

$$\begin{aligned}
&= \int dp p^2 \frac{1}{\sqrt{2\pi\hbar}} e^{i p x / \hbar} \langle p | \psi \rangle \\
&= \left(\frac{\hbar}{i}\right)^2 \frac{d^2}{dx^2} \int dp \frac{1}{\sqrt{2\pi\hbar}} e^{i p x / \hbar} \langle p | \psi \rangle \\
&= \left(\frac{\hbar}{i}\right)^2 \frac{d^2}{dx^2} \int dp \langle \pi | p \rangle \langle p | \psi \rangle \\
&= \left(\frac{\hbar}{i}\right)^2 \frac{d^2}{dx^2} \langle \pi | \psi \rangle
\end{aligned}$$

□

(1.27) 在解 $\langle \pi | \psi \rangle$

$$\psi_n(x) = \sqrt{\frac{1}{2^n n!}} \sqrt{\frac{m\omega}{\pi\hbar}} H_n\left(\sqrt{\frac{m\omega}{\hbar}} x\right) e^{-\frac{m\omega}{2\hbar} x^2}$$

* $H_n(x)$: 厄米多项式

$$E_n = \hbar\omega\left(n + \frac{1}{2}\right)$$

↑ 位置, 动量: 连续

← E : 离散



\hat{H} 的量子化

$$\hat{H} = \frac{1}{2} m \omega^2 \hat{x}^2 + \frac{1}{2m} \hat{p}^2$$

$$= \hbar\omega \left(\frac{m\omega}{2\hbar} \hat{x}^2 + \frac{1}{2m\hbar\omega} \hat{p}^2 \right)$$

$$= \hbar\omega \left(\left(\sqrt{\frac{m\omega}{2\hbar}} \hat{x} - i \sqrt{\frac{1}{2m\hbar\omega}} \hat{p} \right) \left(\sqrt{\frac{m\omega}{2\hbar}} \hat{x} + i \sqrt{\frac{1}{2m\hbar\omega}} \hat{p} \right) \right)$$

$$-i \sqrt{\frac{m\omega}{2\hbar}} \sqrt{\frac{\hbar}{2m\omega}} [\hat{x}, \hat{p}] \quad i\hbar$$

$$= \hbar\omega \left(\hat{a}^\dagger \hat{a} + \frac{1}{2} \right)$$

(準備)

$$|\psi(x)\rangle := \langle x | \psi \rangle$$

$$|\psi\rangle = \int dx \psi(x) |x\rangle$$

$$\langle x | \hat{p} | \psi \rangle = -i\hbar \frac{d}{dx} \psi(x)$$

└

$$\hat{H} |\psi\rangle = E |\psi\rangle$$

$\langle x | \downarrow$

$$\langle x | \hat{H} | \psi \rangle = E \psi(x)$$

↖ $\psi(x)$ の表現!

Step 1 $\langle x | \hat{p} | x \rangle = -i\hbar \frac{d}{dx} \psi(x)$

Step 2 $\langle x | p \rangle = \frac{1}{\sqrt{2\pi\hbar}} e^{\frac{ixp}{\hbar}}$

Step 3. $\langle x | \hat{p}^2 | x \rangle$

$$\text{step 0} \quad f(x) \delta'(x) = -f'(x) \delta(x) \quad \lim_{|x| \rightarrow \infty} f(x) \rightarrow 0$$

(証明)

$$\int_{-\infty}^{\infty} \underline{f(x)} \delta'(x) dx = [f(x) \delta(x)]_{-\infty}^{\infty} - \int \underline{f(x)'} \delta(x) dx$$

$$f(x) = x \rightarrow \frac{\delta(x)}{x} = -\delta'(x)$$

step 1.

$$\begin{aligned} \langle x | [\hat{x}, \hat{p}] | x' \rangle &= \langle x | \hat{x} \hat{p} - \hat{p} \hat{x} | x' \rangle \\ &= x \langle x | \hat{p} | x' \rangle - x' \langle x | \hat{p} | x' \rangle \\ &= (x - x') \langle x | \hat{p} | x' \rangle \\ &= i\hbar \delta(x - x') \end{aligned}$$

$$\langle x | \hat{p} | x' \rangle = i\hbar \frac{\delta(x - x')}{x - x'}$$

$$= i\hbar \frac{d}{dx'} \delta(x - x')$$

$$\langle x | \hat{p} | \psi \rangle = \langle x | \hat{p} \hat{1} | \psi \rangle$$

