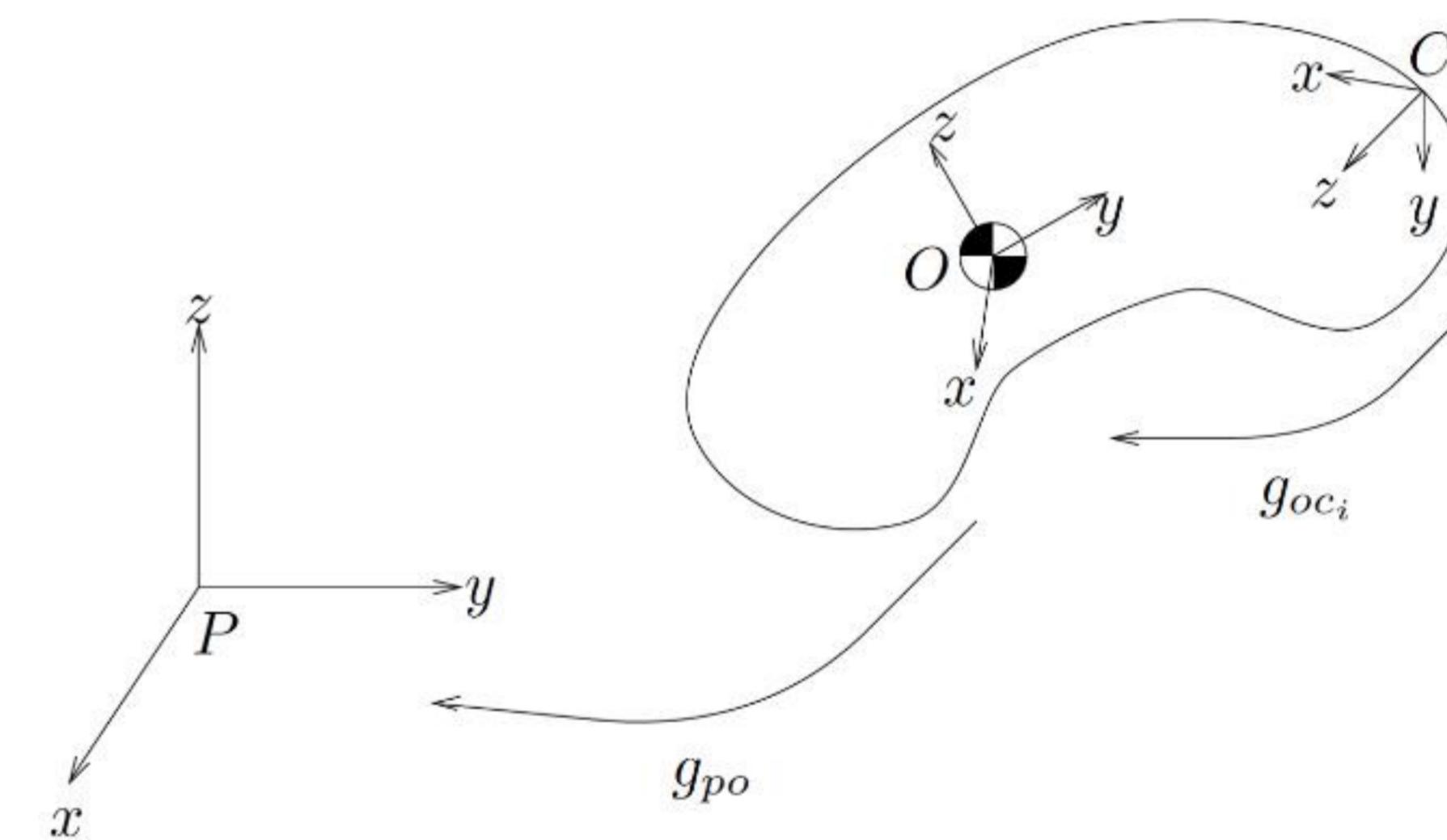


Figure 5.2: Coordinate frames for contact and object forces.

Step 2: Adding Constraints

- While assigning the additional generalized coordinates, we also need to add constraints.
 - In our example, ξ^{C_i} must conform to two sets of constraints:
 - All contact frames must move in accordance with the movement of the object;
 - All contact frames must move in accordance with the movement of the arms.

