



## On the Computational Power of $QAC^0$ with Barely Superlinear Ancillae



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- Basic concepts of approximation
  - Boolean Fourier expansion

- Ancilla models and channels
- Comparison between  $AC^0$  model



Overview



Pauli/Fourier  
degree



Boolean function



Pauli degree  
growth



Applications



- Major Theorems and concepts
- Basic  $\varepsilon$  – *error* approximations

- Embedding Booleans as operators
- Approximate degree conservation

- Comparison between prev. works
- Available applications (lack of insight)



Quantum Arithmetic Circuit의 가장 아래 변환층

$\text{QAC}^0$  is the family of constant-depth polynomial-size quantum circuits consisting of arbitrary single qubit unitaries and multi-qubit Toffoli gates. It was introduced by Moore as a quantum counterpart of  $\text{AC}^0$ , along with the conjecture that  $\text{QAC}^0$  circuits cannot compute PARITY. In this work, we make progress on this long-standing conjecture: we show that any depth- $d$   $\text{QAC}^0$  circuit requires  $n^{1+3^{-d}}$  ancillae to compute a function with approximate degree  $\Theta(n)$ , which includes PARITY, MAJORITY and  $\text{MOD}_k$ . We further establish superlinear lower bounds on quantum state synthesis and quantum channel synthesis. This is the first lower bound on the super-linear sized  $\text{QAC}^0$ . Regarding PARITY, we show that any further improvement on the size of ancillae to  $n^{1+\exp(-o(d))}$  would imply that  $\text{PARITY} \notin \text{QAC}^0$ .

These lower bounds are derived by giving low-degree approximations to  $\text{QAC}^0$  circuits. We show that a depth- $d$   $\text{QAC}^0$  circuit with  $a$  ancillae, when applied to low-degree operators, has a degree  $(n+a)^{1-3^{-d}}$  polynomial approximation in the spectral norm. This implies that the class  $\text{QLC}^0$ , corresponding to linear size  $\text{QAC}^0$  circuits, has an approximate degree  $o(n)$ . This is a quantum generalization of the result that  $\text{LC}^0$  circuits have an approximate degree  $o(n)$  by Bun, Kothari, and Thaler. Our result also implies that  $\text{QLC}^0 \neq \text{NC}^1$ .

$$f(n) = O(g(n))$$

$$f(n) = \Theta(g(n))$$

$$f(n) = o(g(n))$$

$$f(n) = \Omega(g(n))$$



For any  $2^n \times 2^n$  operator  $A$  with degree  $\ell$ , and any unitary  $U$  implemented by a depth- $d$  QAC<sup>0</sup> circuit, the approximate degree of  $UAU^\dagger$  is upper bounded by  $\tilde{O}\left(n^{1-3^{-d}}\ell^{3^{-d}}\right)$ .

Let  $f: \{0,1\}^n \rightarrow \{0,1\}$  be a Boolean function with approximate degree  $\Omega(n)$ . Suppose  $U$  is a depth  $d$  QAC0 circuit with  $n$  input qubits and  $a = \tilde{O}(n^{1+3^{-d}})$  ancillae initialized in any quantum state. Then  $U$  cannot compute  $f$  with the worst-case error strictly below  $1/2$ . And  $U$  can't approximate  $\text{Parity}_n$  nor  $\text{Majority}_n$  over uniform inputs.

**Theorem 1.4** (informal of Corollary 5.3). *If any QAC<sup>0</sup> circuit with  $n^{1+\exp(-o(d))}$  ancillae, where  $d$  is the depth of this circuit family, can not compute  $\text{Parity}_n$  with the worst-case error  $\text{negl}(n)$ , then any QAC<sup>0</sup> circuit family with arbitrary polynomial ancillae can not compute  $\text{Parity}_n$  with the worst-case error  $\text{negl}(n)$ .*

**Theorem 1.6** (informal of Theorem 7.2). *Suppose  $\mathcal{E}_{U,\psi}$  is a quantum channel from  $n$  qubits to  $k$  qubits, implemented by a depth- $d$  QAC<sup>0</sup> circuit  $U$  with  $n$  input qubits and  $a$  ancillae. The upper bound of approximate degree of the Choi representation  $\Phi_{U,\psi}$  of  $\mathcal{E}_{U,\psi}$  is then given by  $\tilde{O}\left((n+a)^{1-3^{-d}}k^{3^{-d}/2}\right)$ .*

$$C(\Phi) = \sum_{i,j} |i\rangle\langle j|_{A'} \otimes \Phi(|i\rangle\langle j|_A).$$



