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### One-Way Quantum Computation

### Wang Kun

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December 27, 2014

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## QCM v.s. 1WQC

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- Quantum computation is usually described in terms of quantum circuit model (QCM).
- One-way quantum computation (1WQC) is a very different way of describing quantum computation.
- Different models of quantum computation give us
  - New approaches to quantum computer design.
  - Different ways of understanding quantum computation.

### General framework for 1WQC

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■ General framework for 1WQC

Step1. A special a special entangled state called a *Cluster State* or *Graph State* is prepared.



## General framework for 1WQC

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■ General framework for 1WQC

Step1. A special a special entangled state called a *Cluster State* or *Graph State* is prepared.



Step2. The qubits are then measured *one* at a time (single-qubit measurements), with local measurements, in a sequence of *different measurement bases*.

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## Single qubit rotations

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• A single qubit rotation through angle  $\phi$  around the X, Y, or Z axis corresponds to the operator

$$R_x(\phi) = \exp[-i\frac{\phi}{2}X]$$

# Single qubit rotations

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Reference:

■ A single qubit rotation through angle  $\phi$  around the X, Y, or Z axis corresponds to the operator

$$R_x(\phi) = \exp[-i\frac{\phi}{2}X]$$

■ Euler Theorem: Any rotation on a sphere can be formed from three successive rotations about mutually orthogonal axes.

# Single qubit rotations

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• A single qubit rotation through angle  $\phi$  around the X, Y, or Z axis corresponds to the operator

$$R_x(\phi) = \exp[-i\frac{\phi}{2}X]$$

- **Euler Theorem:** Any rotation on a sphere can be formed from three successive rotations about mutually orthogonal axes.
- Thus any single qubit rotation can be expressed as

$$R_z(\alpha)$$
  $R_z(\beta)$   $R_z(\gamma)$ 

# Single qubit measurements

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■ Single qubit measurements are important in 1WQC.

# Single qubit measurements

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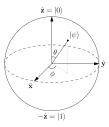
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- Single qubit measurements are important in 1WQC.
- We can use a Bloch sphere representation:



# Single qubit measurements

### 1WQC

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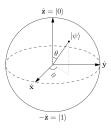
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- Single qubit measurements are important in 1WQC.
- We can use a Bloch sphere representation:



- Single qubit measurements applied in 1WQC:
  - **Z** measurement projects on  $|0\rangle$ ,  $|1\rangle$ .
  - **X** measurement projects on  $|\pm\rangle = [|0\rangle + |1\rangle]/\sqrt{2}$ .
  - "X-Y plane" measurement  $\cos(\alpha)X + \sin(\alpha)Y$  projects onto  $[|0\rangle \pm e^{-i\alpha}|1\rangle]/\sqrt{2}$ .

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# Graph states

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- Graph states are entangled states over many qubits, with a very simple description.
- A collection of vertices, some of whom are linked by an edge. This leads to their name "graph states".
- Each vertex represents a qubit, each edge represents an entangle operation.
- The graph structure gives a method of constructing the graph states.

# Constructing the graph states

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- With every vertex we associate a qubit, initially prepared in state  $|+\rangle = [|0\rangle + |1\rangle]/\sqrt{2}$ .
- The graph state is obtained by applying a CZ gate (entangle operation) to each edge on the graph.
- A CZ gate has the following truth table and is order-independent:

$$|00\rangle \mapsto |00\rangle, \ |01\rangle \mapsto |01\rangle, \ |10\rangle \mapsto |10\rangle, \ |00\rangle \mapsto -|11\rangle$$

$$CZ|i\rangle|j\rangle = (-1)^{ij}|i\rangle|j\rangle, \ i, j \in \{0, 1\}$$

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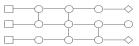
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■ The term *cluster state* is reserved for a graph state on a square lattice:



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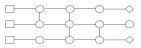
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■ The term *cluster state* is reserved for a graph state on a square lattice:



In one-way quantum computation, by simply making single-qubit measurements on cluster states of sufficient size, any quantum circuit can be simulated.

