Quantum Mechanics - Lecture 7

Brajesh Kumar Mani



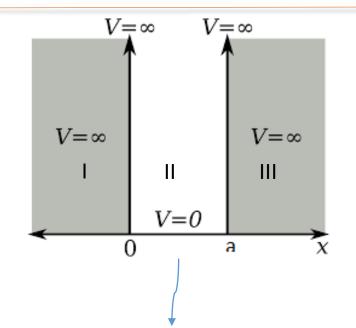
Infinite square well/Particle in a box

"bound states"

The particle is in the potential well, such that $V(x) = \begin{cases} 0, & \text{if } 0 \le x \le a \\ \infty, & \text{otherwise} \end{cases}$

Find:

- (a) Wave function of the particle, $\psi(x)$
- (b) Energy of the particle, *E*
- (c) Time evolved wave function, $\Psi(x, t)$



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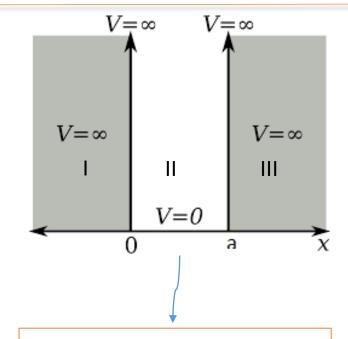
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- The time independent Schrodinger equation in region II is

$$-\frac{\hbar^2}{2m}\,\frac{d^2\psi}{dx^2} = E\psi$$

$$\Rightarrow \frac{d^2\psi}{dx^2} = -k^2\psi$$
, where $k = \frac{\sqrt{2mE}}{\hbar}$ Eq.(1)



In region I and III wave function is zero because infinite potential

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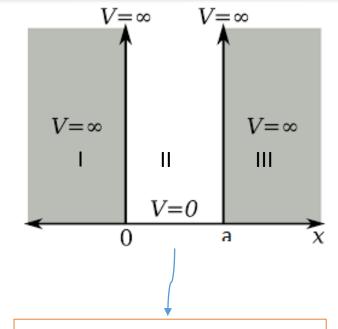
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• Let us assume the solution of this equation in the form $\psi(x) = A \sin kx + B \cos kx, \text{ where A and B are the constants}$

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Using BC(2) we get: $\psi(a) = A \sin ka = 0 \Rightarrow ka = \pm m\pi \Rightarrow k = \pm \frac{m\pi}{a}$, where m = 0, 1, 2, 3

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Then we can write the wave function and energy of the particle as

$$\psi_n(x) = A \sin kx = A \sin \left(\frac{n\pi x}{a}\right)$$
 and $E_n = \frac{k^2 \hbar^2}{2m} = \frac{n^2 \pi^2 \hbar^2}{2ma^2}$

We can calculate the normalization constant A as

$$\int_0^a \psi^*(x)\psi(x)dx = 1 \Rightarrow \int_0^a |A|^2 \sin^2 kx \, dx = 1 \Rightarrow |A|^2 \times \frac{a}{2} = 1 \Rightarrow A = \sqrt{\frac{2}{a}}$$

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$$\psi_{1}(x) = \sqrt{\frac{2}{a}} \sin\left(\frac{\pi x}{a}\right); \quad E_{1} = \frac{\pi^{2} \hbar^{2}}{2ma^{2}}$$

$$\psi_{2}(x) = \sqrt{\frac{2}{a}} \sin\left(\frac{2\pi x}{a}\right); \quad E_{2} = \frac{4\pi^{2} \hbar^{2}}{2ma^{2}}$$

$$\psi_{3}(x) = \sqrt{\frac{2}{a}} \sin\left(\frac{3\pi x}{a}\right); \quad E_{3} = \frac{9\pi^{2} \hbar^{2}}{2ma^{2}}, \text{ and so on.}$$

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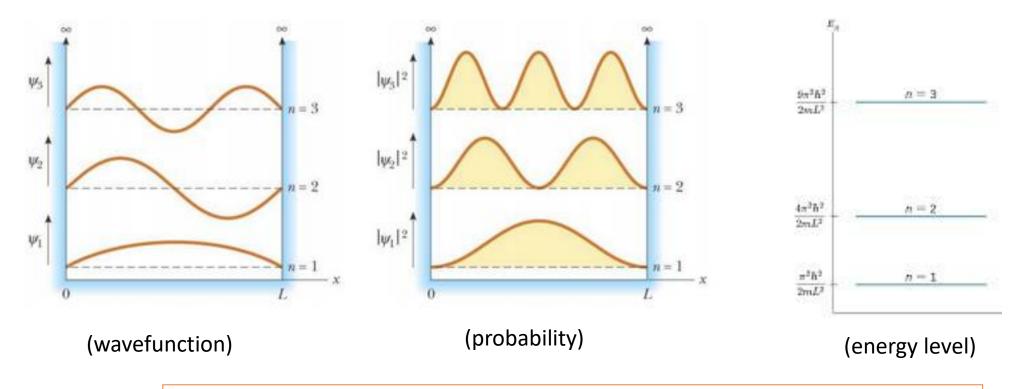
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" $\rightarrow \psi_1, \psi_2, \psi_3$ represent the ground state, first excited, second excited state,"

Let us sketch the wave functions and corresponding energy



- → Wave functions are continuous
- → Confined quantum mechanical systems have discrete energies
- → Represents an example of "bound" quantum system.

