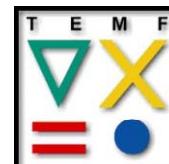


# Technical Electrodynamics for iCE

(in English language)

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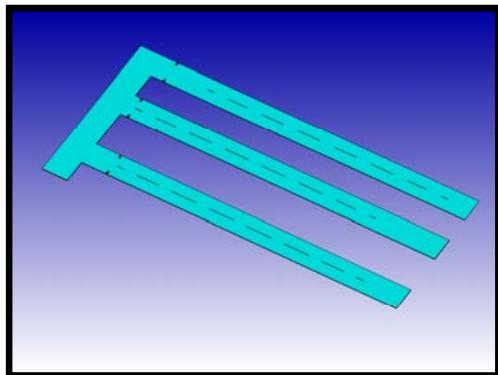


Graduate School of Computational Engineering (GSC)

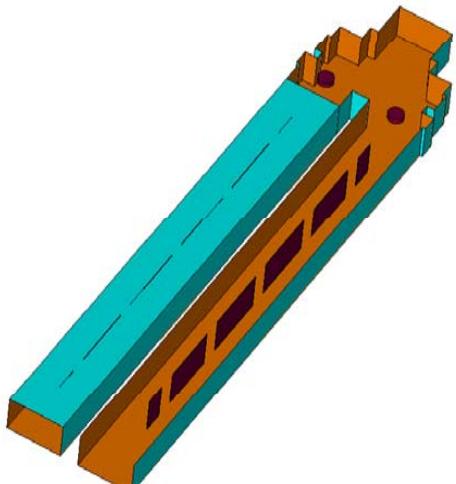
## Waveguides

# Waveguide structures: applications

Multiplexer

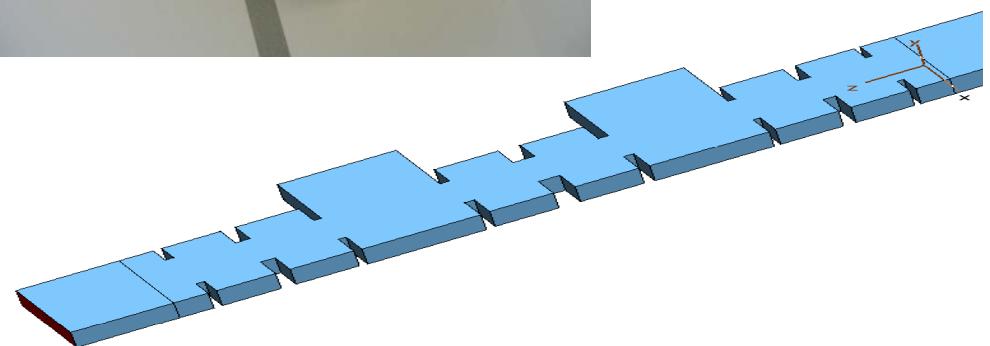
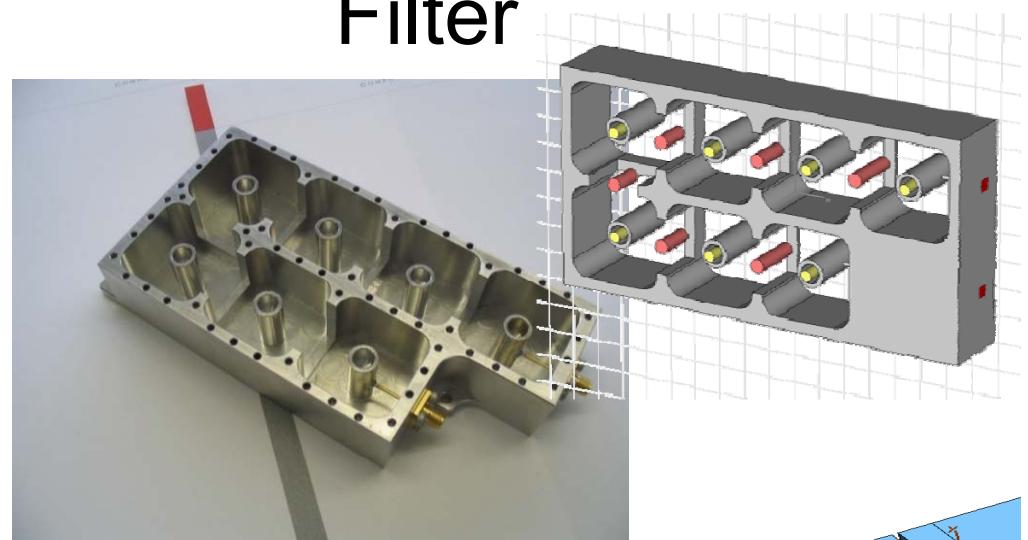


Diplexer

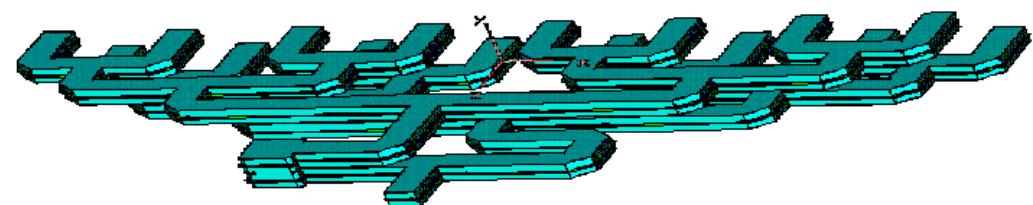


Source: CST AG

Filter



16-way power divider



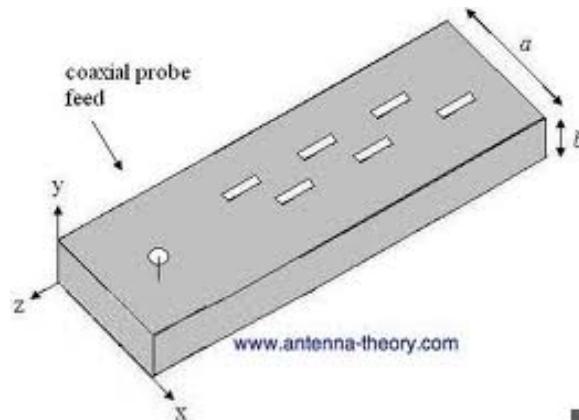
# Waveguide structures: applications



Source: Microwave Engineering Services Corp.



Source: Quinstar.com

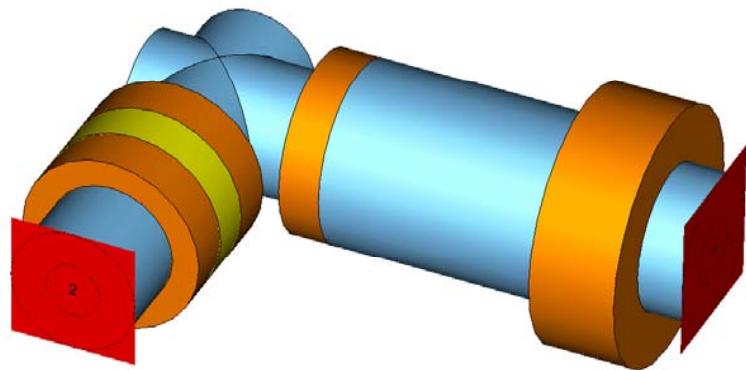


Source: antennatheory.com



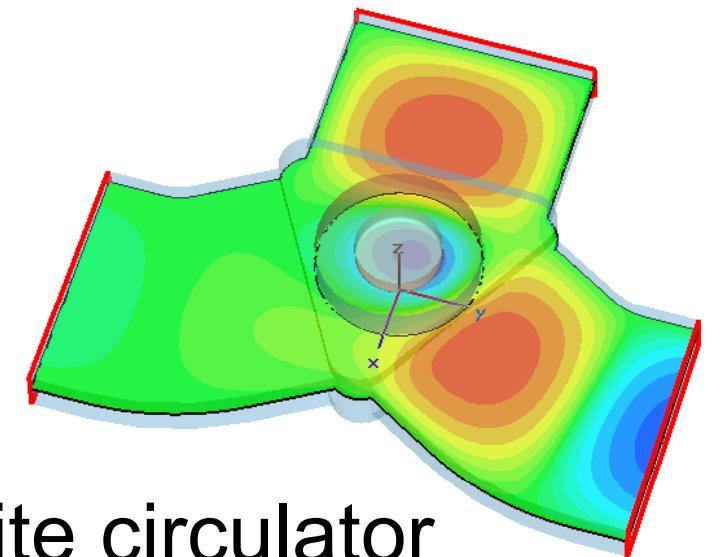
Source: fairviewmicrowave.com

# Waveguide structures: applications



Coaxial connectors

Source: CST AG



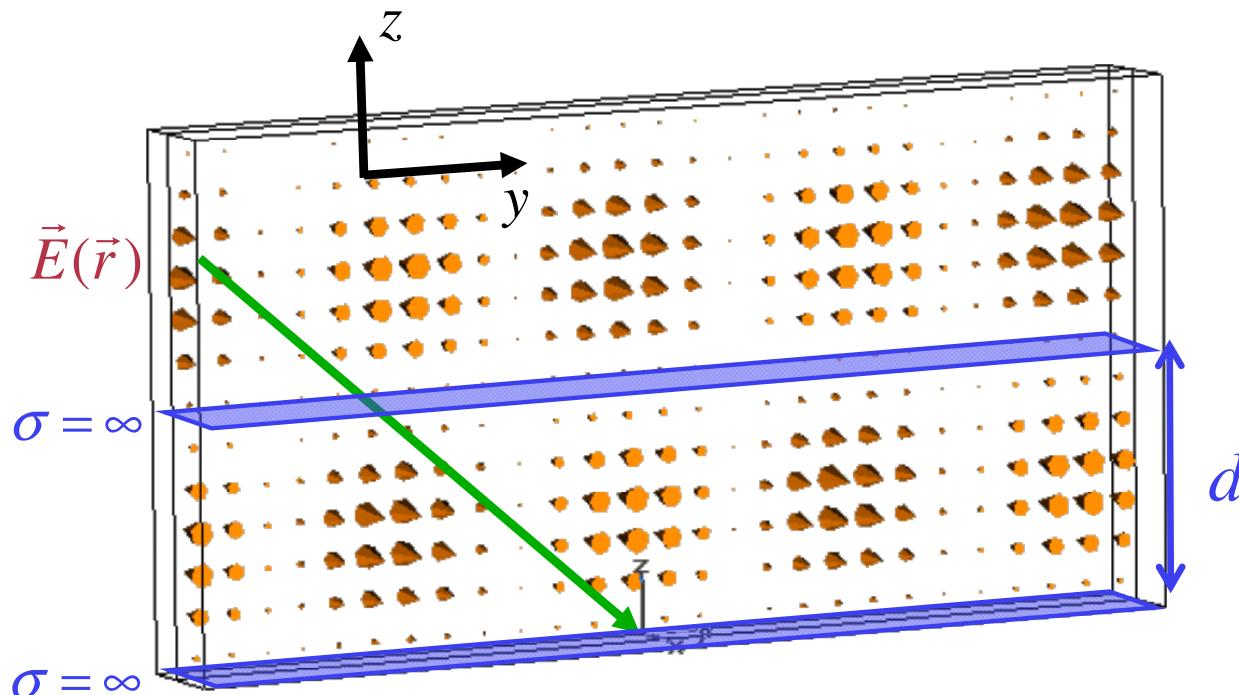
Ferrite circulator

Source: CST AG

- **Transmit power from a source to a load for antennas, magnetrons, power division, ...**
- **Able to transfer large powers → e.g. microwave oven**
- **Mechanically robust → used in safety-critical applications (e.g. space applications such as satellites)**



# Parallel Plate Waveguides



$$Z_2 = 0 \rightarrow r_s = -1$$

$$E_x(y, z) \propto e^{-jk_1 \sin \phi_i y}.$$

$$\sin(k_1 \cos \phi_i z)$$

$$k_1 = \frac{\omega}{c}$$

transversal:

$$\sin(k_1 \cos \phi_i d) = 0 \Rightarrow \frac{\omega}{c} \cos \phi_i d = n\pi$$

longitudinal:

$$\beta_y = k_1 \sin \phi_i = \frac{\omega}{c} \sin \phi_i$$

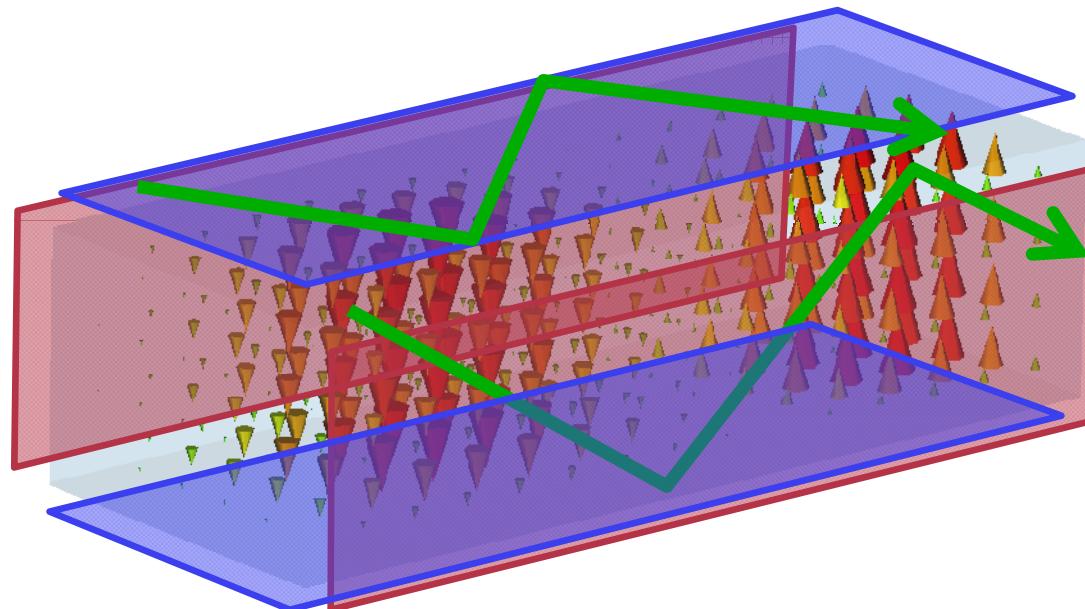


$$\beta_y^2 + \left( \frac{n\pi}{d} \right)^2 = \left( \frac{\omega}{c} \right)^2$$

dispersion relation of a parallel plate waveguide

## ● From plane waves to guided waves

- ☞ Oblique incidence of plane waves at metallic half plane
  - standing wave part transversal to the metallic plane
- ☞ Repeat this principle in 2<sup>nd</sup> direction:



→ (Rectangular)  
Hollow  
Waveguide

- ☞ Modified dispersion characteristics

- **Rectangular waveguide:**  $a>b$ , perfectly conducting walls

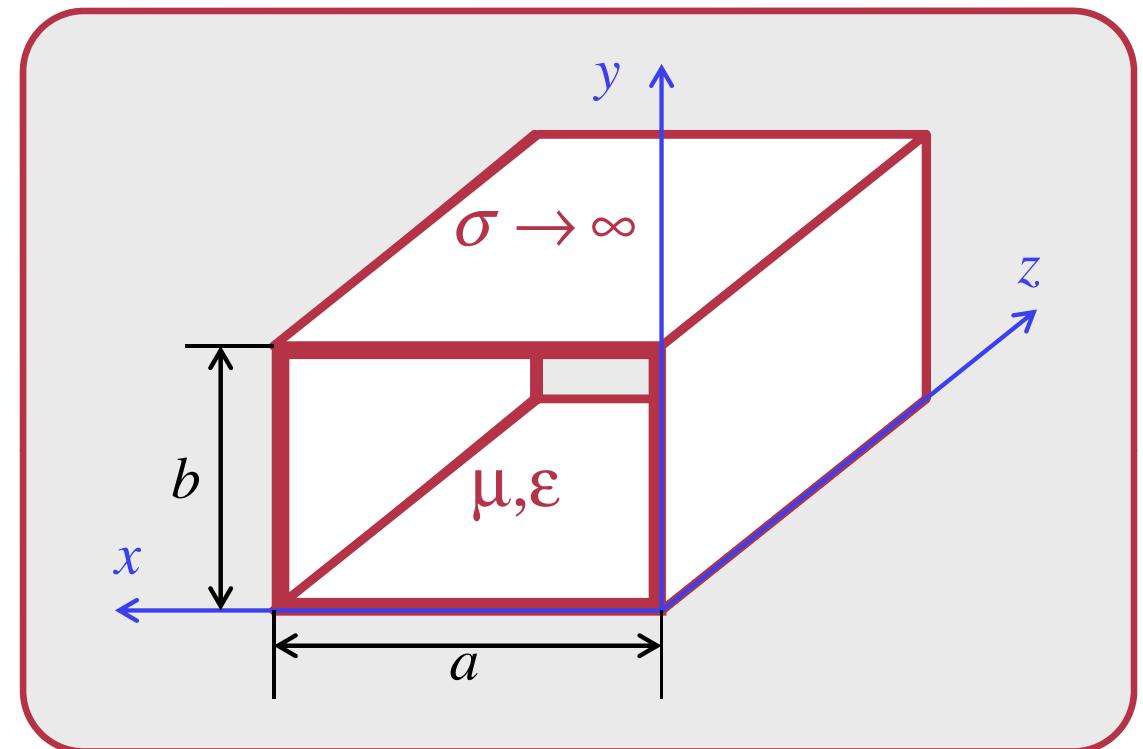
☞ **2 Wave types:**

- **TE (‘transversal electric’)**

$$E_z = 0$$

- **TM (‘transversal magnetic’)**

$$H_z = 0$$



(Lecture Notes pp. 49-57)

- Wave equation for vector potential (homogeneous medium):

$$\operatorname{div} \vec{E} = 0 \Rightarrow \vec{E} = \operatorname{curl} \vec{A} \quad (TE\text{-case, TM: equivalent})$$

$$\operatorname{curl} \vec{E} = \operatorname{curl} \operatorname{curl} \vec{A} = \operatorname{grad} \operatorname{div} \vec{A} - \Delta \vec{A} = -j\omega\mu \vec{H} \quad (\text{Faraday's law})$$

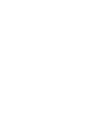
$$\operatorname{curl} \vec{H} = j\omega\epsilon \vec{E} = j\omega\epsilon \operatorname{curl} \vec{A} \quad (\text{Ampere's law})$$

$$\Rightarrow \vec{H} = j\omega\epsilon \vec{A} - \operatorname{grad} \phi$$

$$\operatorname{grad} \operatorname{div} \vec{A} - \Delta \vec{A} = \omega^2 \mu \epsilon \vec{A} + j\omega\mu \operatorname{grad} \phi$$

$$\text{set: } \operatorname{div} \vec{A} = j\omega\mu \phi$$

(Integration constant)



,Lorentz gauge'

$$\Rightarrow \boxed{\Delta \vec{A} + k^2 \vec{A} = 0} \quad k^2 = \omega^2 \mu \epsilon$$

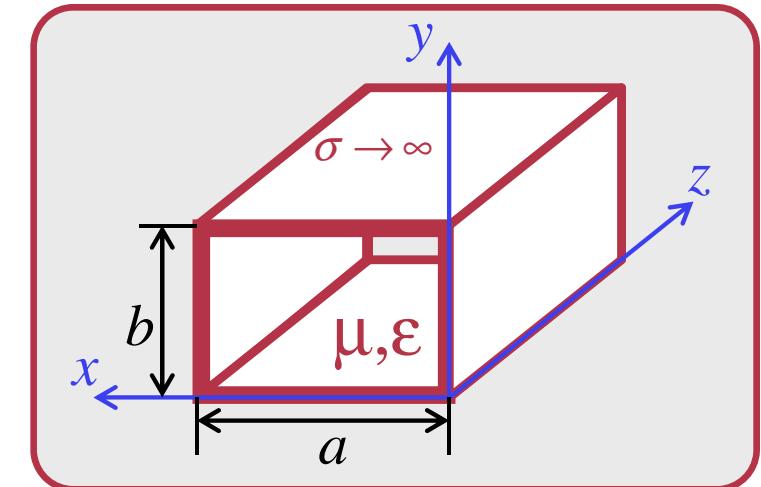
- 2 Wave types:

  - ☞ **TE:**  $E_z = 0$

  - ☞ **TM:**  $H_z = 0$

  - ☞ **Vector potential approach:**

$$\vec{A}(\vec{r}) = A_z(x, y, z) \vec{e}_z$$



(Lecture Notes p. 49)

$$\vec{E} = \text{curl } \vec{A} = \text{curl} \begin{pmatrix} 0 \\ 0 \\ A_z \end{pmatrix} = \begin{pmatrix} * \\ * \\ 0 \end{pmatrix} \rightarrow \text{TE-case}$$

$$\vec{H} = \text{curl } \vec{A} = \text{curl} \begin{pmatrix} 0 \\ 0 \\ A_z \end{pmatrix} = \begin{pmatrix} * \\ * \\ 0 \end{pmatrix} \rightarrow \text{TM-case}$$

} formulate and solve wave equation for  $\vec{A}(\vec{r})$

- **Solutions of the wave equation (Summary):** Rectangular WG

$$\Delta \vec{A} + k^2 \vec{A} = 0 \quad k^2 = \omega^2 \mu \epsilon$$

$$\rightarrow \Delta A_z + k^2 A_z = 0 \quad \vec{A} = \vec{e}_z A_z$$



☞ product approach:  $A_z(x, y, z) = f(x) \cdot g(y) \cdot h(z)$

$$\rightarrow (...) \quad A_z(x, y, z) = \underbrace{\begin{cases} \sin(k_x x) \\ \cos(k_x x) \end{cases}}_{\text{superposition}} \cdot \underbrace{\begin{cases} \sin(k_y y) \\ \cos(k_y y) \end{cases}}_{\text{superposition}} \cdot \underbrace{\begin{cases} e^{+jk_z z} \\ e^{-jk_z z} \end{cases}}_{\text{superposition}}$$

superposition:  $(C_1 \sin(k_x x) + C_2 \cos(k_x x)) \cdot \dots$

☞ *separation relation:*  $k_x^2 + k_y^2 + k_z^2 = k^2$  (insert  $A_z$  in wave eqn.)

☞ To determine unknown coefficients: boundary conditions for fields !

- Derivation of the solutions of the wave equation:

Rectangular WG

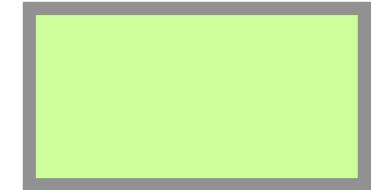


- Solutions of the wave equation:

Rectangular WG

$$\Delta \vec{A} + k^2 \vec{A} = 0 \quad k^2 = \omega^2 \mu \epsilon$$

$$\rightarrow \Delta A_z + k^2 A_z = 0 \quad \vec{A} = \vec{e}_z A_z$$



$$\frac{\partial^2 A_z}{\partial x^2} + \frac{\partial^2 A_z}{\partial y^2} + \frac{\partial^2 A_z}{\partial z^2} + k^2 A_z = 0$$

( $A_z$  complex variable)

1. Product approach:  $A_z(x, y, z) = f(x) \cdot g(y) \cdot h(z)$

2. Replace in the equation and divide by  $A_z$  :

$$\frac{1}{f} \underbrace{\frac{\partial^2 f}{\partial x^2}}_{-k_x^2} + \frac{1}{g} \underbrace{\frac{\partial^2 g}{\partial y^2}}_{-k_y^2} + \frac{1}{h} \underbrace{\frac{\partial^2 h}{\partial z^2}}_{-k_z^2} + k^2 = 0$$

$$k_x^2 + k_y^2 + k_z^2 = k^2$$

!

## ● Solutions of the wave equation:

Rectangular WG

- ☞ Remember the solution of the wave equation:

$$\frac{\partial^2 \underline{E}_x}{\partial z^2} + k^2 \underline{E}_x = 0 \Leftrightarrow \frac{1}{\underline{E}_x} \frac{\partial^2 \underline{E}_x}{\partial z^2} + k^2 = 0$$



$$\Rightarrow \underline{E}_x = \underline{C}_1 e^{-jkz} + \underline{C}_2 e^{jkz}$$

- ☞ By using  $e^{j\alpha} = \cos(\alpha) + j \sin(\alpha)$  the solution can be rewritten as

$$\underline{E}_x = \underline{C}_1 (\cos kz - \sin kz) + \underline{C}_2 (\cos kz + \sin kz) = \underline{D}_1 \cos kz + \underline{D}_2 \sin kz$$

with  $\underline{D}_1 = \underline{C}_1 + \underline{C}_2$ ,  $\underline{D}_2 = \underline{C}_1 - \underline{C}_2$

- ☞ The form with exponential is useful for describing a propagation, the form with cos/sin for describing standing waves

## ● Solutions of the wave equation:

Rectangular WG



$$\text{separation relation: } k_x^2 + k_y^2 + k_z^2 = k^2$$



3. All equations are of the same form and have the solutions:

$$\frac{1}{f} \frac{\partial^2 f}{\partial x^2} = -k_x^2 \Rightarrow f(x) = A_1 \cos(k_x x) + A_2 \sin(k_x x)$$

$$\frac{1}{g} \frac{\partial^2 g}{\partial y^2} = -k_y^2 \Rightarrow g(y) = B_1 \cos(k_y y) + B_2 \sin(k_y y)$$

$$\frac{1}{h} \frac{\partial^2 h}{\partial z^2} = -k_z^2 \Rightarrow h(z) = C_1 e^{jk_z z} + C_2 e^{-jk_z z}$$

- **Solutions of the wave equation:**

5. Final solution:

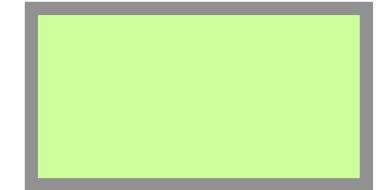
$$f(x) = A_1 \cos(k_x x) + A_2 \sin(k_x x)$$

$$g(y) = B_1 \cos(k_y y) + B_2 \sin(k_y y)$$

$$h(z) = C_1 e^{jk_z z} + C_2 e^{-jk_z z}$$

$$A_z(x, y, z) = f(x) \cdot g(y) \cdot h(z)$$

Rectangular WG



☞ Compact notation:



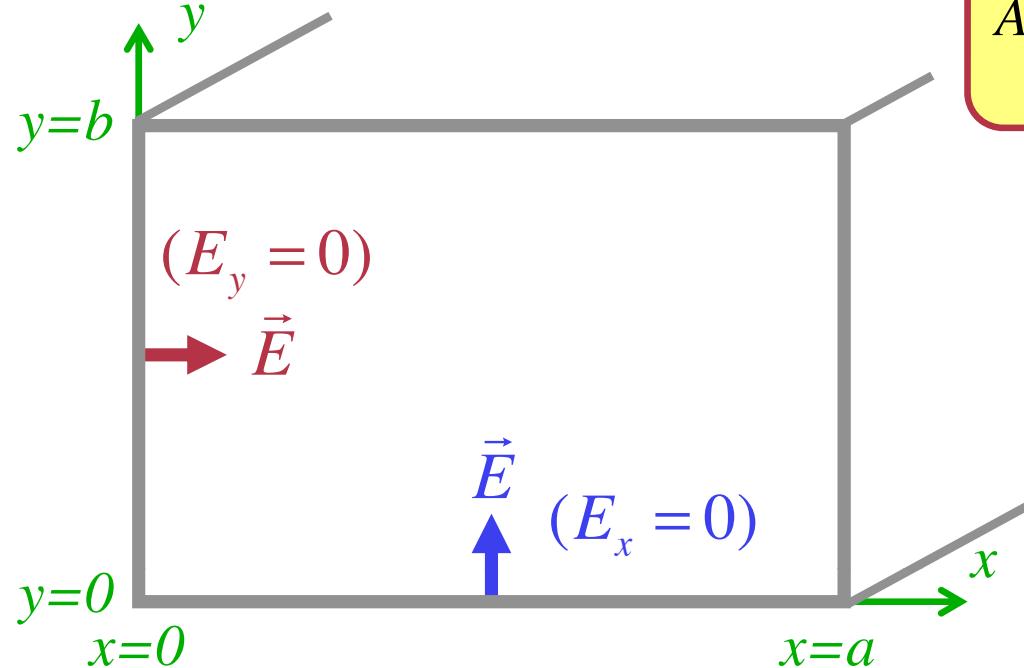
$$A_z(x, y, z) = \underbrace{\begin{cases} \sin(k_x x) \\ \cos(k_x x) \end{cases}}_{\text{superposition: } (A_2 \sin(k_x x) + A_1 \cos(k_x x)) \cdots} \cdot \underbrace{\begin{cases} \sin(k_y y) \\ \cos(k_y y) \end{cases}}_{\text{superposition: } (B_2 \sin(k_y y) + B_1 \cos(k_y y)) \cdots} \cdot \underbrace{\begin{cases} e^{+jk_z z} \\ e^{-jk_z z} \end{cases}}_{\text{superposition: } (C_2 e^{jk_z z} + C_1 e^{-jk_z z}) \cdots}$$

superposition:  $(A_2 \sin(k_x x) + A_1 \cos(k_x x)) \cdots$

☞ To determine unknown coefficients: boundary conditions for fields !