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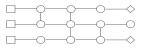
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■ The term *cluster state* is reserved for a graph state on a square lattice:



- In one-way quantum computation, by simply making single-qubit measurements on cluster states of sufficient size, any quantum circuit can be simulated.
- Extra flexibility of graph states require far fewer qubits when implementing the computation.

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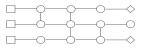
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■ The term *cluster state* is reserved for a graph state on a square lattice:



- In one-way quantum computation, by simply making single-qubit measurements on cluster states of sufficient size, any quantum circuit can be simulated.
- Extra flexibility of graph states require far fewer qubits when implementing the computation.
- Regular layout of cluster states means that they can be generated efficiently in physical implementation.

## One-qubit teleportation

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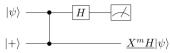
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■ The key idea in 1WQC is known as *one-bit teleportation* 



# One-qubit teleportation

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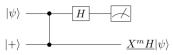
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■ The key idea in 1WQC is known as *one-bit teleportation* 



 $\blacksquare$  *m* is the outcome (0 or 1) of the computational basis measurement on the first qubit.

# One-qubit teleportation

### 1WQC

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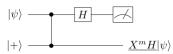
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■ The key idea in 1WQC is known as *one-bit teleportation* 



- $\blacksquare$  *m* is the outcome (0 or 1) of the computational basis measurement on the first qubit.
- Let  $|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$ , then the evolution is

$$\begin{split} |\psi\rangle|+\rangle & \xrightarrow{\text{CZ}} & \alpha|0\rangle|+\rangle + \beta|1\rangle|-\rangle \\ & \xrightarrow{\text{H}\otimes\text{I}} & \alpha|+\rangle|+\rangle + \beta|-\rangle|-\rangle \\ & = & [|0\rangle\otimes\text{H}|\psi\rangle + |1\rangle\otimes\text{XH}|\psi\rangle]/\sqrt{2} \\ & \xrightarrow{\text{M}} & \text{X}^m\text{H}|\psi\rangle \end{split}$$

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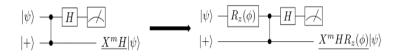
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■ Imagine now that a Z-rotation  $R_z(\phi)$  is added at the start of this circuit



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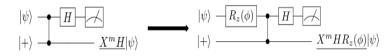
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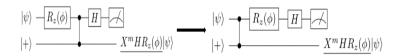
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■ Imagine now that a Z-rotation  $R_z(\phi)$  is added at the start of this circuit



Note that  $[R_z(\phi), CZ] = 0$ , thus we can rearrange the order of the gates



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- Introduce a new symbol which explicitly states the measurement's basis states  $X^m|0\rangle$ 
  - If measurement result is m = 0, then the measured computational basis is  $|0\rangle$ ;

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■ Introduce a new symbol which explicitly states the measurement's basis states  $X^m|0\rangle$ 

- If measurement result is m = 0, then the measured computational basis is  $|0\rangle$ ;
- If measurement result is m=1, then the measured computational basis is  $X|0\rangle=|1\rangle$ .

$$|\psi\rangle \longrightarrow R_z(\phi) - H \longrightarrow X^m H R_z(\phi) |\psi\rangle \longrightarrow |\psi\rangle \longrightarrow R_z(\phi) - H \longrightarrow X^m H R_z(\phi) |\psi\rangle$$

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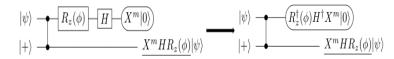
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Unitary operators before a measurement can be "absorbed" into the measurement with result reserved.



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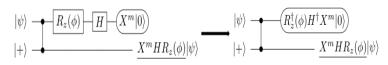
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Unitary operators before a measurement can be "absorbed" into the measurement with result reserved.