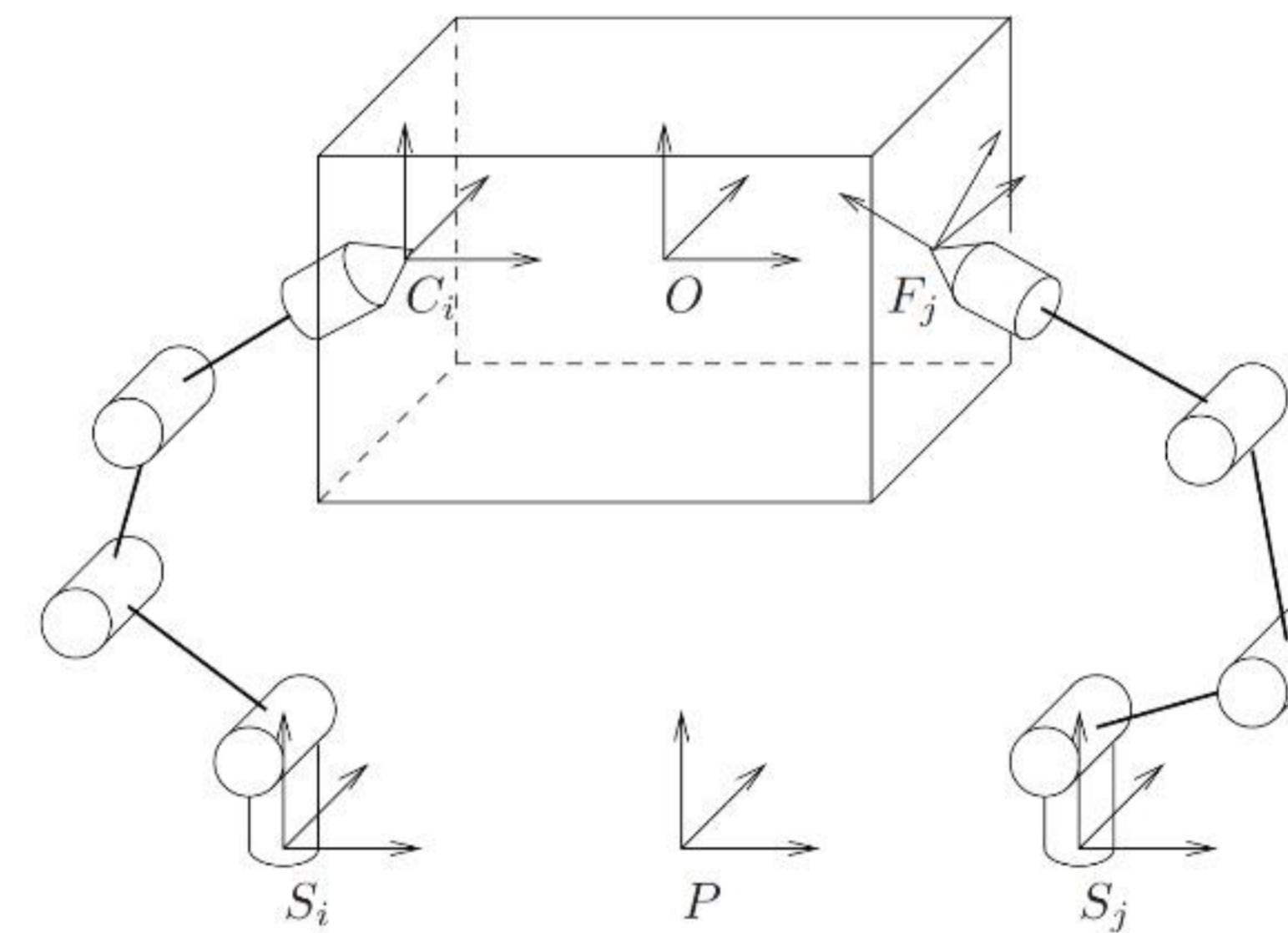


Figure 5.14: Grasp coordinate frames.

Step 2: Adding Constraints

"All contact frames must move in accordance with the movement of the object."

- Assume the body frame twist of the box is ξ^b .
- Then, $\xi^{C_i} = [\text{Ad}_{C_i \rightarrow b}] \xi^b \quad \forall i$.

