

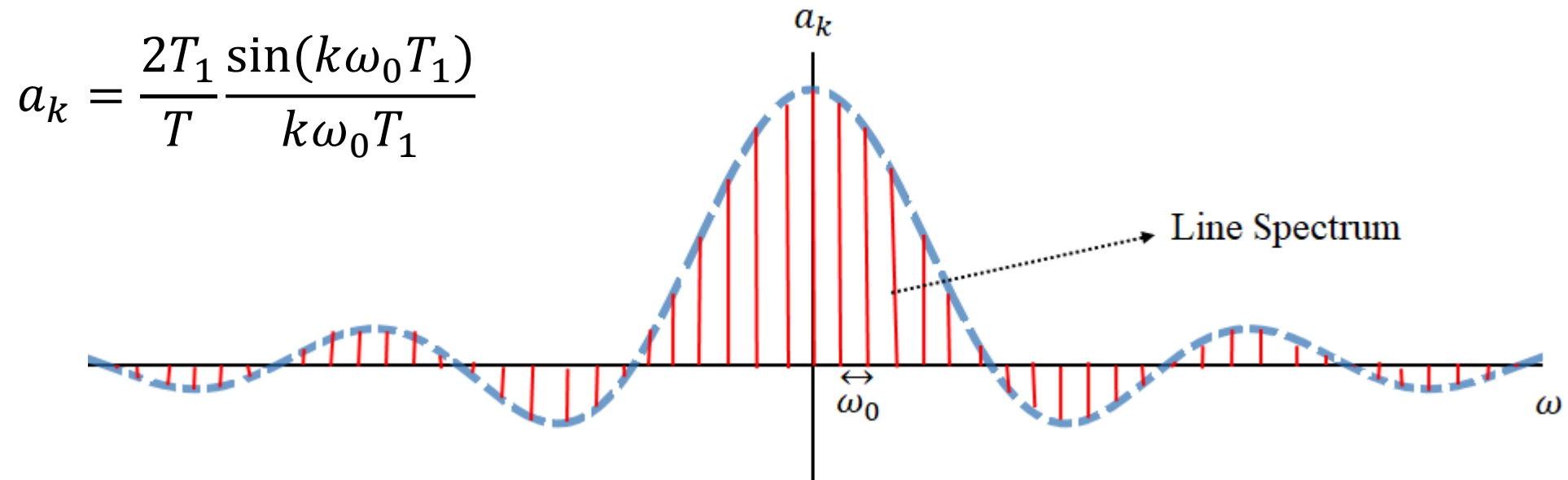
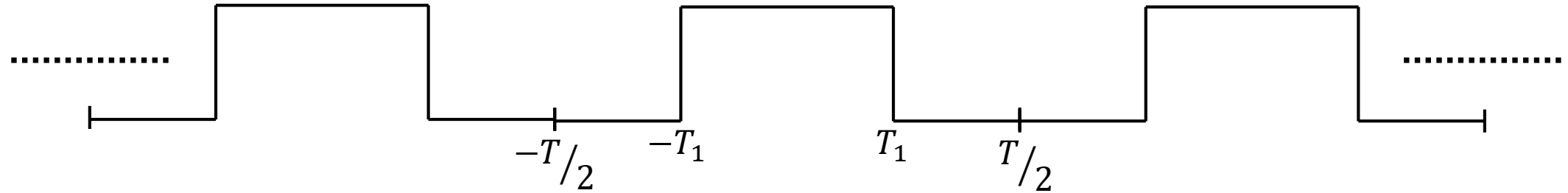
# Lecture 21

## Signals and Systems (ELL205)

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# Lecture 21: Introduction to Fourier Transforms

# Continuous-Time Fourier Series



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$$a_k = \frac{2T_1}{T} \frac{\sin(k\omega_0 T_1)}{k\omega_0 T_1}$$

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By substituting  $k\omega_0 = \omega$

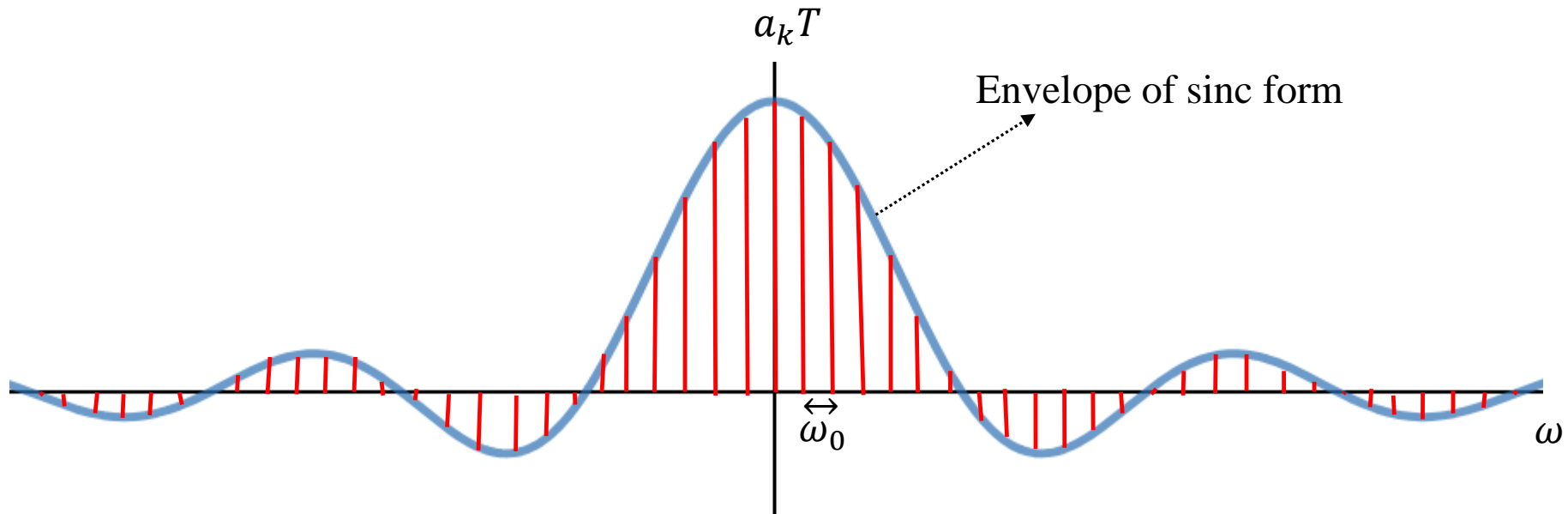
$$a_k T = 2T_1 \frac{\sin(\omega T_1)}{\omega T_1}$$