 $\blacksquare H^{\dagger}X^m|0\rangle = Z^m|+\rangle$ ,  $R_z^{\dagger}(\phi) = R_z(-\phi)$ , thus

$$|\psi\rangle \longrightarrow \underbrace{\left(R_z^\dagger(\phi)H^\dagger X^m|0\rangle\right)}_{|+\rangle} \longrightarrow \underbrace{\left|\psi\rangle}_{|+\rangle} \xrightarrow{\left(R_z(-\phi)Z^m|+\right)\right)}_{|+\rangle}$$

### Kernel pattern

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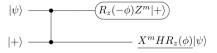
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• From the one-bit teleportation protocol, we derive the kernel pattern in 1WQC:



# Kernel pattern

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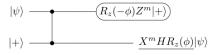
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Reference:

■ From the one-bit teleportation protocol, we derive the kernel pattern in 1WQC:



■ In the Bloch sphere picture, the measurement with basis  $R_z(-\phi)|+\rangle$  and  $R_z(-\phi)|-\rangle$  lies on the X-Y plane.

## Kernel pattern

### 1WQC

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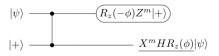
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■ From the one-bit teleportation protocol, we derive the kernel pattern in 1WQC:



- In the Bloch sphere picture, the measurement with basis  $R_z(-\phi)|+\rangle$  and  $R_z(-\phi)|-\rangle$  lies on the X-Y plane.
- Actually, the first qubit is measured in the basis

$$\frac{1}{\sqrt{2}}[|0]\rangle \pm e^{i\phi}|1\rangle]$$

## Kernel pattern (cont.)

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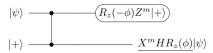
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■ The kernel pattern is implementing  $HR_z(\phi)$  modulo a Pauli correction  $X^m$ 

$$\underbrace{X^m}_{\text{Random "by-product" Implemented unitary operator}} \underbrace{HR_z(\phi)}_{\text{Random unitary operator}}$$

■ *m* is the measurement result of first qubit, that's why we say measurement based!

## Cluster states in kernel pattern

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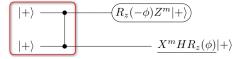
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■ Imagine we set the input state to  $|+\rangle$ 



## Cluster states in kernel pattern

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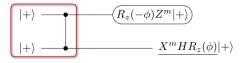
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In this case, the state before measurement is a two-qubit cluster state.

## Cluster states in kernel pattern

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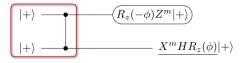
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■ Imagine we set the input state to  $|+\rangle$ 



- In this case, the state before measurement is a two-qubit cluster state.
- The circuit consists of a single qubit measurement on a cluster state.

### Cluster states in kernel pattern

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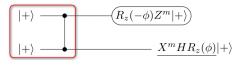
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■ Imagine we set the input state to  $|+\rangle$ 



- In this case, the state before measurement is a two-qubit cluster state.
- The circuit consists of a single qubit measurement on a cluster state.
- How can we connect the kernel patterns to implement more complicated unitary operators?

#### 1WOC

Connecting Kernel Patterns

**Euler's Theorem:** any single-qubit rotation U can be decomposed as a product of three rotations

$$\forall U \in SU(2), \exists \alpha, \beta, \gamma \text{ such that } U = R_z(\gamma)R_x(\beta)R_z(\alpha).$$

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**Euler's Theorem:** any single-qubit rotation U can be decomposed as a product of three rotations

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• Consider a sequence of three  $HR_z(\phi)$  operations

$$HR_z(\alpha)$$
  $HR_z(\beta)$   $HR_z(\gamma)$ 

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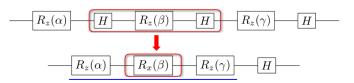
**Euler's Theorem:** any single-qubit rotation U can be decomposed as a product of three rotations

$$\forall U \in SU(2), \exists \alpha, \beta, \gamma \text{ such that } U = R_z(\gamma)R_x(\beta)R_z(\alpha).$$

■ Consider a sequence of three  $HR_z(\phi)$  operations

$$HR_z(\alpha)$$
  $HR_z(\beta)$   $HR_z(\gamma)$ 

 $lue{}$  This gives a circuit for simulating the rotation U



# Concatenating the kernel patterns

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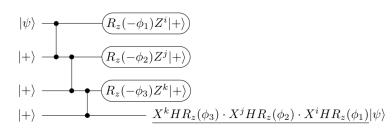
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• We can simulate the above circuit by concatenating three copies of kernel pattern



which implements the following unitary operator:

$$- \boxed{HR_z(\phi_1)} - \boxed{X^i} - \boxed{HR_z(\phi_2)} - \boxed{X^j} - \boxed{HR_z(\phi_3)} - \boxed{X^k} -$$

### By-product Pauli operators

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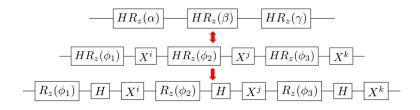
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• Measurement-dependent by-product operators  $X^i$  are sandwiched among the H and  $R_z$  gates.

