



B

a) Step 1 with EF

$$\frac{z(t+\Delta t) - z(t)}{\Delta t} = 0, \text{ since } \dot{z}(0) = 0$$

$$\frac{z_1 - z_0}{\Delta t} = 0$$

$$\underline{z_1 = z_0}$$

General formula

$$\frac{z(t+\Delta t) - 2z(t) + z(t-\Delta t))}{\Delta t^2} = - \frac{\gamma M}{(z(t) + R_E)^2}$$

$$z_{t+1} = \frac{-\gamma M \Delta t^2}{(z(t) + R_E)^2} + 2z(t) - z(t-\Delta t)$$

$$b) \frac{d^2 z}{dt^2} = \frac{-\gamma M}{(z + R_E)^2}$$

$$u_1 = z$$

$$u_2 = \frac{dz}{dt}$$

$$\frac{du_1}{dt} = u_2$$

$$\frac{du_2}{dt} = \frac{dz}{dt^2} = \frac{-\gamma M}{(z + R_E)^2}$$

$$\begin{pmatrix} u_{1,t+\Delta t} \\ u_{2,t+\Delta t} \end{pmatrix} = \begin{pmatrix} u_{1,t} \\ u_{2,t} \end{pmatrix} + \Delta t \begin{pmatrix} u_{2,t} \\ \frac{-\gamma M}{(u_{1,t} + R_E)^2} \end{pmatrix}$$

$$\frac{du_1}{dt} = u_2$$

$$\frac{u_1(t+\Delta t) - u_1(t)}{\Delta t} = u_2(t)$$

$$\frac{du_2}{dt} = \frac{-\gamma M}{(u_1 + R_E)^2}$$

$$\frac{u_2(t+\Delta t) - u_2(t)}{\Delta t} = \frac{-\gamma M}{(u_1 + R_E)^2}$$

$$u_1(t+\Delta t) = u_1(t) + u_2(t) \Delta t$$

$$u_2(t+\Delta t) = u_2(t) - \frac{\gamma M}{(u_1(t) + R_E)^2} \Delta t$$

14. $\frac{d^2 T(x)}{dx^2} = h(x)$

$$h(x) = 0.1x \quad 1 \quad 0 \leq x \leq 10$$

$$T(0) = 3 \quad 1 \quad \left. \frac{dT}{dx} \right|_{x=0} = -2$$

a) $\underline{C} = h_1 - \frac{T_0}{\Delta x^2}, h_2, \dots$

left side with dirichlet

$$\frac{T_{j+1} - 2T_j + T_{j-1}}{\Delta x^2} = h_j$$

$$\underline{j=1} \quad T_2 - 2T_1 = \left(h_1 - \frac{T_0}{\Delta x^2} \right) \Delta x^2$$

$$\underline{j=2} \quad \dots = h_j$$

$$\underline{j=N-1}$$

Right side with Neumann

$$\underline{j=N} \quad \frac{T_{N+1} - 2T_N + T_{N-1}}{\Delta x^2} = h_N \rightarrow \frac{T_{N+1} - T_N - T_N + T_{N-1}}{\Delta x^2} = h_N$$

$$\frac{dT}{dx} \Delta x + \frac{-T_N + T_{N-1}}{\Delta x^2} = h_N$$

$$\frac{-T_N + T_{N-1}}{\Delta x^2} = h_N - \frac{dT}{dx} \Delta x$$

$$\rightarrow \underline{C} = \left(h_1 - \frac{T_0}{\Delta x^2}, h_2, \dots, h_{N-1}, h_N - \frac{dT}{dx} \Delta x \right)$$

first - Dirichlet

$$\underline{M} = \begin{pmatrix} -\frac{2}{\Delta x^2} & \frac{1}{\Delta x^2} & 0 & \dots & 0 \\ \frac{1}{\Delta x^2} & -\frac{2}{\Delta x^2} & \frac{1}{\Delta x^2} & 0 & \dots & 0 \\ 0 & \frac{1}{\Delta x^2} & -\frac{2}{\Delta x^2} & \frac{1}{\Delta x^2} & 0 & \dots & 0 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & \dots & \dots & \dots & \frac{1}{\Delta x^2} & -\frac{2}{\Delta x^2} & \frac{1}{\Delta x^2} \\ 0 & \dots & \dots & \dots & \frac{1}{\Delta x^2} & -\frac{1}{\Delta x^2} \end{pmatrix}$$

last row Neumann

14.