- **Imposing the boundary conditions**

- **TE-solution:**



Continuity (1):  $E_x(x \quad , y = 0) = 0$

$E_y(x = 0, y \quad ) = 0$

(Lecture Notes pp. 51-52)

$$A_z(x, y, z) = \begin{Bmatrix} \sin(k_x x) \\ \cos(k_x x) \end{Bmatrix} \cdot \begin{Bmatrix} \sin(k_y y) \\ \cos(k_y y) \end{Bmatrix} \cdot \begin{Bmatrix} e^{+jk_z z} \\ e^{-jk_z z} \end{Bmatrix}$$



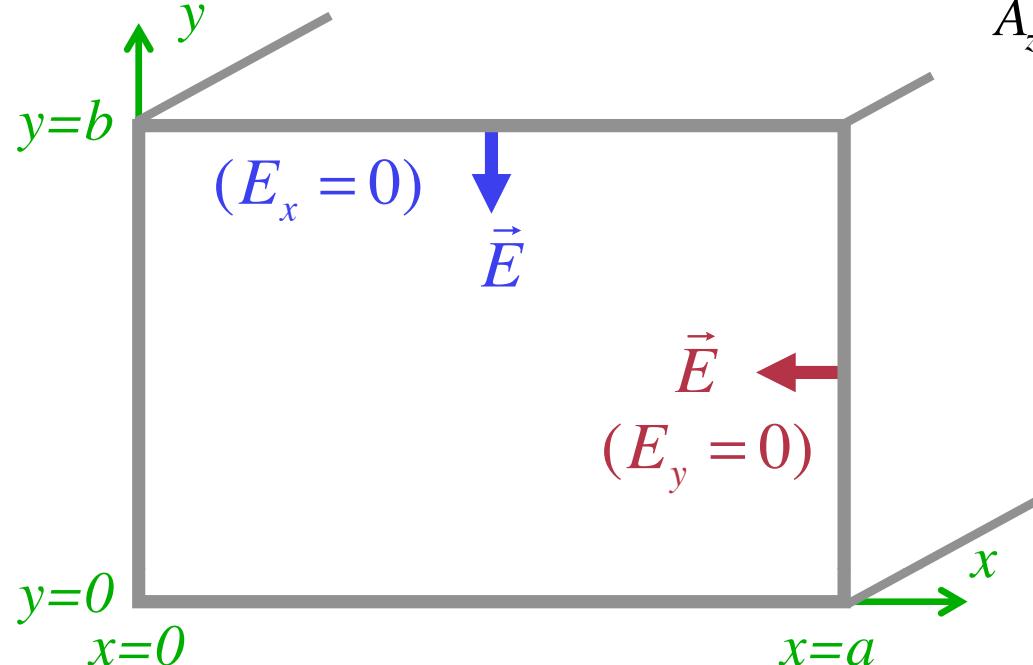
$$\vec{E} = \text{curl } \vec{A} = (\dots E_x \dots) \mathbf{e}_x + (\dots E_y \dots) \mathbf{e}_y + (\dots 0 \dots) \mathbf{e}_z$$

$$\downarrow \quad \vec{E}_x \sim \frac{\partial A}{\partial y}, \vec{E}_y \sim \frac{\partial A}{\partial x}$$

(....)

$$A_z(x, y, z) = \begin{Bmatrix} \cancel{\sin(k_x x)} \\ \cancel{\cos(k_x x)} \end{Bmatrix} \cdot \begin{Bmatrix} \cancel{\sin(k_y y)} \\ \cancel{\cos(k_y y)} \end{Bmatrix} \cdot \begin{Bmatrix} e^{+jk_z z} \\ e^{-jk_z z} \end{Bmatrix}$$

- **TE-solution (...):**



$$A_z(x, y, z) = \begin{Bmatrix} \sin(k_x x) \\ \cos(k_x x) \end{Bmatrix} \cdot \begin{Bmatrix} \sin(k_y y) \\ \cos(k_y y) \end{Bmatrix} \cdot \begin{Bmatrix} e^{+jk_z z} \\ e^{-jk_z z} \end{Bmatrix}$$



$$E_x(x, y) \propto \cos(k_x x) \cdot \sin(k_y y) k_y$$

$$E_y(x, y) \propto \sin(k_x x) \cdot \cos(k_y y) k_x$$

Continuity (2):  $E_x(x = a, y = b) = 0$

$$E_y(x = a, y = b) = 0$$

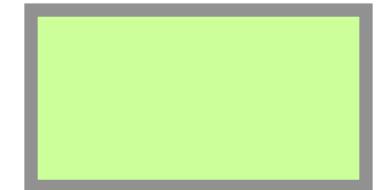
→ (...)  $k_y = \frac{n \pi}{b}$   
 $k_x = \frac{m \pi}{a}$

(transversal separation constants)

(Lecture Notes p. 52)

- Calculation of the field components – TE modes

Rectangular WG



$$\vec{E} = \operatorname{curl} \vec{A}$$

$$\vec{H} = -\frac{1}{j\omega\mu} \operatorname{curl} \vec{E}$$

## ● Field components

### ☞ TE modes

$$E_x = -C \frac{n\pi}{b} \cos(m\pi x/a) \sin(n\pi y/b) e^{\pm jk_z z}$$

$$E_y = +C \frac{m\pi}{a} \sin(m\pi x/a) \cos(n\pi y/b) e^{\pm jk_z z}$$

$$E_z = 0$$

### Rectangular WG



$$H_x = \pm C \frac{k_z}{\omega\mu} \frac{m\pi}{a} \sin(m\pi x/a) \cos(n\pi y/b) e^{\pm jk_z z}$$

$$H_y = \pm C \frac{k_z}{\omega\mu} \frac{n\pi}{b} \cos(m\pi x/a) \sin(n\pi y/b) e^{\pm jk_z z}$$

$$H_z = -C \frac{k^2 - k_z^2}{j\omega\mu} \cos(m\pi x/a) \cos(n\pi y/b) e^{\pm jk_z z}$$

### ☞ TM modes - similar

$$E_x = \pm C \frac{k_z}{\omega\epsilon} \frac{m\pi}{a} \cos(m\pi x/a) \sin(n\pi y/b) e^{\pm jk_z z}$$

$$E_y = \pm C \frac{k_z}{\omega\epsilon} \frac{n\pi}{b} \sin(m\pi x/a) \cos(n\pi y/b) e^{\pm jk_z z}$$

$$E_z = +C \frac{k^2 - k_z^2}{j\omega\epsilon} \sin(m\pi x/a) \sin(n\pi y/b) e^{\pm jk_z z}$$

$$H_x = +C \frac{n\pi}{b} \sin(m\pi x/a) \cos(n\pi y/b) e^{\pm jk_z z}$$

$$H_y = -C \frac{m\pi}{a} \cos(m\pi x/a) \sin(n\pi y/b) e^{\pm jk_z z}$$

$$H_z = 0$$

- Infinite set of solutions (,Waveguide Modes':)**

$$A_z(x, y, z) = A_0 \cos(k_x x) \cdot \cos(k_y y) \cdot \begin{cases} e^{+jk_z z} \\ e^{-jk_z z} \end{cases}$$

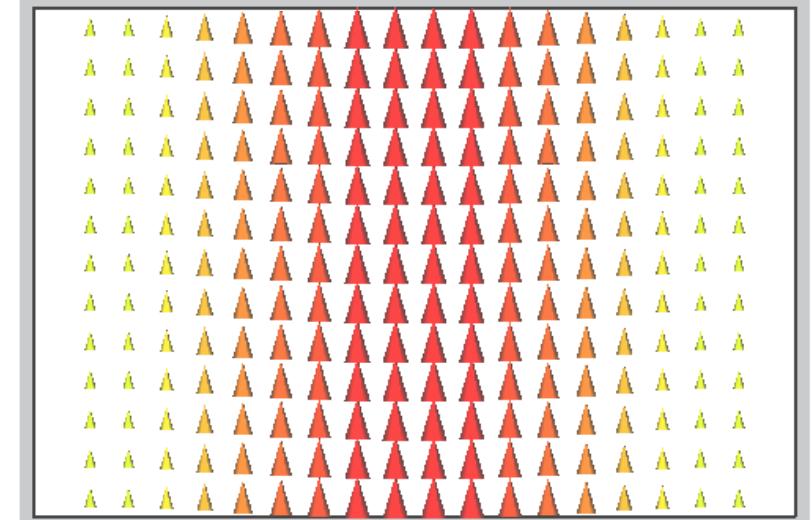
$$E_x(x, y) \propto \cos(k_x x) \cdot \sin(k_y y)$$

$$E_y(x, y) \propto \sin(k_x x) \cdot \cos(k_y y)$$

$$k_x = \frac{m \pi}{a} \quad m = 0, 1, 2, \dots$$

$$k_y = \frac{n \pi}{b} \quad n = 0, 1, 2, \dots$$

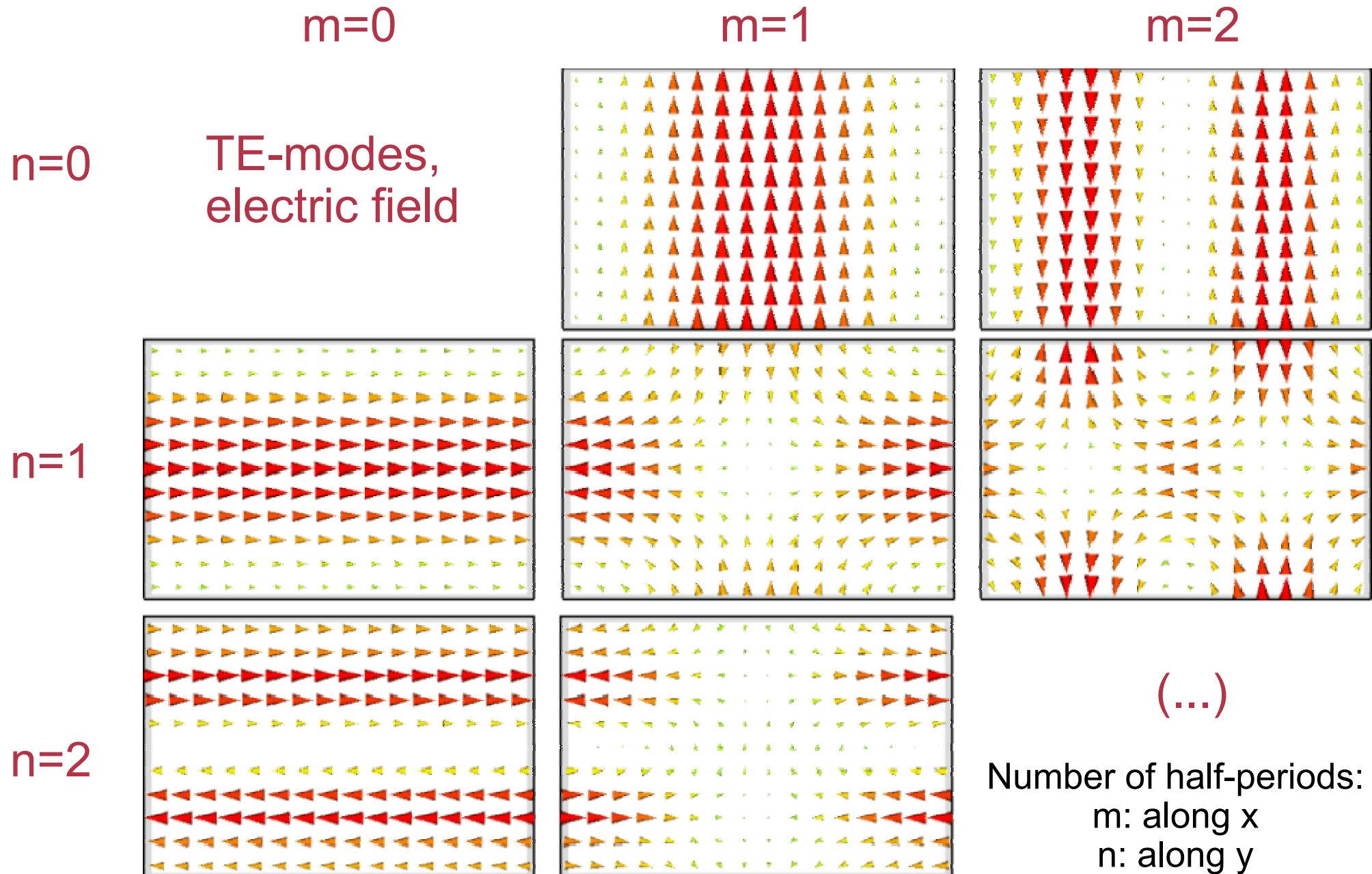
$$\begin{aligned} k_z^2 &= k^2 - (k_x^2 + k_y^2) \\ &= \omega^2 \mu \epsilon - (k_x^2 + k_y^2) \end{aligned}$$



$$(m, n) = \{ \quad (1,0), (2,0), (3,0), \dots, \\ (0,1), (1,1), (2,1), (3,1), \dots, \\ (0,2), (1,2), (2,2), (3,2), \dots, \\ \dots \}$$

*(m = n = 0 : all zero)*

# Rectangular Waveguides



## ● Waveguide modes

- ☞ Solution of the wave equation in a hollow guide define field patterns for the electric (and the magnetic) field
- ☞ This pattern propagates in longitudinal direction
- ☞ Total field may be a superposition of:
  - different modes
  - modes propagating in +/- z-direction
- ☞ Two types of modes:
  - TE (‘transversal electric’) and TM (‘transversal magnetic’)
  - sometimes also referred to as  $H_z$  and  $E_z$ -modes
- ☞ In rectangular guides:
  - modes are determined by lengths a and b

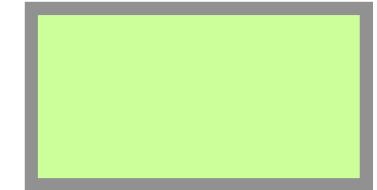
# Waveguides (dispersion relation)

- Dispersion relation:

Rectangular WG

$$\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2 + k_z^2 = \omega^2 \mu \epsilon$$

↑      ↑      ↑  
 periodicity      periodicity      wave propagation in waveguide  
 in  $x$             in  $y$             depends on mode type !



$$f_c = \frac{c}{2} \sqrt{\left(\frac{m}{a}\right)^2 + \left(\frac{n}{b}\right)^2}$$

$f_c$  = cutoff-frequency

(hint: see degenerated modes)

$$\begin{aligned}
 k_z &= \sqrt{\omega^2 \mu \epsilon - \left( \frac{(m\pi)^2}{a^2} + \frac{(n\pi)^2}{b^2} \right)} \\
 &= \frac{2\pi}{c} \sqrt{f^2 - f_c^2}
 \end{aligned}$$

$$k_z = \begin{cases} \text{real (wave propagation)} & f > f_c \\ \text{imag. (attenuation)} & f < f_c \end{cases}$$

# Two very important formulas

- Cutoff frequency

$$f_c = \frac{c}{2} \sqrt{\left(\frac{m}{a}\right)^2 + \left(\frac{n}{b}\right)^2}$$

- Propagation constant in the waveguide ( $\neq \omega\sqrt{\mu\varepsilon}$  !)

$$\begin{aligned} k_z &= \sqrt{\omega^2 \mu \varepsilon - \left( \frac{(m\pi)^2}{a^2} + \frac{(n\pi)^2}{b^2} \right)} = \sqrt{\left( \frac{2\pi f}{c} \right)^2 - \left( \frac{(m\pi)^2}{a^2} + \frac{(n\pi)^2}{b^2} \right)} = \\ &= \frac{2\pi}{c} \sqrt{f^2 - f_c^2} \quad \text{at cutoff } k_z = 0 \end{aligned}$$

- Fundamental mode: the mode with the smallest cutoff frequency – TE10

- Dispersion diagram:

  - ☞ Plot  $k_z^2$  as a function of  $f^2$

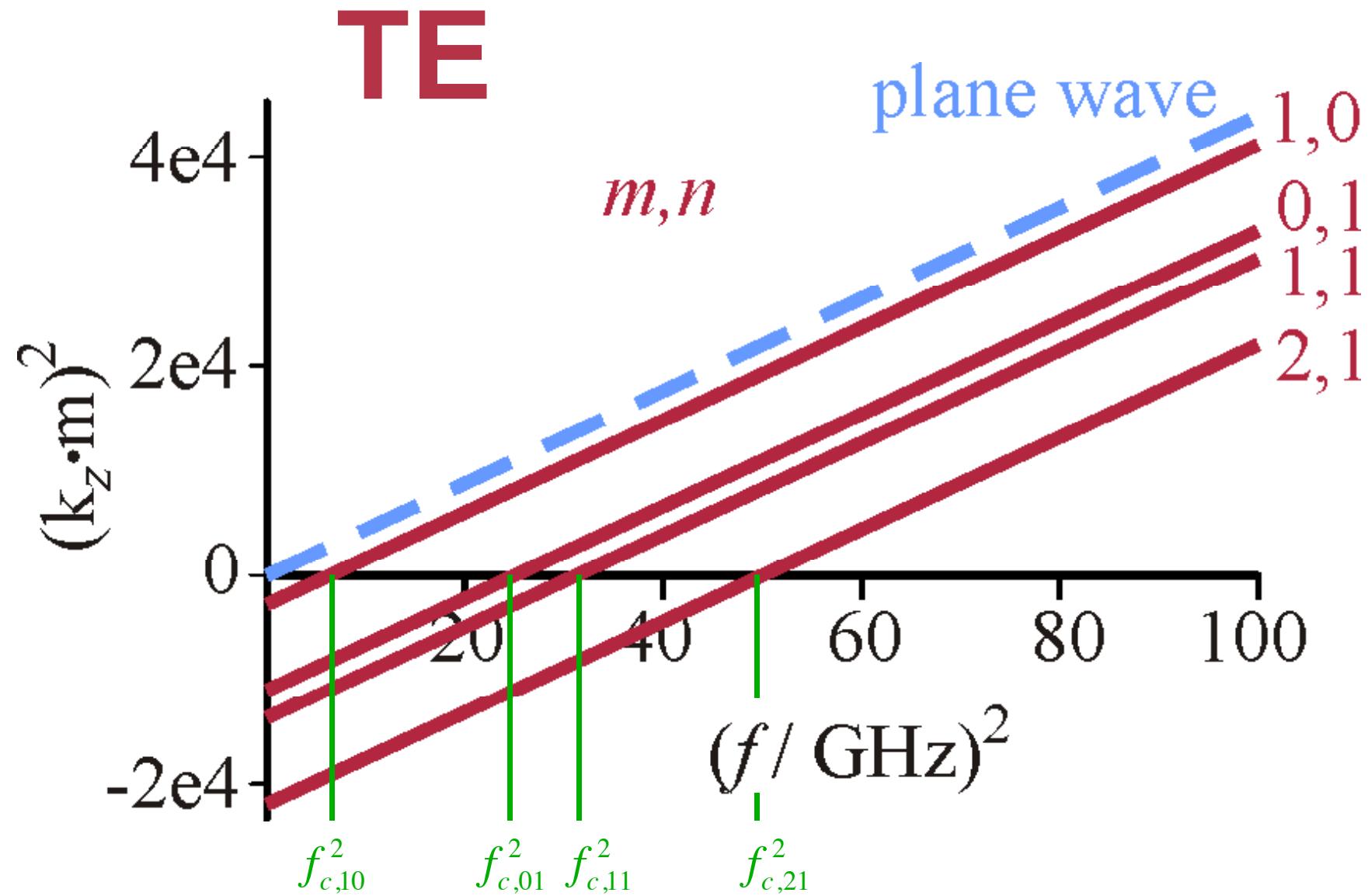
- Formula:

$$k_z^2(f^2) = (2\pi/c)^2 - (f^2 - f_c^2)$$

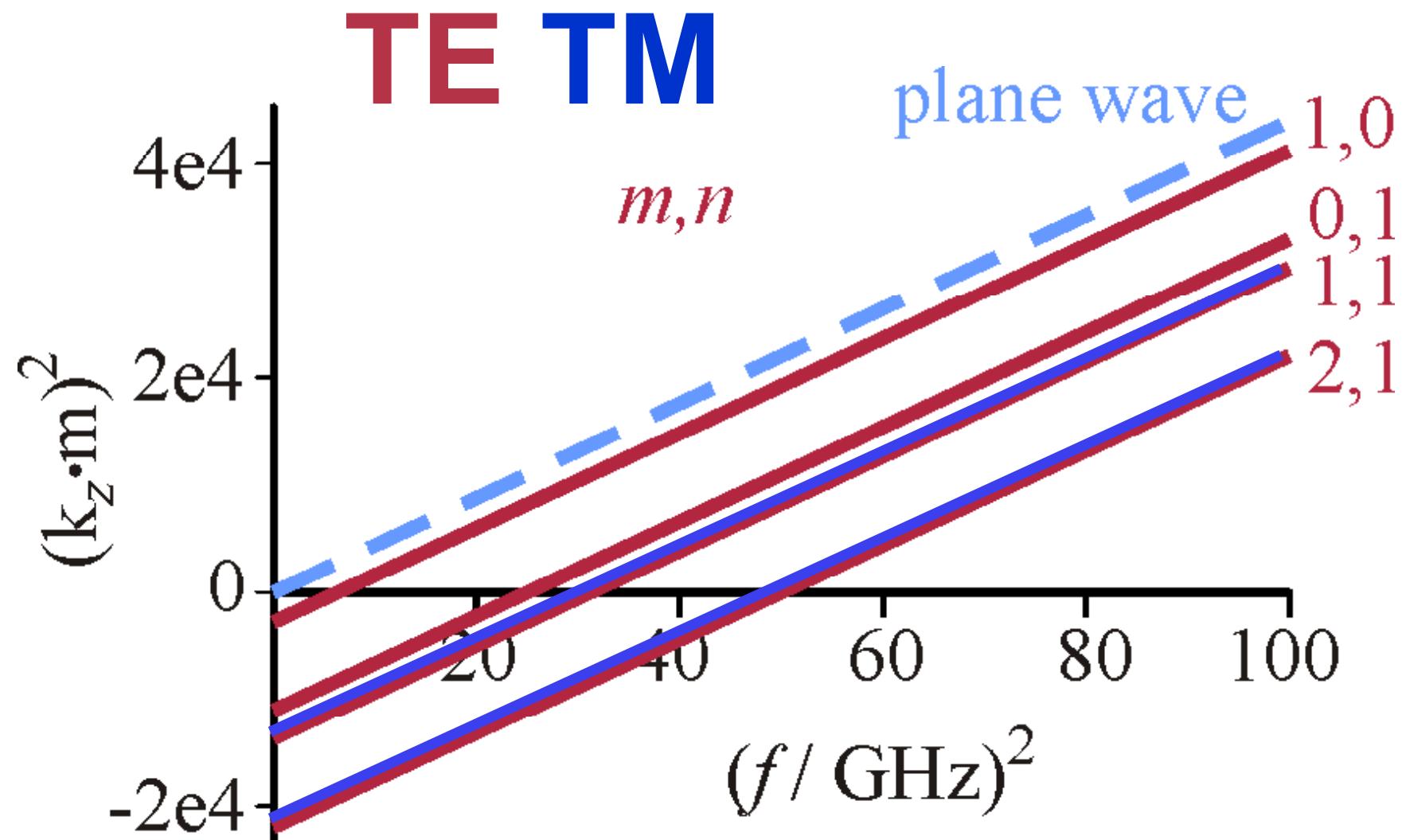
- Alternative:

  - ☞ Plot  $|k_z|$  as a function of  $f$

- Dispersion diagram:  $k_z^2(f^2) = (2\pi/c)^2 - (f^2 - f_c^2)$



- Dispersion diagram:  $k_z^2(f^2)$



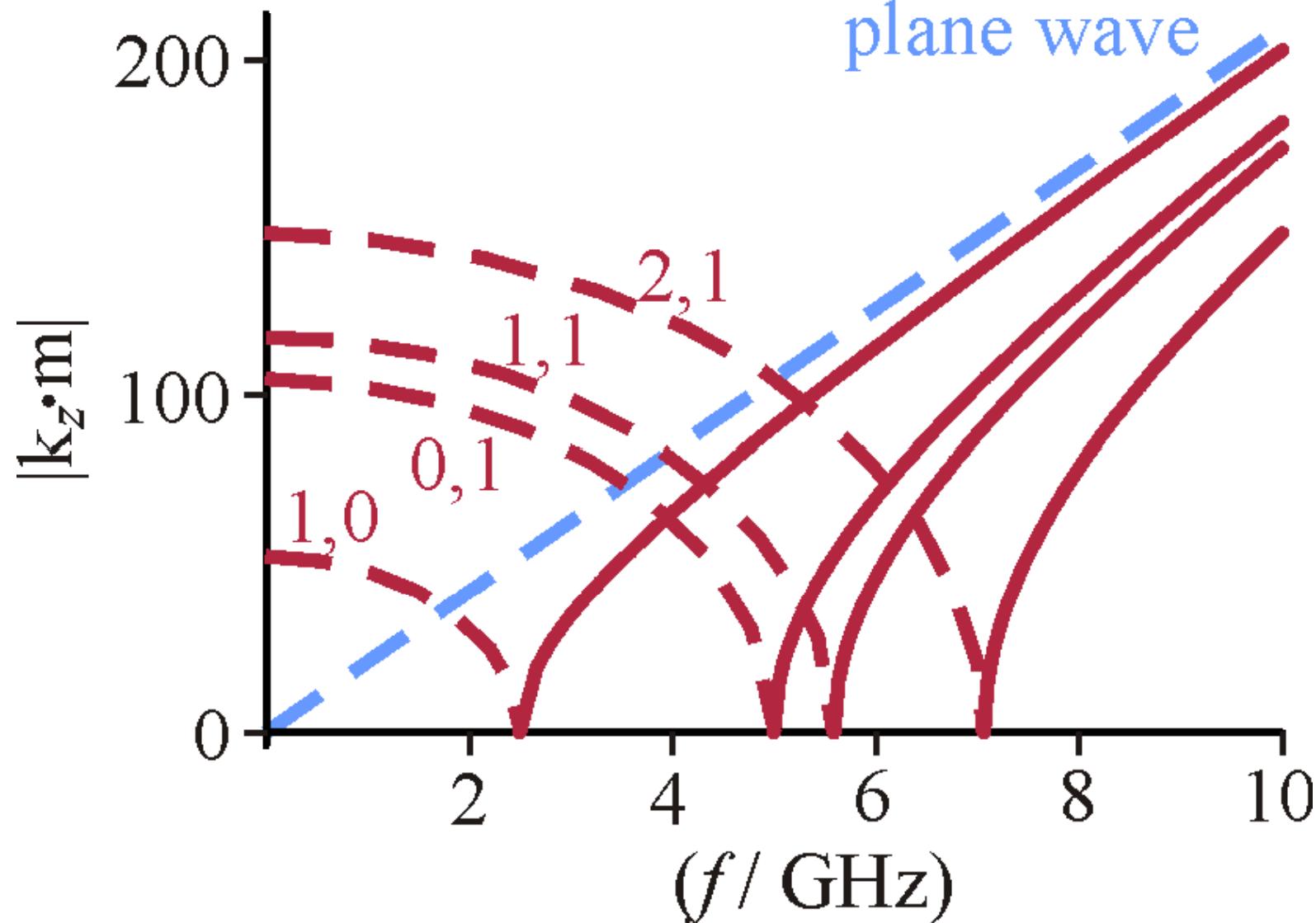
- Some modes can have the same cutoff frequency, although they have completely different mode patterns
- These are called *degenerated modes*
- Above cutoff  $k_z^2(f^2)$  is positive → *propagating modes*