- We can represent  $\rho_{AB}$  as an ensemble of pure states

$$\rho_{AB} = \sum_{i=1}^N p_i |\psi_i\rangle\langle\psi_i|.$$

- We can also view  $\rho_{AB}$  as part of a pure state (its purification) as

$$\rho_{AB} = \text{Tr}_E [|\psi_{ABE}\rangle\langle\psi_{ABE}|].$$

**Theorem 4.1** (Choi and Kraus). *For a linear map  $T : B(\mathcal{H}_A) \rightarrow B(\mathcal{H}_B)$  the following are equivalent:*

1. *The map  $T$  is completely positive.*
2. *There exist operators  $\{K_i\}_{i=1}^R \subset B(\mathcal{H}_A, \mathcal{H}_B)$  and some  $R \in \mathbb{N}$  such that*

$$T = \sum_{i=1}^R \text{Ad}_{K_i}. \quad (4.1)$$

$$C(\Phi) = \sum_{i,j} |i\rangle\langle j|_{A'} \otimes \Phi(|i\rangle\langle j|_A).$$

*Proof.* Since the maps  $\text{Ad}_{K_i} : B(\mathcal{H}_A) \rightarrow B(\mathcal{H}_B)$  are completely positive, it is clear that 2. implies 1.. To see the other direction, consider a completely positive map  $T : B(\mathcal{H}_A) \rightarrow B(\mathcal{H}_B)$ . By Theorem 3.8, the Choi matrix

$$C_T = \dim(\mathcal{H}_A) (\text{id}_{A'} \otimes T) (\omega_{\mathcal{H}_A})$$

is positive semidefinite. Using the spectral decomposition, we can write

$$C_T = \sum_{i=1}^R |\psi_i\rangle\langle\psi_i|,$$

for  $R = \text{rk}(C_T)$  and some (unnormalized) vectors  $|\psi_i\rangle \in \mathcal{H}_A \otimes \mathcal{H}_B$ . Next, we use the inverse of the vectorization isomorphism and Theorem 2.14 to show that

$$|\psi_i\rangle = \text{vec}(K_i) = (\mathbb{1}_{\mathcal{H}_A} \otimes K_i)|\Omega_{\mathcal{H}_A}\rangle,$$

for some operators  $K_i \in B(\mathcal{H}_A, \mathcal{H}_B)$ . Combining the previous equations, we find that

$$C_T = \sum_{i=1}^R (\mathbb{1}_{d_A} \otimes K_i) \omega_{d_A} (\mathbb{1}_{d_A} \otimes K_i)^\dagger = \sum_{i=1}^R C_{\text{Ad}_{K_i}} = C_{\sum_{i=1}^R \text{Ad}_{K_i}}.$$

Finally, we use that the Choi-Jamiolkowski isomorphism is indeed an isomorphism and the last equation implies

$$T = \sum_{i=1}^R \text{Ad}_{K_i}.$$

□



$$\mathbb{E}_{\mathbf{x}}[f(\mathbf{x})g(\mathbf{x})]$$

$$\chi_S(x) = (-1)^{\sum_{i \in S} x_i}$$

$$f = \sum_{S \subseteq [n]} \widehat{f}(S) \chi_S$$

$$\|f\|_2^2 = \sum_{S \subseteq [n]} \widehat{f}(S)^2 \quad \deg(A) = \max_{\sigma: \widehat{A}(\sigma) \neq 0} |\sigma|$$

$$\|M\|_p = \left( \frac{1}{n} \operatorname{Tr} [|M|^p] \right)^{1/p} \quad \text{Normalized Schatten p-norm}$$

$$F(\rho, \sigma) = \operatorname{Tr} \left[ \sqrt{\sqrt{\rho}\sigma\sqrt{\rho}} \right] \quad \text{Fidelity}$$

$$1 - \frac{1}{2} \|\rho - \sigma\|_{TD} \leq F(\rho, \sigma) \leq \sqrt{1 - \frac{1}{4} \|\rho - \sigma\|_{TD}^2}$$

$$M_f = \sum_x f(x) \cdot |x\rangle\langle x| = \begin{bmatrix} f(0^n) & & \\ & \ddots & \\ & & f(1^n) \end{bmatrix} = \sum_{\sigma \in \{0,1\}^n} \widehat{f}(S_\sigma) \cdot \mathcal{B}_\sigma$$