c) Ansatz $T(x) = Ax^3 + Bx^2 + Cx + D$

$$T'(x) = 3Ax^2 + 2Bx + C$$

$$T''(x) = 6Ax + 2B$$

Sheet: $T''(x) = 0,1x$ Coefficients

$$\rightarrow 0,1x = 6Ax + 2B$$

no x in front of B , therefore $B \stackrel{!}{=} 0$

$$0,1x = 6Ax \rightarrow A = \frac{0,1}{6}$$

$$T'(10) = -2$$

$$\rightarrow -2 = \cancel{3} \cdot \frac{0,1}{6} \cdot 10^2 + C$$

$$C = -7$$

$$T(0) = 3$$

$$\rightarrow 3 = D$$

$$\boxed{T(x) = \frac{0,1}{6}x^3 - 7x + 3}$$

sheet9

October 31, 2017

1 13. Free fall in Earth's inhomogeneous gravitational field

1.1 a)

```
In [1]: using PyPlot
        g= 6.67408e-11 #[m^3/(kg*s^2)]
        re=6371000 #mean radius [m]
        m=5.97237e24 #[kg]
        dt=10 #[s]

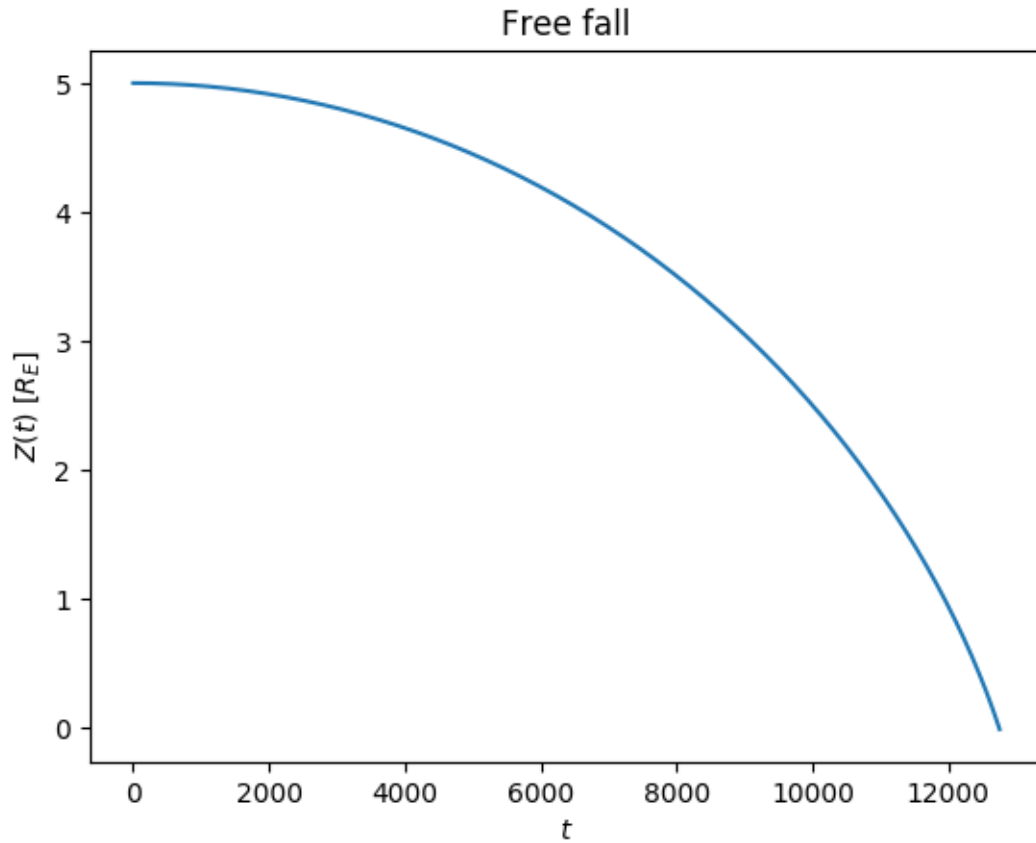
        z=Float64[]
        dz=Float64[]
        t=Float64[]

        push!(z,5*re)
        push!(dz,0)
        push!(t,0)

        #first value
        push!(z,z[1])
        push!(t,t[1]+dt)
        #do until z>surface (z=0)
        i=2
        while (z[i]>0)
            push!(z,-g*m*dt^2/(z[i]+re)^2+2*z[i]-z[i-1])
            push!(t,t[i]+dt)
            i+=1
        end

        title("Free fall")
        plot(t,z/re)
        ylabel(L"$Z(t)$ $[R_E]$")
        xlabel(L"$t$")

        sleep(1)
        println("Time to reach the surface: ",t[length(t)],"s")
```