We can write the initial wave function as a linear combination

$$\Psi(x,0) = \sum_{n=1}^{\infty} c_n \, \psi_n(x) = \sum_{n=1}^{\infty} c_n \, \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi x}{a}\right)$$

Where,
$$c_n = \langle \psi_n(x) | \Psi(x,0) \rangle = \sqrt{\frac{2}{a}} \int_0^a \sin\left(\frac{n\pi x}{a}\right) \Psi(x,0) dx$$

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The stationary states are then written as

$$\Psi_n(x,t) = \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi x}{a}\right) e^{-i\left(\frac{n^2\pi^2\hbar}{2ma^2}\right)t}$$

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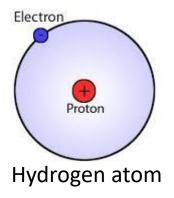
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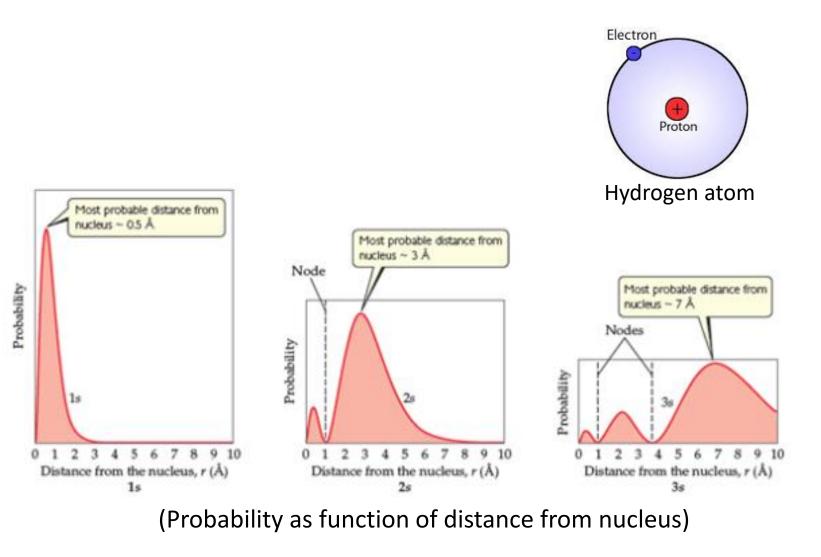
The most general solution of the time-dependent Schrodinger equation is then

$$\Psi(x,t) = \sum_{n=1}^{\infty} c_n \Psi_n(x,t) = \sum_{n=1}^{\infty} c_n \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi x}{a}\right) e^{-i\left(\frac{n^2\pi^2\hbar}{2ma^2}\right)t}$$

Other examples of bound system: Hydrogen atom

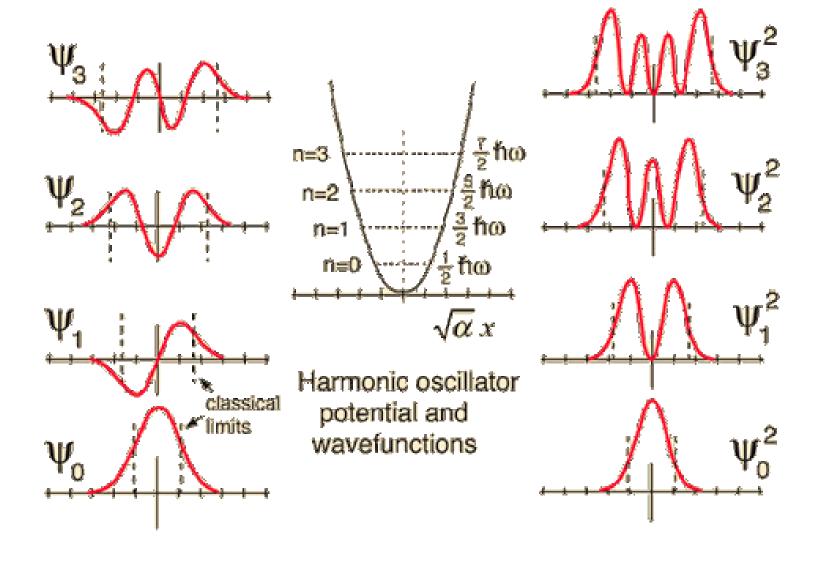


Other examples of bound system: Hydrogen atom



-0.28 eVLyman Paschen Pfund n = 6-0.38 eVseries series series n = $-0.54 \, eV$ -0.85 eV n =n = 3 $-1.51 \, eV$ Brackett series n=2 $-3.40 \, eV$ Balmer series $-13.60 \, eV$ (Energy level diagram)