$$E_x = f(x, y) e^{\pm j k_z z}$$

- Below cutoff,  $k_z^2(f^2)$  is negative → *evanescent modes* (cannot propagate, field attenuation along the waveguide)

$$E_x = f(x, y) e^{-k_z z}$$

- Dispersion diagram:  $|k_z(f)|$

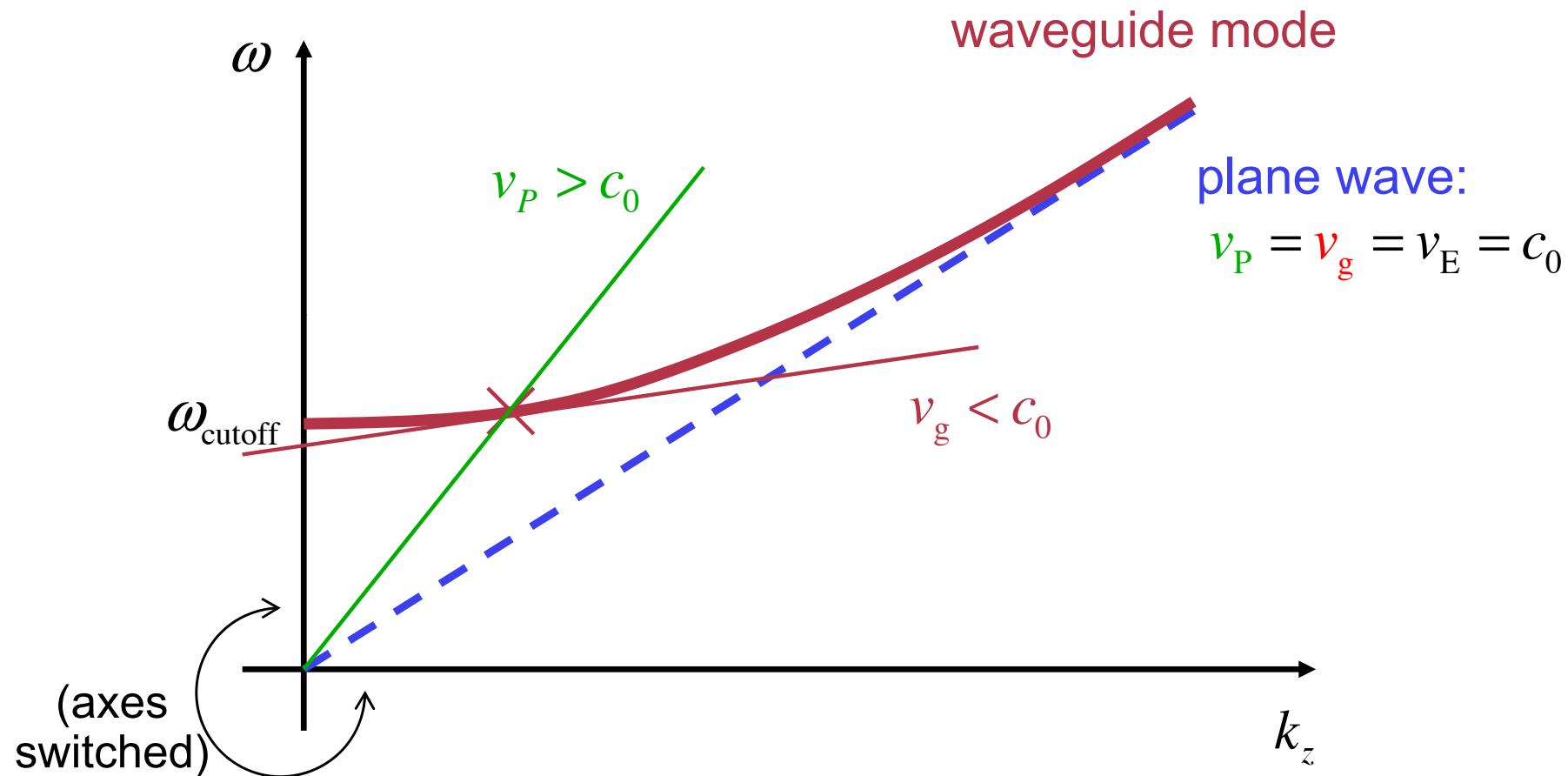


- Phase and group velocity in waveguides:

$$v_P = \frac{\omega}{k_z}$$

$$v_g = \frac{d\omega(k_z)}{dk_z}$$

$$\vec{v}_E = \frac{\text{Re}\{\underline{S}(\vec{r})\}}{\bar{w}(\vec{r})}$$



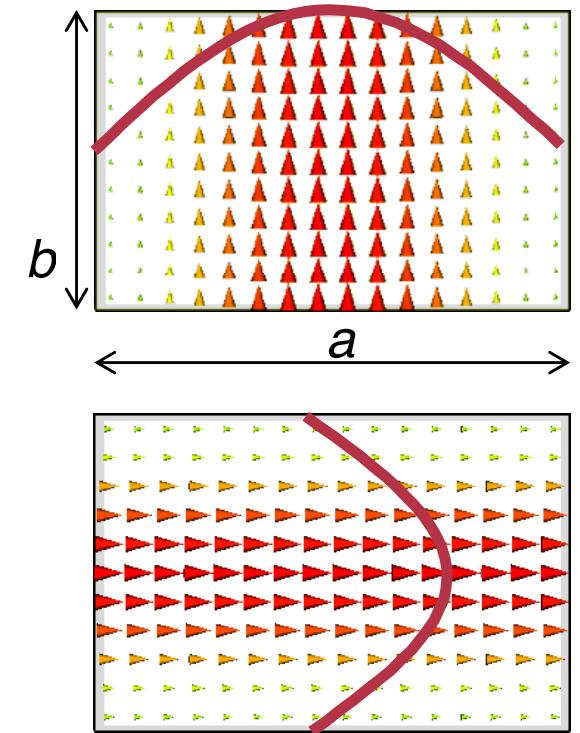
- **Size of Rectangular Waveguides:**  
smallest cutoff-frequencies (cutoff-wavelength):

$$f_{c,10} = \frac{c}{2} \sqrt{\left(\frac{1}{a}\right)^2 + \left(\frac{0}{b}\right)^2} = \frac{c}{2a}$$

$$\lambda_{c,10} = \frac{c}{f_{c,10}} = 2a$$

$$f_{c,01} = \frac{c}{2} \sqrt{\left(\frac{0}{a}\right)^2 + \left(\frac{1}{b}\right)^2} = \frac{c}{2b}$$

$$\lambda_{c,01} = \frac{c}{f_{c,01}} = 2b$$

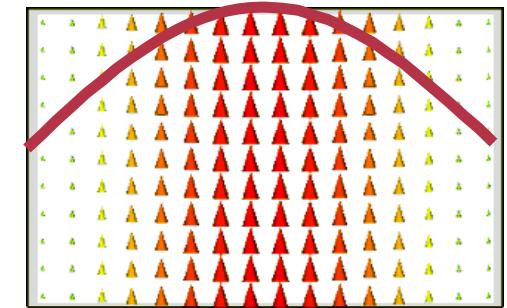


mode with  
smallest cutoff:  $\frac{\lambda_c}{2} = \max(a,b)$

$$f_c = \frac{c}{\lambda_c} = \frac{c}{2 \max(a,b)}$$

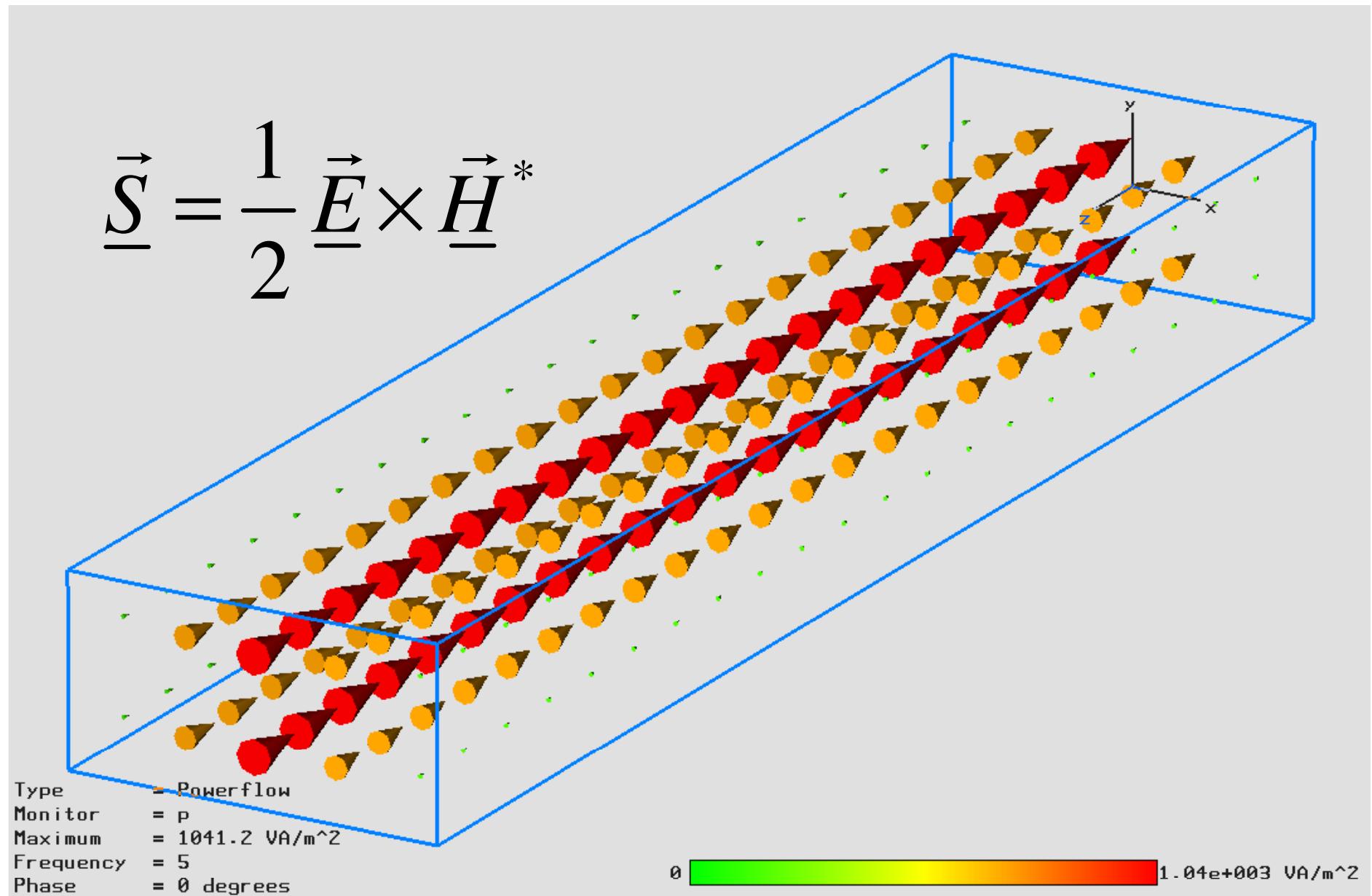
$$\frac{\lambda_c}{2} = \max(a, b)$$

$$f_c = \frac{c}{\lambda_c} = \frac{c}{2 \max(a, b)}$$

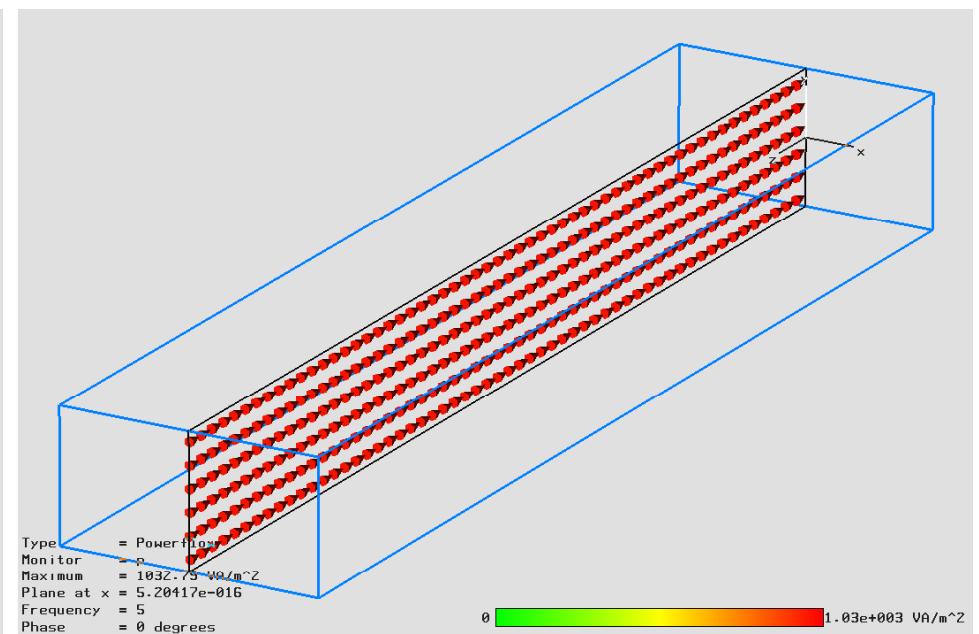
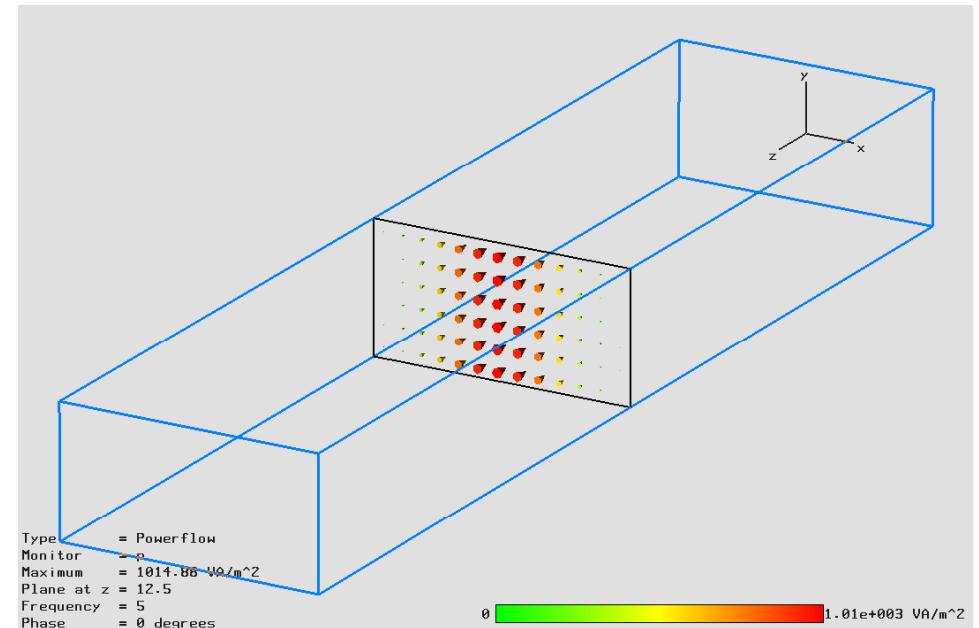
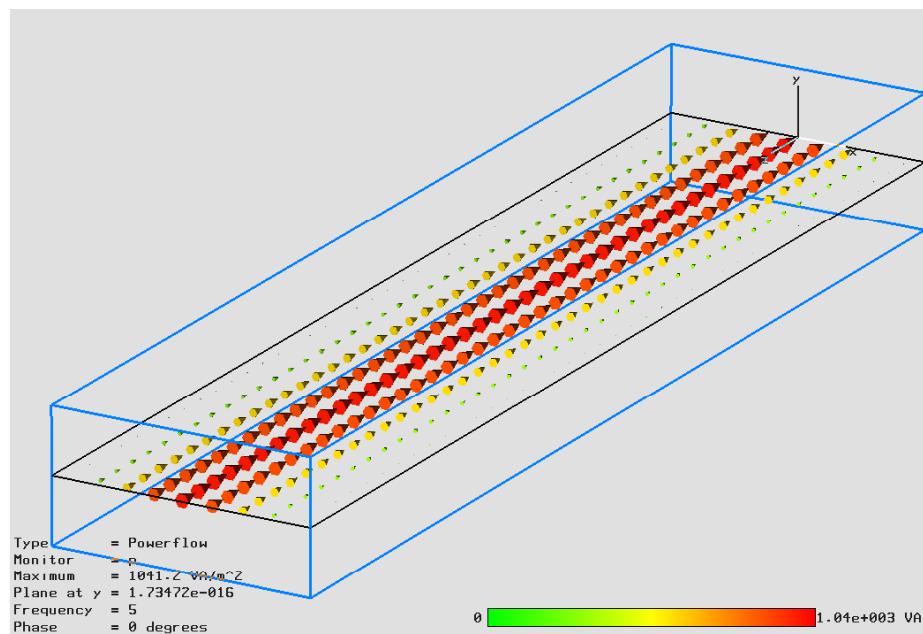


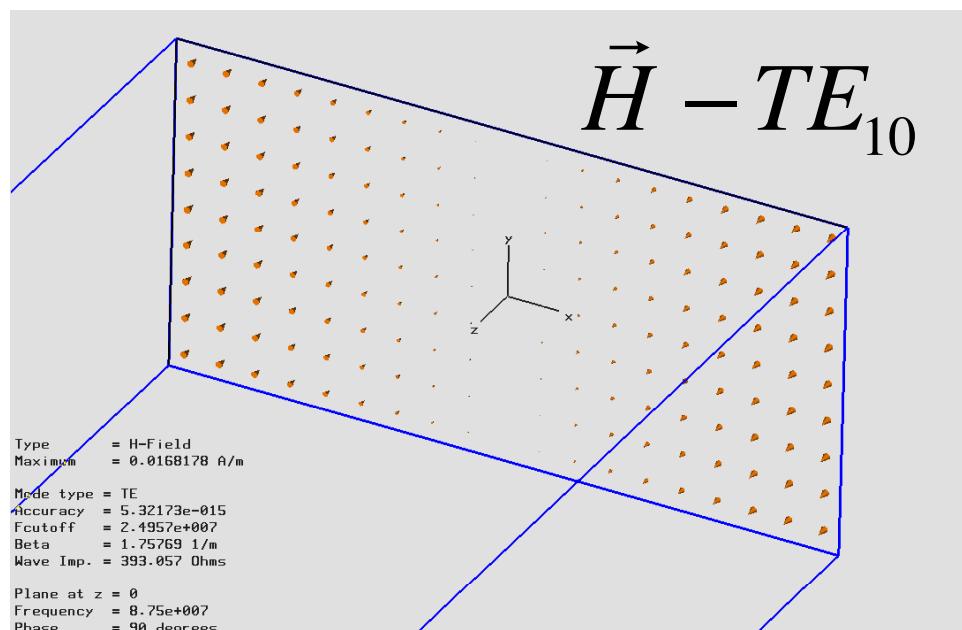
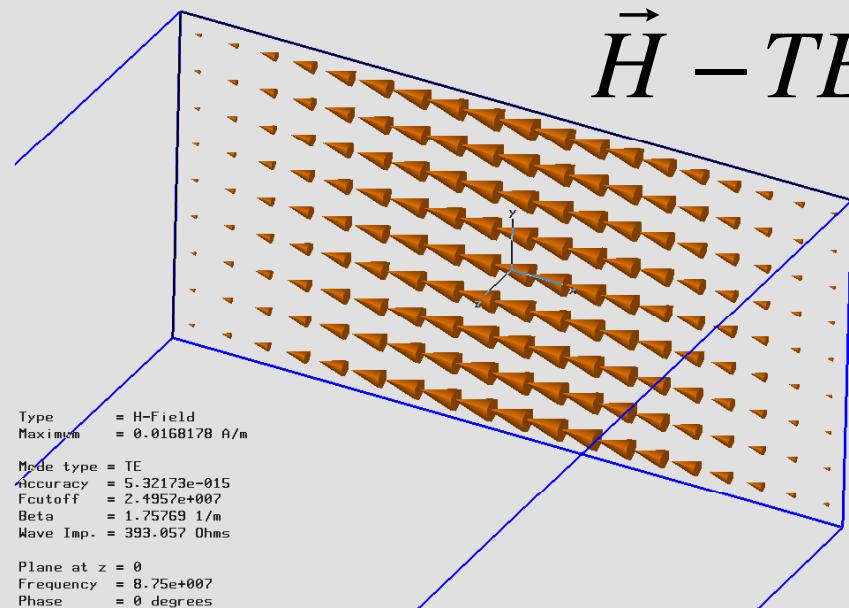
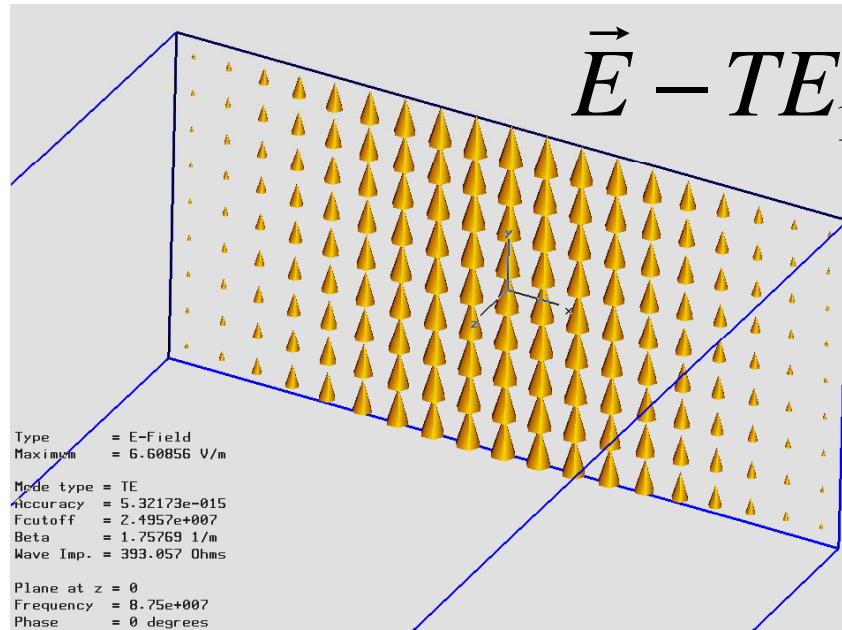
- Examples:

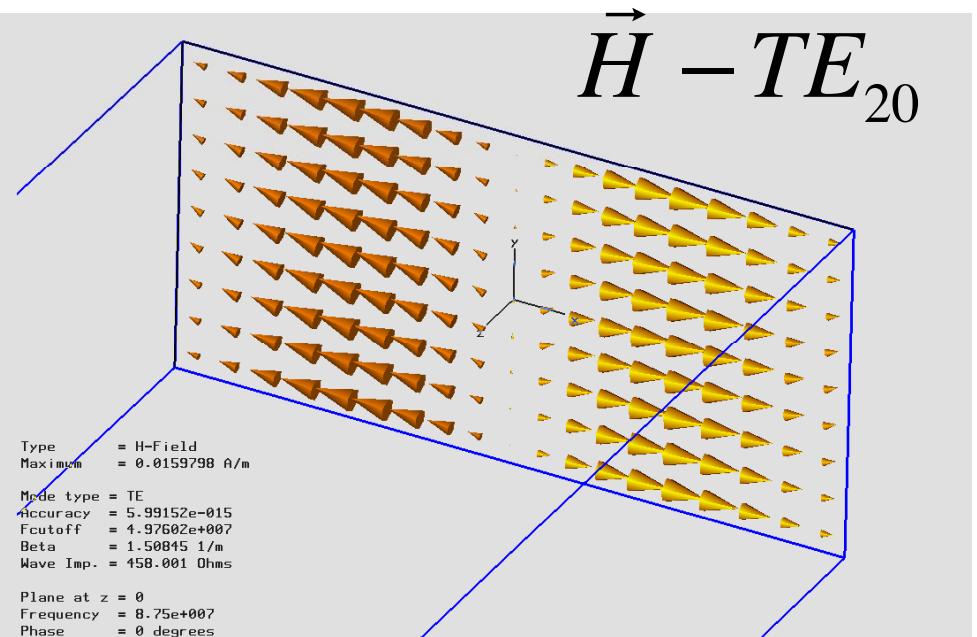
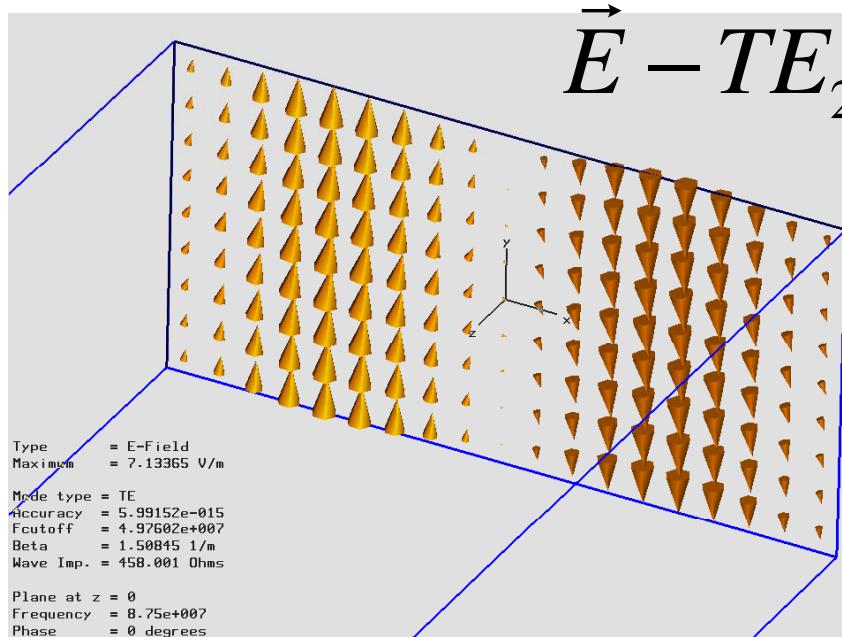
$f_{\text{cutoff}}$	$a$	
1 MHz	150 m	
10 MHz	15 m	
100 MHz	1.5 m	
0.91 GHz	165.1 mm x 82.55 mm („WR 650“)	too large for most applications !
1 GHz	15 cm	
10 GHz	1.5 cm	



$$\underline{\vec{S}} = \frac{1}{2} \underline{\vec{E}} \times \underline{\vec{H}}^*$$