Time to reach the surface: 12760.0s

1.2 b)

```
In [2]: #initialization
z=Float64[] #u1
dz=Float64[] #u2
t=Float64[]

push!(z,5*re)
push!(dz,0)
push!(t,0)

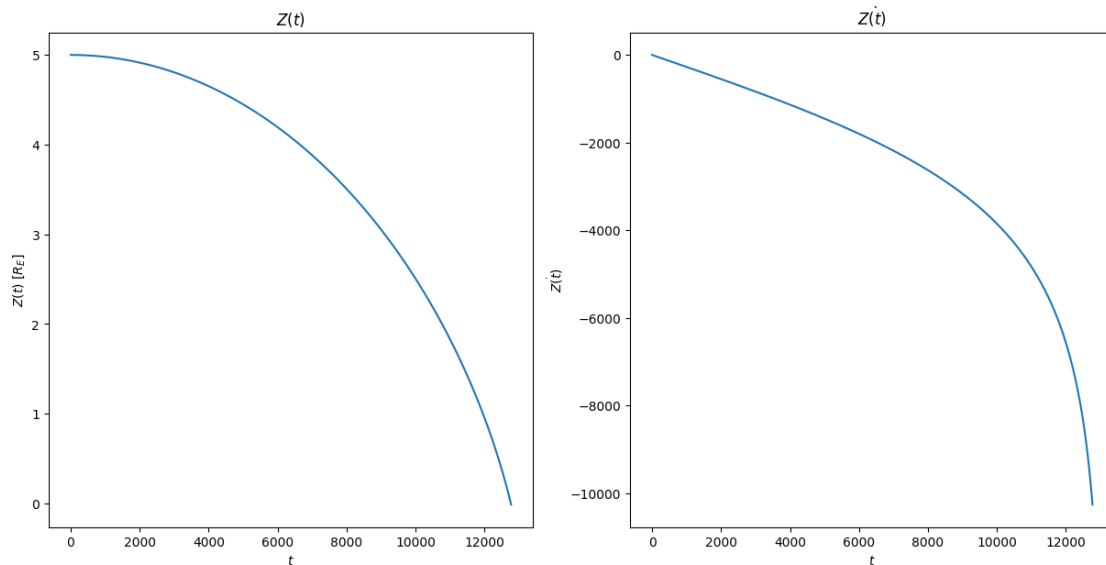
i=1
while (z[i]>0)
    push!(z,z[i]+dz[i]*dt)
    push!(dz,dz[i]-g*m*dt/(z[i]+re)^2)
    push!(t,t[i]+dt)
    i+=1
end
```

```

end

#plotting commands
figure(1,figsize=(15,7))
#2figs in line, linenumber=1, rownumber=2,number of figure=1
subplot(121)
title(L"$Z(t)$")
plot(t,z/re)
ylabel(L"$Z(t)$ $[R_E]$")
xlabel(L"$t$")
legend()
println("Time to reach the surface: ",t[length(t)],"s")
subplot(122)
title(L"$\dot{Z}(t)$")
plot(t,dz)
ylabel(L"$\dot{Z}(t)$")
xlabel(L"$t$")
legend()