# Continuous-Time Fourier Series

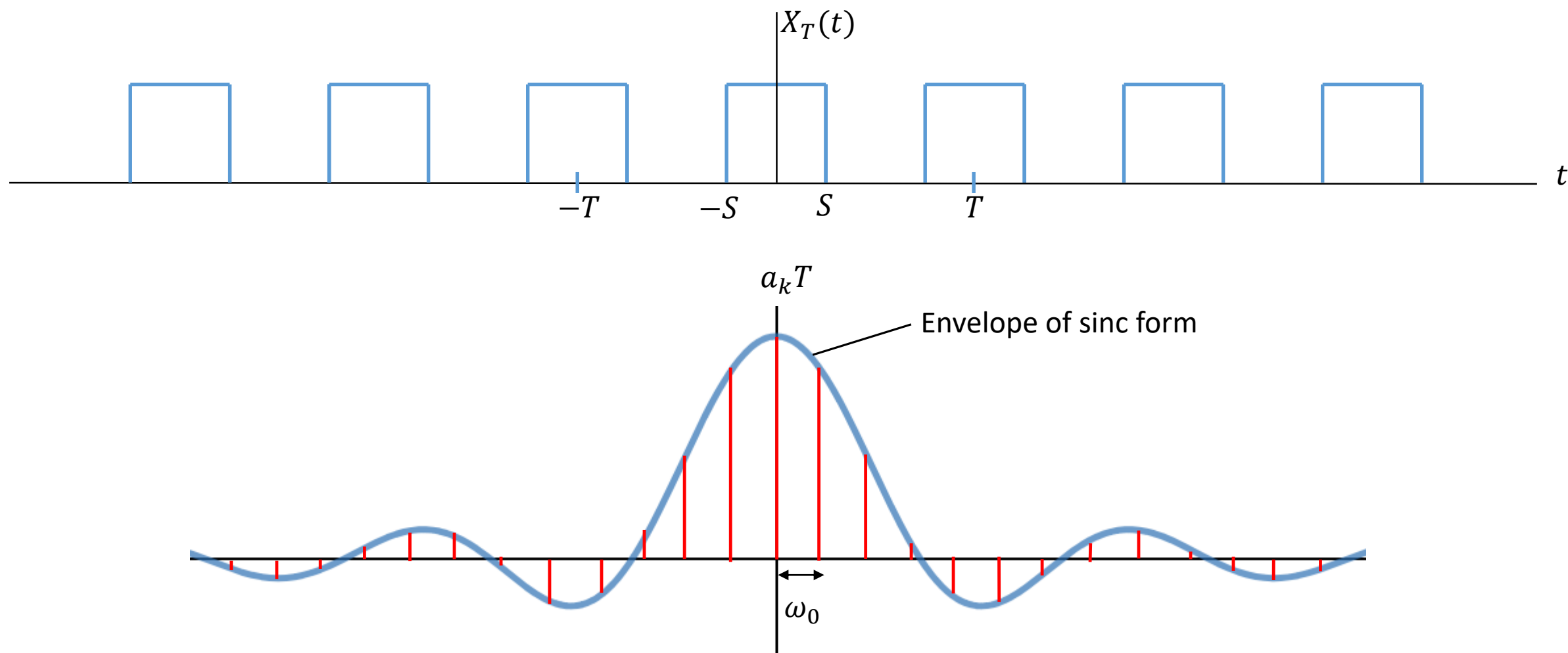
$$a_k = \frac{2T_1}{T} \frac{\sin(k\omega_0 T_1)}{k\omega_0 T_1}$$

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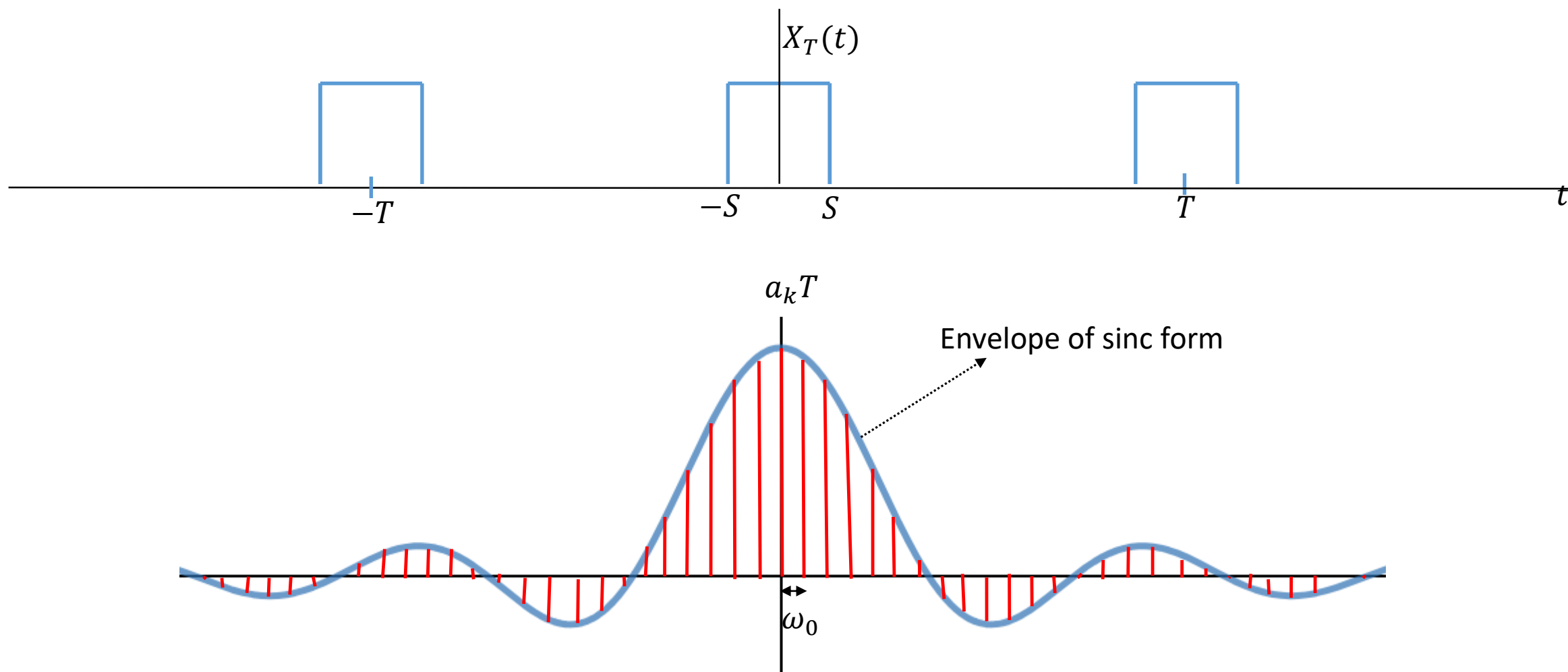
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# Effect of changing period