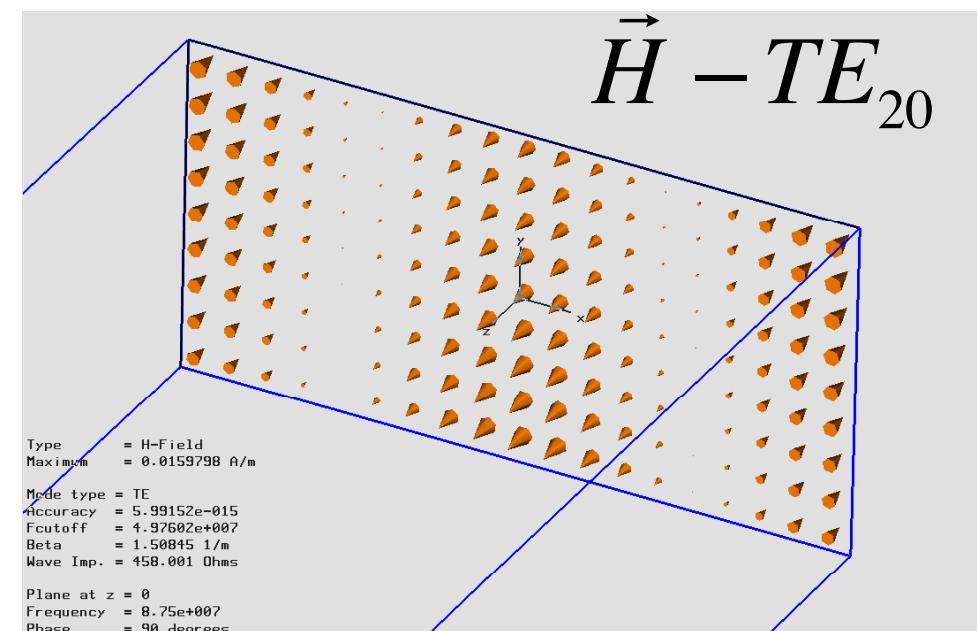


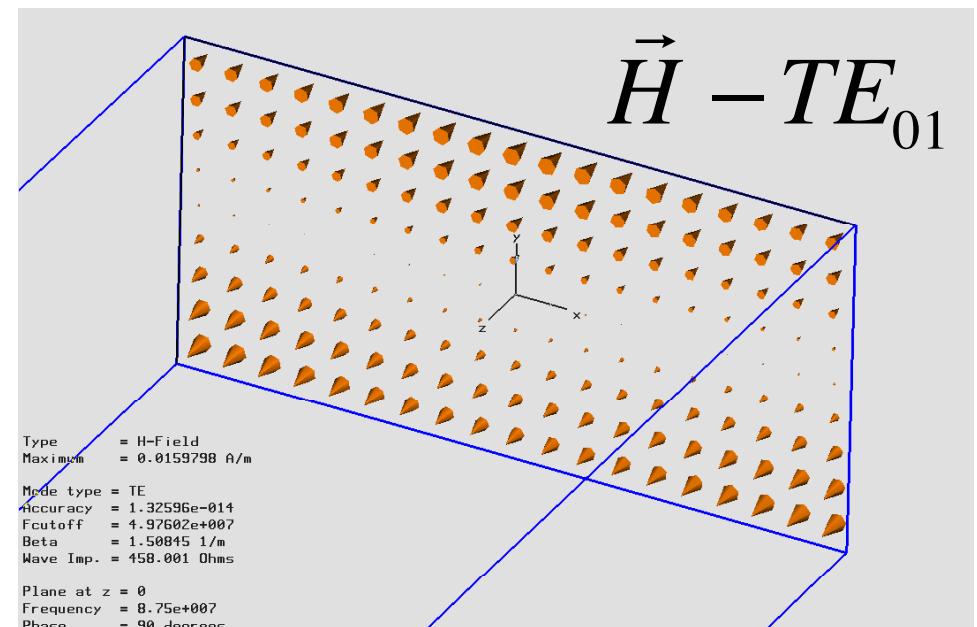
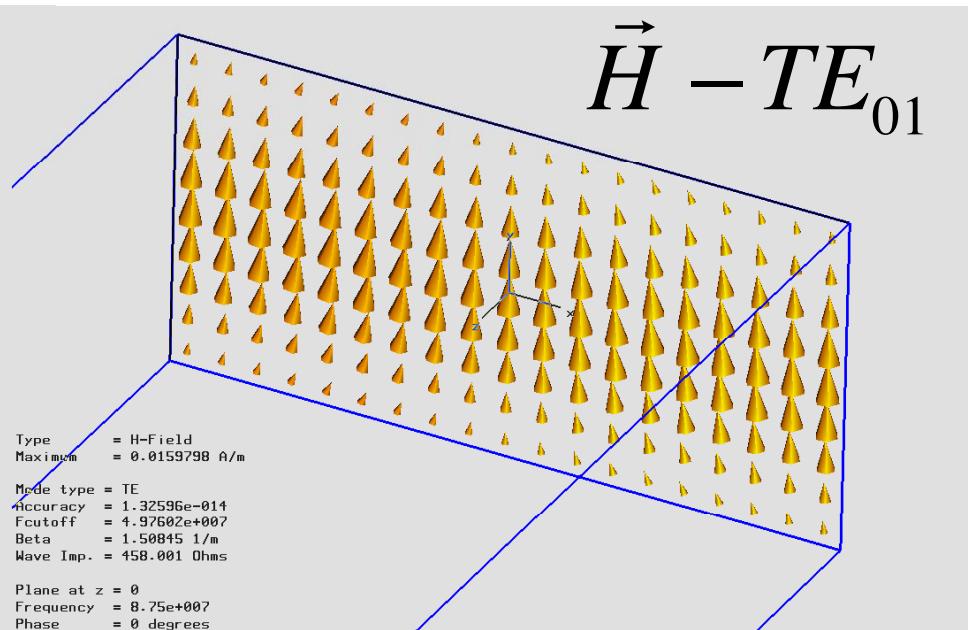
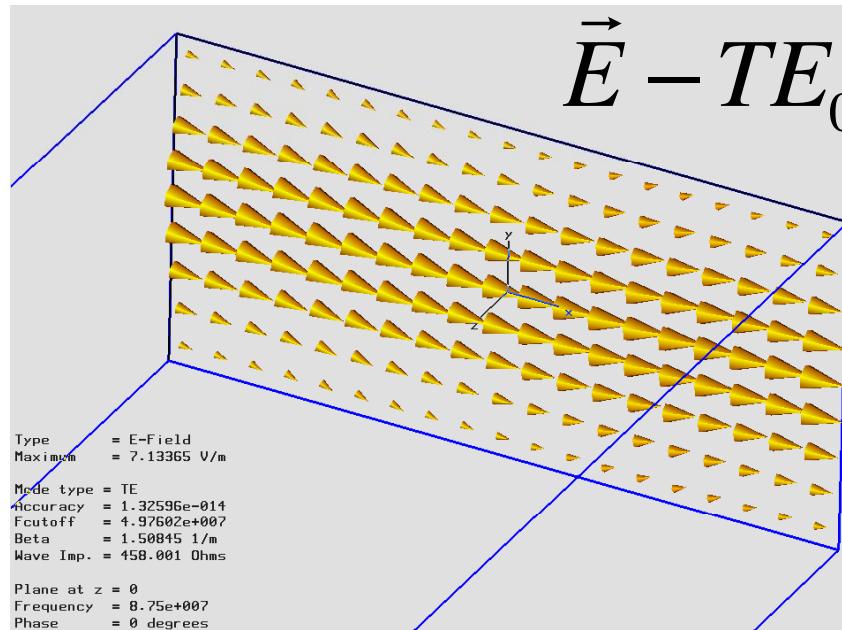


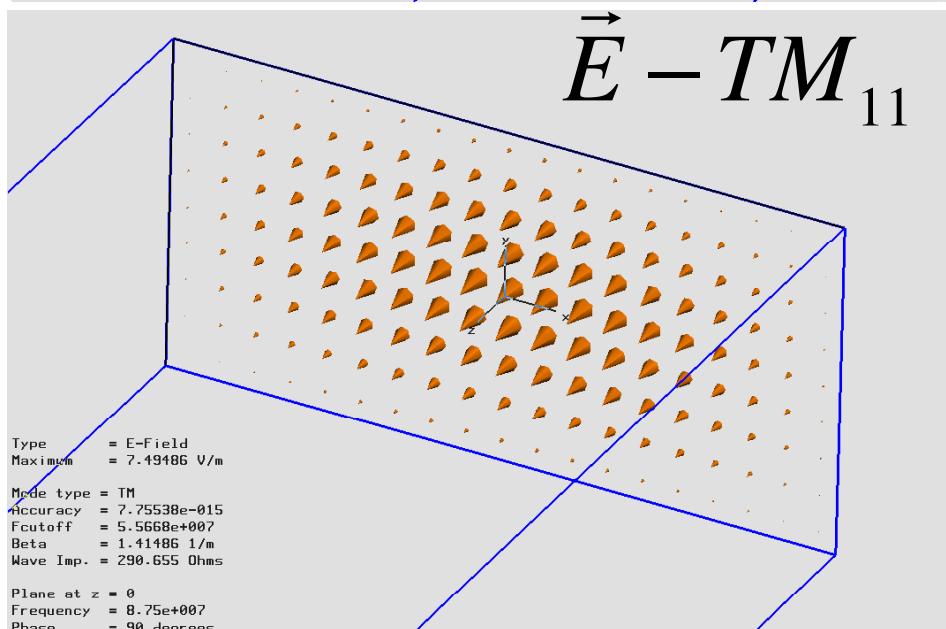
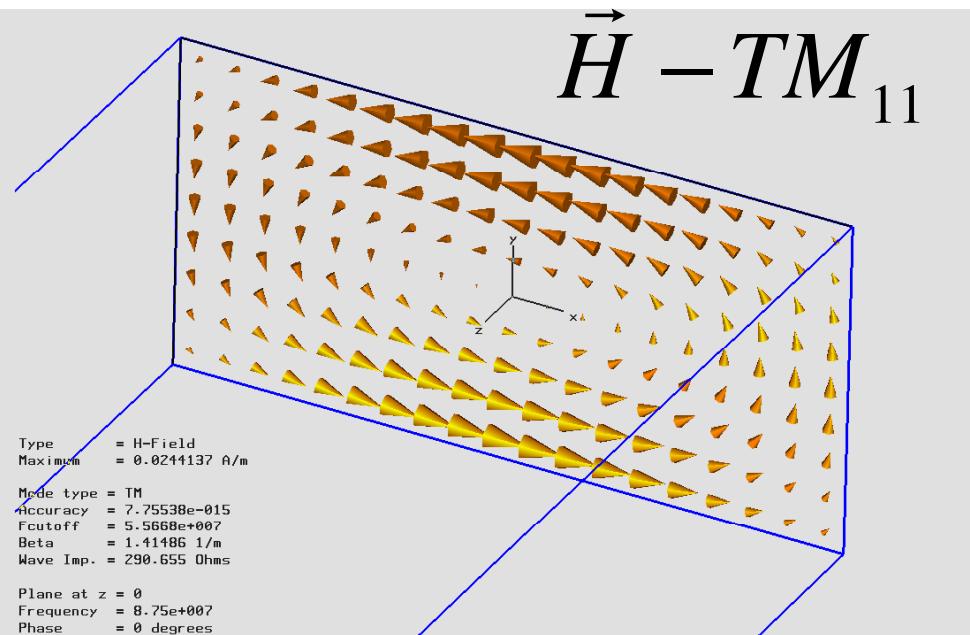
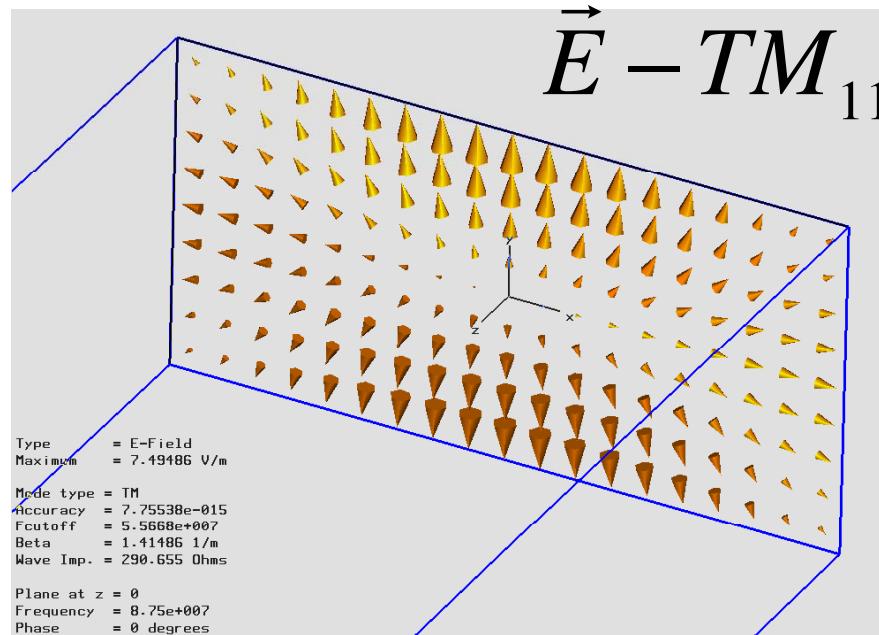


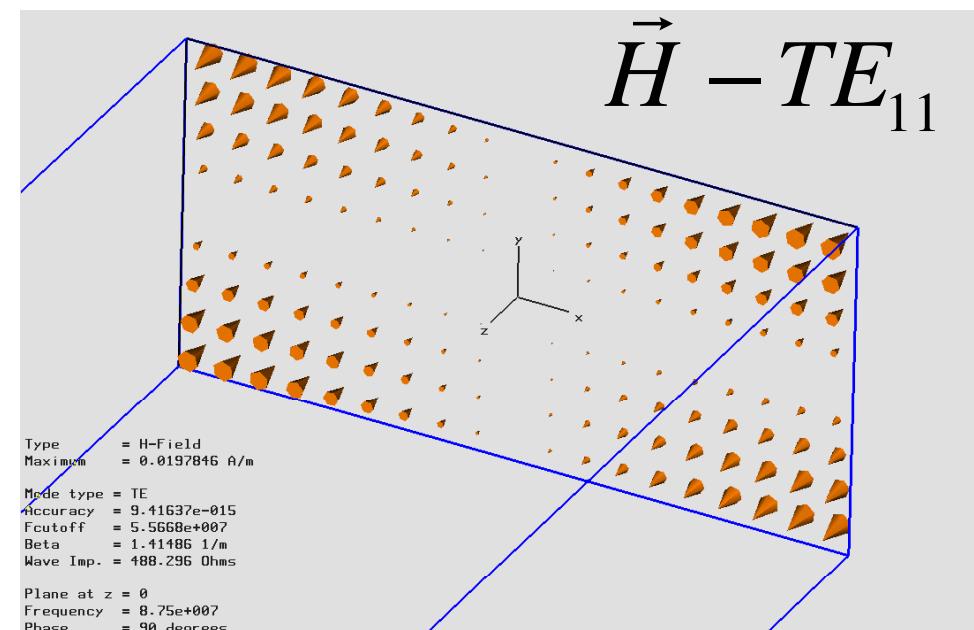
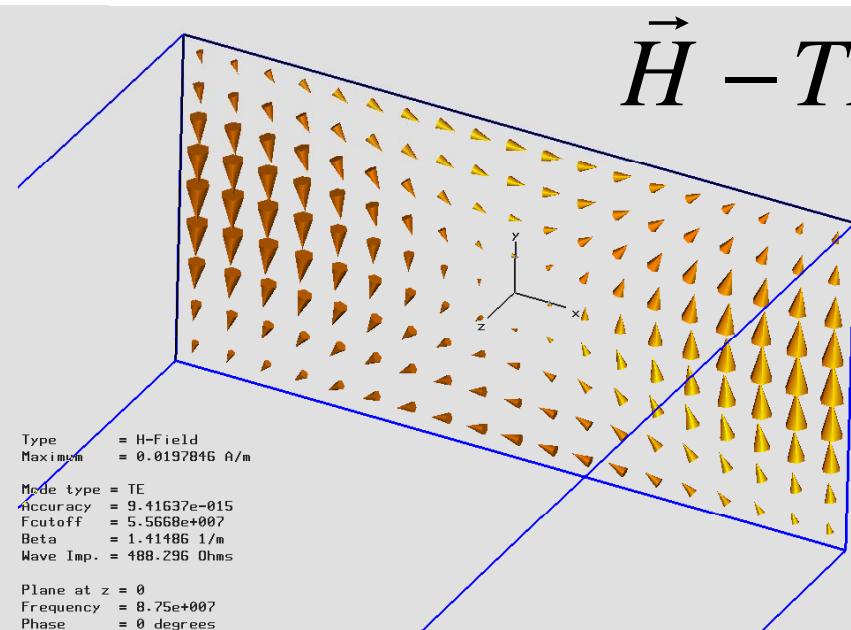
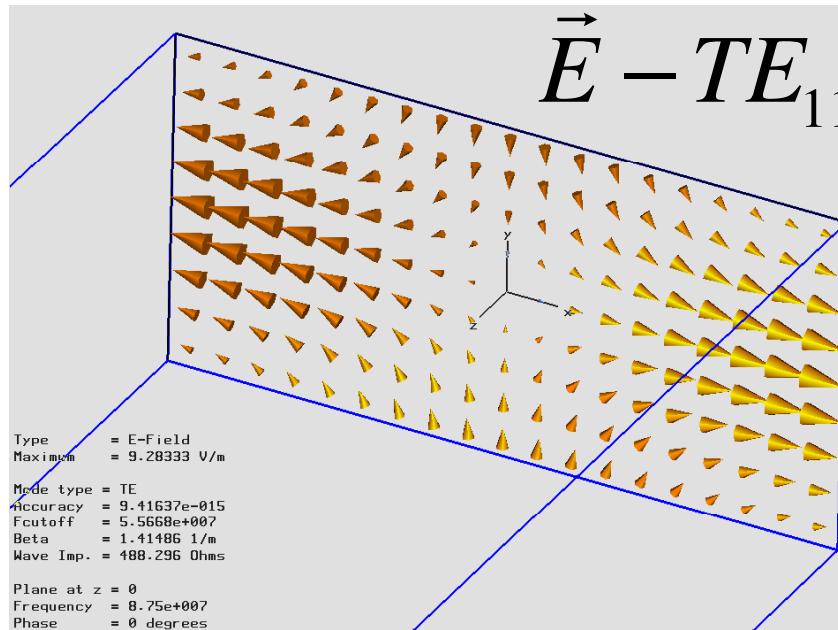
$m=2, n=0$

Number of half-periods:  
 m: along x  
 n: along y

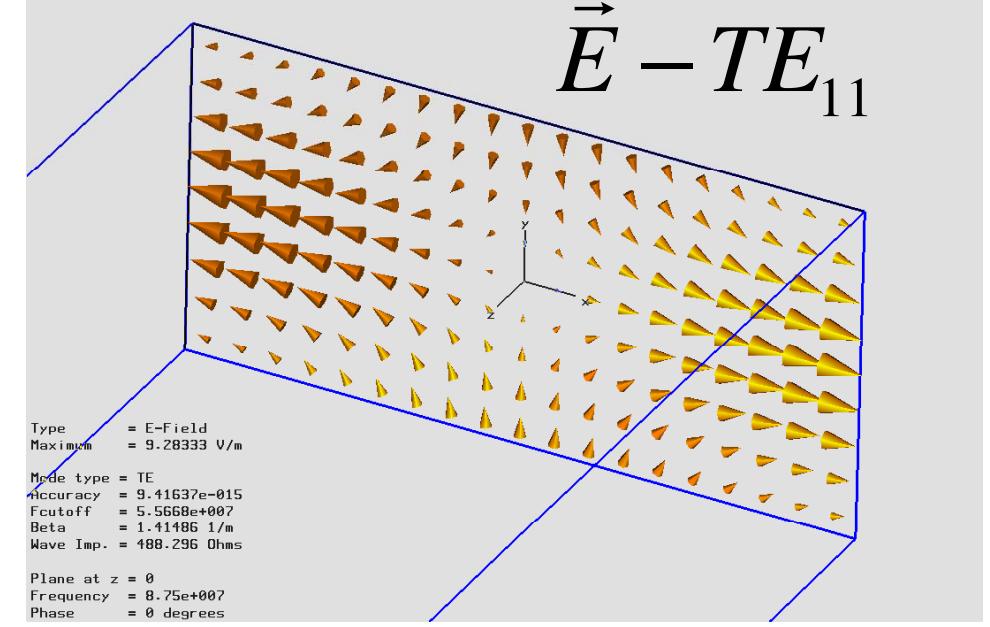
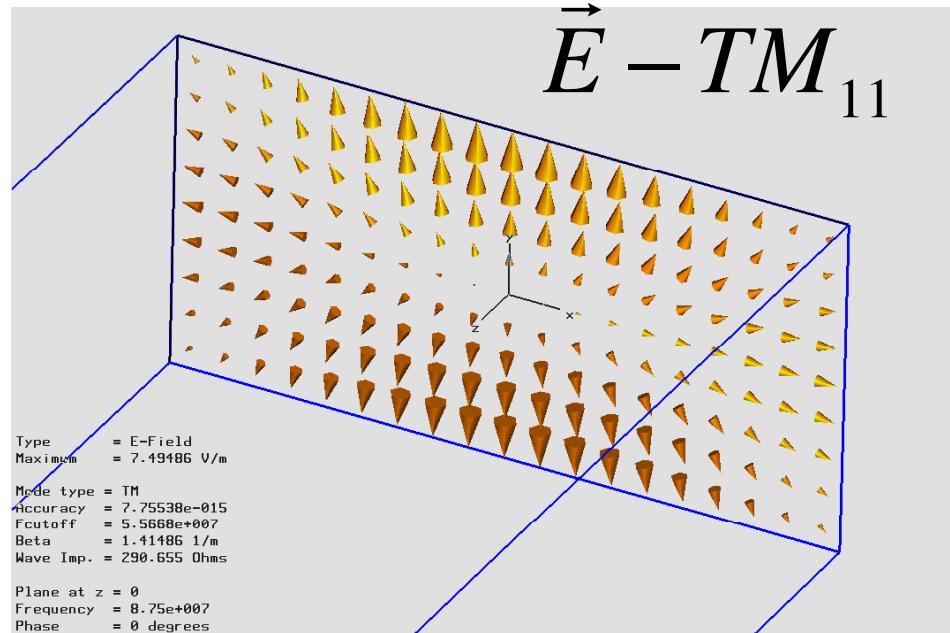








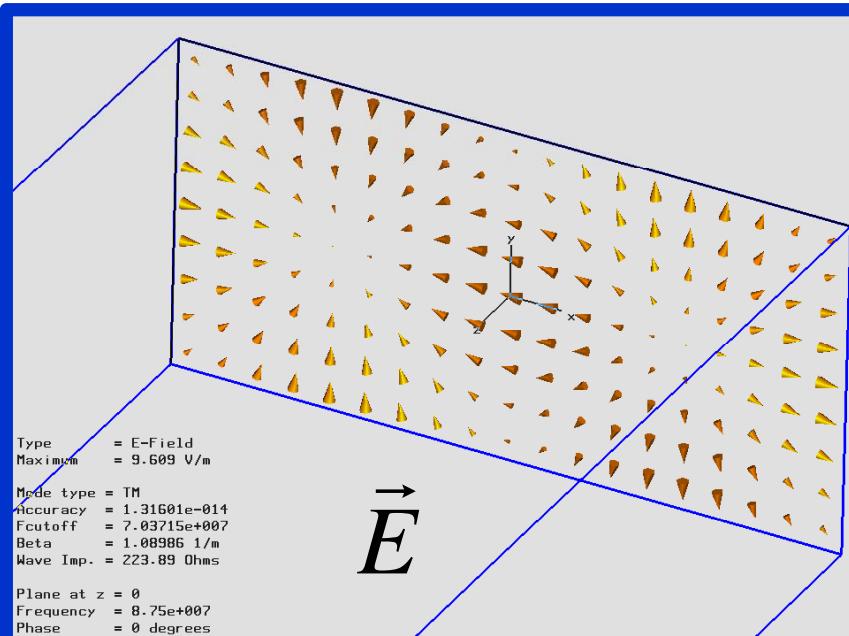
# Degenerated modes



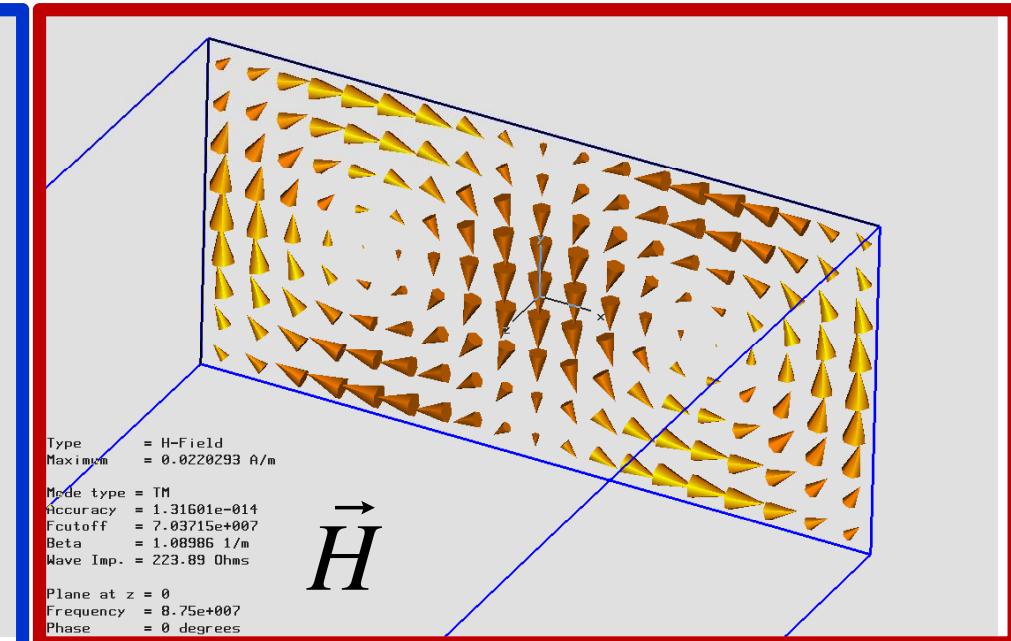
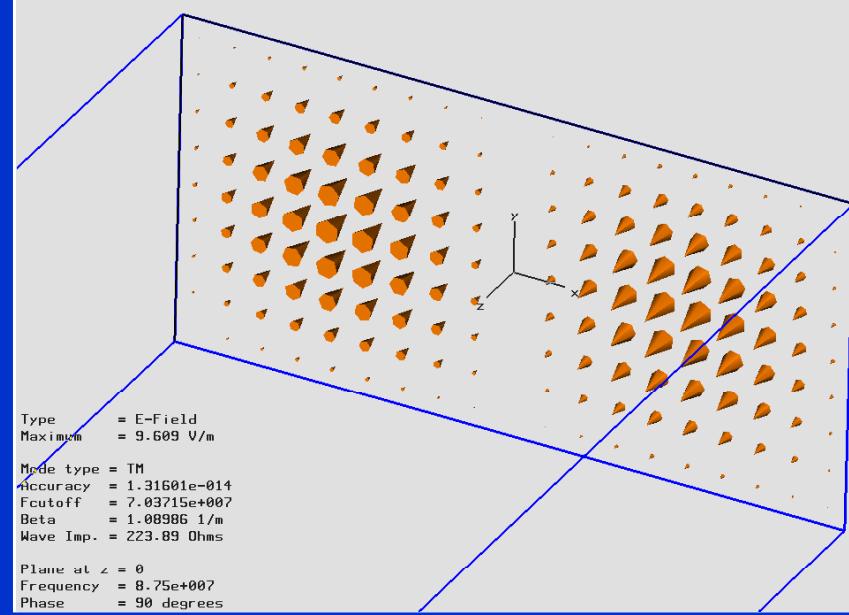
$$f_c = \frac{c}{2} \sqrt{\left(\frac{m}{a}\right)^2 + \left(\frac{n}{b}\right)^2}$$