```



Time to reach the surface: 12770.0s

/usr/local/lib/python2.7/dist-packages/matplotlib/axes/_axes.py:545: UserWarning: No labelled objects found.
warnings.warn("No labelled objects found. ")

There's a 10s difference between both methods for $\Delta t = 1s$. Since the method used in b) is of first order accuracy, compared to 2nd order accuracy in a), I would stick with the value from a)

2 14. Steady-state diffusion equation

2.1 b)

```
In [3]: #algorithm from sheet 7
function thomasalgo(d,a,b,y)
    N=length(d)
    #A=a' and Y=y''
    A=Array{Float64}(N)
    Y=Array{Float64}(N)
    x=Array{Float64}(N)
    #timestep1
    A[1]=a[1]/d[1]
    Y[1]=y[1]/d[1]
    #steps 2 to N-1
    for i in 2:(N-1)
        A[i]=a[i]/(d[i]-b[i]*A[i-1])
        Y[i]=(y[i]-b[i]*Y[i-1])/(d[i]-b[i]*A[i-1])
    end
    #N-th value
    Y[N]=(y[N]-b[N]*Y[N-1])/(d[N]-b[N]*A[N-1])

    #calculating the x-Vector
    x[N]=Y[N]
    for i in (N-1):-1:1
        x[i]=Y[i]-A[i]*x[i+1]
    end
    return(x)
end
```

Out[3]: thomasalgo (generic function with 1 method)

```
In [4]: dx=0.05
xmax=10.
T0=3
dtx=-2

#making M
#dx:xmax
leng=floor(Int,xmax/dx)
diagonalelements=ones(leng)*(-2)/dx^2
#last element is different
diagonalelements[leng]+=1/dx^2
#offdiagonal quite easy
offdiagonala=ones(leng)*1/dx^2
offdiagonalb=ones(leng)*1/dx^2
offdiagonala[leng]=0.
offdiagonalb[1]=0
```

```
#M=
```

```
M=diagm(diagonalelements,0)+diagm(offdiagonala[1:leng-1],1)+diagm(offdiagonalb[2:leng]
```

```
Out[4]: 200x200 Array{Float64,2}:
```

```
-800.0  400.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0
 400.0 -800.0  400.0   0.0   0.0   0.0   0.0   0.0   0.0
   0.0  400.0 -800.0  400.0   0.0   0.0   0.0   0.0   0.0
   0.0   0.0  400.0 -800.0  400.0   0.0   0.0   0.0   0.0
   0.0   0.0   0.0  400.0 -800.0   0.0   0.0   0.0   0.0
   0.0   0.0   0.0   0.0  400.0   0.0   0.0   0.0   0.0
   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0
   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0
   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0
   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0
   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0
   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0
   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0
   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0
   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0
   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0
   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0
   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0   0.0
   0.0   0.0   0.0   0.0   0.0   400.0   0.0   0.0   0.0
   0.0   0.0   0.0   0.0   0.0  -800.0  400.0   0.0   0.0
   0.0   0.0   0.0   0.0   0.0  400.0 -800.0  400.0   0.0
   0.0   0.0   0.0   0.0   0.0   0.0  400.0 -800.0  400.0
   0.0   0.0   0.0   0.0   0.0   0.0   0.0  400.0 -400.0
```

```
In [5]: #making C
```

```
h(x)=0.1*x
```

```
C=Array{Float64}(leng)
for (i,x) in enumerate(dx:dx:10)
    C[i]=h(x)
end
C[1]-=T0/dx^2
C[leng]-=dtx/dx
```

```
C
```

```
Out[5]: 200-element Array{Float64,1}:
```

```
-1199.99
  0.01
  0.015
  0.02
```