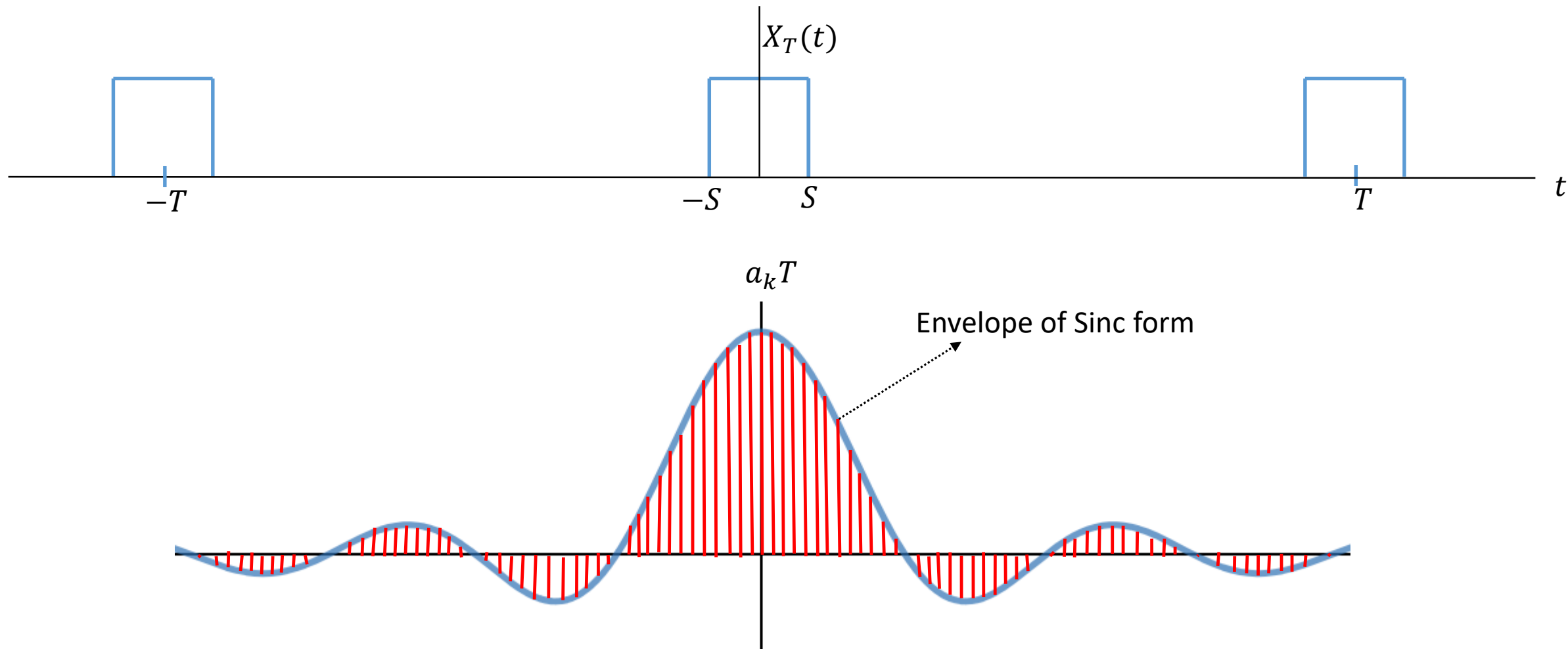


# Effect of changing period



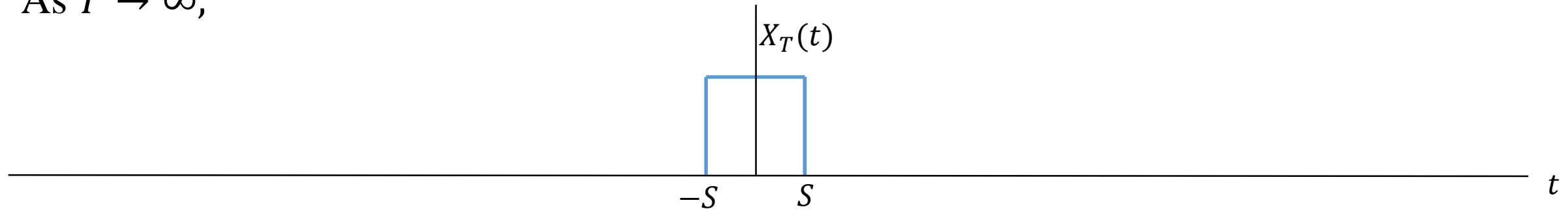


# Effect of changing period

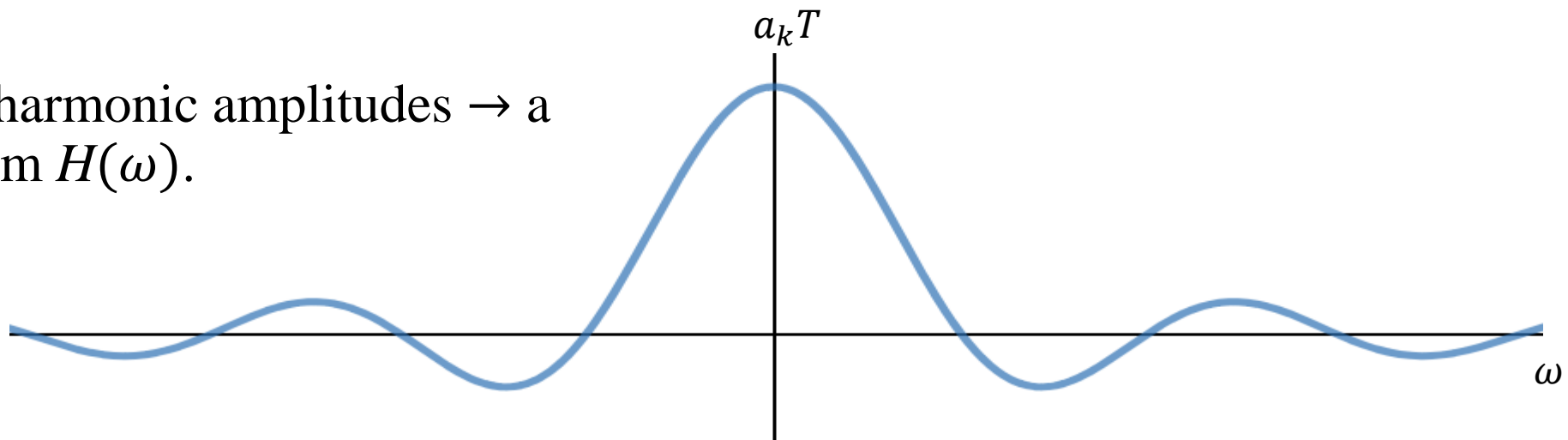


# Effect of changing period

As  $T \rightarrow \infty$ ,



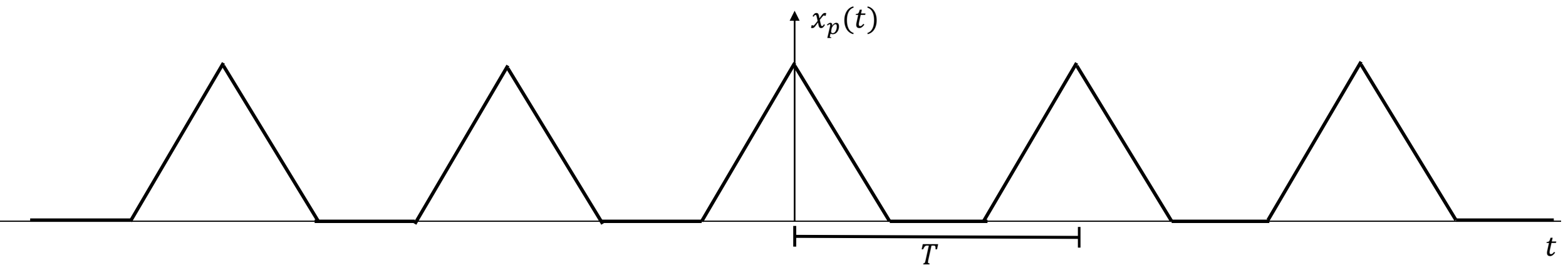
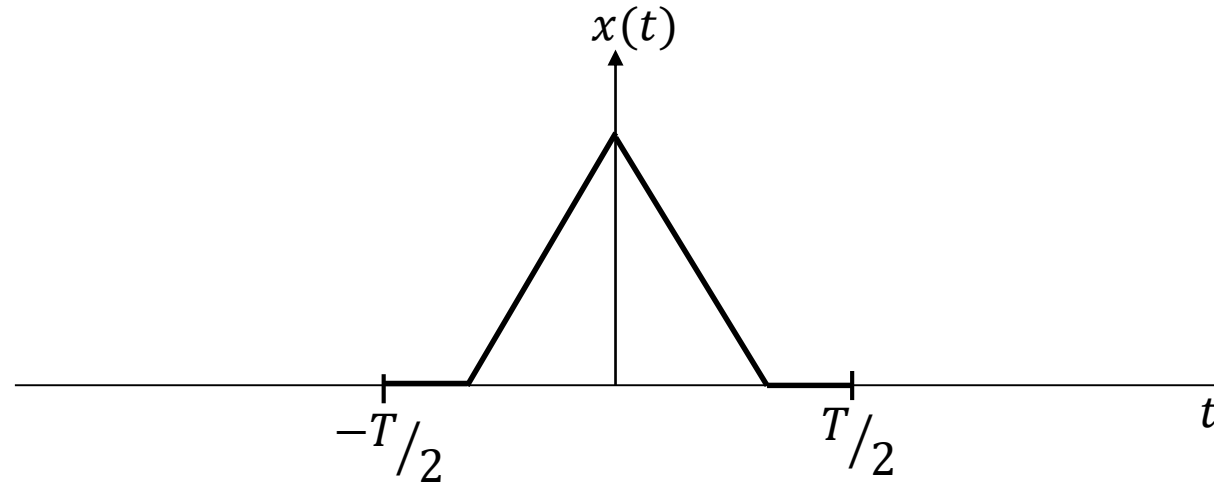
discrete harmonic amplitudes  $\rightarrow$  a  
continuum  $H(\omega)$ .



# Summary

Period	Spectrum	Analysis tool
Aperiodic	Continuous	Fourier Transforms
Periodic	Discrete (line)	Fourier Series

# Analysis and Synthesis Eq. of Fourier Transform



$x_p(t) \rightarrow$  periodic version of  $x(t)$  with period  $T$  (periodic extension)

# Analysis & Synthesis

$$x_p(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t}$$

$$a_k = \frac{1}{T} \int_{-\frac{T}{2}}^{\frac{T}{2}} x_p(t) e^{-jk\omega_0 t} dt$$

# Analysis & Synthesis

Step 1:  $a_k = \frac{1}{T} \int_{-\frac{T}{2}}^{\frac{T}{2}} x_p(t) e^{-jk\omega_0 t} dt$

Step 2: Since  $x(t)$  and  $x_p(t)$  are same in  $-\frac{T}{2}$  to  $+\frac{T}{2}$  so  $x(t)$  replaces  $x_p(t)$