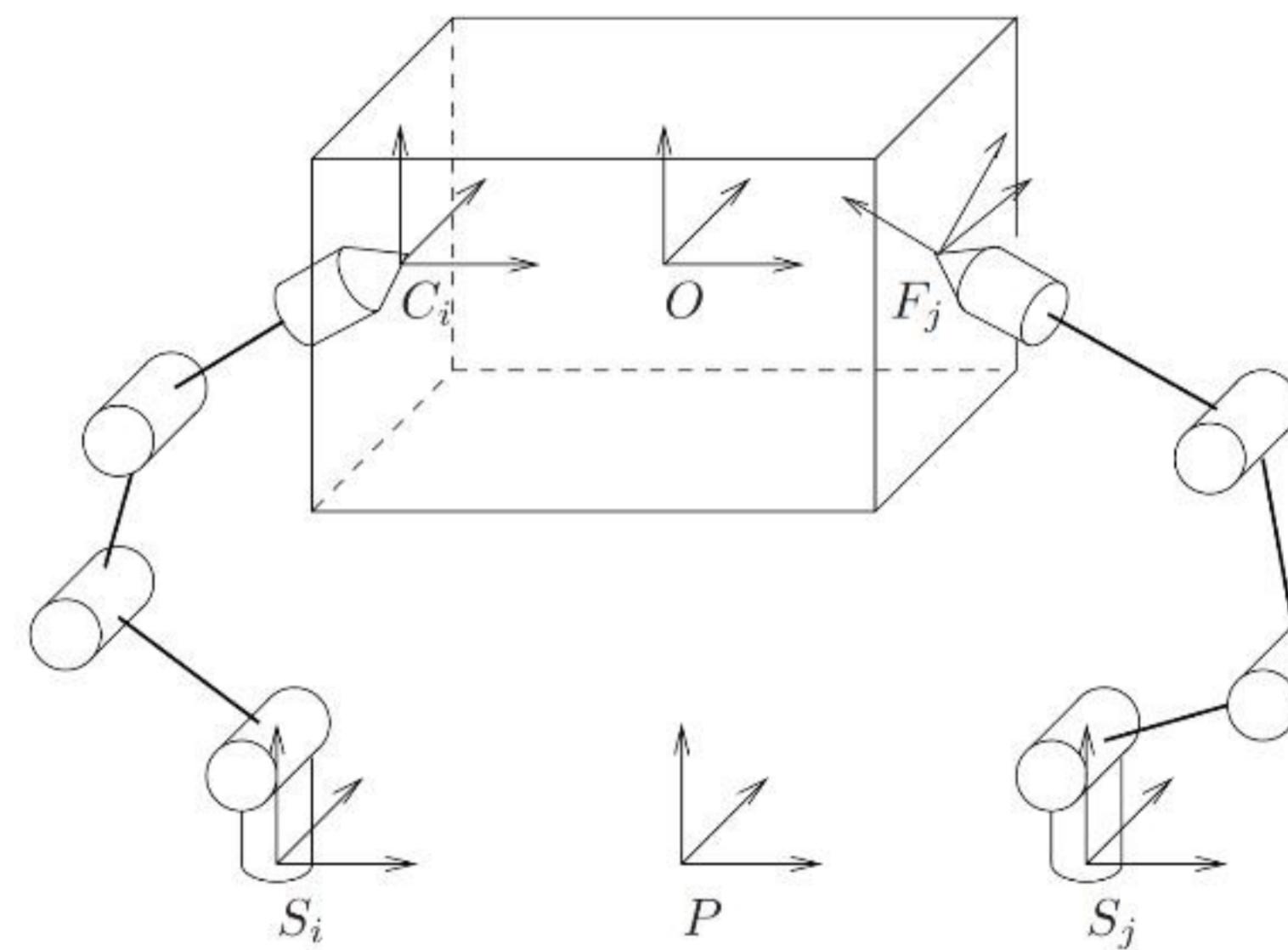


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Step 2: Adding Constraints

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- Assume the body frame twist of the box is ξ^b .
- Then, $\xi^{C_i} = [\text{Ad}_{C_i \rightarrow b}] \xi^b \quad \forall i$.

"All contact frames must move in accordance with the movement of the arms."

- The body frame twist of the contact frame is ξ^{C_i} .
- We can develop the hand Jacobian as $\xi^{C_i} = J^{C_i}(\theta) \dot{\theta} \quad \forall i$.

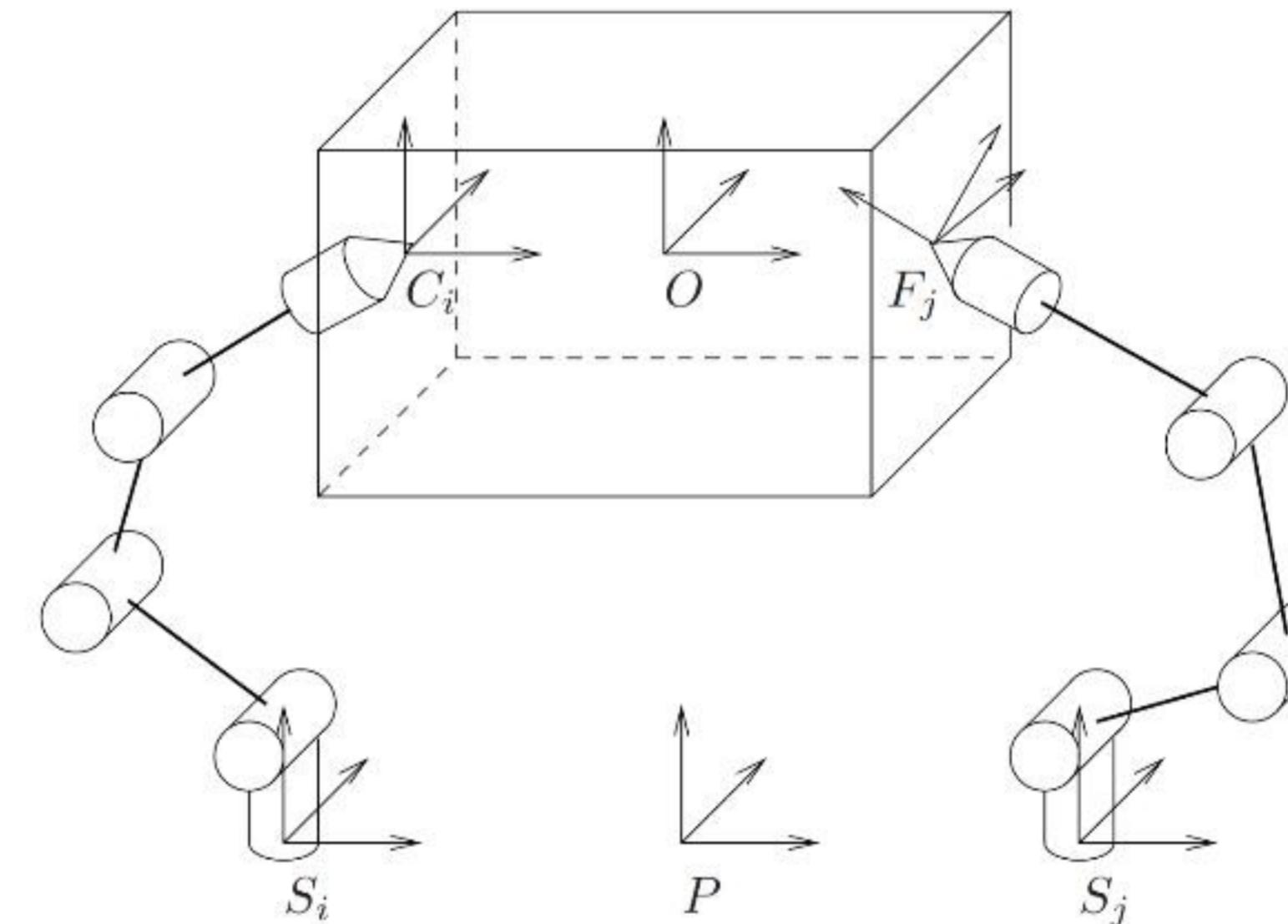


Figure 5.14: Grasp coordinate frames.