- For this particular waveguide, these modes have the same cutoff frequency – “degenerated modes”

xy

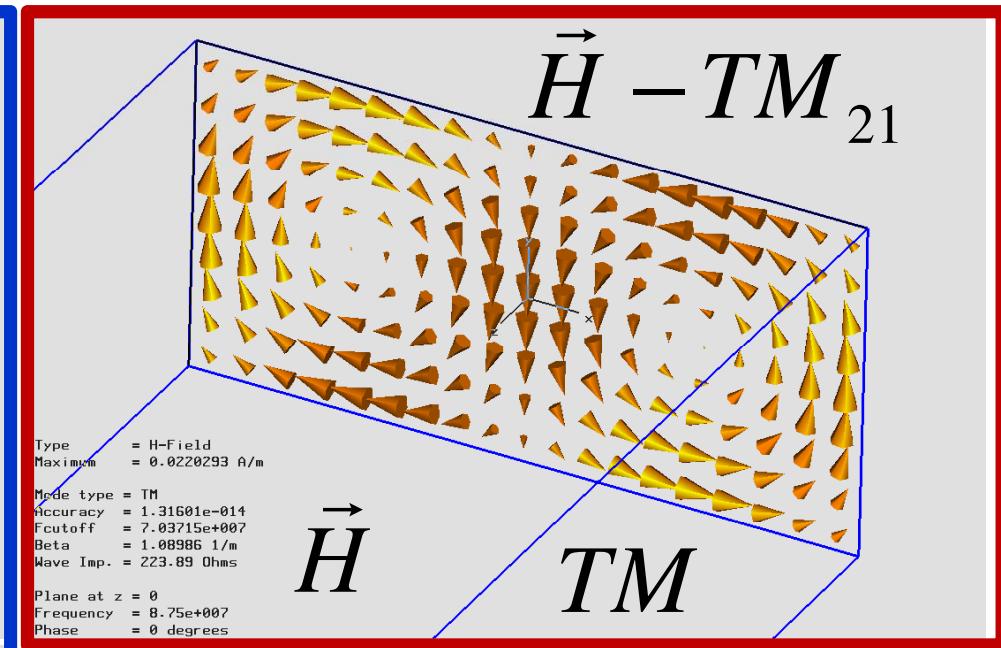
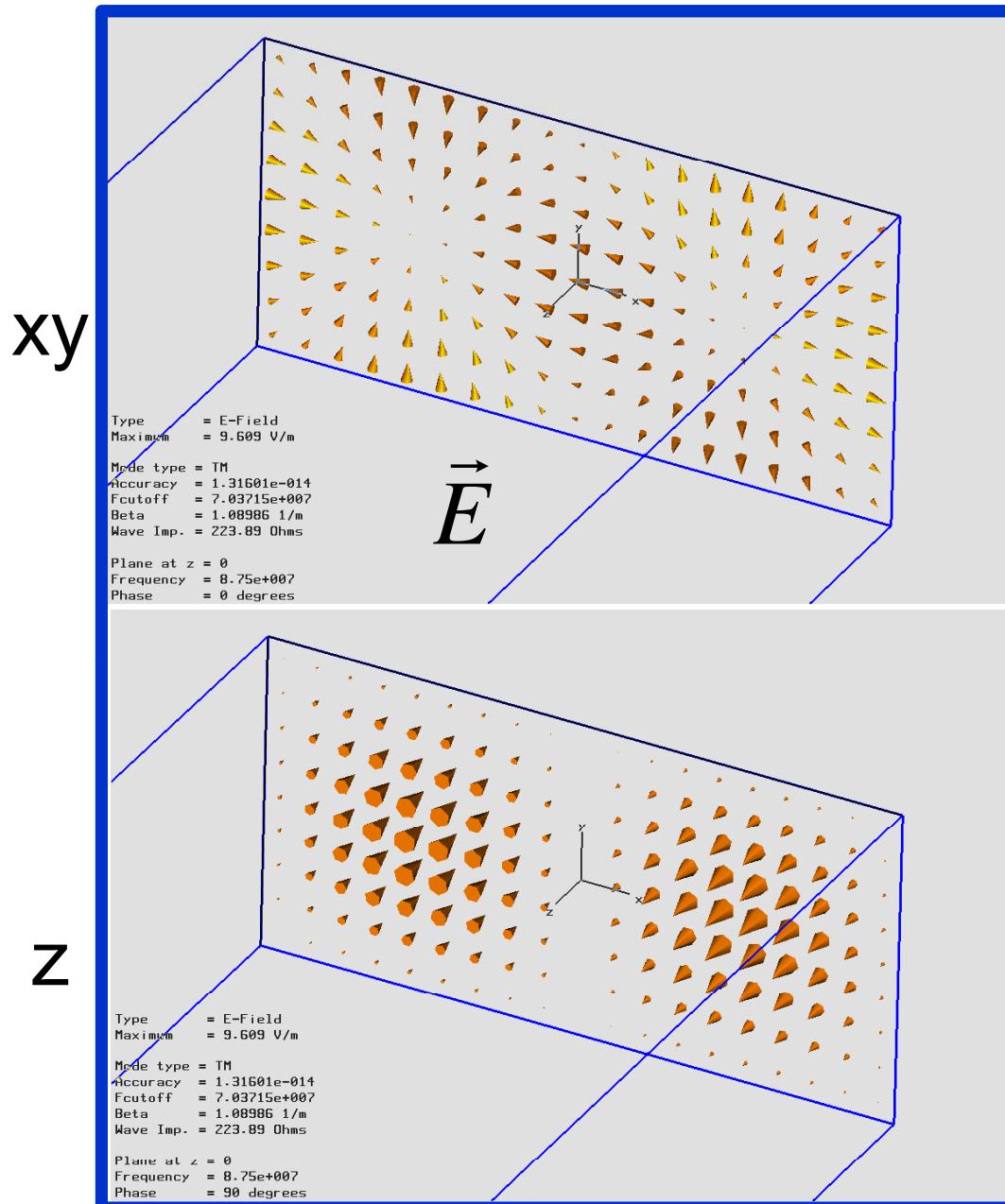


Z



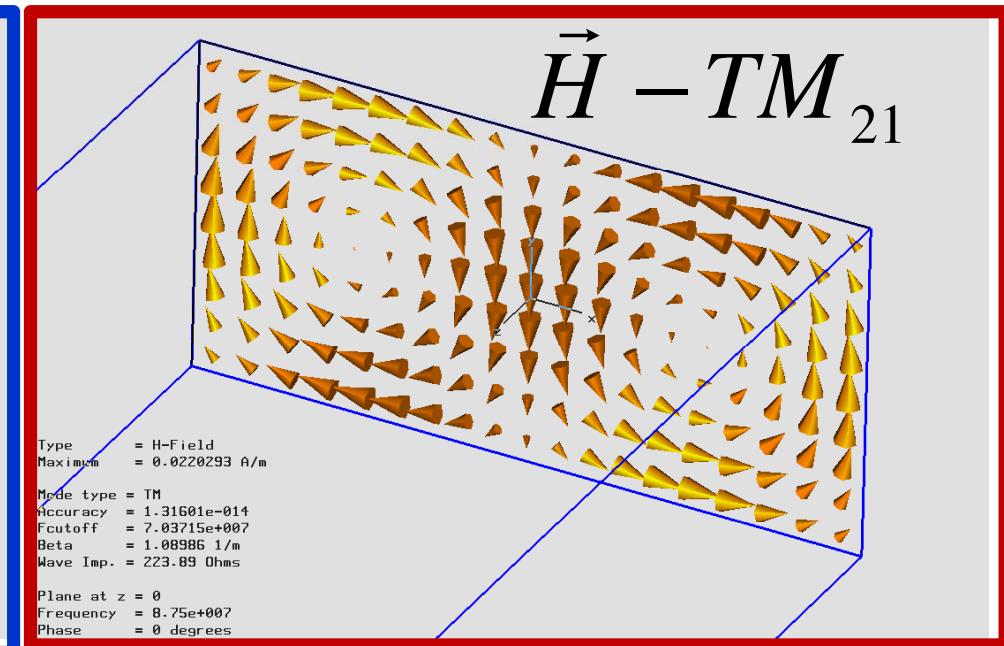
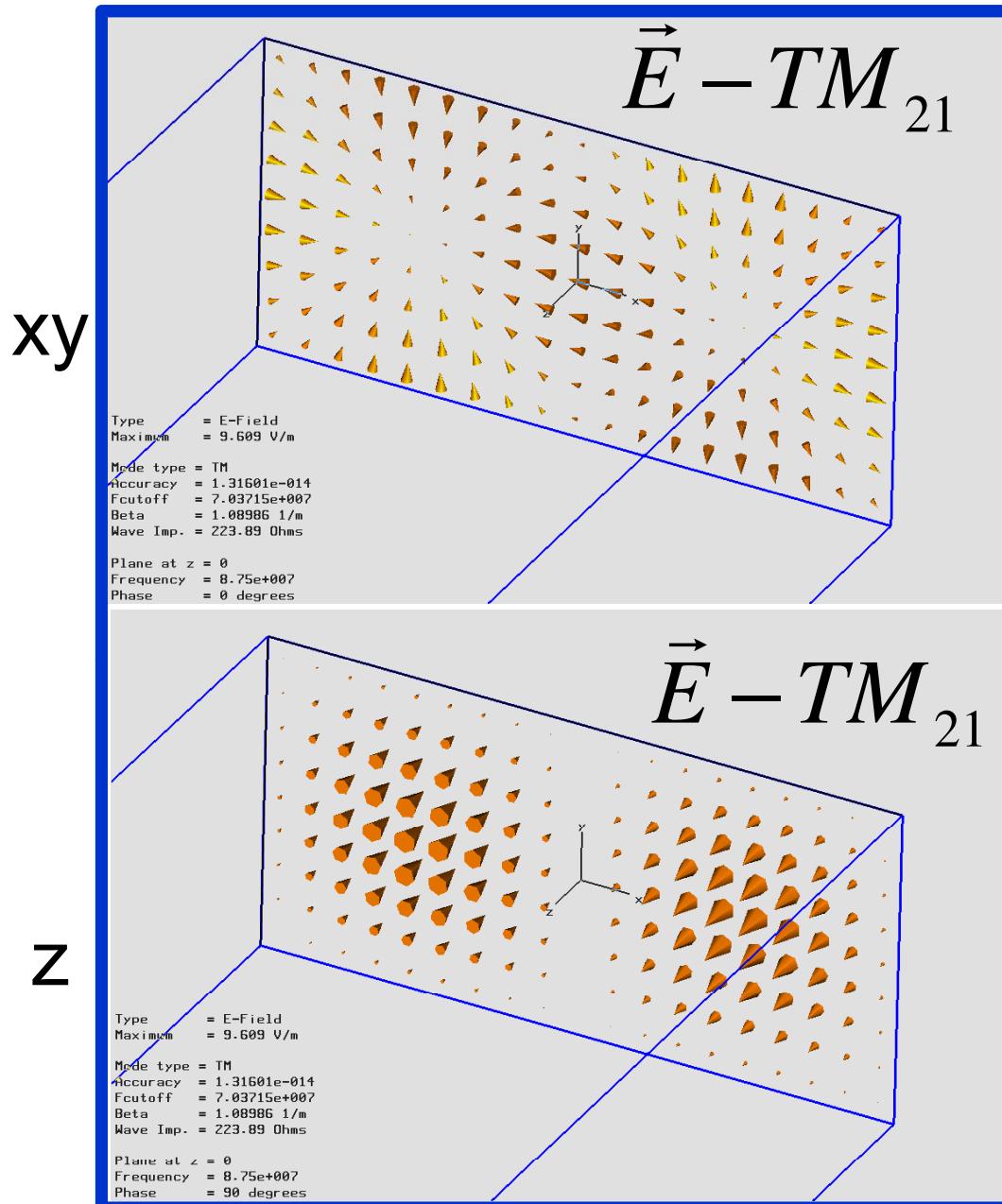
These field plots all correspond to the same mode

Which mode is this?



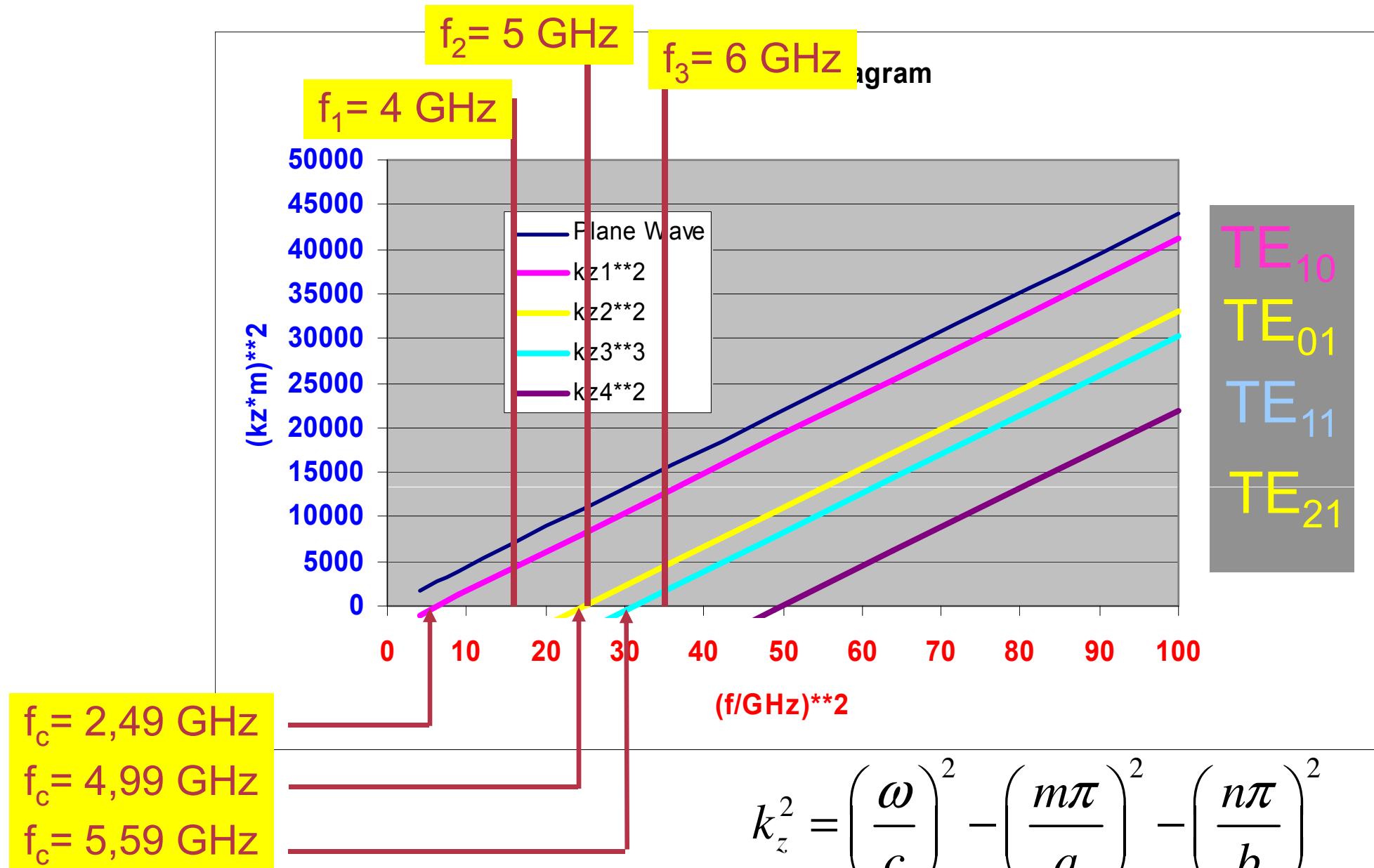
These field plots all correspond to the same mode

Which mode is this?



These field plots all correspond to the same mode

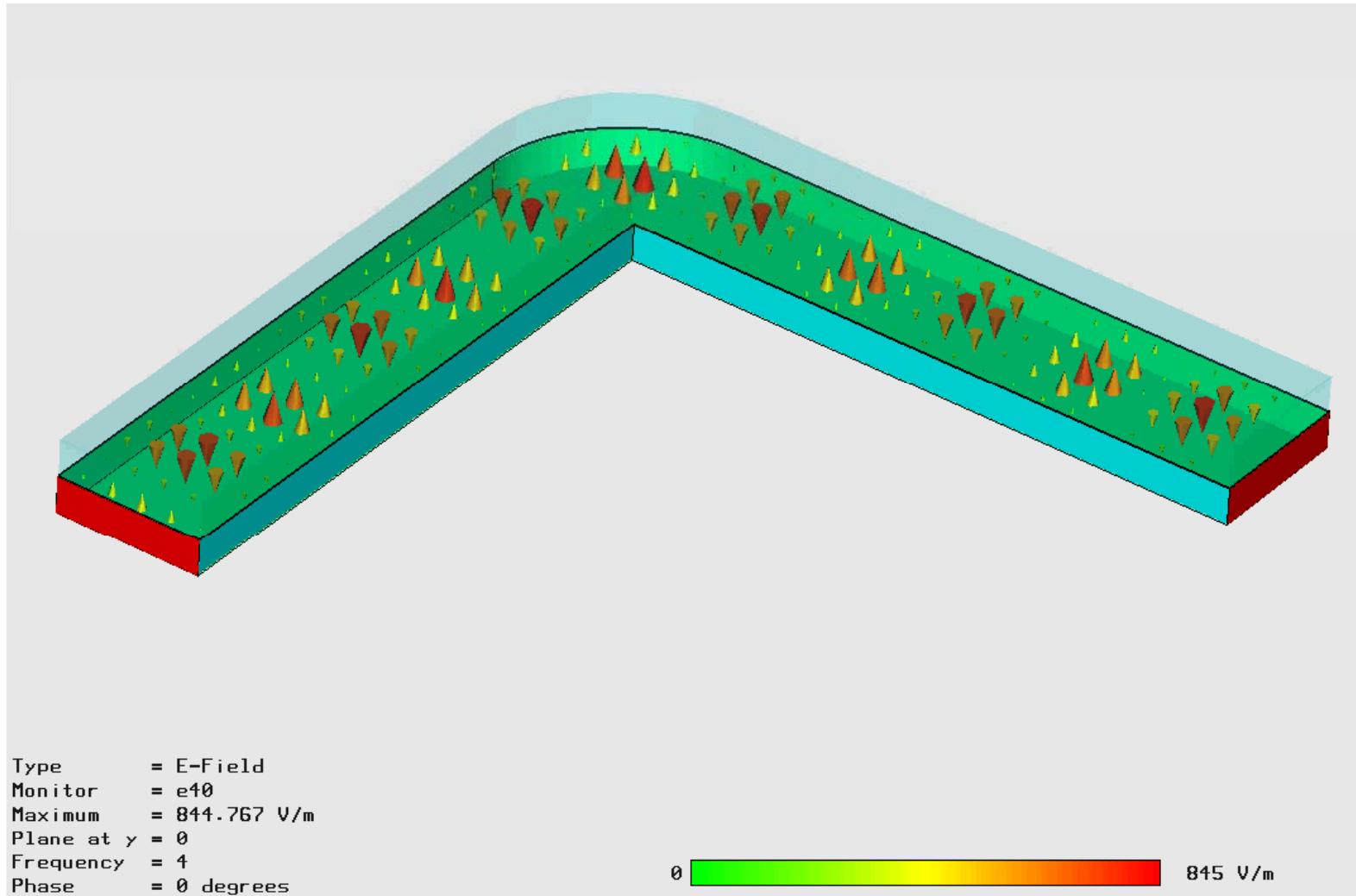
Which mode is this?