$$a_k = \frac{1}{T} \int_{-\frac{T}{2}}^{\frac{T}{2}} x(t) e^{-jk\omega_0 t} dt$$

Step 3: Since  $x(t)$  exists only in  $-\frac{T}{2}$  to  $+\frac{T}{2}$  so  $a_k$  is:

$$a_k = \frac{1}{T} \int_{-\infty}^{\infty} x(t) e^{-jk\omega_0 t} dt$$

# Analysis & Synthesis

Step 4:  $a_k = \frac{1}{T} \int_{-\infty}^{\infty} x(t) e^{-jk\omega_0 t} dt \quad \Rightarrow \quad a_k = \frac{1}{T} X(k\omega_0)$

$$H(\omega) \triangleq \int_{-\infty}^{\infty} h(\tau) e^{-j\omega\tau} d\tau,$$

and thus  $X(\omega) \triangleq \int_{-\infty}^{\infty} x(\tau) e^{-j\omega\tau} d\tau$

# Analysis & Synthesis

Step 5:  $x_p(t) = \sum_{k=-\infty}^{\infty} \frac{1}{T} X(k\omega_0) e^{jk\omega_0 t}$  using  $a_k = \frac{1}{T} X(k\omega_0)$

Step 6: Substituting  $T = \frac{2\pi}{\omega_0}$ ,  $x_p(t) = \frac{1}{2\pi} \sum_{k=-\infty}^{\infty} X(k\omega_0) e^{jk\omega_0 t} \omega_0$

Step 7: As  $T \rightarrow \infty$ ,  $x_p(t) = x(t)$ ,  $k\omega_0 \rightarrow \omega$ ,  $\omega_0 \rightarrow d\omega$ ,  $\sum \rightarrow \int$

$$x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(\omega) e^{j\omega t} d\omega$$