Constraints for Augmented Generalized Velocities

- We put all constraints over generalized velocities in a linear equation:

$$A(q)\dot{q} = 0, \quad \text{where } A(q) \in \mathbb{R}^{m \times n}$$

- m : number of constraints
- n : dimension of generalized coordinates
- After some infinitesimal time internal δt , the displacement of the generalized coordinates is $\delta q = \dot{q} \delta t$. We rewrite the constraints as

$$A(q)\delta q = 0$$

A Few More Words on Constraints

- A constraint on a mechanical system restricts the motion of the system by limiting the set of paths that the system can follow.
- **Holonomic constraints:**
 - Roughly speaking, the constraint $A(q)\delta q = 0$ does not depend on \dot{q} (e.g., our grasping example)
 - Strictly speaking, a constraint is said to be **holonomic** if it restricts the motion of the system to a smooth hyper-surface in the (unconstrained) space \mathcal{C} .
- Example of **Nonholonomic constraints:**
 - No-side-slip constraint of wheeled robot: Due to *rolling friction*, a car can only move along the velocity direction, thus the constraint is velocity-based ($A(q, \dot{q})\delta q = 0$), and cannot be written as purely position-based.

Variational Method

Some Tricky (but Sloppy) Math is Coming!



source: *Plants vs Zombies V3*

Background: Unconstrained Function Optimization

- For some $f : \mathbb{R}^n \rightarrow \mathbb{R}$, let us solve the optimization problem

$$\underset{x}{\text{minimize}} \quad f(x)$$

- The first-order Taylor's expansion of $f(x)$ is

$$f(x + \delta x) - f(x) = \nabla f(x)^T \delta x + o(\delta x)$$

- The optimal point must be a *stationary point*, i.e., *the first-order term diminishes*:

$$\nabla f(x^*)^T \delta x = 0$$

- Because the problem is unconstrained, δx can vary along any direction.
- Therefore, $\nabla f(x^*) = 0$ is a necessary condition.

Background: Constrained Function Optimization

- For some $f : \mathbb{R}^n \rightarrow \mathbb{R}$, let us solve the optimization problem

$$\begin{array}{ll}\text{minimize}_x & f(x) \\ \text{subject to} & Ax = 0, \quad A \in \mathbb{R}^{m \times n}\end{array}$$

- The first-order Taylor's expansion of $f(x)$ is

$$f(x + \delta x) - f(x) = \nabla f(x)^T \delta x + o(\delta x)$$

- The optimal point must be a *stationary point*, i.e., *the first-order term diminishes for variations in the constraint set*:

$$\forall \delta x : A\delta x = 0, \nabla f(x^*)^T \delta x = 0$$

- $\delta x \in \text{Null}(A)$ and $\nabla f(x^*)^T \perp \delta x$ implies that $\nabla f(x^*) = -A^T \lambda$ for some $\lambda \in \mathbb{R}^m$

Background: Constrained Function Optimization

- Optimality Condition:

$$\begin{cases} \nabla f(x) + A^T \lambda = 0 & \text{(stationarity)} \\ Ax = 0 & \text{(feasibility)} \end{cases}$$

Principle of Stationary Action (Unconstrained)

- Given some $q : [t_1, t_2] \rightarrow \mathcal{C}$, assume a small perturbation function $\delta q : [t_1, t_2] \rightarrow \mathcal{C}$ such that

$$\delta q(t_1) = \delta q(t_2) = 0$$

- For the functional $S[q] = \int_{t_1}^{t_2} L(q, \dot{q}) dt$, the first-order variation is

$$\begin{aligned}\delta S &= S[q + \delta q] - S[q] = \int_{t_1}^{t_2} L(q + \delta q, \dot{q} + \delta \dot{q}) - L(q, \dot{q}) \\&= \int_{t_1}^{t_2} L(q, \dot{q}) + \left(\frac{\partial L}{\partial q} \right)^T \delta q + \left(\frac{\partial L}{\partial \dot{q}} \right)^T \delta \dot{q} - L(q, \dot{q}) + o(q) \\&= \int_{t_1}^{t_2} \left(\frac{\partial L}{\partial q} \right)^T \delta q + \left(\frac{\partial L}{\partial \dot{q}} \right)^T \delta \dot{q} + o(q) = \int_{t_1}^{t_2} \left(\frac{\partial L}{\partial q} \right)^T \delta q + \left(\frac{\partial L}{\partial \dot{q}} \right)^T \frac{d}{dt} \delta q + o(q) \\&= \int_{t_1}^{t_2} \left(\frac{\partial L}{\partial q} \right)^T \delta q + \left. \frac{\partial L}{\partial \dot{q}} \delta q \right|_{t_1}^{t_2} - \left(\frac{d}{dt} \frac{\partial L}{\partial \dot{q}} \right)^T \delta q + o(t) = \int_{t_1}^{t_2} \left(\frac{\partial L}{\partial q} - \frac{d}{dt} \frac{\partial L}{\partial \dot{q}} \right)^T \delta q + o(t)\end{aligned}$$

- Because the problem is unconstrained, δq can vary along any direction.
- Therefore, $\frac{d}{dt} \frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q} = 0$ is the condition for stationary solution.