$$= \langle x | \hat{p} \int dx' | x' \rangle \langle x' | \psi \rangle$$

$$= \int dx' \langle x | \hat{p} | x' \rangle \langle x' | \psi \rangle$$

$$= i\hbar \int dx \left[\frac{d}{dx} \delta(x-x') \right] \psi(x')$$

$$= i\hbar \left\{ \left[\delta(x-x') \psi(x') \right]_{-\infty}^{\infty} - \int dx \psi'(x) \delta(x-x) \right\}$$

$$= -i\hbar \frac{d}{dx} \psi(x)$$

Step 2.

$$\langle x | \hat{p} | p \rangle = p \langle x | p \rangle$$

$$= p \cdot p(x)$$

$$\langle x | \hat{p} | p \rangle = -i\hbar \frac{d}{dx} p(x) \quad (\because \text{Step 1})$$

$$p(x) = C \exp\left(\frac{ixp}{\hbar}\right)$$

$$\delta(x-x') = \langle x | x' \rangle$$

$$= \langle x | \int dp | p \rangle \langle p | x' \rangle = \int dp p(x) p^*(x')$$

$$= |C|^2 \int \exp\left(\frac{i(x-x')p}{\hbar}\right) dp$$

$$1 = \int_{-\infty}^{\infty} dt \delta(t) e^{-i\omega t}$$

$$\delta(t) = \frac{1}{2\pi} \int d\omega 1 \cdot e^{i\omega t}$$

$$\left\{ \begin{array}{l} \omega \rightarrow \frac{p}{\hbar} \\ t \rightarrow x - x' \end{array} \right.$$

$$\delta(x - x') = \frac{1}{2\pi\hbar} \int dp e^{i \frac{x-x'}{\hbar} p}$$

$$|C|^2 = \frac{1}{2\pi\hbar} \quad \therefore C = \frac{1}{\sqrt{2\pi\hbar}}$$

$$\therefore \langle x|p \rangle = \frac{1}{\sqrt{2\pi\hbar}} \exp\left(\frac{ixp}{\hbar}\right)$$

Step 3.

$$\langle x | \hat{p}^2 | \psi \rangle$$

$$= \langle x | \hat{p}^2 \int dp |p\rangle \langle p | \psi \rangle$$

$$= \int dp p^2 \langle x | p \rangle \langle p | \psi \rangle$$

$$= \int dp p^2 \cdot \underbrace{\frac{1}{\sqrt{2\pi\hbar}} \exp\left(\frac{ixp}{\hbar}\right)}_{\left(\frac{d^2}{dx^2} \langle x | p \rangle \cdot \left(\frac{\hbar}{i}\right)^2 = p^2 \langle x | p \rangle\right)} \langle p | \psi \rangle$$

$$\left(\frac{d^2}{dx^2} \langle x | p \rangle \cdot \left(\frac{\hbar}{i}\right)^2 = p^2 \langle x | p \rangle \right)$$

$$= \left(\frac{\hbar}{i}\right)^2 \int dx \frac{d^2}{dx^2} \langle x|p \rangle \langle p|x \rangle$$

$$= \underbrace{\left(\frac{\hbar}{i}\right)^2}_{\frac{d^2}{dx^2}} \langle x| \left[\int dp |p\rangle \langle p| \right] |x\rangle = \left(\frac{\hbar}{i}\right)^2 \frac{d^2}{dx^2} \psi(x)$$

$$\left(\frac{1}{2} m \omega^2 \hat{x}^2 + \frac{1}{2m} \hat{p}^2 \right) |\psi\rangle = E |\psi\rangle$$

$\langle x|$
 $\psi(x)$

$$\frac{1}{2} m \omega^2 x^2 \psi(x) + \frac{1}{2m} \left(\frac{\hbar}{i}\right)^2 \frac{d^2}{dx^2} \psi(x) = E \psi(x)$$

$$\left(-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + \frac{1}{2} m \omega^2 x^2 \right) \psi(x) = E \psi(x)$$