# Analysis & Synthesis

**Synthesis:**

$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t}$$

**Analysis:**

$$a_k = \frac{1}{T} \int_{-\frac{T}{2}}^{\frac{T}{2}} x(t) e^{-jk\omega_0 t} dt$$

**Fourier  
Series**

$$x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(\omega) e^{j\omega t} d\omega$$

$\omega \rightarrow t$  Inverse Fourier Transform

$$X(\omega) = \int_{-\infty}^{\infty} x(\tau) e^{-j\omega \tau} d\tau$$

$t \rightarrow \omega$  Fourier Transform

**Fourier  
Transform**

# Physical meaning of $X(\omega)$

$$a_k = \frac{1}{T} X(k\omega_0)$$

$$a_k = \frac{\omega_0}{2\pi} X(k\omega_0)$$

When  $T \rightarrow \infty$ ,

$$a_\omega = \frac{X(\omega)d\omega}{2\pi}$$

$$2\pi a_\omega = X(\omega)d\omega$$

$X(\omega)$  is amplitude spectral density

# Convergence of Fourier Transforms

$$X(\omega) = \int_{-\infty}^{\infty} x(\tau) e^{-j\omega\tau} d\tau$$

$$\tilde{x}(t) = \frac{1}{2\pi} \int_{-W}^W X(\omega) e^{j\omega t} d\omega$$

$$e(t) = x(t) - \tilde{x}(t)$$

- $x(\tau)$  square integrable,

$$\text{If } \int_{-\infty}^{\infty} |x(\tau)|^2 d\tau < \infty \quad \text{Then } \int_{-\infty}^{\infty} \lim_{W \rightarrow \infty} |e(t)|^2 dt \rightarrow 0$$

# Convergence of Fourier Transforms

- Dirichlet conditions

If  $\int_{-\infty}^{\infty} |x(\tau)| d\tau < \infty$  and  $x(t)$  is “well-behaved”

Then  $\lim_{W \rightarrow \infty} e(t) \rightarrow 0$  except at discontinuities

Well-behaved  $\triangleq$

- 1) Finite no. of maxima and minima in a finite time period
- 2) Finite no. of finite discontinuities in a finite time period

These are sufficient conditions not necessary conditions. Ex.  $u(t)$  is neither absolutely integrable nor an energy signal but has a Fourier Transform.

# Try yourself:

Fourier Transform of  $e^{-at}u(t)$  is  $X(\omega)$  where  $a > 0$

How many statements are correct?

1) $X(\omega) = \frac{1}{a+j\omega}$	2) $ X(\omega) $ is an even function
3) $\angle X(\omega)$ is an odd function	4) $X(\omega)$ is a low pass filter

Try yourself:

$$X(\omega) = \int_{-\infty}^{\infty} e^{-at} e^{-j\omega t} dt u(t)$$

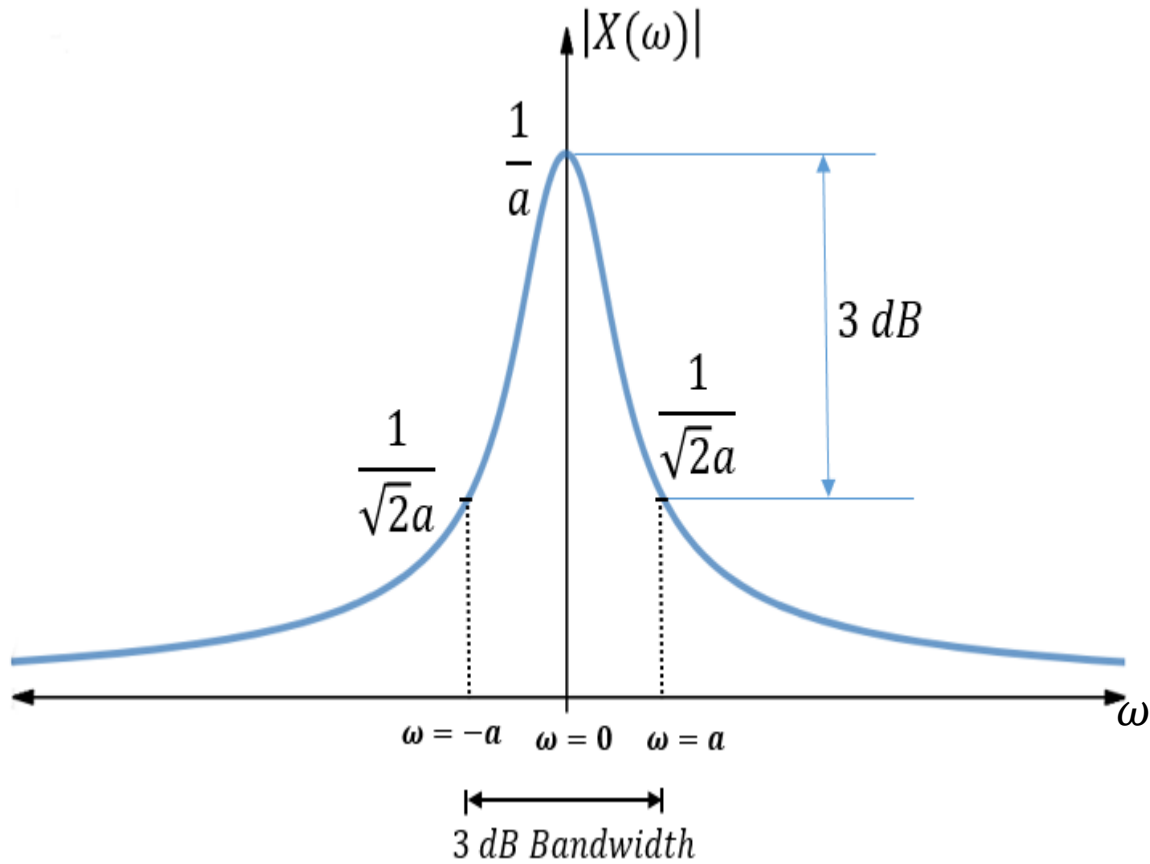
$$= \int_0^{\infty} e^{-(a+j\omega)t} dt$$