Principle of Stationary Action (Holonomically Constrained)

- Given some $q : [t_1, t_2] \rightarrow \mathcal{C}$, assume a small perturbation function $\delta q : [t_1, t_2] \rightarrow \mathcal{C}$ such that

$$\delta q(t_1) = \delta q(t_2) = 0, \quad A(q)\delta q = 0$$

- For the functional $S[q] = \int_{t_1}^{t_2} L(q, \dot{q}) dt$, the first-order variation is

$$\delta S = S[q + \delta q] - S[q] = \int_{t_1}^{t_2} \left(\frac{\partial L}{\partial q} - \frac{d}{dt} \frac{\partial L}{\partial \dot{q}} \right)^T \delta q + o(t)$$

- The optimal point must be a *stationary point*, i.e., *the first-order term diminishes for variations in the constraint set*:

$$\forall \delta q : A(q)\delta q = 0, \int_{t_1}^{t_2} \left(\frac{\partial L}{\partial q} - \frac{d}{dt} \frac{\partial L}{\partial \dot{q}} \right)^T \delta q = 0$$

- $\forall \delta q \in \text{Null}(A(q))$, $\frac{d}{dt} \frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q} \perp \delta q$, which implies that $\frac{d}{dt} \frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q} = -A(q)^T \lambda$ for some $\lambda \in \mathbb{R}^m$

Principle of Stationary Action (Holonomically Constrained)

- Stationary point condition (no external force):

$$\begin{cases} \frac{d}{dt} \frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q} + A(q)^T \lambda = 0 & \text{(stationarity)} \\ A(q)\dot{q} = 0 & \text{(feasibility)} \end{cases}$$

- When external force exists,

$$\begin{cases} \frac{d}{dt} \frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q} + A(q)^T \lambda = \mathbf{F} & \text{(stationarity)} \\ A(q)\dot{q} = 0 & \text{(feasibility)} \end{cases}$$

- **Constraint Force:** $\mathbf{F}_{cons} = A(q)^T \lambda$
- Note that $\mathbf{F}_{cons}^T \dot{q} = 0$, thus internal force does no work.

Principle of Stationary Action (Holonomically Constrained)

- The feasibility constraint can be rewritten as:

$$A(\dot{q})\dot{q} + A(q)\ddot{q} = 0$$

- The system of equations is thus:

$$\begin{cases} \frac{d}{dt} \frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q} + A(q)^T \lambda = \mathbf{F} & \text{(stationarity)} \\ A(\dot{q})\dot{q} + A(q)\ddot{q} = 0 & \text{(feasibility)} \end{cases}$$

Principle of Stationary Action (Holonomically Constrained)

- Our object manipulation example:

$$\begin{cases} M(q)\ddot{q} + C(q, \dot{q})\dot{q} + g(q, \dot{q}) + A(q)^T \lambda = \mathbf{F} & \text{(stationarity)} \\ \dot{A}(q)\dot{q} + A(q)\ddot{q} = 0 & \text{(feasibility)} \end{cases}$$

therefore,

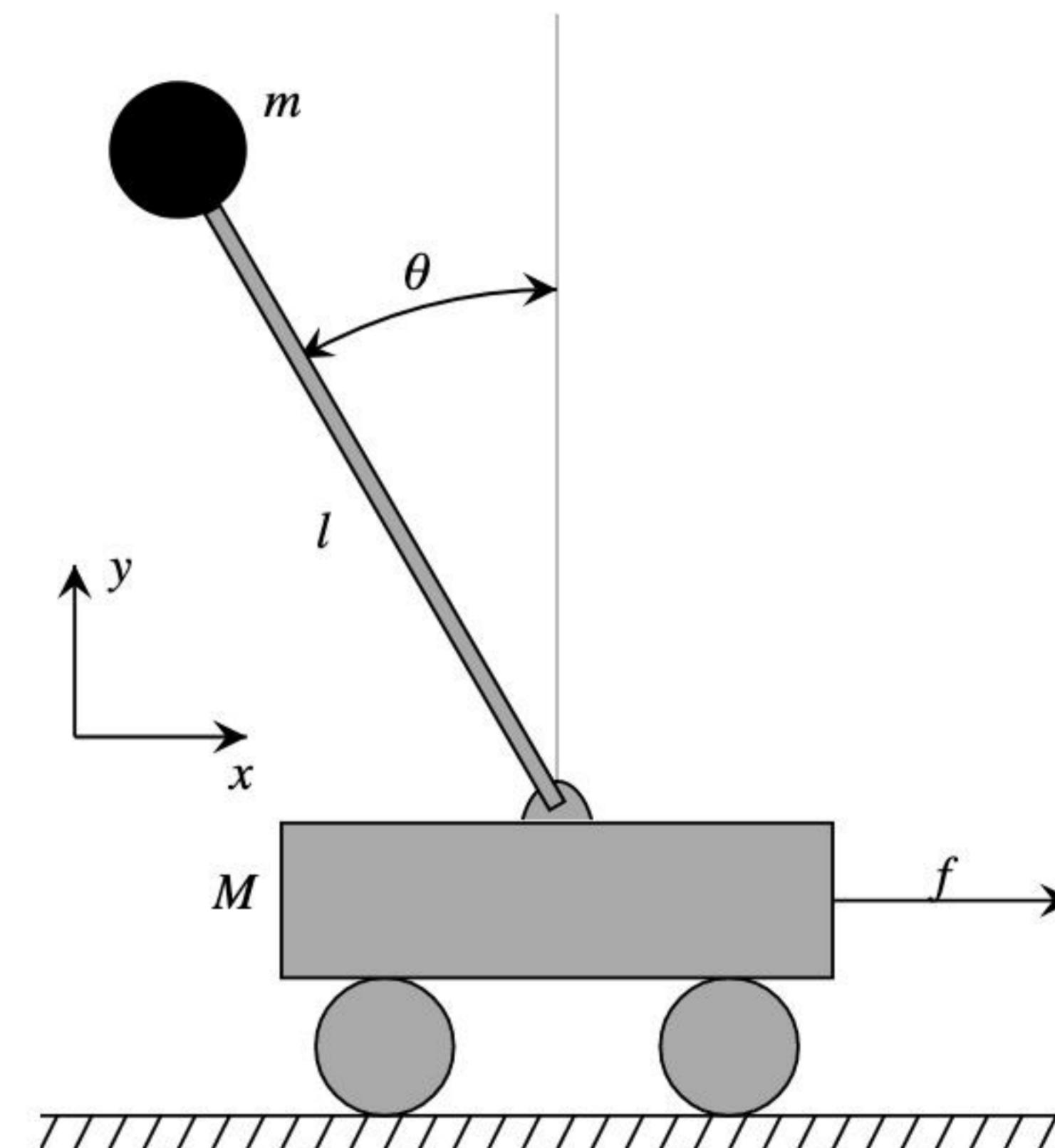
$$\lambda = (AM^{-1}A^T)^{-1}(AM^{-1}(F - C\dot{q} - N) + \dot{A}\dot{q})$$