$f = 4 \text{ GHz}$

Input-Mode =  $\text{TE}_{10}$       Reflection = -26 dB

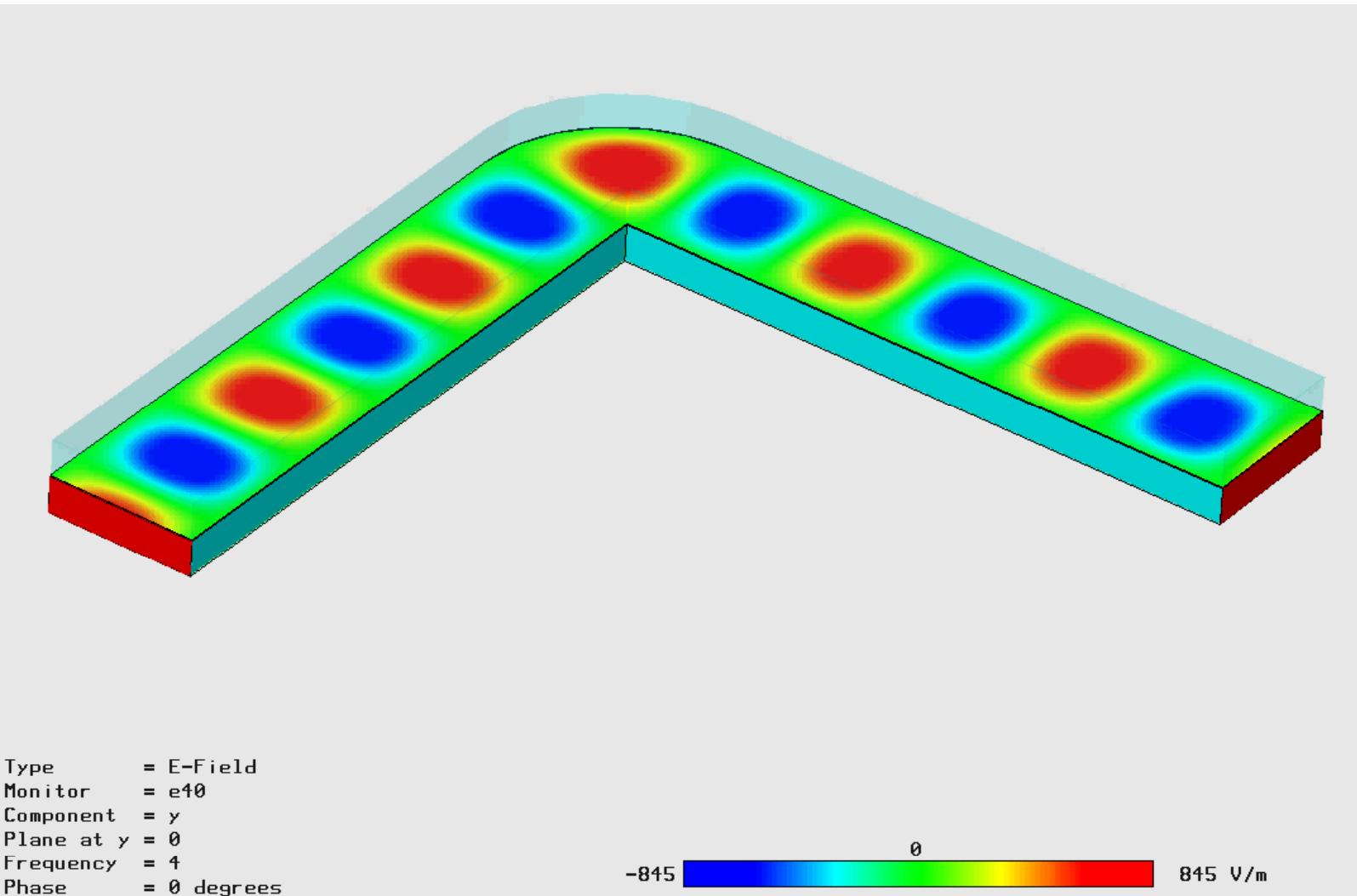
$\vec{E}$



$f = 4 \text{ GHz}$

Input-Mode =  $\text{TE}_{10}$       Reflection = -26 dB

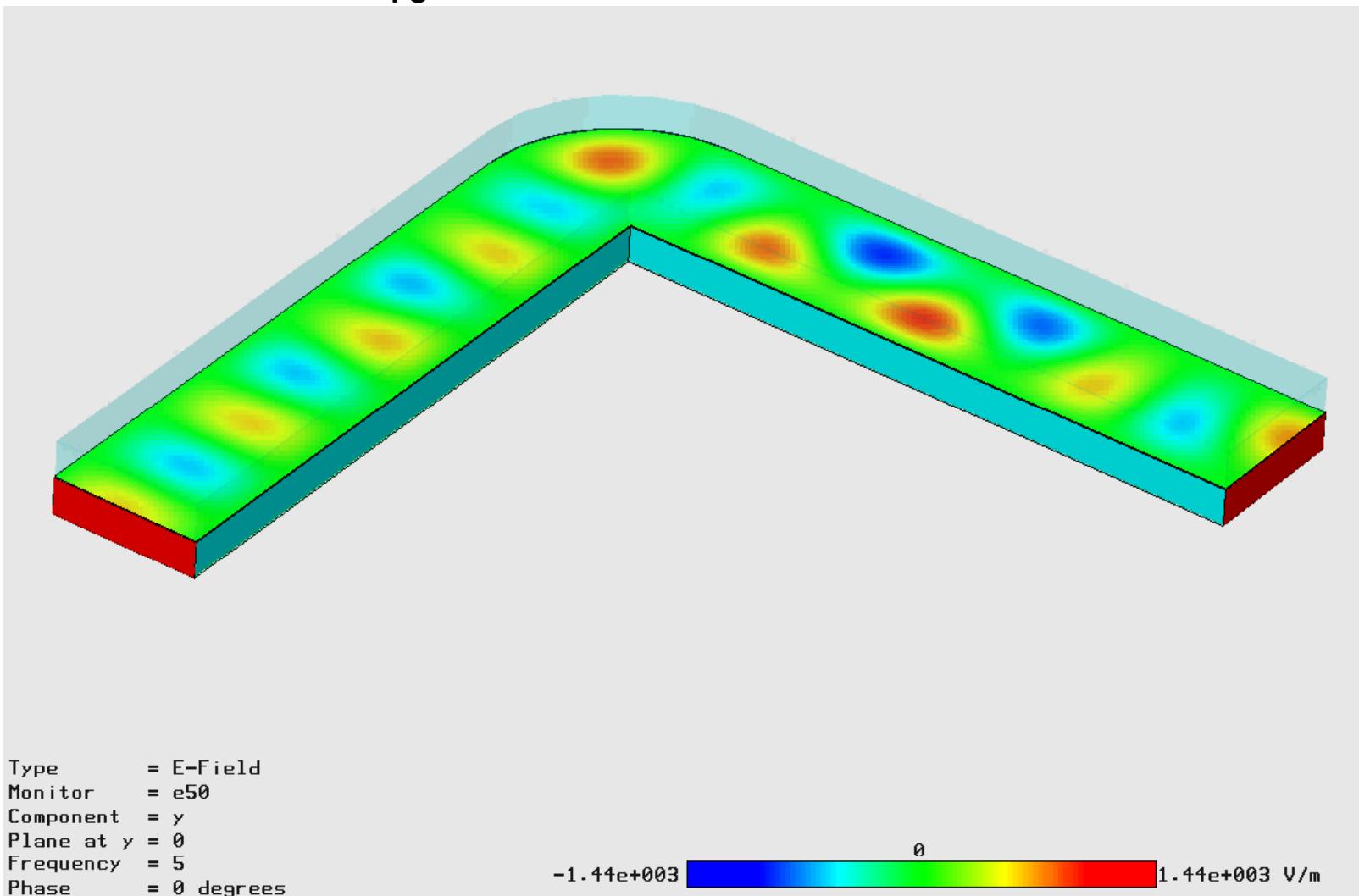
$E_y$



$f = 5 \text{ GHz}$

Input-Mode = TE<sub>10</sub> - 15% of energy in wrong mode

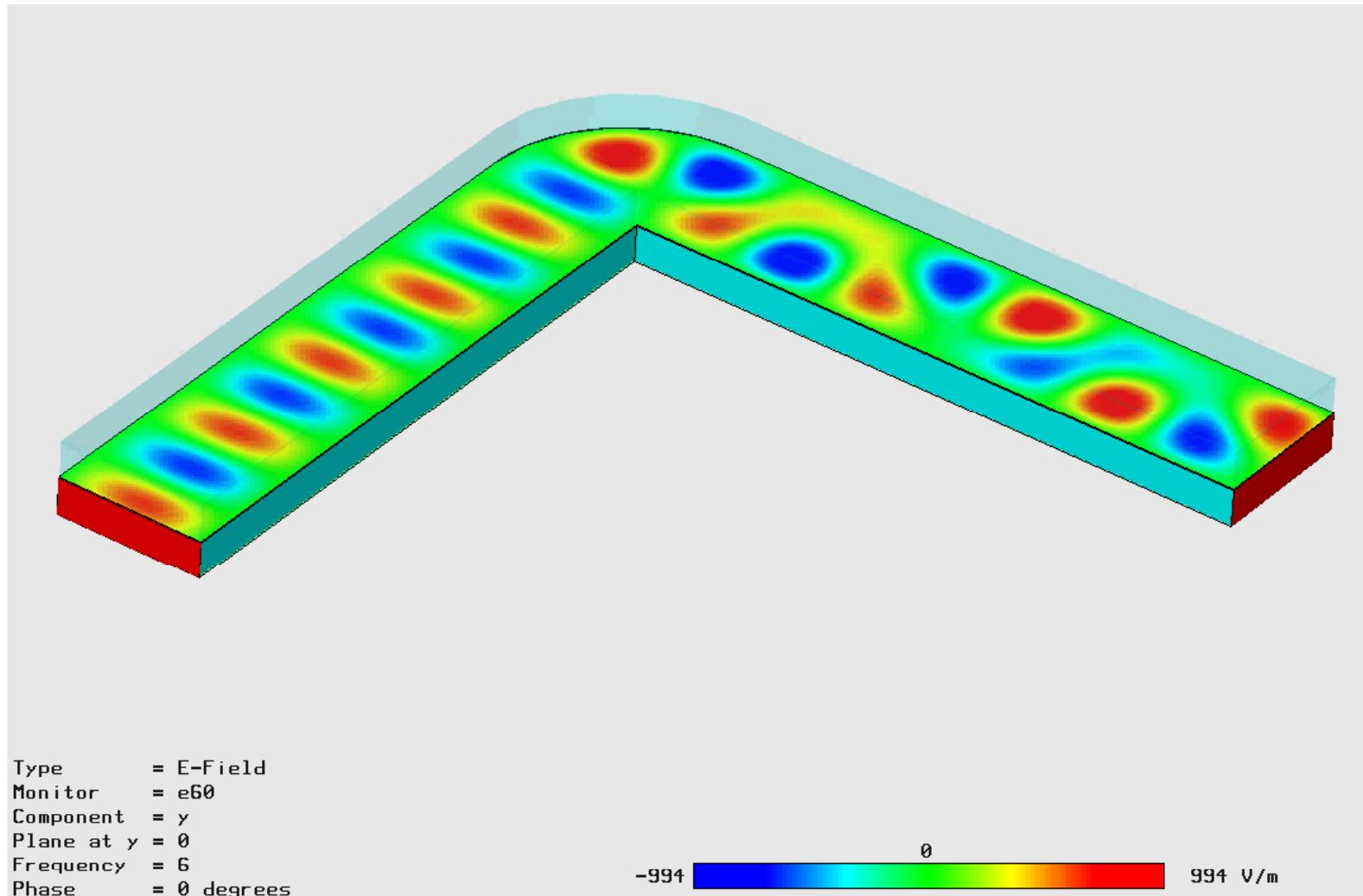
$E_y$

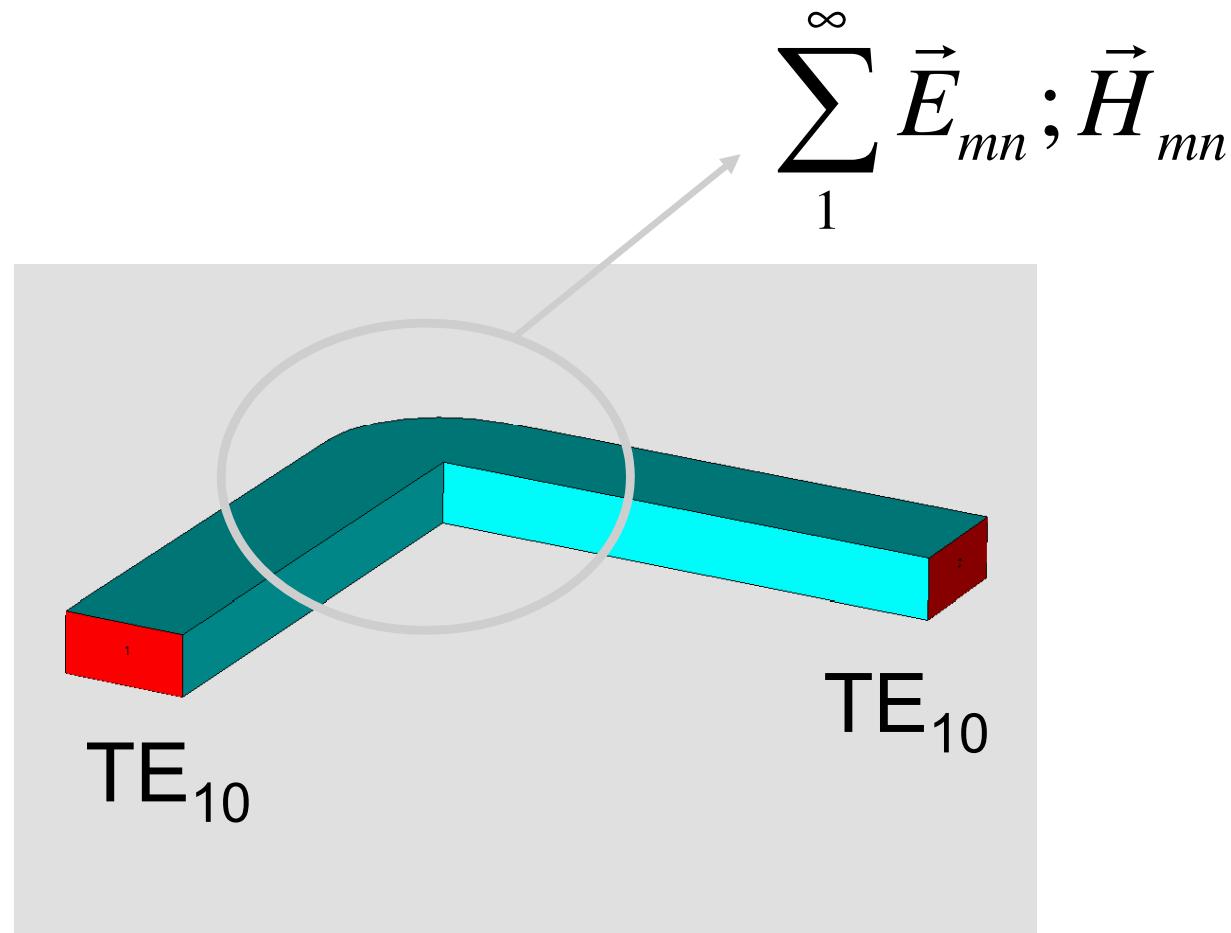


$f = 6 \text{ GHz}$

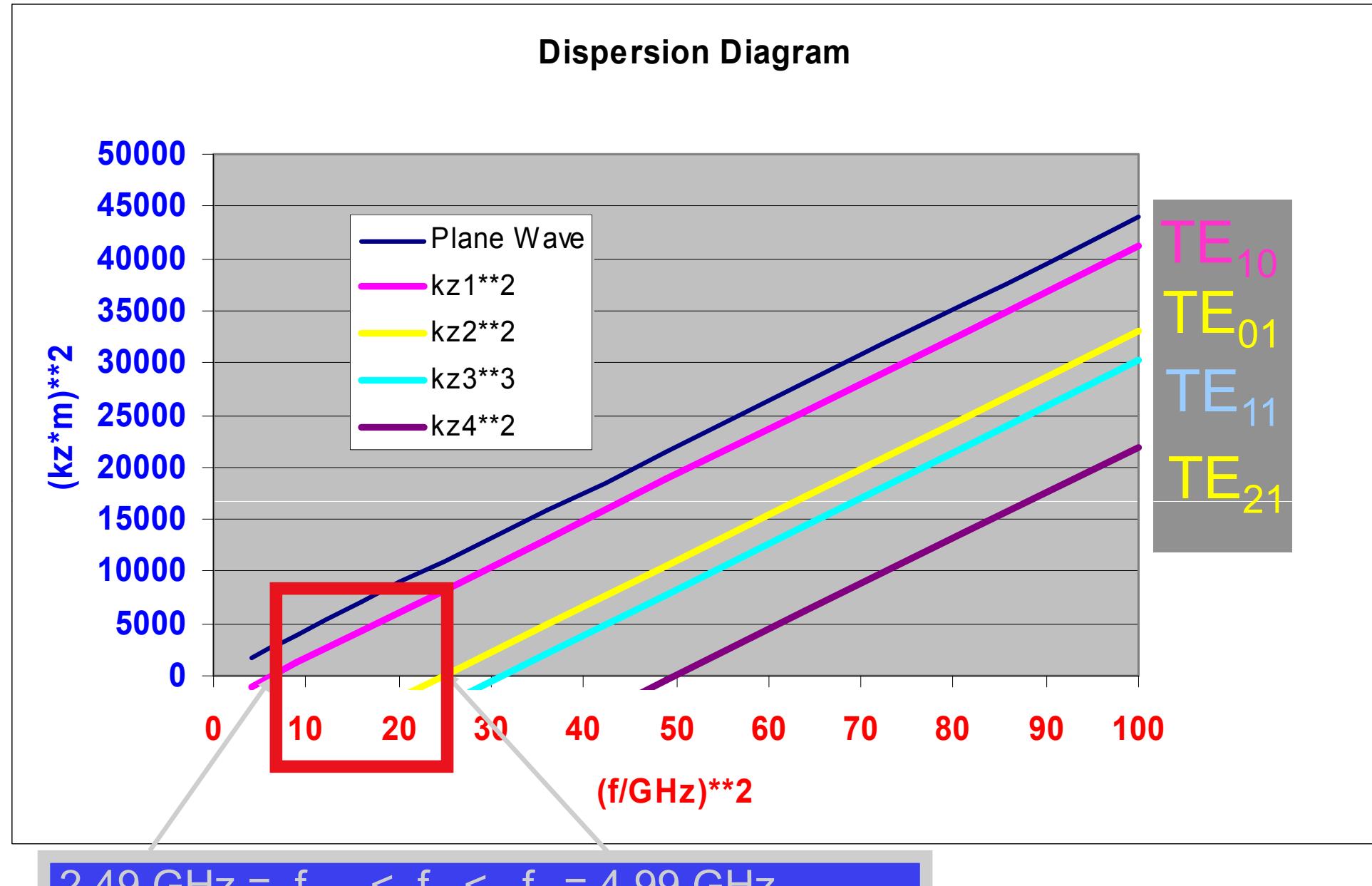
Input-Mode =  $\text{TE}_{10}$  - 75% of energy in wrong mode

$E_y$





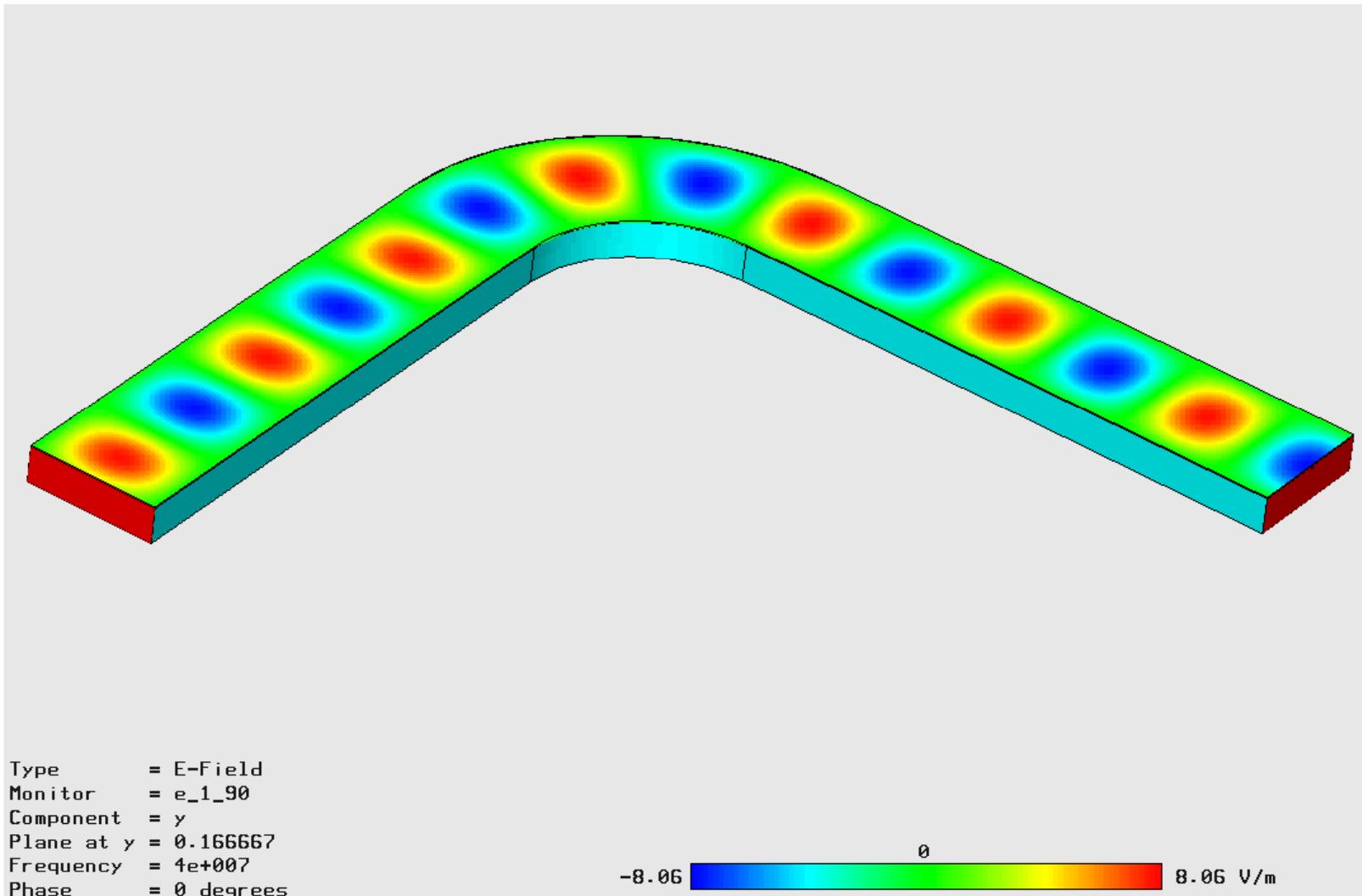
At the *sharp* corner, many higher modes can appear



$f = 4 \text{ GHz}$

Input-Mode =  $\text{TE}_{10}$       Reflection = -50 dB

$E_y$

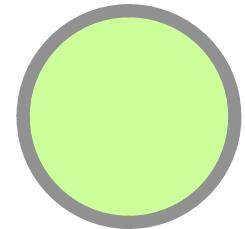


## Circular Waveguides

- **Solution of the wave equation (Summary):**

Circular WG

$$\Delta \vec{A} + k^2 \vec{A} = 0 \quad k^2 = \omega^2 \mu \epsilon$$



- ☞ Vector potential:  $\vec{A} = \vec{e}_z A_z$
- ☞  $\frac{\partial^2 A_z}{\partial r^2} + \frac{1}{r} \frac{\partial A_z}{\partial r} + \frac{1}{r^2} \frac{\partial^2 A_z}{\partial \varphi^2} + \frac{\partial^2 A_z}{\partial z^2} + k^2 A_z = 0 \quad (\text{cylindrical coordinates})$
- ☞ Product approach:  $A_z(r, \varphi, z) = f(r) \cdot g(\varphi) \cdot h(z)$

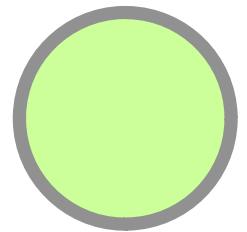


$$A_z(r, \varphi, z) = \begin{cases} J_\mu(Kr) \\ N_\mu(Kr) \end{cases} \cdot \begin{cases} \sin(\mu\varphi) \\ \cos(\mu\varphi) \end{cases} \cdot \begin{cases} e^{+jk_z z} \\ e^{-jk_z z} \end{cases}$$

- ☞ Separation relation:  $K^2 + k_z^2 = k^2$  (Lecture Notes pp. 57-61)

- Derivation of the solutions of the wave equation:

Circular WG

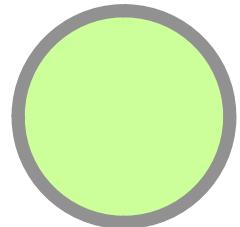


- **Solution of the wave equation:**

$$\frac{\partial^2 A_z}{\partial r^2} + \frac{1}{r} \frac{\partial A_z}{\partial r} + \frac{1}{r^2} \frac{\partial^2 A_z}{\partial \varphi^2} + \frac{\partial^2 A_z}{\partial z^2} + k^2 A_z = 0$$

☞ Product approach:  $A_z(r, \varphi, z) = f(r) \cdot g(\varphi) \cdot h(z)$

Circular WG



☞ Insert this into the equation, divide by  $A_z$

$$\frac{1}{f} \frac{\partial^2 f}{\partial r^2} + \frac{1}{fr} \frac{\partial f}{\partial r} + \frac{1}{r^2} \underbrace{\frac{1}{g} \frac{\partial^2 g}{\partial \varphi^2}}_{-\mu^2} + \underbrace{\frac{1}{h} \frac{\partial^2 h}{\partial z^2}}_{-k_z^2} + k^2 = 0 \quad (\text{cylindrical coordinates})$$



$$g(\varphi) = A_1 \sin(\mu\varphi) + A_2 \cos(\mu\varphi)$$

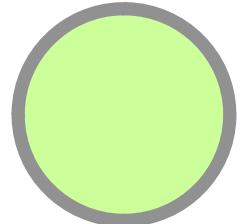
$$h(z) = C_1 e^{jk_z z} + C_2 e^{-jk_z z}$$

(Lecture Notes pp. 57-61)

- Solution of the wave equation:

$$\underbrace{\frac{1}{f} \frac{\partial^2 f}{\partial r^2} + \frac{1}{fr} \frac{\partial f}{\partial r} - \frac{1}{r^2} \mu^2}_{-K^2} - k_z^2 + k^2 = 0$$

Circular WG



☞ Separation relation:  $K^2 + k_z^2 = k^2$

☞ After multiplication with  $f$ :

$$\frac{\partial^2 f}{\partial r^2} + \frac{1}{r} \frac{\partial f}{\partial r} + \left( k^2 - k_z^2 - \frac{\mu^2}{r^2} \right) f = 0$$

$$\frac{\partial^2 f}{\partial r^2} + \frac{1}{r} \frac{\partial f}{\partial r} + \left( K^2 - \frac{\mu^2}{r^2} \right) f = 0$$

Bessel differential equation



$$f(r) = B_1 J_\mu(Kr) + B_2 N_\mu(Kr)$$

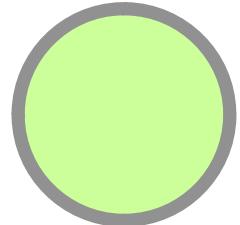
(Lecture Notes pp. 57-61)

- Solution of the wave equation:

Circular WG



$$A_z(r, \varphi, z) = \begin{cases} J_\mu(Kr) \\ N_\mu(Kr) \end{cases} \cdot \begin{cases} \sin(\mu\varphi) \\ \cos(\mu\varphi) \end{cases} \cdot \begin{cases} e^{+jk_z z} \\ e^{-jk_z z} \end{cases}$$

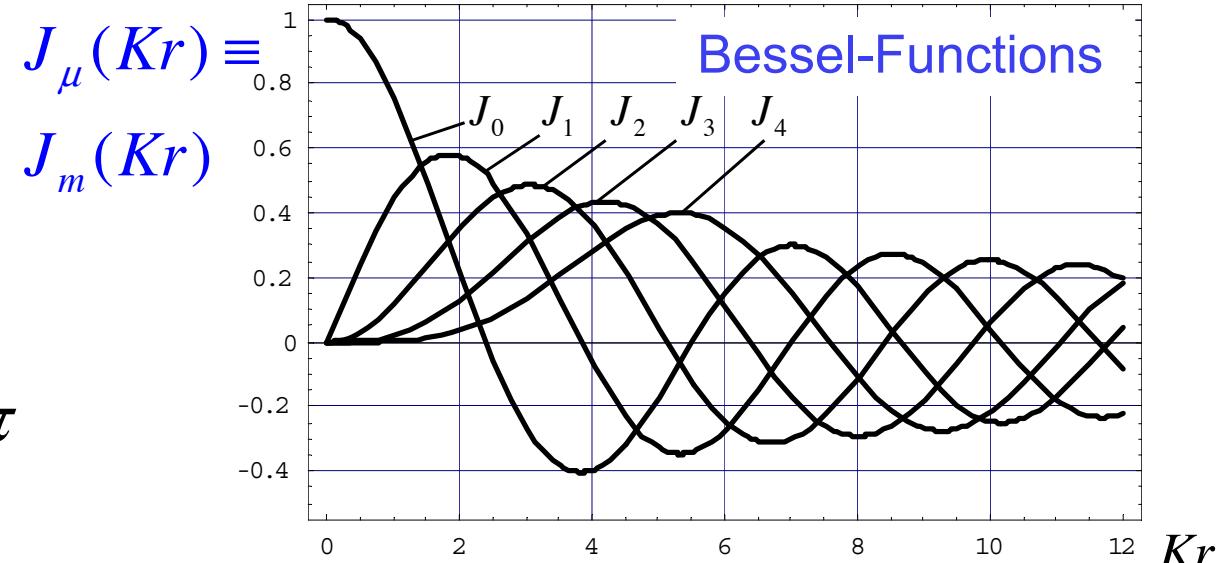


*Separation relation:*  $K^2 + k_z^2 = k^2$

(Lecture Notes pp. 57-61)

- **Cylindrical functions:**

- ☞  $\mu$  couples azimuthal and radial dependency
- ☞  $\mu \in \mathbb{N}_0$  in cylindrical systems with  $0 \leq \varphi < 2\pi$

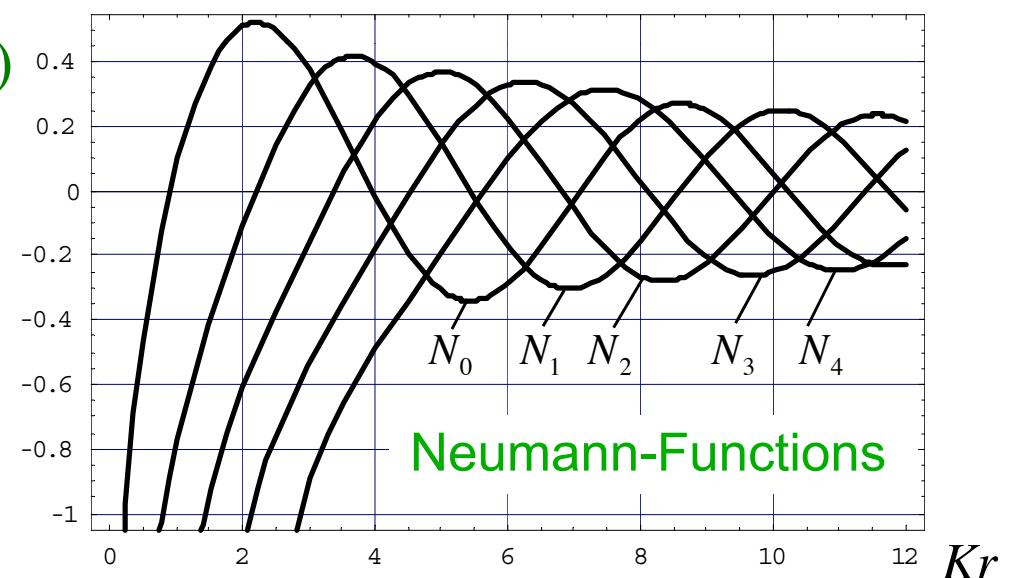


$$N_\mu(Kr)$$

- ☞ If the z-axis is included only the functions of the first form (Bessel) may be considered

$$\lim_{r \rightarrow 0} N_\mu(Kr) = -\infty \quad \forall \mu$$

$$\Rightarrow A \sim J_\mu(Kr)$$



$$\mu = m \quad A_z(r, \varphi, z) = \begin{Bmatrix} J_m(Kr) \\ N_m(Kr) \end{Bmatrix} \cdot \begin{Bmatrix} \sin(m\varphi) \\ \cos(m\varphi) \end{Bmatrix} \cdot \begin{Bmatrix} e^{+jk_z z} \\ e^{-jk_z z} \end{Bmatrix}$$

- $\varphi$  dependency: linear combination of sine and cosine functions

- ☞ Rotational symmetry: choose origin for  $\varphi$  such that one single function is sufficient
- ☞ Sine function is 0 for  $m = 0$

$$\Rightarrow A \sim \cos(m\varphi)$$

- z dependency: consider first forward-propagating waves

$$\Rightarrow A \sim e^{-jk_z z}$$

$$A_z(r, \varphi, z) = C J_m(Kr) \cos(m\varphi) e^{-jk_z z}$$

## ● TE-modes

☞ Field vectors:

$$\vec{E} = \text{curl} \vec{A} \quad \vec{H} = \frac{-1}{j\omega\mu} \text{curl} \vec{E}$$

$$A_z(r, \varphi, z) = C J_m(Kr) \cos(m\varphi) e^{-jk_z z}$$

☞ Components:

$$E_r = \frac{1}{r} \frac{\partial A_z}{\partial \varphi} = -C \frac{m}{r} J_m(Kr) \sin m\varphi e^{-jk_z z}$$

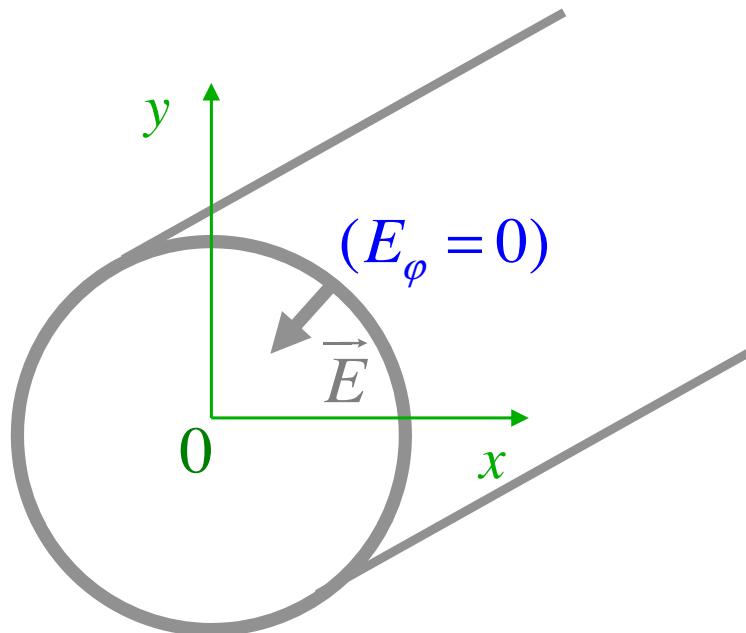
$$E_\varphi = -\frac{\partial A_z}{\partial r} = -C K J'_m(Kr) \cos m\varphi e^{-jk_z z}$$

$$H_r = \frac{1}{j\omega\mu} \frac{\partial E_\varphi}{\partial z} = -C \frac{k_z}{\omega\mu} K J'_m(Kr) \cos m\varphi e^{-jk_z z}$$

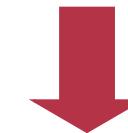
$$H_\varphi = \frac{-1}{j\omega\mu} \frac{\partial E_r}{\partial z} = -C \frac{k_z}{\omega\mu} \frac{m}{r} J_m(Kr) \sin m\varphi e^{-jk_z z}$$

$$H_z = -\frac{K^2}{j\omega\mu} A_z = -C \frac{K^2}{j\omega\mu} J_m(Kr) \cos m\varphi e^{-jk_z z}$$

- TE-solution:



$$E_\phi = -C K J_m(Kr) \cos m\varphi e^{-jk_z z} = 0$$



$$J_m(Kr) = 0 \quad \text{for } r = a \rightarrow K = \frac{j_{mn}}{a}$$

$j_{mn}$  : n-th root of  $J_m(Kr)$

$j_{mn}$	$m$			
	0	1	2	3
1	3.832	1.841	3.054	4.201
2	7.016	5.331	6.706	8.015
3	10.173	8.536	9.969	11.346

Continuity:

$$E_\phi(r = a, \varphi, z) = 0$$

- Fundamental mode for the circular waveguide: TE11
- Cutoff angular frequency:

$$k_z = \sqrt{k^2 - K^2} = 0 \quad \Rightarrow K = k = \frac{\omega}{c}$$
$$\Rightarrow \omega = Kc$$

- for TE modes:

$$\omega_{c_{mn}} = Kc = \frac{j'_{mn}}{a} c$$

- Cutoff angular frequency for the fundamental mode:

$$\omega_{c_{11}} = \frac{j'_{11}}{a} c$$

## ● TM-modes

☞ Field vectors:

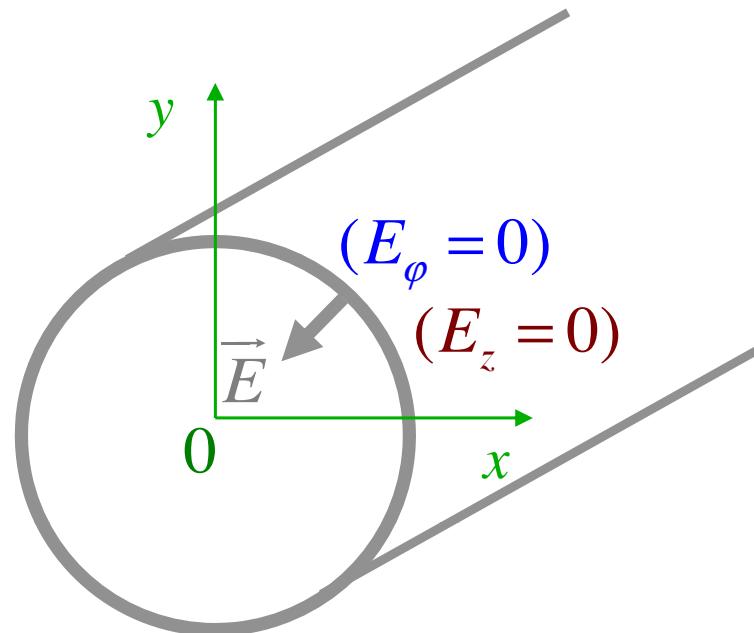
$$\vec{H} = \operatorname{curl} \vec{A}$$

$$\vec{E} = \frac{1}{j\omega\epsilon} \operatorname{curl} \vec{H}$$

☞ Components:

$H_r$	$= \frac{1}{r} \frac{\partial A_z}{\partial \varphi}$	$= -C \frac{m}{r} J_m(Kr) \sin m\varphi e^{-jk_z z}$
$H_\varphi$	$= -\frac{\partial A_z}{\partial r}$	$= -C K J'_m(Kr) \cos m\varphi e^{-jk_z z}$
$E_r$	$= \frac{-1}{j\omega\epsilon} \frac{\partial H_\varphi}{\partial z}$	$= -C \frac{k_z}{\omega\epsilon} K J'_m(Kr) \cos m\varphi e^{-jk_z z}$
$E_\varphi$	$= \frac{1}{j\omega\epsilon} \frac{\partial H_r}{\partial z}$	$= C \frac{k_z}{\omega\epsilon} \frac{m}{r} J_m(Kr) \sin m\varphi e^{-jk_z z}$
$E_z$	$= \frac{K^2}{j\omega\epsilon} A_z$	$= C \frac{K^2}{j\omega\epsilon} J_m(Kr) \cos m\varphi e^{-jk_z z}$

- TM-solution:



$$E_\phi = C \frac{k_z m}{\omega \epsilon r} J_m(Kr) \sin m\varphi e^{-jk_z z} = 0$$

$$E_z = C \frac{K^2}{j\omega \epsilon} J_m(Kr) \cos m\varphi e^{-jk_z z} = 0$$



$$J_m(Kr) = 0 \quad \text{for } r = a \rightarrow K = \frac{j_{mn}}{a}$$

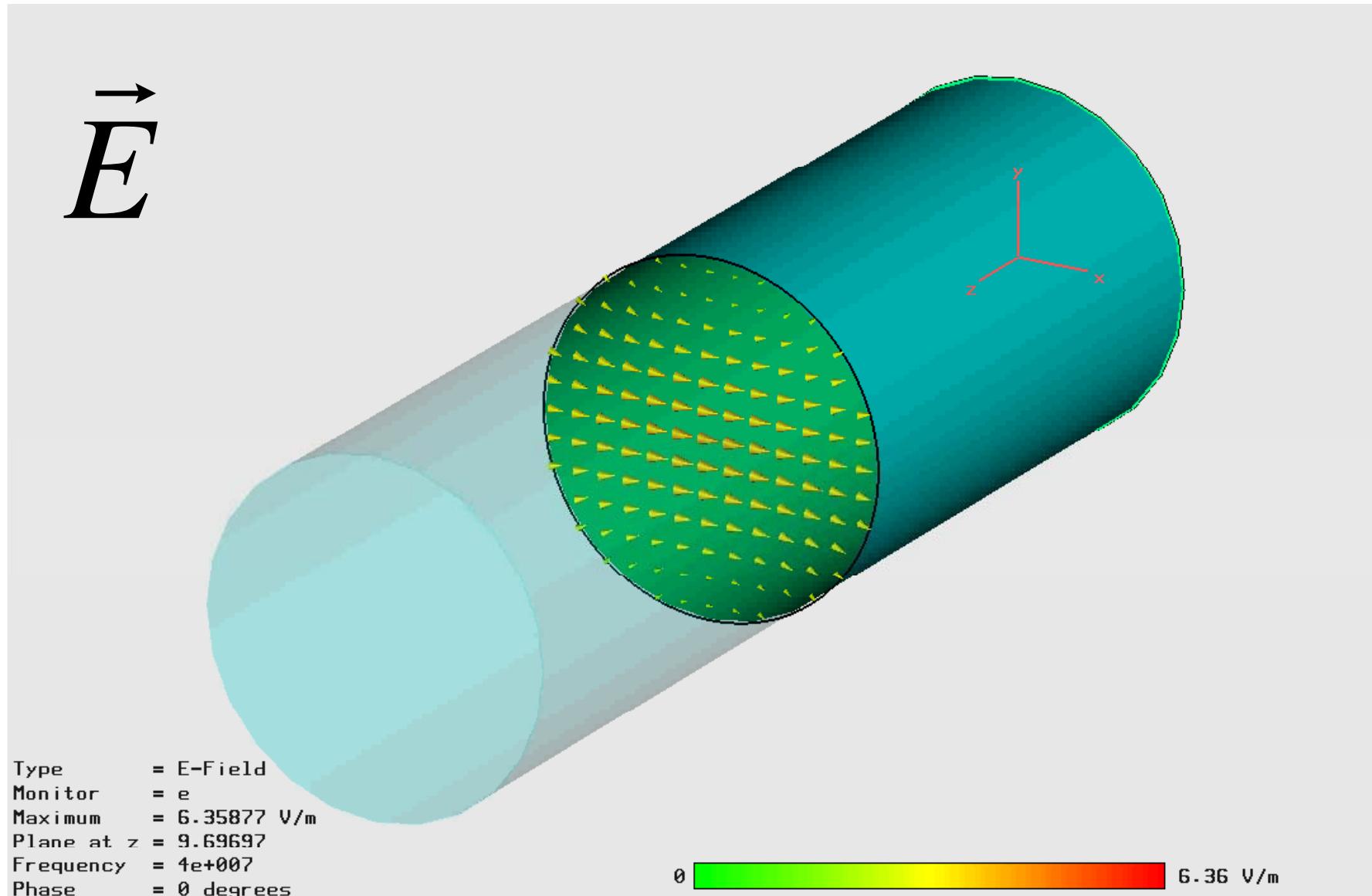
$j_{mn}$	$m$			
	0	1	2	3
1	2.405	3.832	5.136	6.380
2	5.520	7.016	8.417	9.761
3	8.654	10.173	11.620	13.015

Continuity:

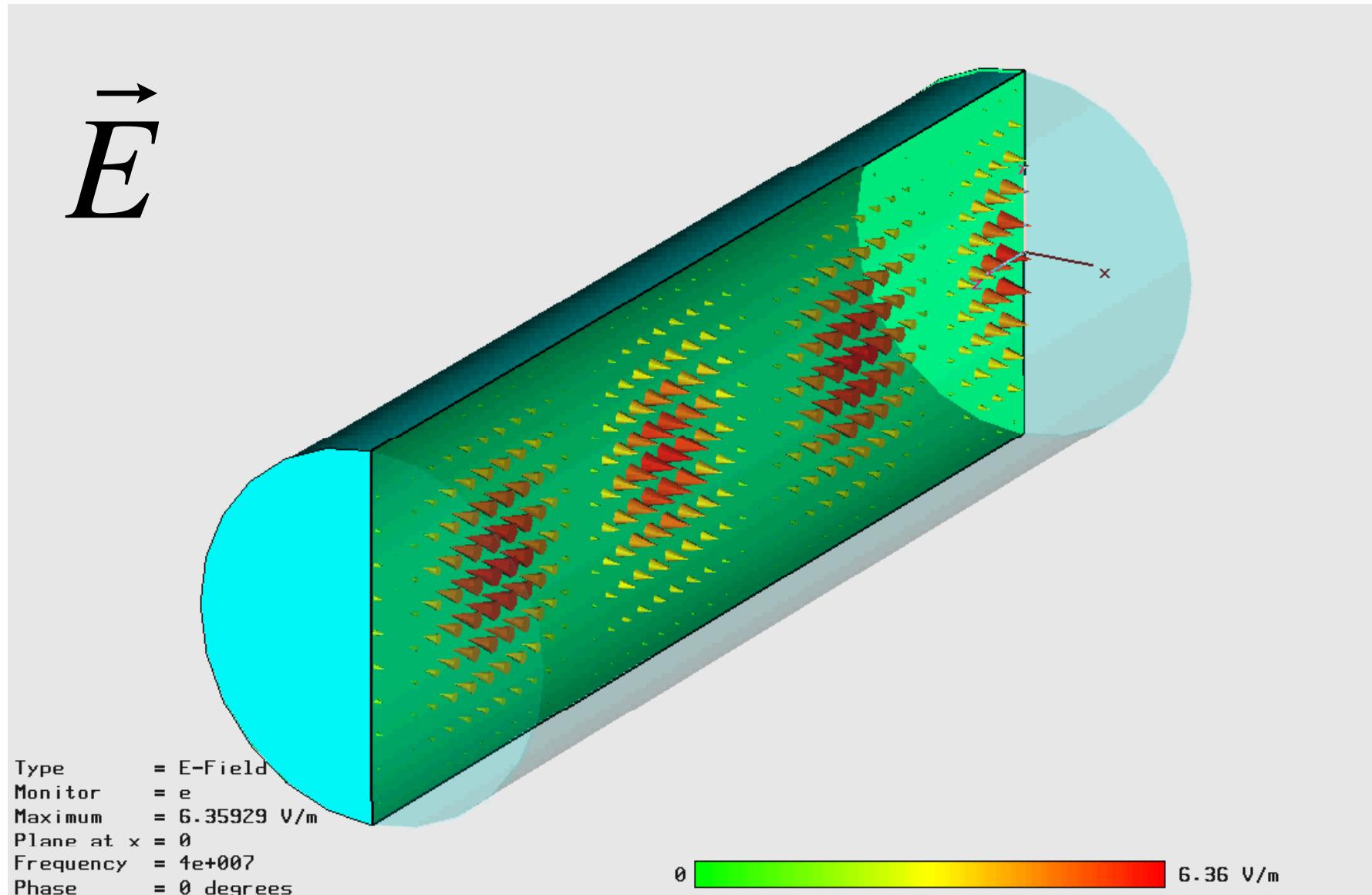
$$E_\phi(r = a, \varphi, z) = 0$$

$$E_z(r = a, \varphi, z) = 0$$

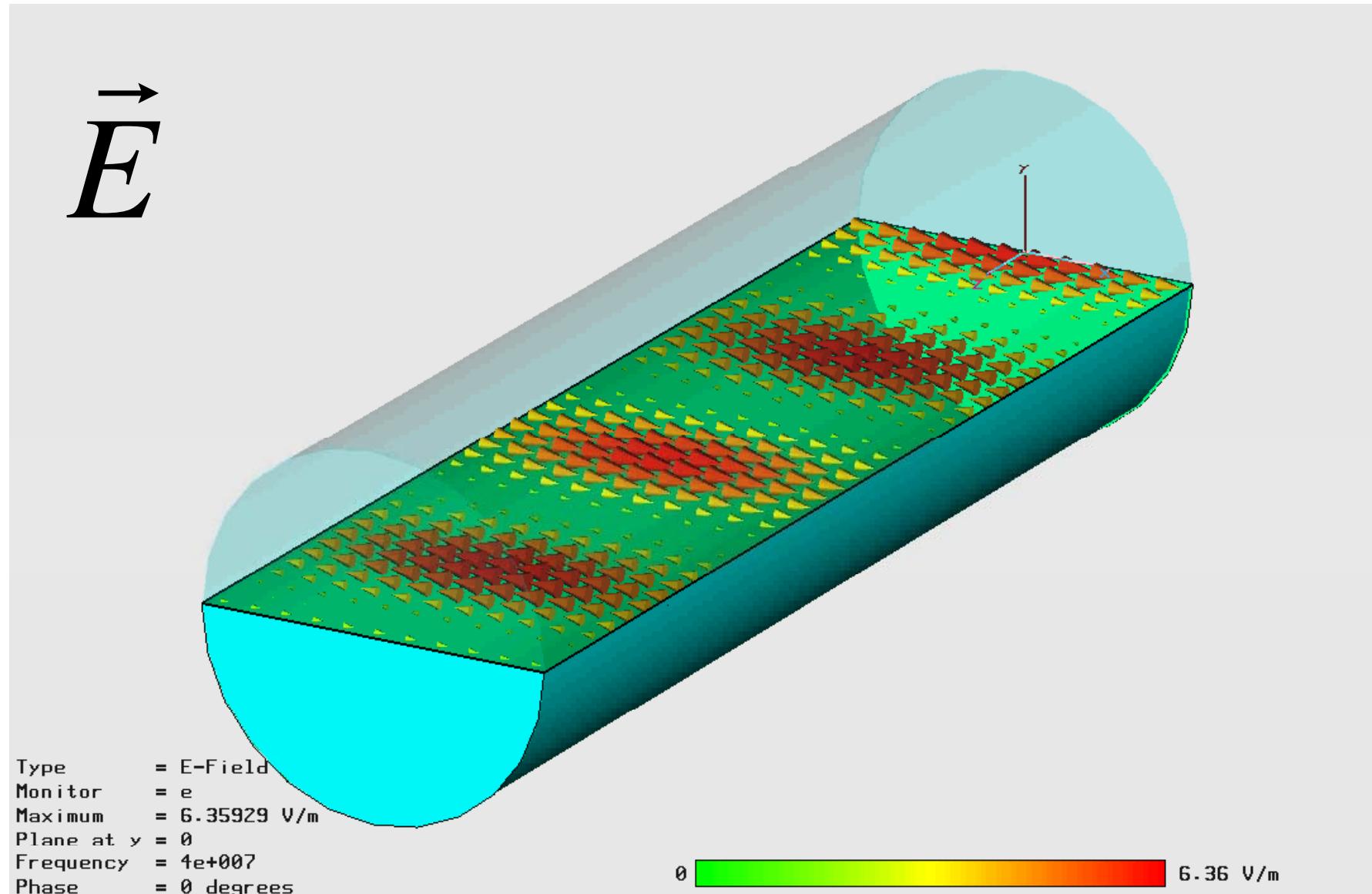
- Electric Field of TE<sub>11</sub>-mode (fundamental mode):



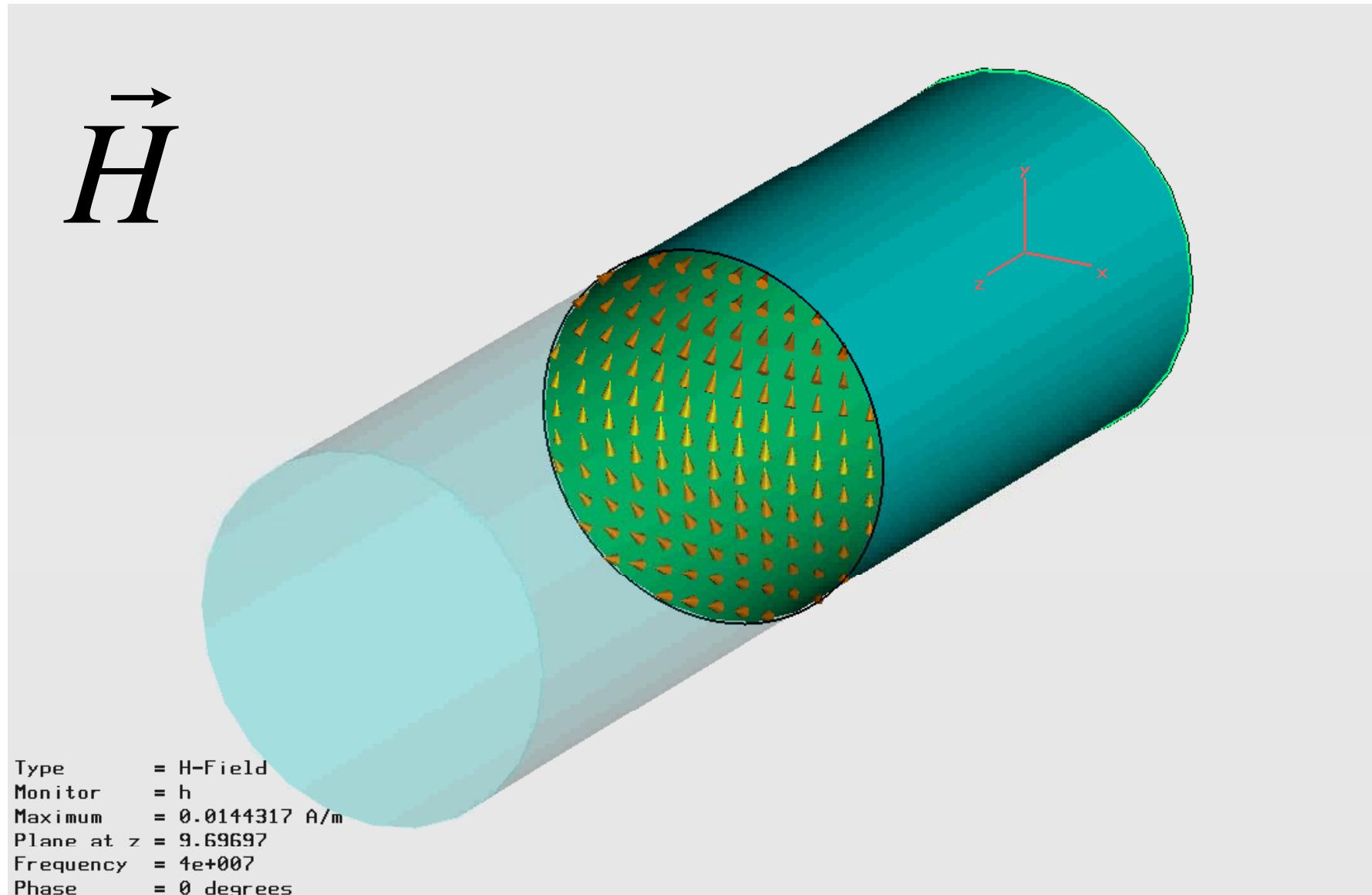
- Electric Field of TE<sub>11</sub>-mode (fundamental mode):



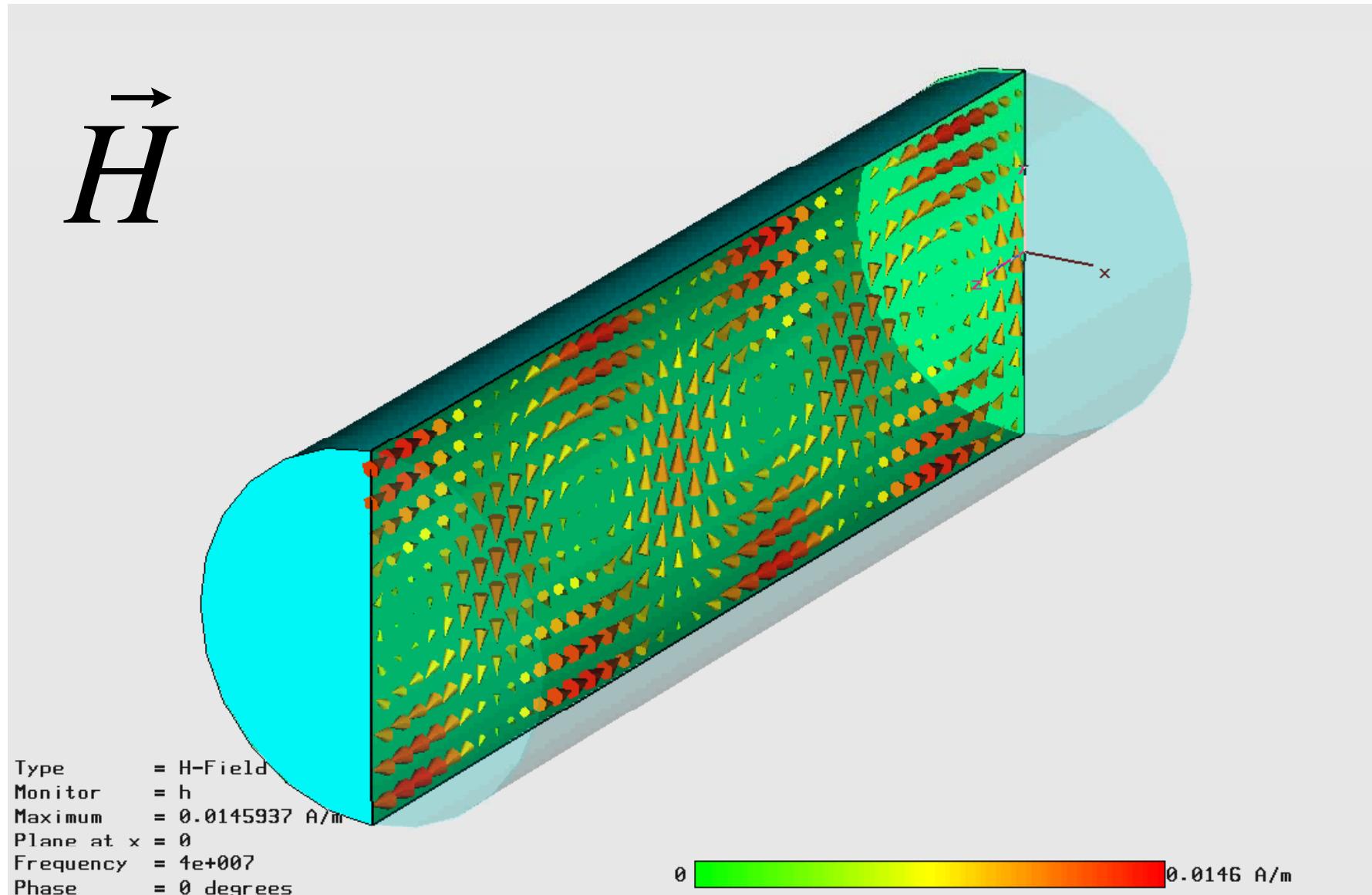
- Electric Field of TE<sub>11</sub>-mode (fundamental mode):



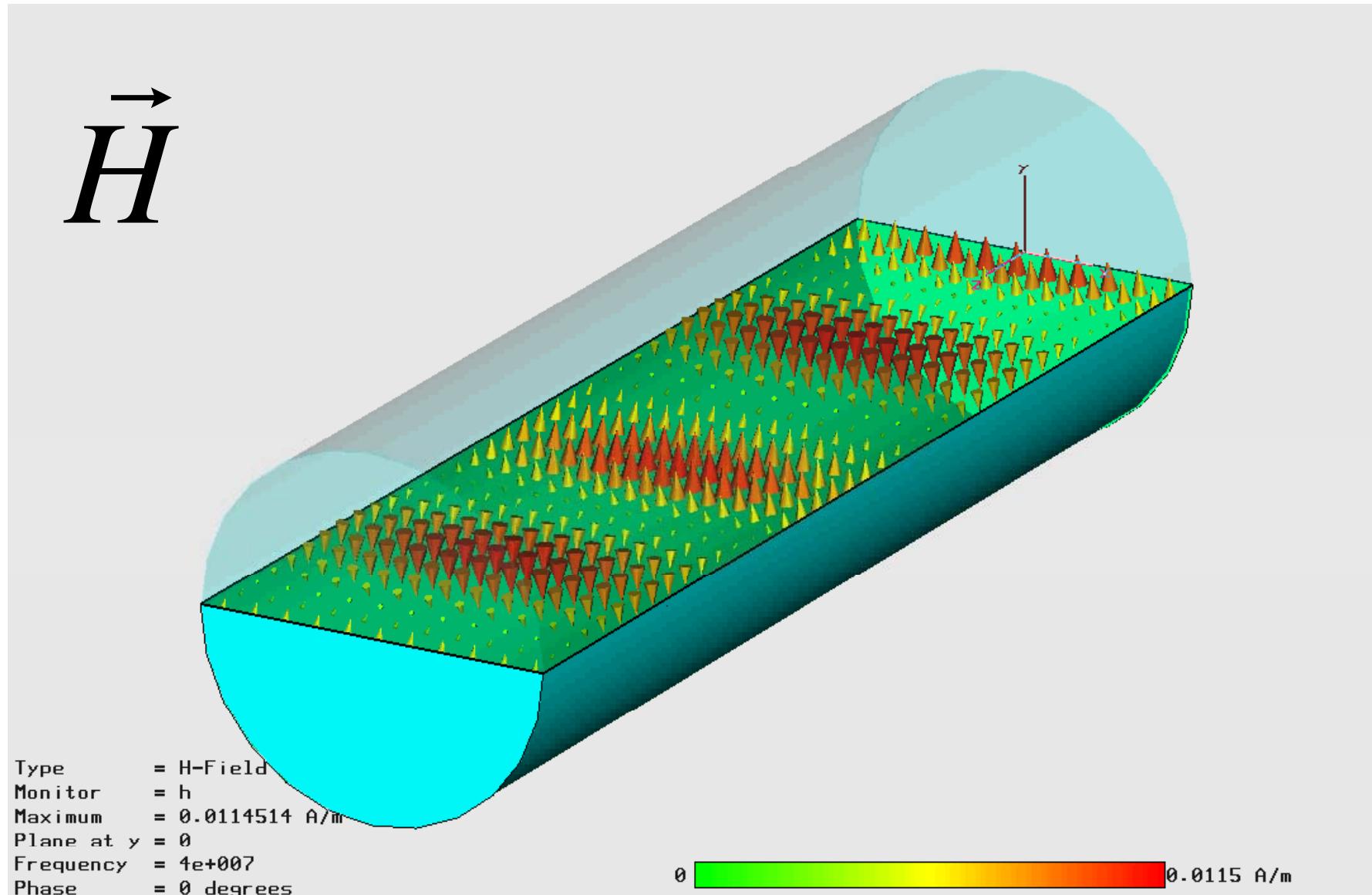
- Magnetic Field of TE<sub>11</sub>-mode (fundamental mode):



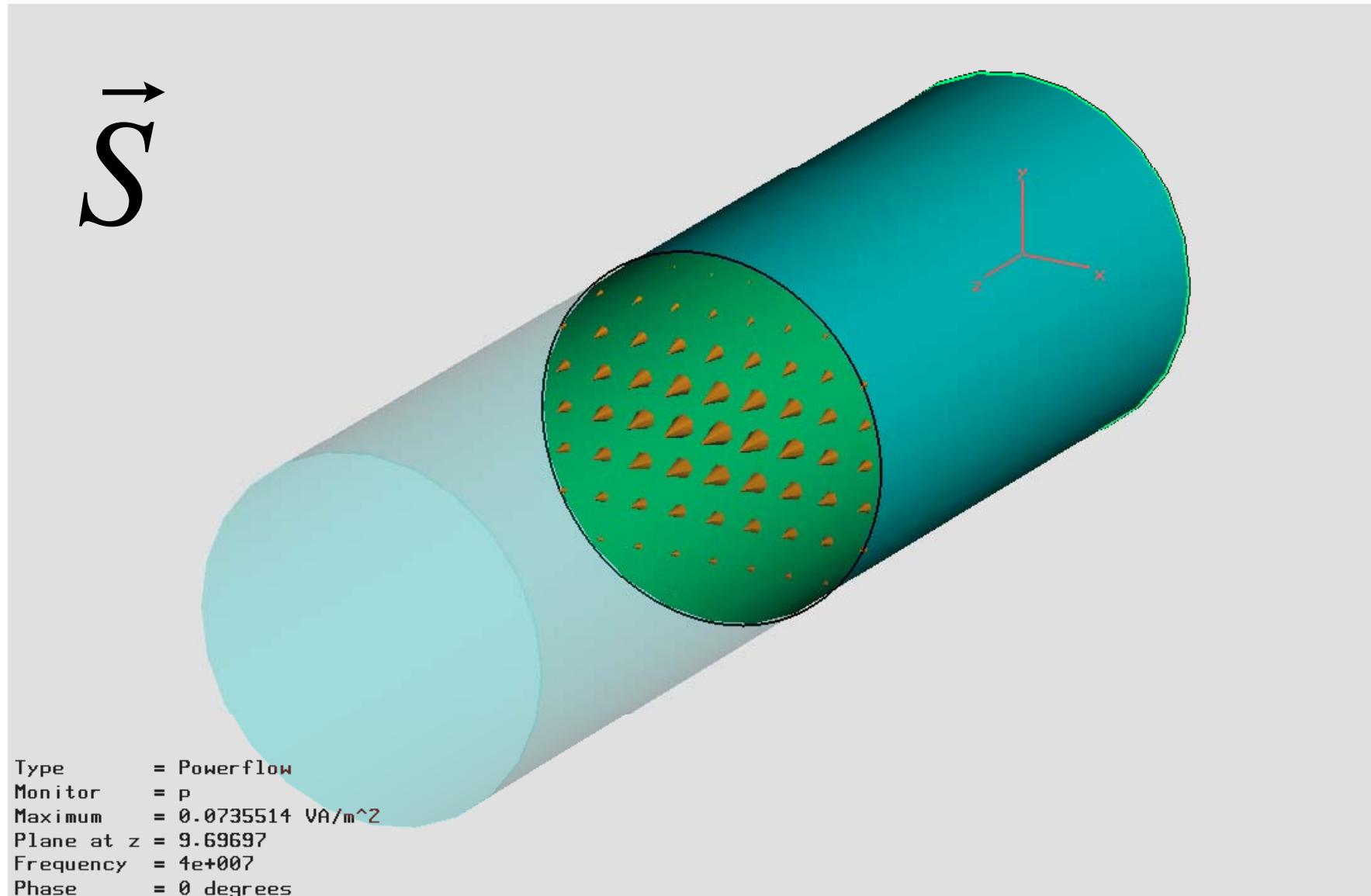
- Magnetic Field of TE<sub>11</sub>-mode (fundamental mode):