**Lemma 2.7.** Let  $A, B, \widetilde{A}, \widetilde{B}$  be operators satisfying

- $\|A\| \leq 1$  and  $\|B\| \leq 1$ .
- $\|\|A - \widetilde{A}\| \leq \varepsilon_0$ .
- $\|\|B - \widetilde{B}\| \leq \varepsilon_1$ .

Then  $\|AB\| \leq 1$  and  $\|AB - \widetilde{A}\widetilde{B}\| \leq \varepsilon = \varepsilon_0 + \varepsilon_1 + \varepsilon_0\varepsilon_1 = (1 + \varepsilon_0)(1 + \varepsilon_1) - 1$ .

The Pauli matrices  $\mathcal{B}_0, \dots, \mathcal{B}_3$  are

$$\mathcal{B}_0 = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}, \mathcal{B}_1 = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, \mathcal{B}_2 = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}, \mathcal{B}_3 = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix},$$

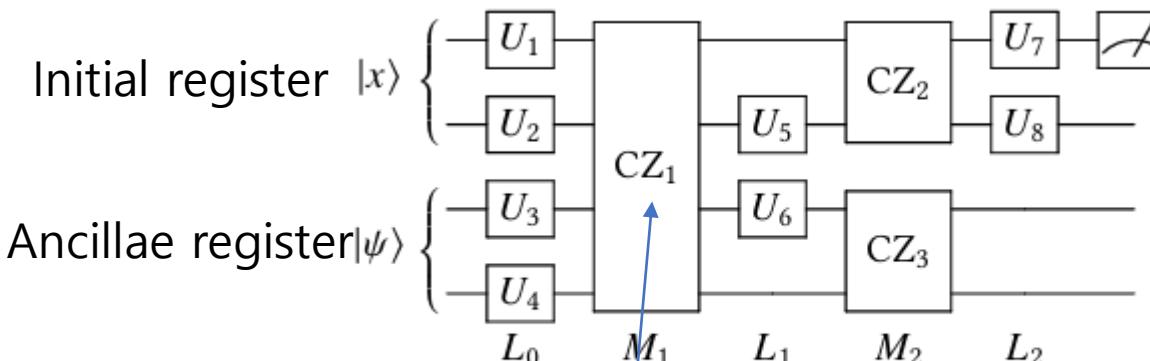
which form an orthonormal basis in  $\mathcal{M}_2$ . For integer  $n \geq 1$  and  $\sigma \in \{0, 1, 2, 3\}^n$ , we define

$$\mathcal{B}_\sigma = \mathcal{B}_{\sigma_1} \otimes \cdots \otimes \mathcal{B}_{\sigma_n}.$$

The set of Pauli matrices  $\{\mathcal{B}_\sigma\}_{\sigma \in \{0,1,2,3\}^n}$  forms an orthonormal basis in  $\mathcal{M}_{2^n}$ . For any  $2^n \times 2^n$  matrix  $A$ , the Pauli expansion of  $A$  is

$$A = \sum_{\sigma \in \{0,1,2,3\}^n} \widehat{A}(\sigma) \cdot \mathcal{B}_\sigma.$$

The coefficients  $\widehat{A}(\sigma)$ 's are called the Pauli coefficients of  $A$ . We can then define the degree and the approximate degree of a matrix.