where $AM^{-1}A^T$ is full rank if the constraints are independent.

- Note: $M(q)$, $C(q, \dot{q})$, and $g(q, \dot{q})$ are computed from the two arms and the object.

Example: Cart Pole

- Kinetic energy: $T = \frac{1}{2}Mv_1^2 + \frac{1}{2}mv_2^2$
- Assume the joint position is $[x_1(t), 0]^T$ and the point mass position is $[x_2(t), y_2(t)]$, then
 - $v_1^2 = \dot{x}_1^2$
 - $v_2^2 = \dot{x}_2^2 + \dot{y}_2^2$
- Potential energy: $V = mgy_2$
- Constraint: $C(x_1, x_2, y_2) = (x_1 - x_2)^2 + y_2^2 - l^2$
 - $A(q)\dot{q} = \nabla_q C(q)\dot{q}$



A schematic drawing of the inverted pendulum on a cart. The rod is considered massless.

https://en.wikipedia.org/wiki/Inverted_pendulum

Example: Cart Pole

```
from sympy import *
from sympy.physics.vector import dynamicsymbols
from sympy import diff, Symbol

M, m, t, g, l, lmd, f = symbols('M m t g l lmd f')

x1 = dynamicsymbols('x1')
x2 = dynamicsymbols('x2')
th = dynamicsymbols('th')
q = Matrix([x1, x2, th])

dqdt = diff(q, t)
dx1dt = dqdt[0]
dx2dt = dqdt[1]
dthdt = dqdt[2]
y2 = l * cos(th)
dy2dt = diff(y2, t)

T = 0.5 * M * dx1dt * dx1dt + 0.5 * m * (dx2dt * dx2dt + dy2dt * dy2dt)
V = m * g * y2
C = (x1-x2)**2+y2**2-l**2
L = T - V + C*lmd

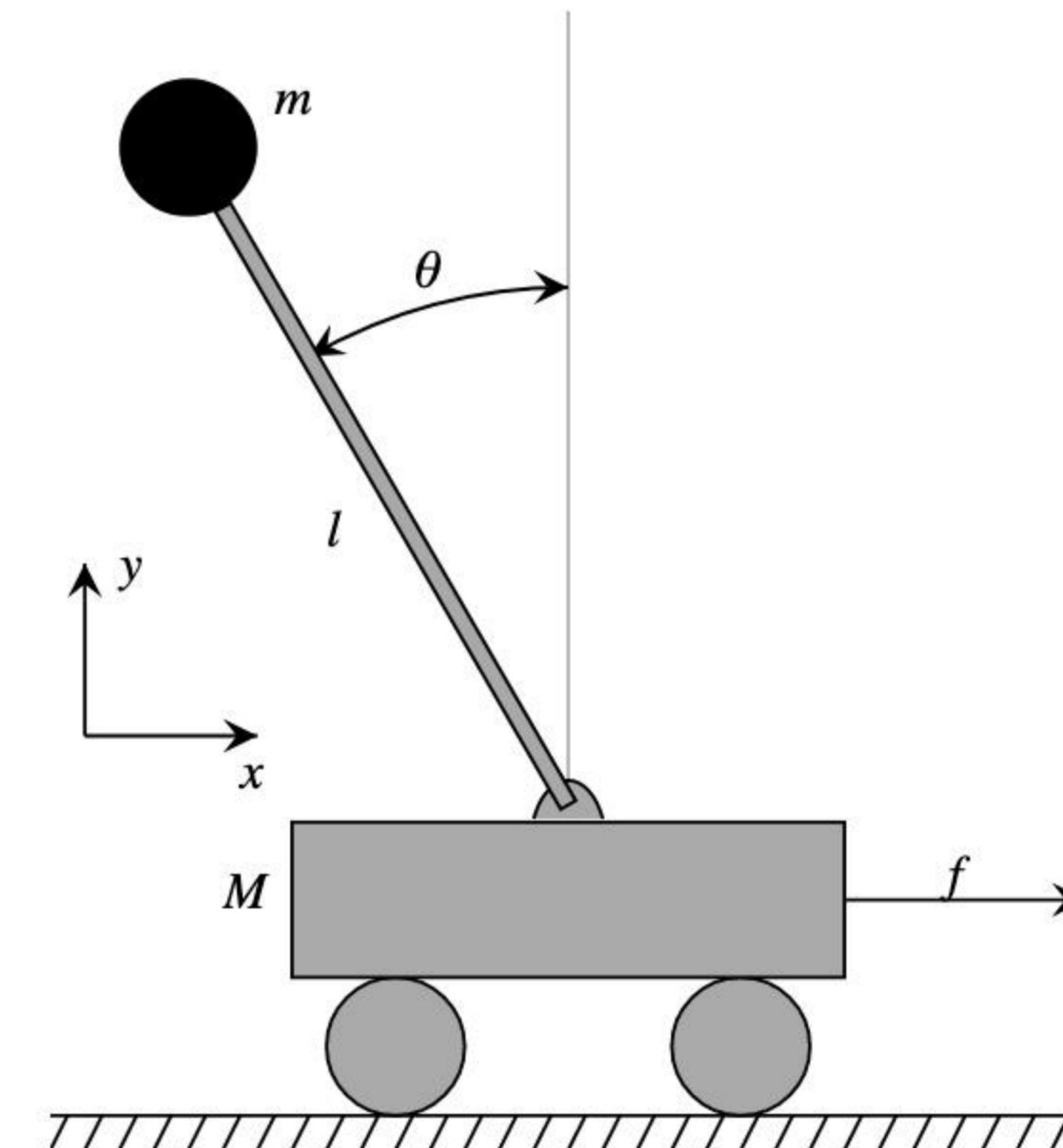
Lagrangian = Lag1 = diff(diff(L, dqdt), t) - diff(L, q)

F = Matrix([f, 0, 0])
stationarity = Lagrangian - F

A = diff(C, q).T

feasibility = diff(A, t)* dqdt + A * diff(dqdt, t)

ddqdt = diff(dqdt, t)
sol = simplify(solve((stationarity[0], stationarity[1], stationarity[2], feasibility), lmd,
ddqdt[0], ddqdt[1], ddqdt[2]))
```