0.025
0.03
0.035
0.04
0.045
0.05
0.055
0.06
0.065

0.945
0.95
0.955
0.96
0.965
0.97
0.975
0.98
0.985
0.99
0.995
41.0

In [6]: T=thomasalgo(diagonalelements,offdiagonala,offdiagonalb,C)

Out[6]: 200-element Array{Float64,1}:

2.64875
2.29751
1.9463
1.59512
1.244
0.892937
0.54195
0.19105
-0.15975
-0.510438
-0.861
-1.21143
-1.5617

-49.3215
-49.4483
-49.5728
-49.6948
-49.8145
-49.9317
-50.0465
-50.1589

```
-50.2688
-50.3763
-50.4813
-50.5838
```

```
In [7]: #just to test, seems like Julia has no problems to invert huge arrays
Xtest=inv(M)*C
```

```
Out[7]: 200-element Array{Float64,1}:
```

```
 2.64875
 2.29751
 1.9463
 1.59512
 1.244
 0.892937
 0.54195
 0.19105
-0.15975
-0.510438
-0.861
-1.21143
-1.5617
```

```
-49.3215
-49.4483
-49.5728
-49.6948
-49.8145
-49.9317
-50.0465
-50.1589
-50.2688
-50.3763
-50.4813
-50.5838
```

```
In [8]: x=Array{Float64}(leng+1)
T2=Array{Float64}(leng+1)
x[1]=0
T2[1]=T0
for (i,x2) in enumerate(dx:dx:10)
    x[i+1]=x2
    T2[i+1]=T[i]
end

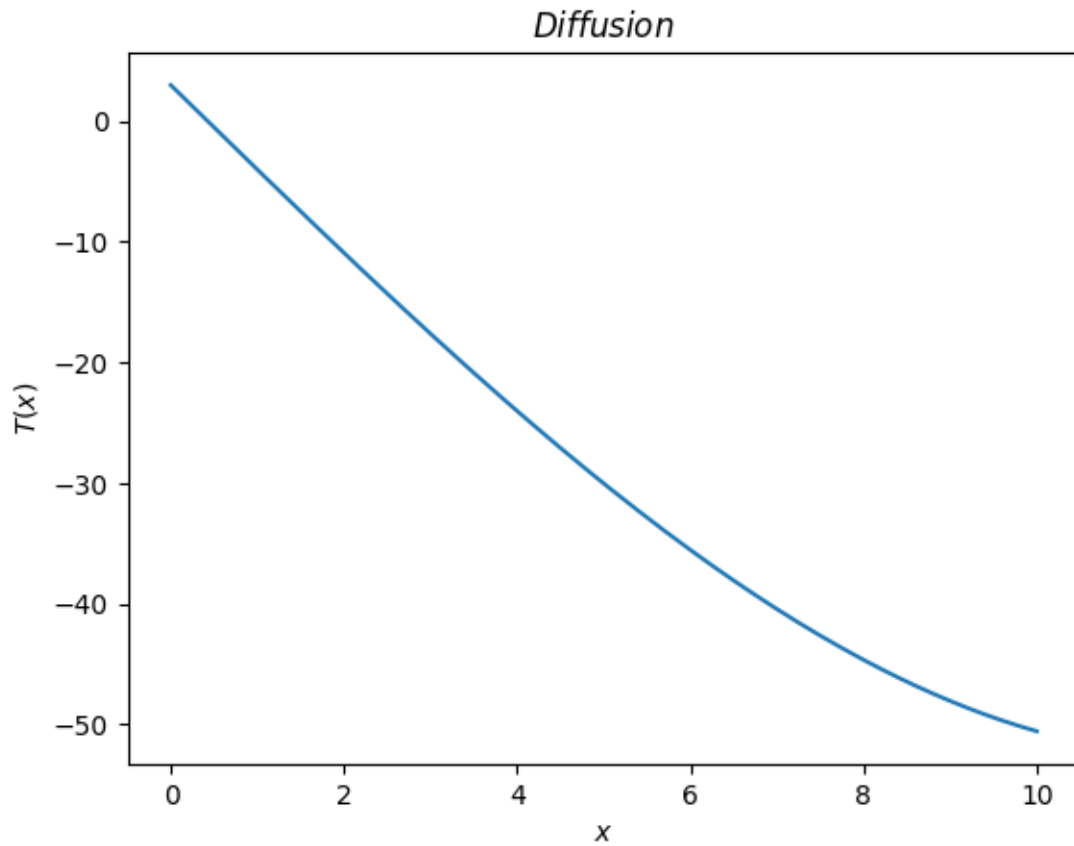
i=1:leng+1
plot(x[i],T2[i])
title(L"Diffusion")
```



```

ylabel(L"$T(x)$")
xlabel(L"$x$")
legend()