$$X(\omega) = \frac{1}{a + j\omega}$$

$$|X(\omega)| = \frac{1}{\sqrt{a^2 + \omega^2}}$$

$$\angle X(\omega) = -\tan^{-1} \left( \frac{\omega}{a} \right)$$

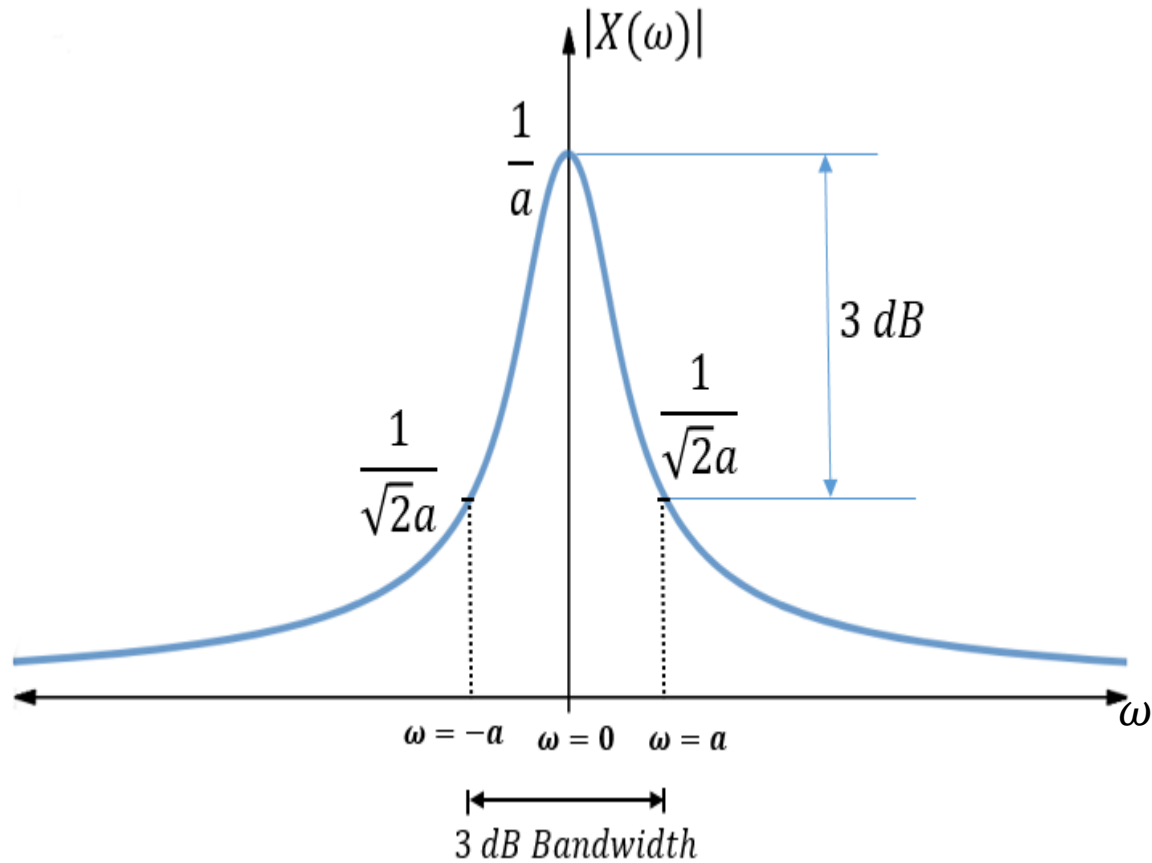
# Try yourself:



$$dB = 20 \log_{10} r$$

Ratio	dB
1	
$\sqrt{2}$	
2	
10	

# Try yourself:

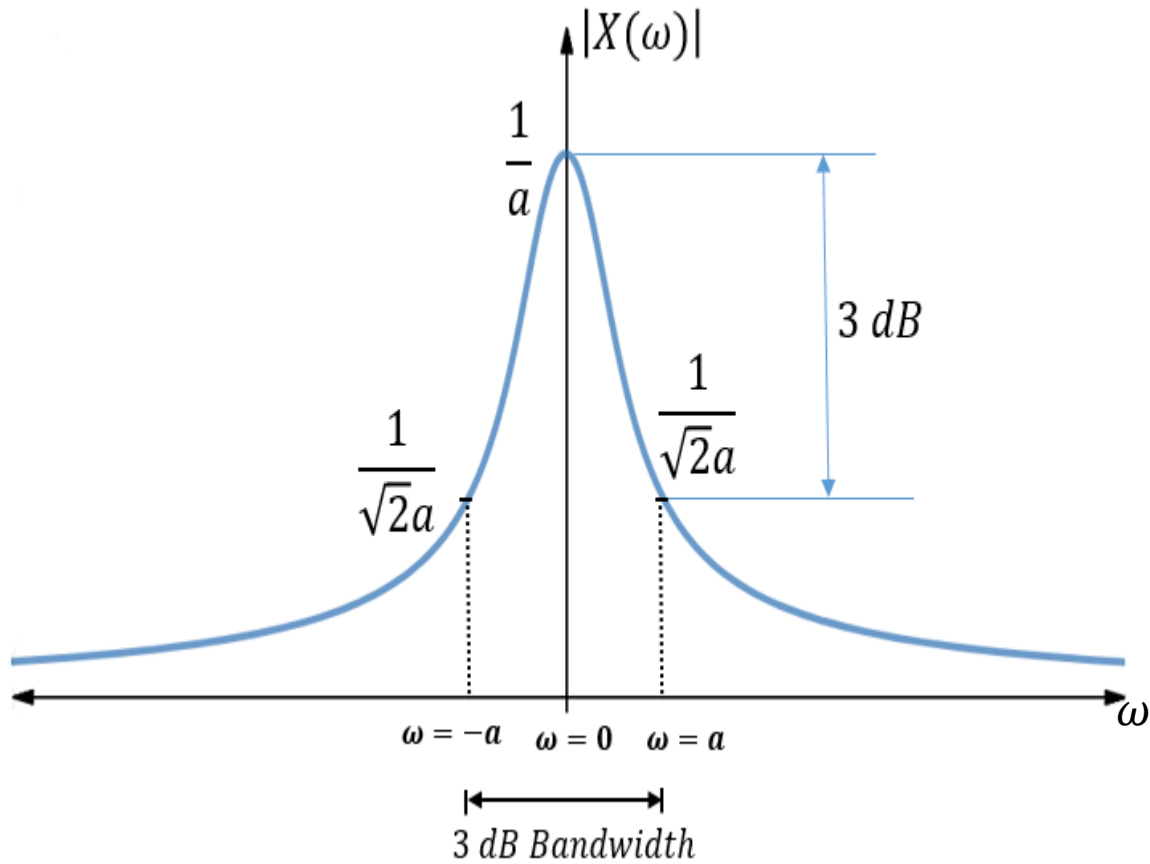


$$dB = 20 \log_{10} r$$

Ratio	dB
1	0
$\sqrt{2}$	
2	
10	



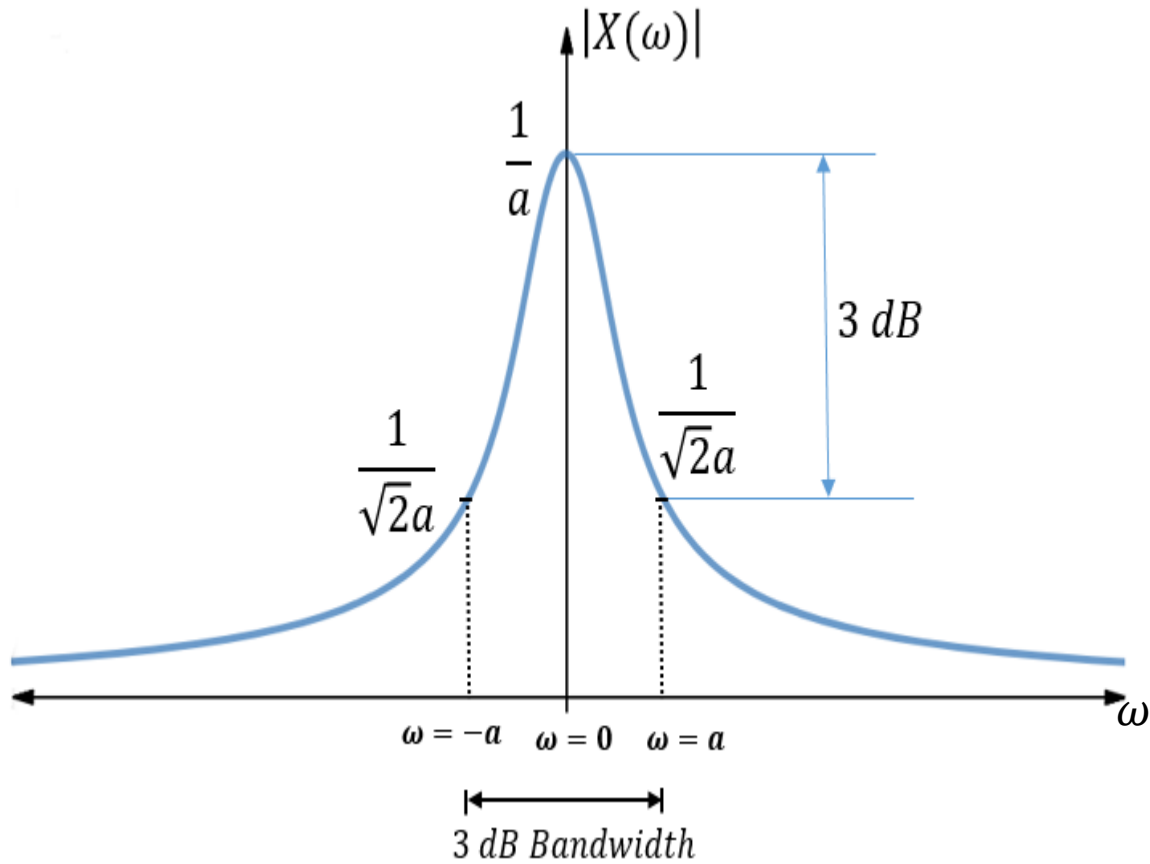
Try yourself:



$$dB = 20 \log_{10} r$$

Ratio	dB
1	0
$\sqrt{2}$	3
2	
10	

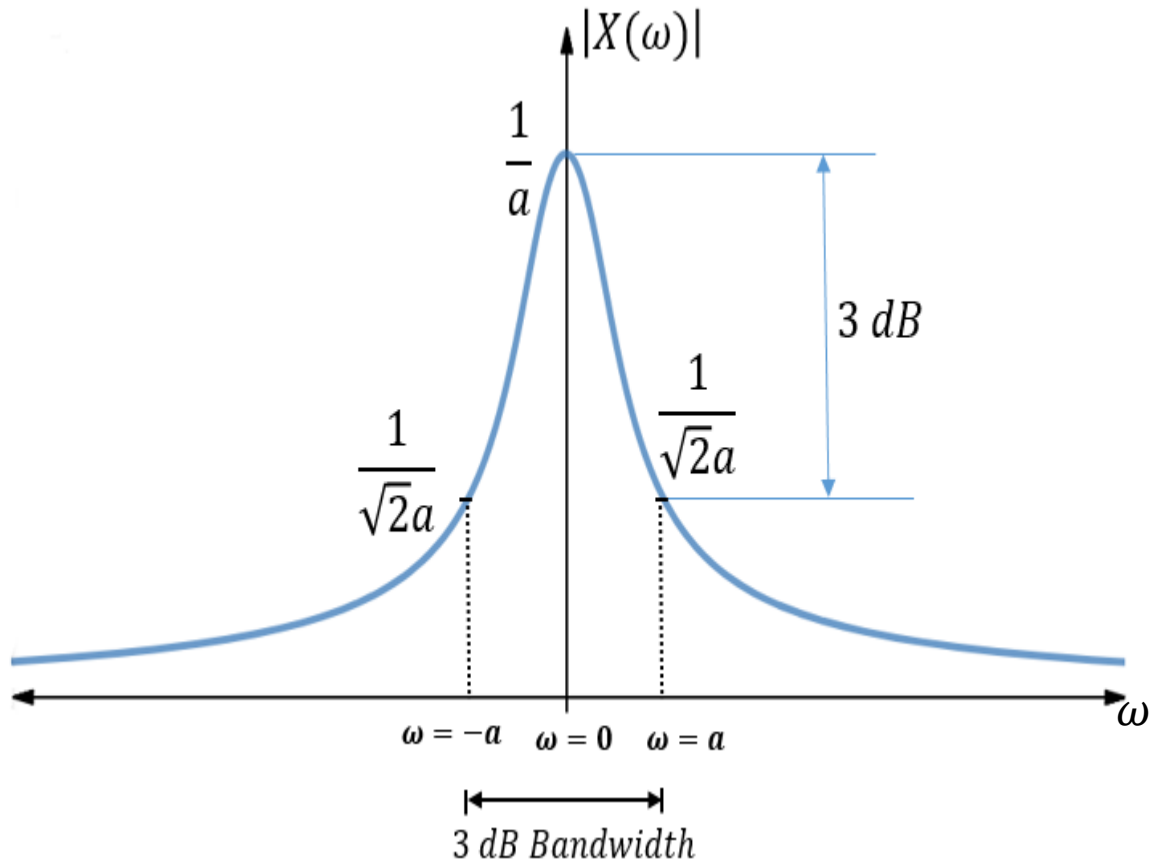
Try yourself:



$$dB = 20 \log_{10} r$$

Ratio	dB
1	0
$\sqrt{2}$	3
2	6
10	

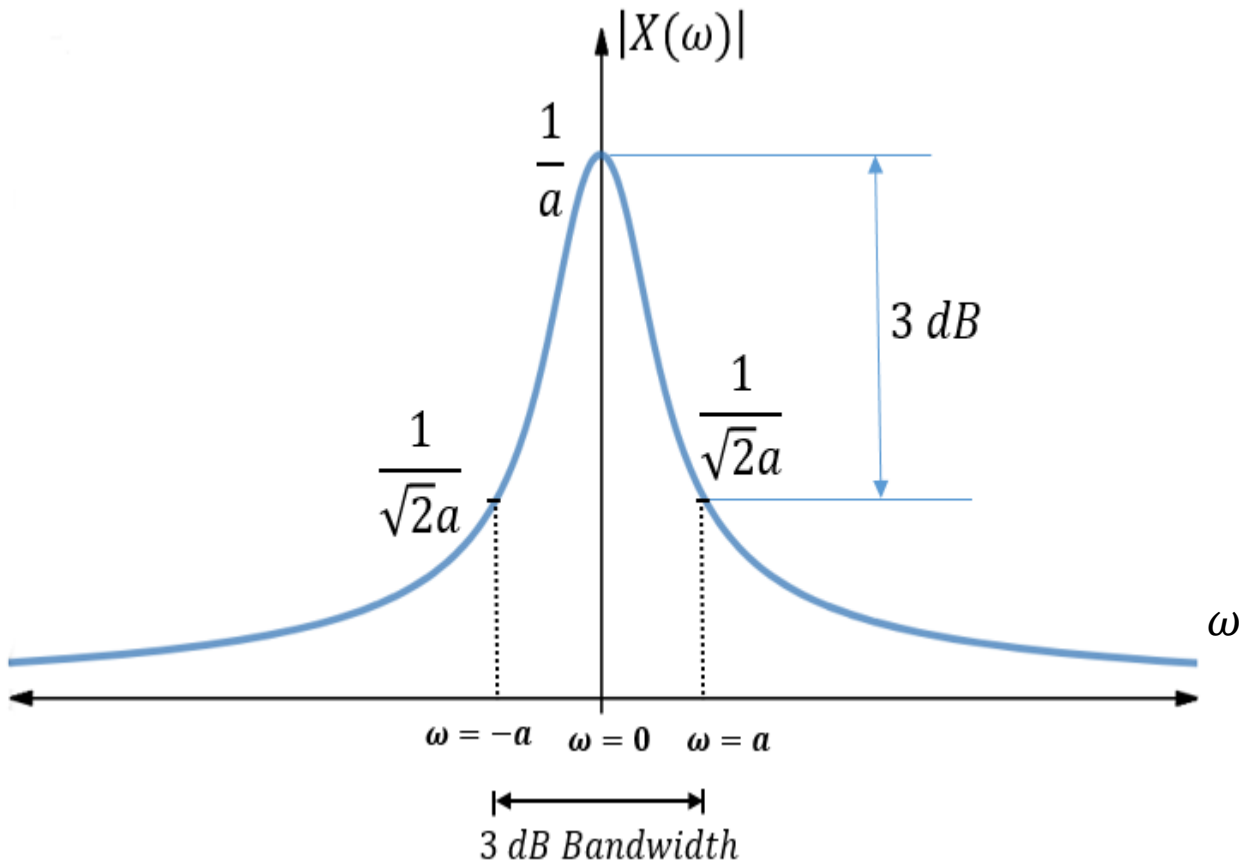
Try yourself:



$$dB = 20 \log_{10} r$$

Ratio	dB
1	0
$\sqrt{2}$	3
2	6
10	20

Try yourself:



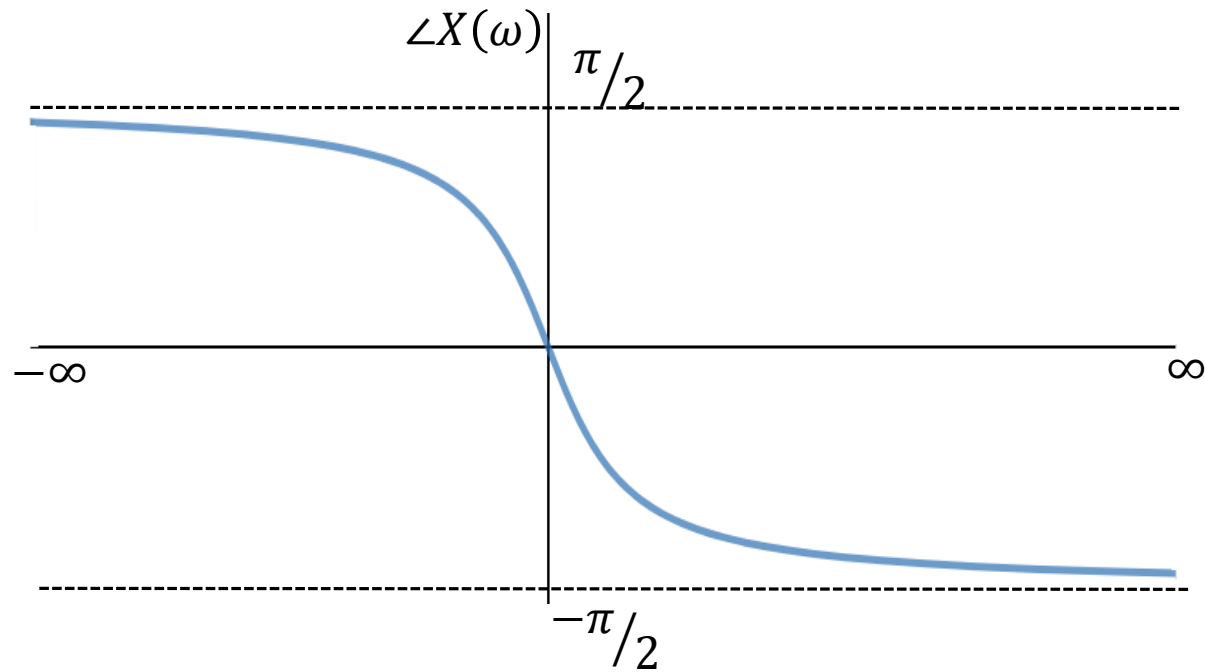
$$\text{dB(decibel)} = 20 \log \left( \frac{1}{a} \times \frac{\sqrt{2}a}{1} \right) \approx 3 \text{ dB}$$

One sided BW =  $a$

Two sided BW =  $2a$

By convention we normally use one-sided BW.

Try yourself:



$$\angle X(\omega) = -\tan^{-1}\left(\frac{\omega}{a}\right)$$