解

準備

$$\left(\frac{d^2}{ds^2} - 2s \frac{d}{ds} + 2n \right) H_n(s) = 0 \quad \leftarrow \text{2111-1.99 模式}$$

$$\int_{-\infty}^{\infty} ds \underbrace{H_m(s) H_n(s)}_{H_m e^{-\frac{s^2}{2}} \quad H_n e^{-\frac{s^2}{2}}} e^{-s^2} = \sqrt{\pi} 2^n n! \delta_n^m$$

$$\left(-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + \frac{1}{2} m \omega^2 x^2 \right) \psi(x) = E \psi(x)$$

$$\left(\frac{d^2}{dx^2} - \frac{m^2 \omega^2}{\hbar^2} x^2 + \frac{2mE}{\hbar^2} \right) \psi(x) = 0$$

(假定)

$$\psi(x) = f(s) e^{-\frac{s^2}{2}}$$

\downarrow
 $H_n(s)$ だといいな...

$$s = \sqrt{\frac{m\omega}{\hbar}} x \quad \sim \text{変数変換}$$

$$\sqrt{\frac{\hbar^2 s^{-1}}{m^2 \omega^2 s^{-1}}} \quad m = 1 \Rightarrow \text{無次元化!}$$

$$\frac{d}{dx} = \frac{ds}{dx} \frac{d}{ds}$$

$$\left(\frac{m\omega}{\hbar} \frac{d^2}{ds^2} - \frac{m^2 \omega^2}{\hbar^2} \frac{\hbar}{m\omega} s^2 + \frac{2mE}{\hbar^2} \right) \psi(x) = 0$$

$$\left(\frac{d^2}{ds^2} - s^2 + \frac{2E}{\hbar\omega} \right) \psi(x) = 0 \quad \leftarrow f(s) e^{-\frac{s^2}{2}} \text{ 代入して}$$

$$\frac{d^2}{ds^2} (f(s) e^{-\frac{s^2}{2}}) = \frac{d}{ds} \left(\frac{df}{ds} e^{-\frac{s^2}{2}} - s f(s) e^{-\frac{s^2}{2}} \right)$$

$$= \frac{d^2 f}{ds^2} e^{-\frac{s^2}{2}} + 2 \frac{df}{ds} (-s) e^{-\frac{s^2}{2}} + f(s) (s^2 - 1) e^{-\frac{s^2}{2}}$$

$$= e^{-\frac{s^2}{2}} \left(\frac{d^2 f}{ds^2} - 2s \frac{df}{ds} + (s^2 - 1) f(s) \right)$$

$$e^{-\frac{s^2}{2}} \left(\frac{d^2}{ds^2} f - 2s \frac{df}{ds} + (s^2 - 1) f(s) \right)$$

$$- \underbrace{e^{-\frac{s^2}{2}} s^2 f} + e^{-\frac{s^2}{2}} \frac{2E}{\hbar \omega} f = 0$$

$$\Leftrightarrow e^{-\frac{s^2}{2}} \left(\frac{d^2}{ds^2} - 2s \frac{d}{ds} + \frac{2E}{\hbar \omega} - 1 \right) f(s) = 0$$

$E_n = \hbar \omega \left(n + \frac{1}{2} \right)$
 $\frac{2E}{\hbar \omega} - 1$
 $2n + 1 = 1$

$\therefore H_n(s)$ 2 個ある!

$$\psi_n(x) \propto H_n(s) e^{-\frac{s^2}{2}}$$

規格化する。

$$1 = \int dx |\psi(x)|^2 \quad x = \sqrt{\frac{\hbar}{m\omega}} s$$

$$= |c|^2 \int ds H_n(s) H_n(s) e^{-s^2}$$

$$= |c|^2 \sqrt{\frac{\hbar}{m\omega}} \int ds H_n(s) H_n(s) e^{-s^2}$$

$$= |c|^2 \sqrt{\frac{\hbar}{m\omega}} \sqrt{\pi} 2^n n!$$

$$\therefore c = \sqrt{\frac{1}{2^n n!}} \sqrt{\frac{m\omega}{\pi \hbar}}$$

$$\therefore \psi(x) = \sqrt{\frac{1}{2^n n!}} \sqrt{\frac{m\omega}{\pi \hbar}} H_n\left(\sqrt{\frac{m\omega}{\hbar}} x\right) \exp\left(-\frac{m\omega}{2\hbar} x^2\right)$$

$$H_0 = 1, H_1 = 2J, H_2 = 4J^2 - J$$

1.3 電磁場の量子化

準備

$$\mathbf{B} = \nabla \times \mathbf{A}$$

$$\mathbf{E} = -\frac{\partial \mathbf{A}}{\partial t}$$