Other examples of bound system: Harmonic oscillator



Example Problem 4: A particle of mass m moves freely inside an infinite potential well of length a. Initially (at t=0), the particle is in the state

$$\Psi(x,0) = \sqrt{\frac{3}{5a}} \sin\left(\frac{3\pi x}{a}\right) + \sqrt{\frac{1}{5a}} \sin\left(\frac{5\pi x}{a}\right)$$

Find $\Psi(x, t)$ at any later time t.

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Solution: For a particle in an infinite box of length a

$$\psi_n(x) = \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi x}{a}\right) \text{ where: } n = 1, 2, 3, 4, \dots$$

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Using this wave function we can write

$$\Psi(x,0) = \sqrt{\frac{3}{5a}} \sin\left(\frac{3\pi x}{a}\right) + \sqrt{\frac{1}{5a}} \sin\left(\frac{5\pi x}{a}\right) = \sqrt{\frac{3}{10}} \psi_3(x) + \sqrt{\frac{1}{10}} \psi_5(x)$$

We can write the time-dependent wavefunction as $\Psi(x,t) = \sum_{n} \Psi_n(x,0) e^{-\frac{iE_nt}{\hbar}}$

$$\Psi(x,t) = \sqrt{\frac{3}{10}} \psi_3(x) e^{-\frac{iE_3t}{\hbar}} + \sqrt{\frac{1}{10}} \psi_5(x) e^{-\frac{iE_5t}{\hbar}}$$

Where,
$$E_3 = \frac{9\pi^2\hbar^2}{2ma^2}$$
 $E_5 = \frac{25\pi^2\hbar^2}{2ma^2}$

Example Problem 5: A particle which is confined to move in an one-dimensional region

$$0 \le x \le a$$
 is represented by the wavefunction $\Psi(x,t) = \sin\left(\frac{\pi x}{a}\right) \exp(-i\omega t)$. Then

- (a) Find the potential V(x)
- (b) Calculate the probability of finding the particle in the interval $\frac{a}{4} \le x \le \frac{3a}{4}$

Solution:

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- (a) Find the potential V(x)
- (b) Calculate the probability of finding the particle in the interval $\frac{a}{4} \le x \le \frac{3a}{4}$

Solution:

(a) We know that the time-dependent Schrodinger equation is

$$i\hbar \frac{\partial \Psi(x,t)}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \Psi(x,t)}{\partial x^2} + V(x)\Psi(x,t)$$

From the given wavefunction, we can calculate: $\frac{\partial \Psi(x,t)}{\partial t} = -i\omega \sin \frac{\pi x}{a} \exp -i\omega t = -i\omega \Psi(x,t),$

and
$$\frac{\partial^2 \Psi(x,t)}{\partial x^2} = -\left(\frac{\pi}{a}\right)^2 \Psi(x,t)$$

We can now write

$$i\hbar (-i\omega)\Psi(x,t) = -\frac{\hbar^2}{2m} \left(-\frac{\pi^2}{a^2}\right) \Psi(x,t) + V(x)\Psi(x,t)$$

$$\Rightarrow V(x) = \hbar\omega - \frac{\hbar^2\pi^2}{2ma^2}$$

(b) We can write the probability as

$$P = \frac{\int_{\frac{a}{4}}^{\frac{3a}{4}} \Psi^{*}(x,t) \Psi(x,t) dx}{\int_{0}^{a} \Psi^{*}(x,t) \Psi(x,t) dx} = \frac{\int_{\frac{a}{4}}^{\frac{3a}{4}} \psi^{*}(x) \psi(x) dx}{\int_{0}^{a} \psi^{*}(x) \psi(x) dx}$$

$$\Rightarrow \frac{\int_{\frac{a}{4}}^{\frac{3a}{4}} \sin^2\left(\frac{\pi x}{a}\right) dx}{\int_{0}^{a} \sin^2\left(\frac{\pi x}{a}\right) dx} = \frac{(\pi + 2)}{2\pi} = 0.82$$

Home work problem:

A particle in the infinite square well has the initial wave function

$$\Psi(x,0) = Ax(a-x); \ 0 \le x \le a.$$

The constant A is a normalization constant. Calculate the wave function $\Psi(x,t)$.