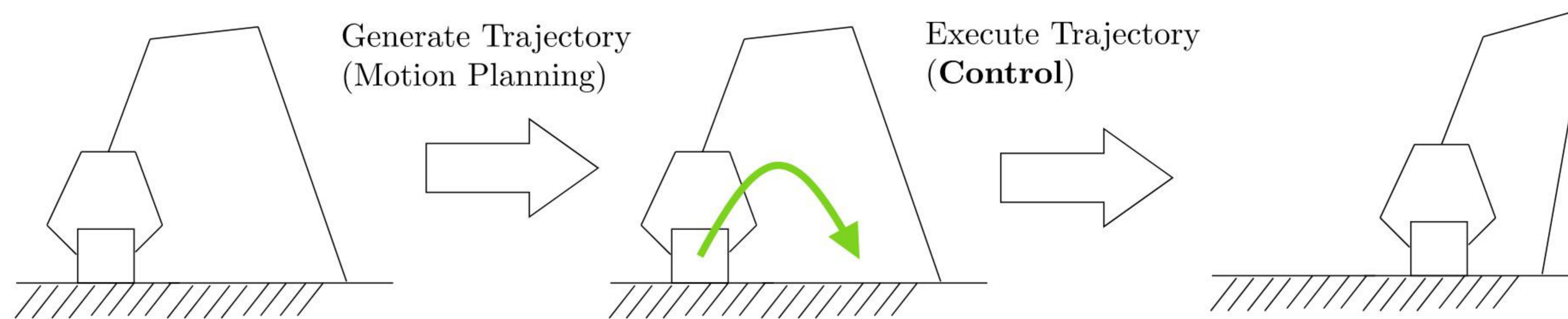
A schematic drawing of the inverted pendulum on a cart. The rod is considered massless.

https://en.wikipedia.org/wiki/Inverted_pendulum

Control

Plan versus Control

- Suppose we have a robot, how do we use it to move objects?



- Motion planning algorithms can generate a trajectory (position, velocity, and acceleration) of the robot.
- Now we need to know how to **control** the robot to follow such a trajectory.

Control

- Let us look at what we have
 - A *desired* trajectory to follow: $(q_d, \dot{q}_d, \ddot{q}_d)$
 - Forward dynamics $\ddot{q} = FD(\mathbf{F}; q, \dot{q})$
 - Inverse dynamics $\mathbf{F} = ID(\ddot{q}; q, \dot{q})$
- Ideally, we just use inverse dynamics to compute \ddot{q} at every moment to match \ddot{q}_d , and we are done!

Control

- Let us look at what we have
 - A *desired* trajectory to follow: $(q_d, \dot{q}_d, \ddot{q}_d)$
 - Forward dynamics $\ddot{q} = FD(\mathbf{F}; q, \dot{q})$
 - Inverse dynamics $\mathbf{F} = ID(\ddot{q}; q, \dot{q})$
- Ideally, we just use inverse dynamics to compute \ddot{q} at every moment to match \ddot{q}_d , and we are done!
- However, the real world is not perfect. There will be many sources of error, and error accumulates if only acceleration is matched.
- We use **control** to deal with delay, overshoot, or steady-state error, and ensure stability.

$$e = q - q_d \quad (\text{steady-state error})$$

PID Controller

Feedforward and Feedback Control

- We need some force to match \ddot{q}_d . This component is called the **feed-forward component**, which comes from $\text{ID}(\cdot)$:

$$\mathbf{F}_{ff} = \text{ID}(\ddot{q}_d; q, \dot{q})$$

- We also need some additional force to correct the steady-state error, which is called the **feedback component**:

$$\mathbf{F}_{fb} = M(q)(-K_v \dot{e} - K_p e)$$

where $M(q)$ is the inertia of the system.

Computed Torque Control Law

- Let us apply the previous control law to our single-arm robot derived last lecture.
- Inverse dynamics equation:

$$\tau = M(\theta)\ddot{\theta} + C(\theta, \dot{\theta})\dot{\theta} + g(\theta)$$

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- Therefore,

$$\begin{aligned}\tau_{ff} &= M(\theta)\ddot{\theta}_d + C(\theta, \dot{\theta})\dot{\theta} + g(\theta) \\ \tau_{fb} &= M(\theta)(-K_v\dot{e} - K_p e)\end{aligned}$$

where $K_v, K_p \in \mathbb{S}^+$ are constant matrices (\mathbb{S}^+ : positive-semidefinite matrices cone).