```

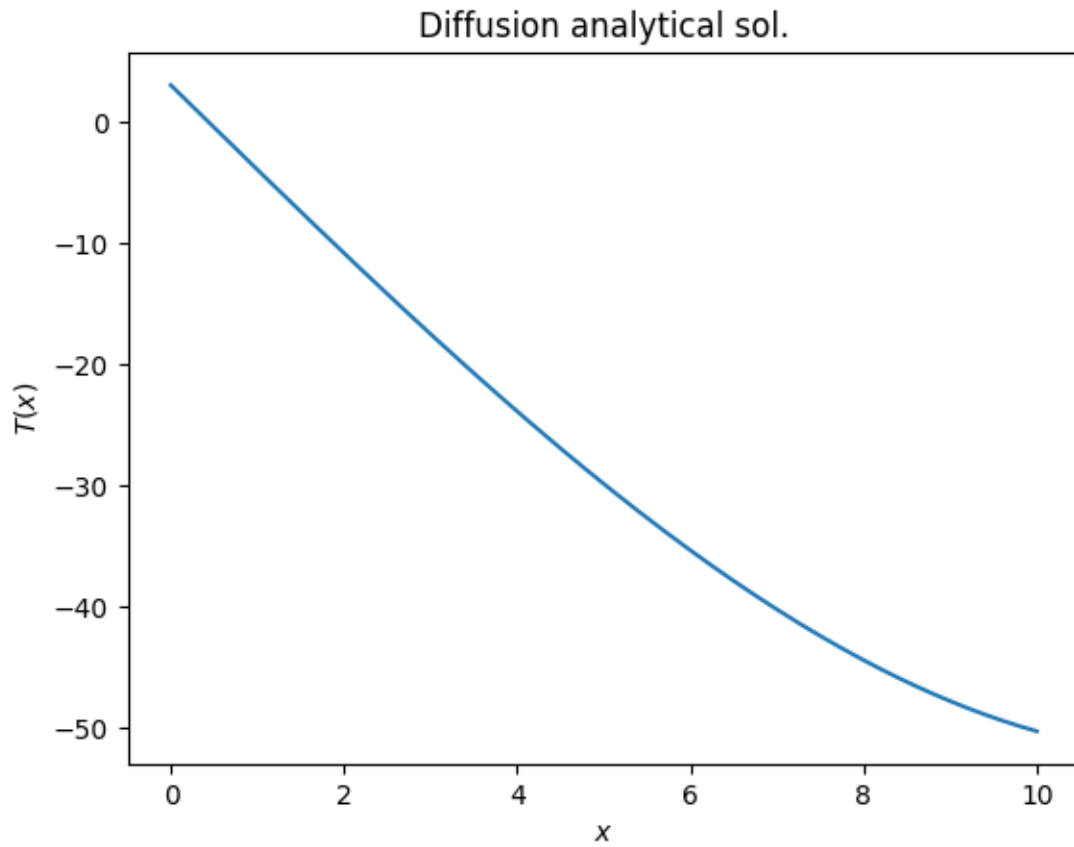


2.2 c)