# By-product Pauli operators

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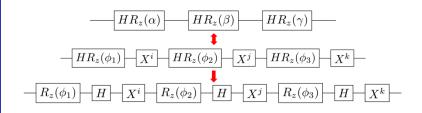
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- Measurement-dependent by-product operators  $X^i$  are sandwiched among the H and  $R_z$  gates.
- We need to change the orders between  $X^i$  and H and  $R_z$ , so that all H and  $R_z$  operations are performed before  $X^i$  operations.

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■ Consider the permutation relation between a rotation  $R_p(\alpha)$  and a Pauli P.

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- Consider the permutation relation between a rotation  $R_p(\alpha)$  and a Pauli P.
- Pauli operators either commute or anti-commute:

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- Consider the permutation relation between a rotation  $R_p(\alpha)$  and a Pauli P.
- Pauli operators either commute or anti-commute:
  - **1** p and P anti-commute. In this case

$$\begin{split} R_p(\alpha)P &= \exp[-i\frac{\alpha}{2}p]P &= [\cos(\frac{\alpha}{2}) - i\sin(\frac{\alpha}{2})p]P \\ &= P[\cos(\frac{\alpha}{2}) + i\sin(\frac{\alpha}{2})p] \\ &= PR_p(-\alpha) \end{split}$$

After the order has been interchanged, P is unchanged, the rotation remains a rotation but the rotation angle  $\alpha$  is reversed!

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After the order has been interchanged, P is unchanged, the rotation remains a rotation but the rotation angle  $\alpha$  is reversed!

**2** p and P commute. The two can be interchanged without change, i.e.  $R_p(\alpha)P = PR_p(\alpha)$ .

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■ Consider the permutation relation between H, CZ and a Pauli. These belong to the **Clifford group**.

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- $lue{}$  Consider the permutation relation between H, CZ and a Pauli. These belong to the **Clifford group**.
- lacktriangle For arbitrary Clifford operator C and Pauli operator  $P_1$

$$\exists P_2 \in \{X, Y, Z\}$$
 such that  $CP_1C^{\dagger} = P_2$ 

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$$CP_1 = P_2C$$

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■ Rearranging this gives us a **permutation relation** 

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- We see that on interchanging the order:
  - **1** The Clifford group operator *C* remains invariant,
  - 2 The Pauli operator remains a (possibly different) Pauli operator.

# Permutation of byproducts

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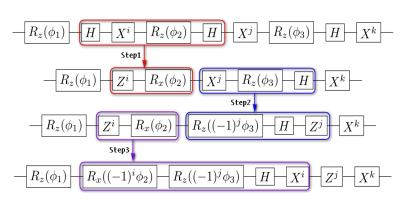
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- We can now apply these permutation rules:
  - **1 Rotation**:  $R_z(\alpha)X = XR_z(-\alpha)$ ,  $R_z(\alpha)Z = ZR_z(\alpha)$ ;
  - **2** Clifford group:  $HX^i = Z^iH$ ,  $HZ^i = X^iH$



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$$R_z(\phi_1) - R_x((-1)^i\phi_2) - R_z((-1)^j\phi_3) - H$$

$$X^i - Z^j - X^k - I$$

■ Choosing  $\phi_1 = \alpha$ ,  $\phi_2 = (-1)^i \beta$ ,  $\phi_3 = (-1)^j \gamma$ , we implement the single-qubit gate  $U = R_z(\gamma) R_x(\beta) R_z(\alpha)$  (up to the Pauli corrections).

$$R_z(\alpha)$$
  $R_z(\beta)$   $R_z(\gamma)$   $H$ 

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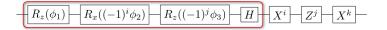
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$$R_z(\alpha)$$
  $R_z(\beta)$   $R_z(\gamma)$   $H$ 

 Hence, we apply kernel patterns to successfully simulate arbitrary single-qubit gates.