Figure 1: QAC<sup>0</sup> Circuit Example

$$U(|\varphi\rangle \otimes |\psi\rangle)$$



$$\Pr[C(x) = 1] = \text{Tr} [(|1\rangle\langle 1| \otimes I) U(|x\rangle\langle x| \otimes |\psi\rangle\langle \psi|) U^\dagger]$$

$$\Pr [C(x) \neq f(x)] \leq \varepsilon \quad \underset{\mathbf{x} \in \{0,1\}}{\mathbb{E}} [\Pr [C(\mathbf{x}) \neq f(\mathbf{x})]] \leq \varepsilon$$

Worst case error

Average case error

$$CZ = \mathbb{1} - 2|1\rangle\langle 1|^n$$

$$CZ = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

$$CZ(X \otimes I)CZ^\dagger = X \otimes Z$$

$$CZ(I \otimes X)CZ^\dagger = Z \otimes X$$

$$|\Psi_n\rangle = \frac{1}{\sqrt{2}}(|0^n\rangle + |1^n\rangle)$$

$$|\nu\rangle = \frac{1}{\sqrt{2}}(|0^n, \psi_0\rangle + |1^n, \psi_1\rangle)$$

The complexity class QAC<sup>0</sup> consists of all languages that can be decided by constant-depth and polynomial-sized QAC quantum circuits. Formally, a language  $L$  is in QAC<sup>0</sup> if there exists a family of constant-depth and polynomial-sized QAC quantum circuits  $\{C_n\}_{n \in \mathbb{N}}$  such that for any  $n \in \mathbb{N}$  and  $x \in \{0, 1\}^n$ , if  $x \in L$  then  $\Pr[C_n(x) = 1] \geq 2/3$ , and if  $x \notin L$ , then  $\Pr[C_n(x) = 0] \geq 2/3$  where  $C_n(x)$  is the measurement outcome on the output qubits of the circuit  $C_n$  on input  $x$ . We also introduce the class of QLC<sup>0</sup> circuits, which consists of QAC<sup>0</sup> circuits with linear-sized ancillae.



**Lemma 3.1** ([AM23, Lemma 3.1], see also [KAAV17]). Let  $H = \sum_{i=1}^n H_i$  be a sum of  $n$  commuting projectors each acting on  $\ell$  qubits, and  $|\psi\rangle$  be the maximum-energy eigenstate of  $H$ . Then, for any  $r \in (\sqrt{n}, n)$ , let  $\varepsilon = 2^{-\frac{r^2}{2^8 n}}$ ,

$$\widetilde{\deg}_\varepsilon(|\psi\rangle\langle\psi|) \leq \ell r.$$

**Corollary 3.2.** Let  $|\psi\rangle$  be an  $\ell$ -qubit pure state. Then for any  $r \in (\sqrt{n}, n)$ , let  $\varepsilon = 2^{-\frac{r^2}{2^8 n}}$ . It holds that

$$\widetilde{\deg}_\varepsilon(|\psi\rangle\langle\psi|^{\otimes n}) \leq \ell r.$$

**Corollary 3.3.** For any CZ-gate CZ acting on  $n$  qubits and real number  $1 < r < n$ , there exists an operator  $\widetilde{\text{CZ}}$  such that