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- Combine them together, and the *computed torque control law* is:

$$\tau = M(\theta)(\ddot{\theta}_d - K_v\dot{e} - K_p e) + C(\theta, \dot{\theta})\dot{\theta} + g(\theta) \quad (\text{computed torque control})$$

Convergence Analysis

- We will control the system by the torque

$$\tau = M(\theta)(\ddot{\theta}_d - K_v \dot{e} - K_p e) + C(\theta, \dot{\theta})\dot{\theta} + g(\theta) \quad (1)$$

- However, inverse dynamics equation tells us that the acceleration from τ is

$$\tau = M(\theta)\ddot{\theta} + C(\theta, \dot{\theta})\dot{\theta} + g(\theta) \quad (2)$$

- Subtracting (2) from (1) and cancel $M(\theta)$, we get the error equation:

$$\ddot{e} + K_v \dot{e} + K_p e = 0$$

Convergence Analysis

$$\ddot{e} + K_v \dot{e} + K_p e = 0$$

- Because $K_v, K_p \in \mathbb{S}^+$, by the theory of ODE, $e(t) = \mathcal{O}(e^{\alpha t})$, $\alpha \leq 0$.
- We say that the computed torque control law has **exponential convergence rate**.

- Sometimes, we do not have the inverse dynamics equation of the system, but we still hope the controller to work
(Isn't that the advertisement of model-free reinforcement learning?)

PD Control

- The PD control law has the form:

$$\tau = -K_v \dot{e} - K_p e \quad (\text{PD control})$$

where $K_v, K_p \in \mathbb{S}^+$ and $e = \theta - \theta_d$.

- Does not compute inverse dynamics at all.
- No theoretical guarantee in general.
- Not practical in general
- e may converge; however, it usually does not converge to 0.

PID Control

- PID Control law has the form

$$\tau = -K_v \dot{e} - K_p e - K_i \int_0^t e(t) dt \quad (\text{PID control})$$

where $K_v, K_p, K_i \in \mathbb{S}^+$ and $e = \theta - \theta_d$.

- K_i term: accumulate errors over all past time steps.
- Inherits all the issues of PD, except
- When e converges, it usually converges to 0.
- Widely used in practice.

Augmented PID Control

- Augmented PID control law has the form:

$$\tau = ID(\ddot{\theta}_d) - K_v \dot{e} - K_p e$$

- Recall the computed torque control law:

$$\tau = ID(\ddot{\theta}_d) - M(\theta)K_v \dot{e} - M(\theta)K_p e$$

- Compared with the computed torque control law, K 's may be harder to tune.
- Does result in exponential convergence for $K_v, K_p \in \mathbb{S}^+$

Tuning PID

Effects of *increasing* a parameter independently

Parameter	Rise time	Overshoot	Settling time	Steady-state error	Stability
K_p	Decrease	Increase	Small change	Decrease	Degrade
K_i	Decrease	Increase	Increase	Eliminate	Degrade
K_d	Minor change	Decrease	Decrease	No effect in theory	Improve if K_d small

Kiam Heong Ang; Chong, G.; Yun Li (2005). "PID control system analysis, design, and technology". *IEEE Transactions on Control Systems Technology*. 13 (4): 559–576.

Jinghua Zhong (Spring 2006). "PID Controller Tuning: A Short Tutorial" (PDF). Archived from the original on 2015-04-21. Retrieved 2011-04-04.

https://en.wikipedia.org/wiki/PID_controller#cite_note-24

L10: Constraint Force Computation, Control

- Constitutive Lagrange Method
- Variational Method
- Contact
- PGT Control

Has Su
Spring 2011

Agenda

- Constitutive Lagrange Method
- Variational Method
- Contact
- PGT Control

Click to return to other modules.

Review: Lagrangian Function

$\dot{q}_1 = \frac{d}{dt} q_1$ for the generalized coordinate q_1

$L(q, \dot{q})$ is the generalized potential function

$L(q, \dot{q}) = T(q) - V(q)$

$T(q)$ is the kinetic energy

$V(q)$ is the potential energy

Review: The Principle of Stationary Action

The principle of stationary action states that the system follows the path that minimizes the action integral:

$$S[q] = \int_{t_1}^{t_2} [L(q, \dot{q}) - V(q)] dt$$

where $q(t_1) = q_1$ and $q(t_2) = q_2$ are the initial and final positions of the system.

For a system with n degrees of freedom, the action integral is:

$$S[q] = \int_{t_1}^{t_2} [T(q) - V(q)] dt$$

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$$\frac{\delta S}{\delta q_i} = 0$$

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Review: Euler-Lagrange Equation

When there are external force components F_x, F_y, F_z acting on the system, the principle of stationary action becomes:

$$F_x = \frac{\partial L}{\partial q_x} - \frac{d}{dt} \frac{\partial L}{\partial \dot{q}_x}$$

(Euler-Lagrange Equation)

Example: Robot Arm

A two-link robot arm is shown below. It consists of two rigid links connected by a joint at the center of mass of the first link. The first link has a length of 1 m and a mass of 2 kg. The second link has a length of 0.5 m and a mass of 1 kg. The center of mass of the first link is located at a distance of 0.5 m from the joint. The center of mass of the second link is located at a distance of 0.25 m from the joint. The joint rotates with an angular velocity of 1 rad/s. The gravitational acceleration is 9.81 m/s². The robot arm is subject to a horizontal force of 10 N applied at its center of mass.

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