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It is convenient to introduce a graphical notation for 1WQC.

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- It is convenient to introduce a graphical notation for 1WQC.
- We are going to represent each qubit by a shape:
  - lacktriangle The "input qubit" by a square  $\Box$
  - lacksquare The "output qubit" by a diamond  $\Diamond$
  - $lue{}$  The other qubits by a circle  $\bigcirc$

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- All qubits are initialized to state  $|+\rangle$ .

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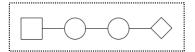
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  - The other qubits by a circle ○
- All qubits are initialized to state  $|+\rangle$ .
- One-dimensional lattice graphical notation for the clusters states of arbitrary single-qubit gates:



### Graph notations (cont.)

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■ The CZ gates (entanglement operation) are represented as lines between these shapes.

### Graph notations (cont.)

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- The CZ gates (entanglement operation) are represented as lines between these shapes.
- Single-qubit measurements will be represented by symbols within the shapes.
  - Greek letters  $\phi$  inside the shape will represent measurements with basis states  $\frac{1}{\sqrt{2}}[|0]\rangle \pm e^{i\phi}|1\rangle]$
  - Letters X or Z inside the shape is used to emphasize the basis states are  $\{|+\rangle, |-\rangle\}$  or  $\{|0\rangle, |1\rangle\}$ .

### Graph notations (cont.)

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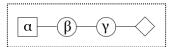
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- One-dimensional lattice graphical notation for simulating the single-qubit gate *U*:



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 For a universal quantum computation, we require two-qubit gates.

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- For a universal quantum computation, we require two-qubit gates.
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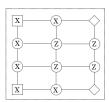
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- For a universal quantum computation, we require two-qubit gates.
- Single qubit rotations plus the *CZ* gate form a universal set.
- A possible two-dimensional lattice graphical notation for simulating the CZ gate:



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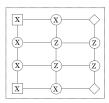
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- For a universal quantum computation, we require two-qubit gates.
- Single qubit rotations plus the CZ gate form a universal set.
- A possible two-dimensional lattice graphical notation for simulating the CZ gate:



Measurement patterns are not unique and subject to the construction approaches.



### Outline

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# Physical implementation of 1WQC

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So far I have described the one-way quantum computation model.

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- So far I have described the one-way quantum computation model.
- It guides us to implement 1WQC in the lab:

Step1. Build a cluster state,

Step2. Measure qubits one by one in specific order.

### Physical implementation of 1WQC

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So far I have described the one-way quantum computation model.

■ It guides us to implement 1WQC in the lab:

Step1. Build a cluster state,

Step2. Measure qubits one by one in specific order.

There are systems where the one-way model has some significant practical advantages.

### Physical implementation of 1WQC (cont.)

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### Optical lattices

- Pro: cluster states can be generated efficiently, already implemented in labs.
- Con: it is (currently!) not possible to perform single-qubit measurements on atoms.

### Physical implementation of 1WQC (cont.)

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### Optical lattices

- Pro: cluster states can be generated efficiently, already implemented in labs.
- Con: it is (currently!) not possible to perform single-qubit measurements on atoms.
- Linear optics
  - **Pro:** efficient strategies for building cluster states.
  - **Con:** single-qubit measurements may fail, thus destroying precious quantum coherence.

### Physical implementation of 1WQC (cont.)

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### Optical lattices

- Pro: cluster states can be generated efficiently, already implemented in labs.
- Con: it is (currently!) not possible to perform single-qubit measurements on atoms.

### Linear optics

- Pro: efficient strategies for building cluster states.
- **Con:** single-qubit measurements may fail, thus destroying precious quantum coherence.

### Cavity QED

- Pro: atoms interact in cavities, implementing CZ gates and cluster state.
- **Con:** hard to implement single-qubit measurements.

# Acknowledgement

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# Thank you!

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