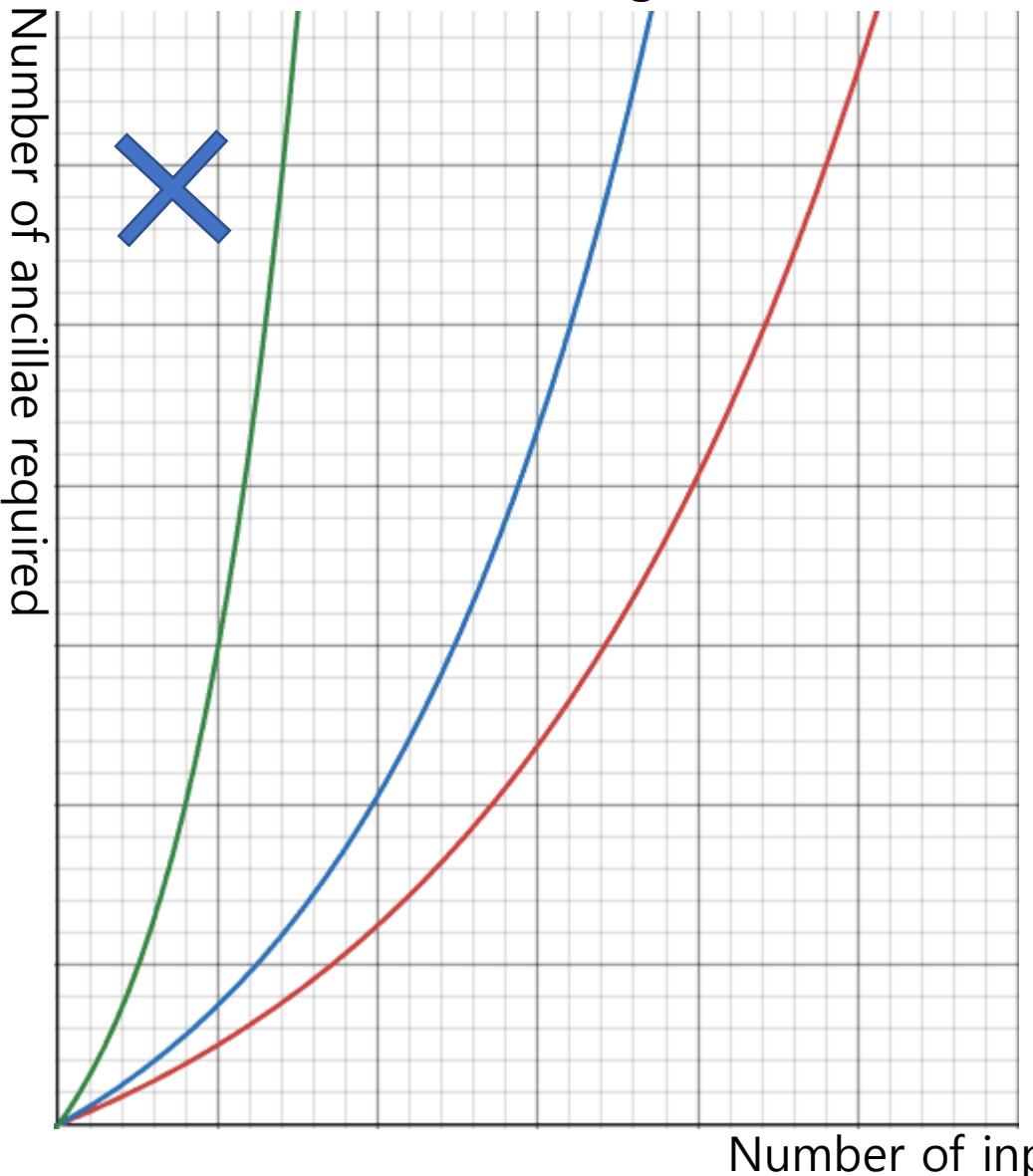
$$\|\text{CZ} - \widetilde{\text{CZ}}\| \leq 2^{1-2^{-8}r}$$

and

$$\deg(\widetilde{\text{CZ}}) \leq \sqrt{nr}.$$



Worst case      Average case



PARITY	$a \geq n2^{-d} - 1$	exact	[FFG <sup>+</sup> 06]
	impossible when $d = 2$	exact	[PFGT20]
	impossible when $d = 2$	average case	[Ros21]
	$a \leq \exp(O(n \log n/\varepsilon))$ when $d = 7$	worst case	[Ros21]
	$a \geq n^{\Omega(1/d)}$	average case	[NPVY24]
	$a \geq n^{1+3^{-d}}$	average/worst case	This work
MAJORITY	$a \geq n^{\Omega(1/d)}$	average case <sup>†</sup>	[NPVY24]
	$a \geq n^{1+3^{-d}}$	average/worst case	This work
MOD <sub>k</sub>	$a \geq n^{1+3^{-d}}$	worst case	This work

$$\deg_\varepsilon(UAU^\dagger) \leq \tilde{O}(n^{1-3^{-d}} \ell^{3^{-d}})$$

Operator spreading 능력의 제한성 수치화

Fan-out과 ancillae를 이용한 Fan-in, 그리고 CZ를 갖지만 Deep entanglement와 high degree operator를 생성할 수 없음

$$\deg_\varepsilon(f) \leq \tilde{O}((n+a)^{1-3^{-d}})$$

Boolean function에서의 계산 가능한 function 경계 디자인

Centre Boolean function 대부분이 계산 불가능  
Classical AC<sup>0</sup>보다는 강하지만 NC<sup>1</sup> – QNC<sup>1</sup> 수준으로는 도당 X

Ancilla가  $n^{1+3^{-d}}$ 보다 작으면 PARITY 불가능  
 $n^{1+t}$  ( $t > 1$ )정도면 PARITY 가능



Thank you  
for listening