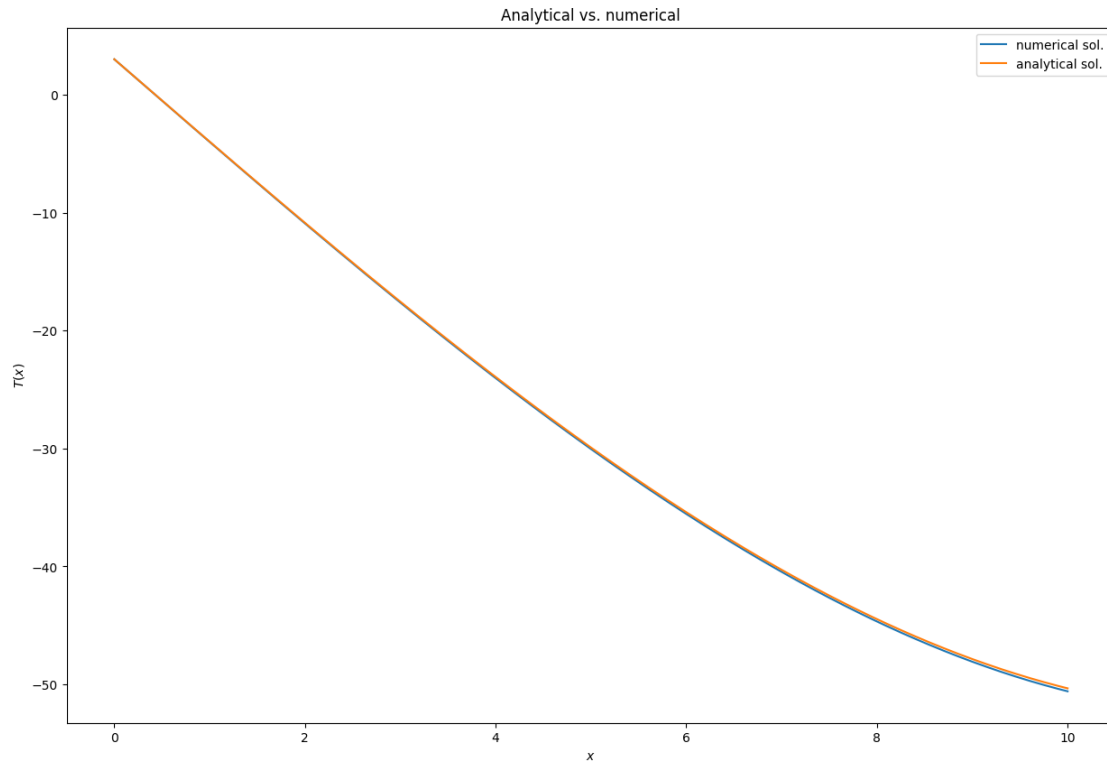
```

In [9]: #analytical sol
T3(x)=0.1/6*x.^3-7*x.+3
x=0.:dx:10.
plot(x,T3(x))
title("Diffusion analytical sol.")
ylabel(L"$T(x)$")
xlabel(L"$x$")
legend()

```



```
In [10]: figure(1,figsize=(15,10))
         plot(x[i],T2[i],label="numerical sol.")
         plot(x,T3(x),label="analytical sol.")
         title("Analytical vs. numerical")
         ylabel(L"$T(x)$")
         xlabel(L"$x$")
         legend()
```



```
Out[10]: PyObject <matplotlib.legend.Legend object at 0x7f3f6fb1dad0>
```

They're almost the same, with a lower stepsize, the numerical approximation would even be better