

Quantum Optics: Homework #1

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Problem 1

Show that

$$\begin{aligned}[a, e^{-\alpha a^\dagger a}] &= (e^{-\alpha} - 1) e^{-\alpha a^\dagger a} a, \\ [a^\dagger, e^{-\alpha a^\dagger a}] &= (e^{\alpha} - 1) e^{-\alpha a^\dagger a} a^\dagger,\end{aligned}$$

where α is a parameter.

Solution

Part One

$$\begin{aligned}[a, e^{-\alpha a^\dagger a}] &= [a, \sum_{n=0}^{\infty} \frac{(-\alpha)^n}{n!} (a^\dagger a)^n] \\ &= \sum_{n=0}^{\infty} \frac{(-\alpha)^n}{n!} [a, (a^\dagger a)^n] \\ &= \sum_{n=1}^{\infty} \frac{(-\alpha)^n}{n!} [a, (a^\dagger a)^n] \\ &= \sum_{n=1}^{\infty} \frac{(-\alpha)^n}{n!} [a(a^\dagger a)^{n-1} a^\dagger a - (a^\dagger a)^n a] \\ &= \sum_{n=1}^{\infty} \frac{(-\alpha)^n}{n!} [a(a^\dagger a)^{n-1} a^\dagger - (a^\dagger a)^n] a \\ &= \sum_{n=1}^{\infty} \frac{(-\alpha)^n}{n!} [(aa^\dagger)^n - (a^\dagger a)^n] a \\ &= (e^{-\alpha aa^\dagger} - e^{-\alpha a^\dagger a}) a \\ &= (e^{-\alpha(1+a^\dagger a)} - e^{-\alpha a^\dagger a}) a \\ &= (e^{-\alpha} - 1) e^{-\alpha a^\dagger a} a;\end{aligned}$$

Part Two

$$\begin{aligned}[a^\dagger, e^{-\alpha a^\dagger a}] &= [a^\dagger, e^{-\alpha(aa^\dagger-1)}] \\ &= e^{\alpha} [a^\dagger, e^{-\alpha aa^\dagger}] \\ &= e^{\alpha} \sum_{n=1}^{\infty} [a^\dagger (aa^\dagger)^n - (aa^\dagger)^n a^\dagger] \\ &= e^{\alpha} \sum_{n=1}^{\infty} [(a^\dagger a)^n - (aa^\dagger)^n] a^\dagger \\ &= e^{\alpha} (e^{-\alpha a^\dagger a} - e^{-\alpha aa^\dagger}) a^\dagger \\ &= e^{\alpha} (e^{-\alpha a^\dagger a} - e^{-\alpha(1+a^\dagger a)}) a^\dagger \\ &= (e^{\alpha} - 1) e^{-\alpha a^\dagger a} a^\dagger.\end{aligned}$$

Problem 2

Show that the free-field Hamiltonian

$$\mathcal{H} = \hbar\nu \left(a^\dagger a + \frac{1}{2} \right)$$

can be written in terms of the number states as

$$\mathcal{H} = \sum_n E_n |n\rangle \langle n|,$$

and hence

$$e^{i\mathcal{H}t/\hbar} = \sum_n e^{iE_n t/\hbar} |n\rangle \langle n|.$$

Proof. In the number representation,

the matrix elements of $\mathcal{H} = \hbar\nu (a^\dagger a + 1/2)$ is

$$\mathcal{H}_{mn} = \langle m | \mathcal{H} | n \rangle = \hbar\nu \left(n + \frac{1}{2} \right) \delta_{mn} = E_n \delta_{mn},$$

and the matrix elements of $\mathcal{H} = \sum_n E_n |n\rangle \langle n|$ is

$$\mathcal{H}_{mk} = \sum_n E_n \langle m | n \rangle \langle n | k \rangle = E_k \delta_{km},$$

so that the free-field Hamiltonian $\mathcal{H} = \hbar\nu (a^\dagger a + 1/2)$ can be written in terms of the number states as $\mathcal{H} = \sum_n E_n |n\rangle \langle n|$.

As for $\exp(i\mathcal{H}t/\hbar)$, we can see that

$$\mathcal{H}^m = \left(\sum_n E_n |n\rangle \langle n| \right)^m = \sum_n E_n^m |n\rangle \langle n|.$$

Therefore,

$$\begin{aligned} e^{i\mathcal{H}t/\hbar} &= \sum_{m=0}^{\infty} \frac{(it/\hbar)^m}{m!} \mathcal{H}^m \\ &= \sum_{m=0}^{\infty} \sum_n \frac{(it/\hbar)^m}{m!} E_n^m |n\rangle \langle n| \\ &= \sum_n e^{iE_n t/\hbar} |n\rangle \langle n|. \end{aligned}$$

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