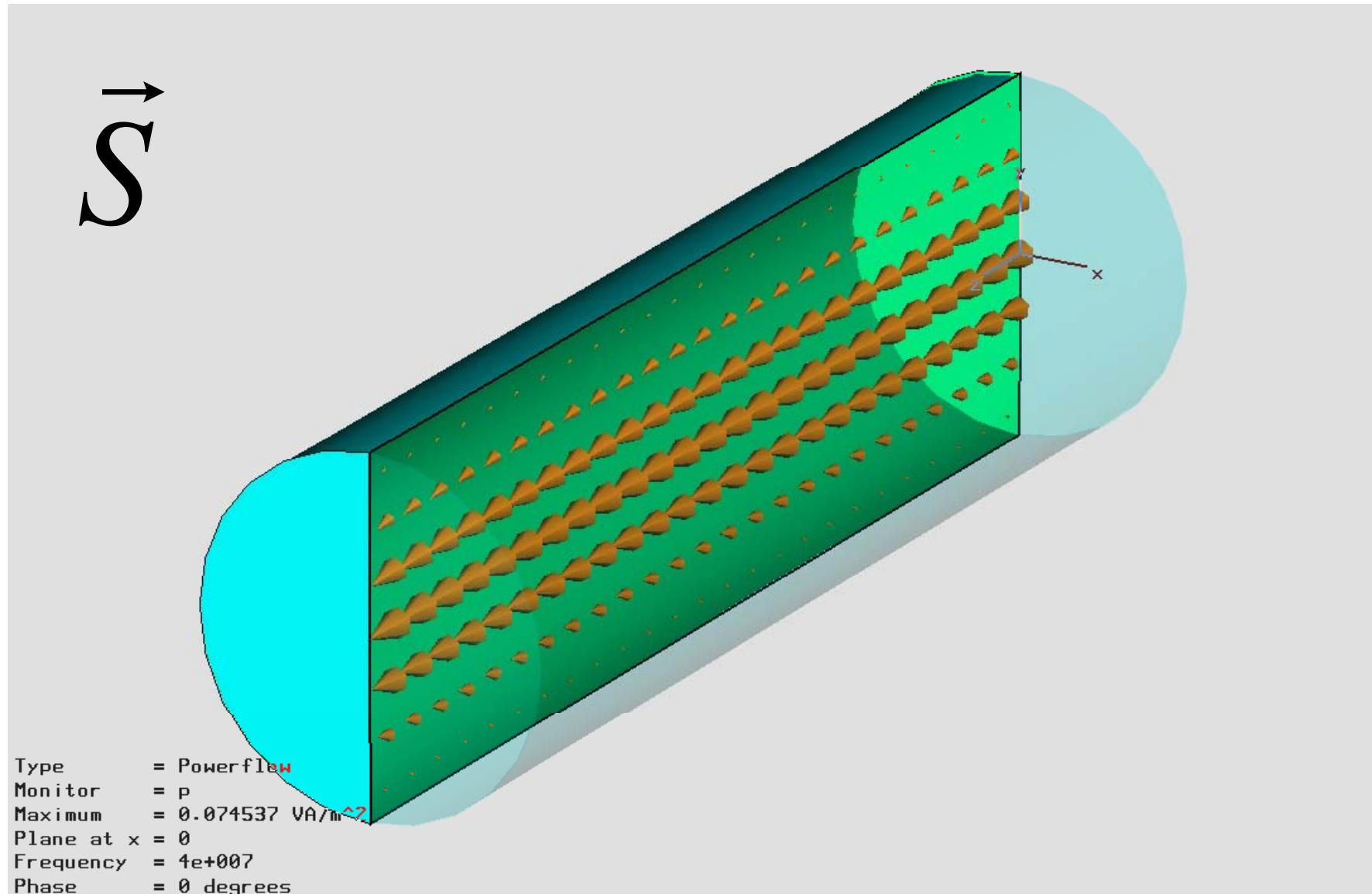
- Magnetic Field of  $\text{TE}_{11}$ -mode (fundamental mode):



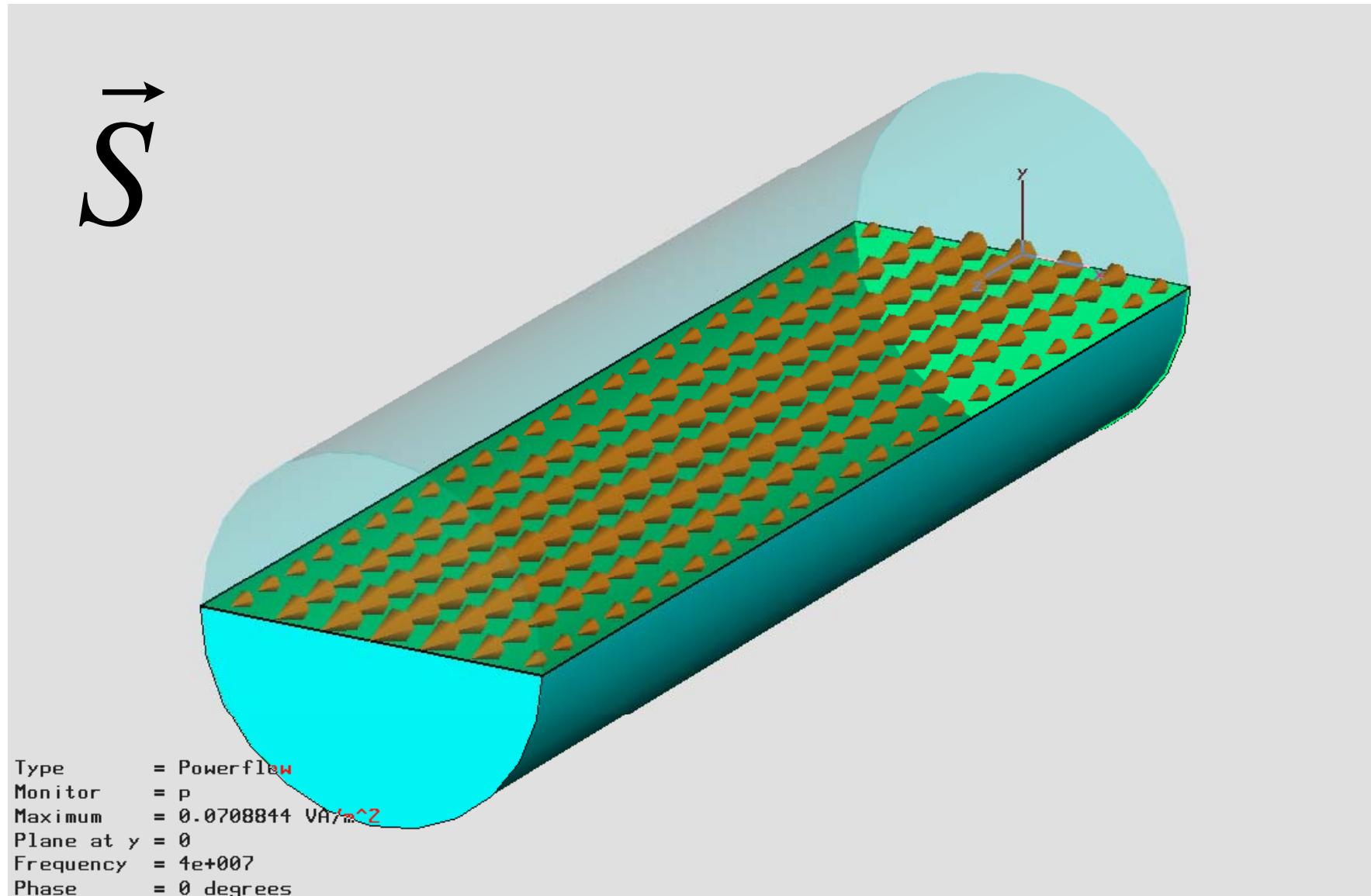
- Poynting vector of TE<sub>11</sub>-mode (fundamental mode):



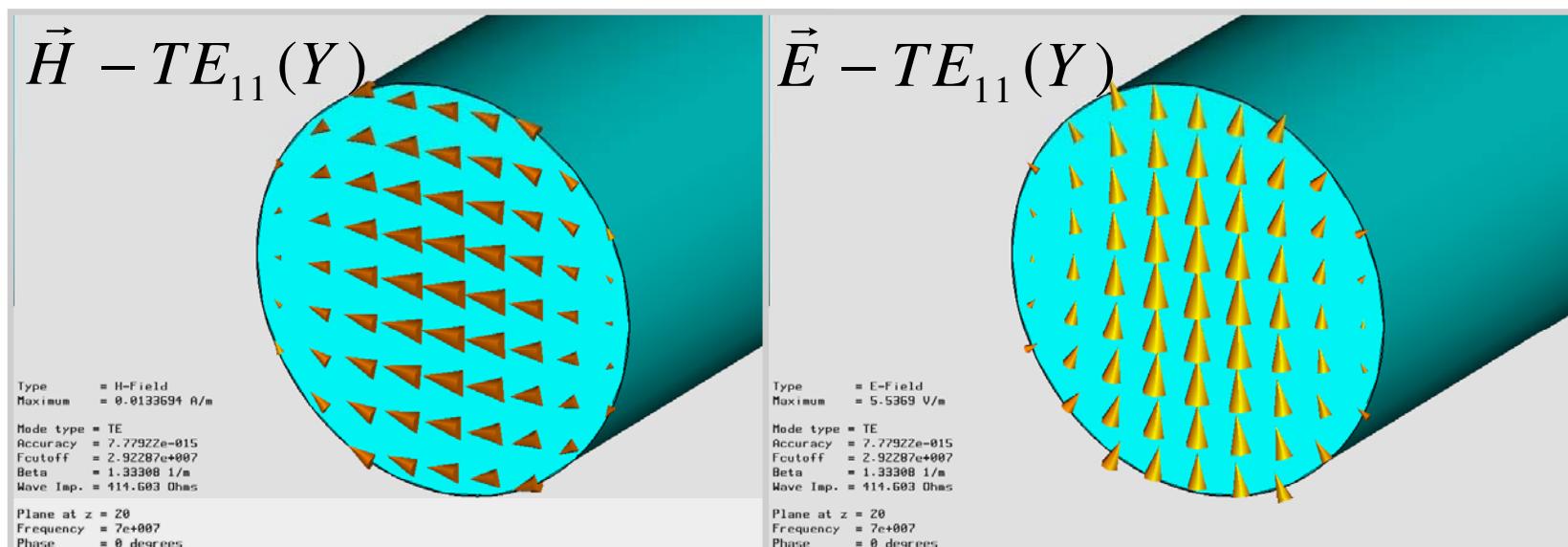
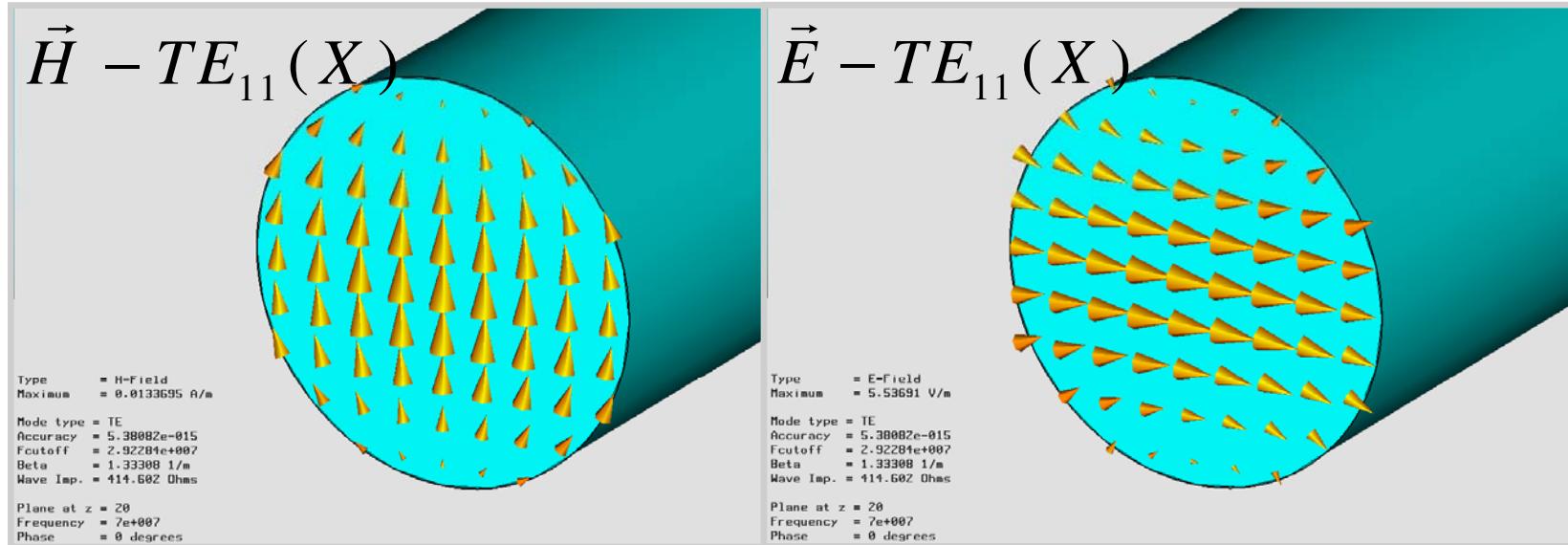
- Poynting vector of TE<sub>11</sub>-mode (fundamental mode):



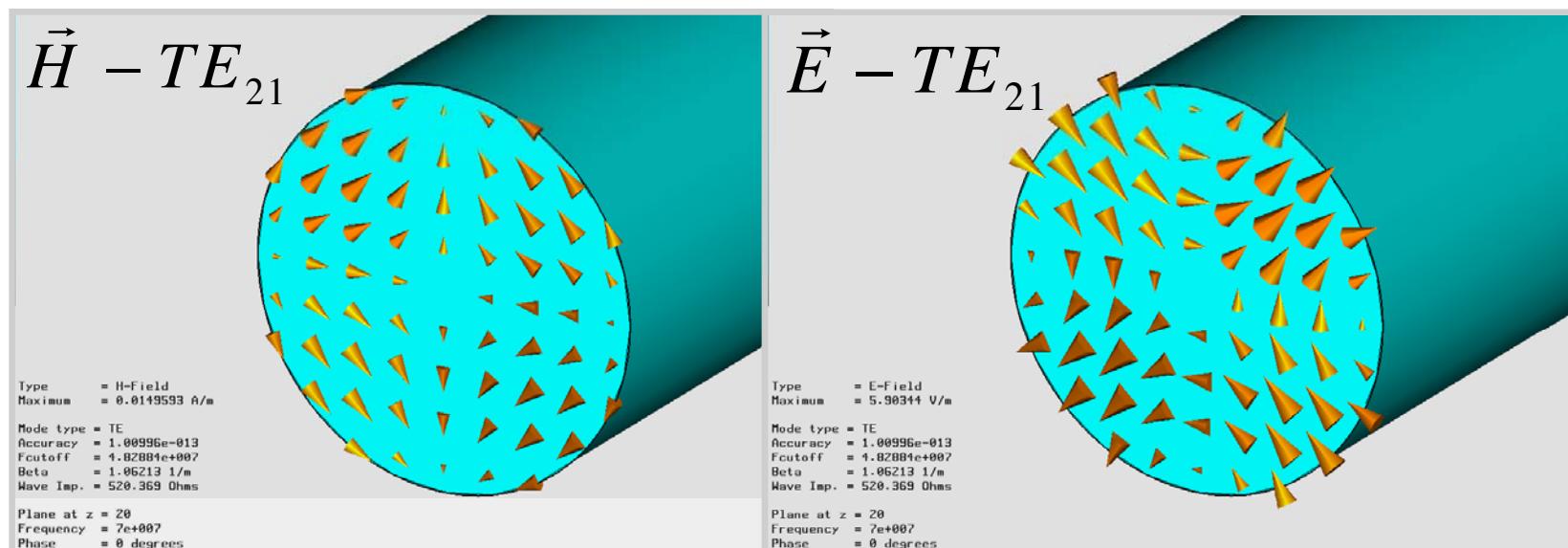
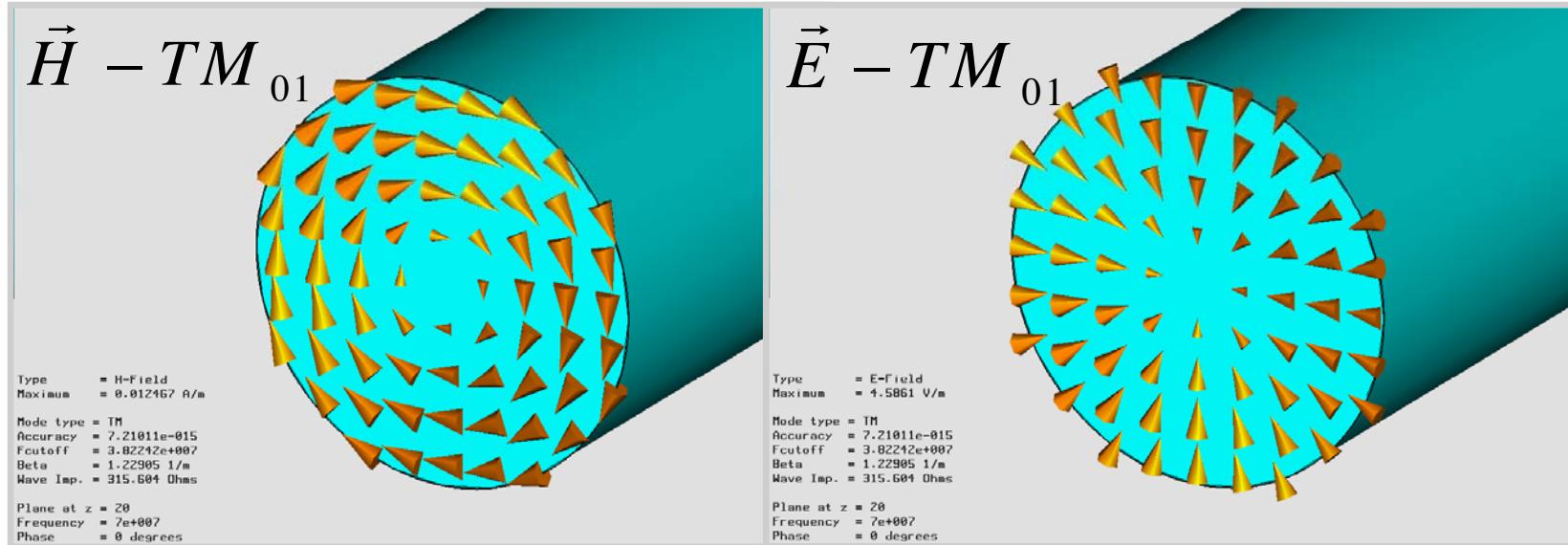
- Poynting vector of TE<sub>11</sub>-mode (fundamental mode):

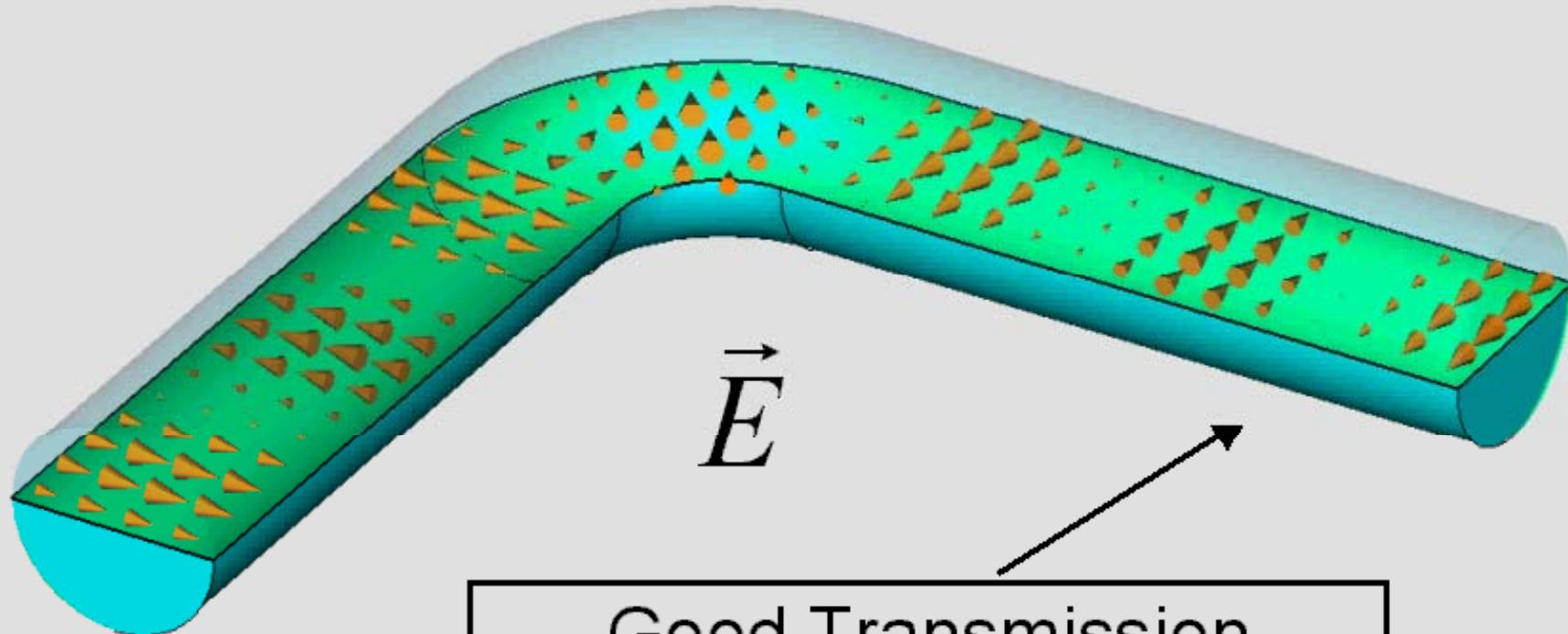


- Magnetic and electric Fields of  $TE_{11}$ -mode:



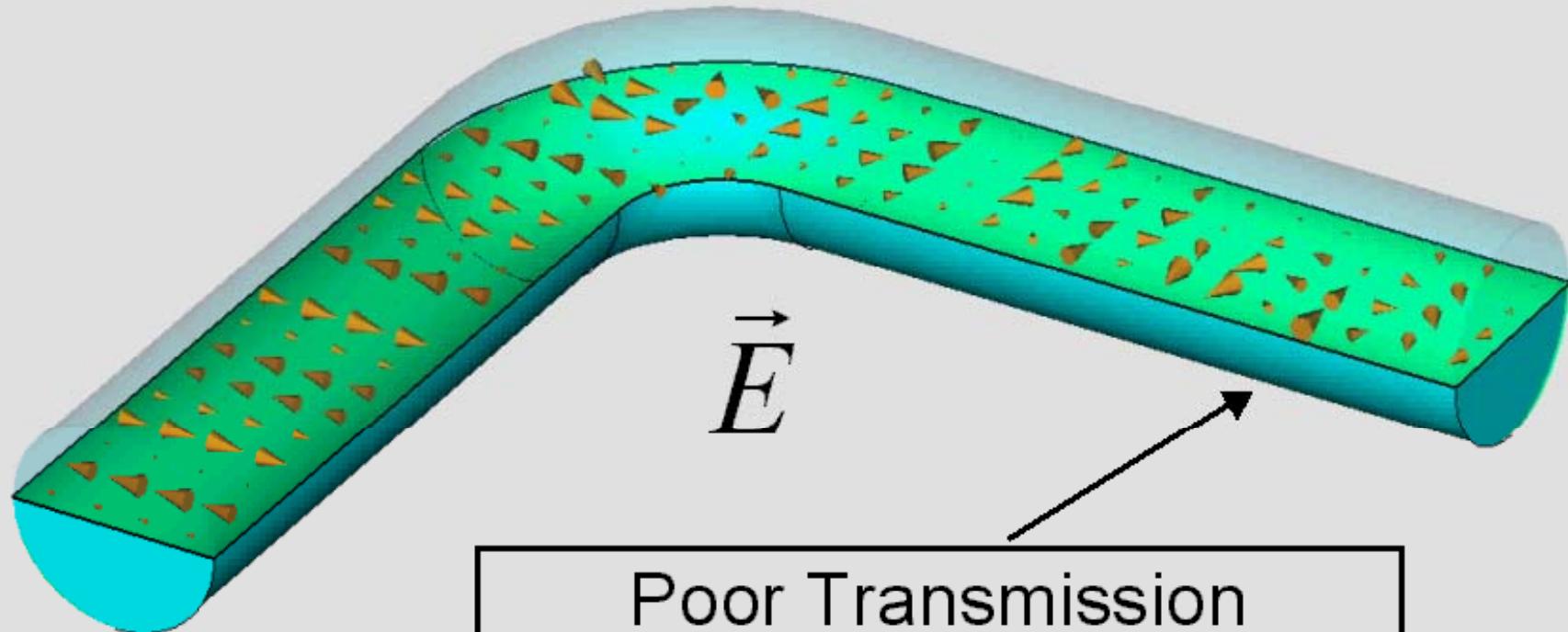
- Magnetic and electric Fields of higher order modes :



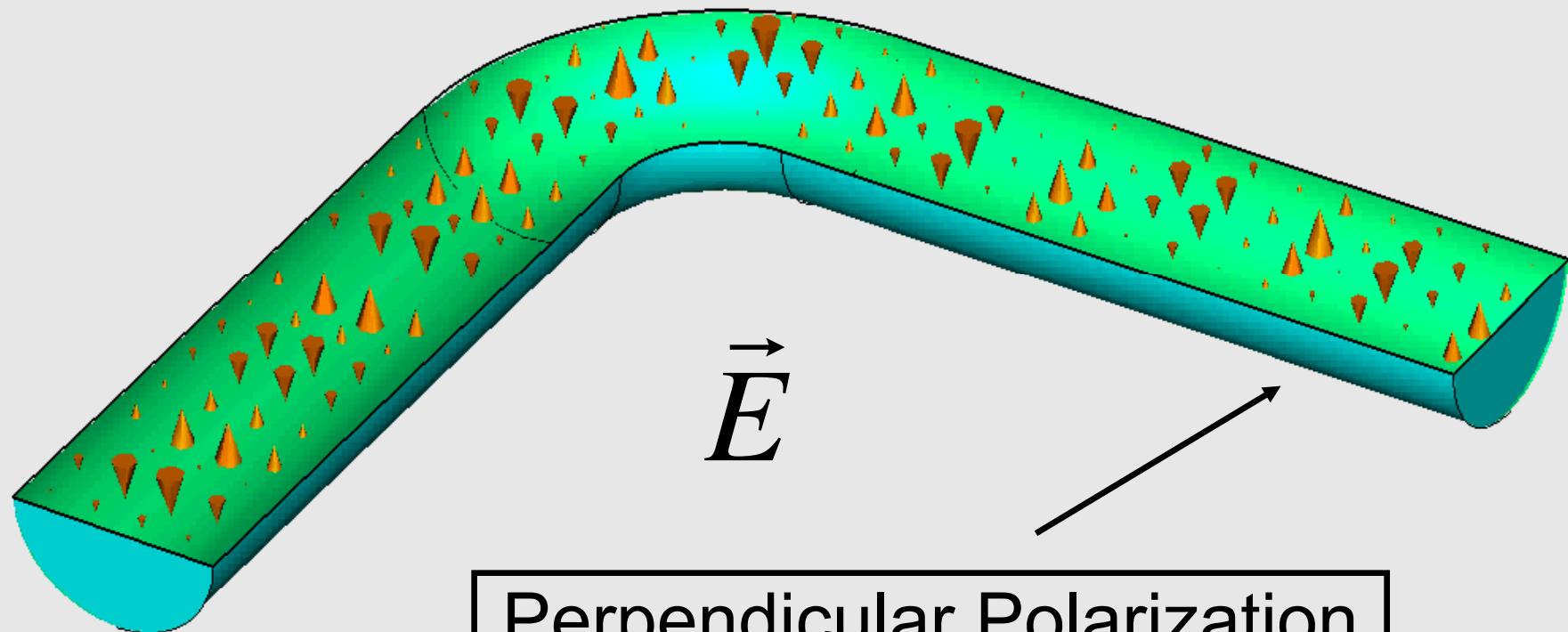


Good Transmission  
below cut-off of next mode

Type	= E-Field
Monitor	= e35
Maximum	= 10.2366 V/m
Plane at y	= 0
Frequency	= 3.5e+007
Phase	= 0 degrees

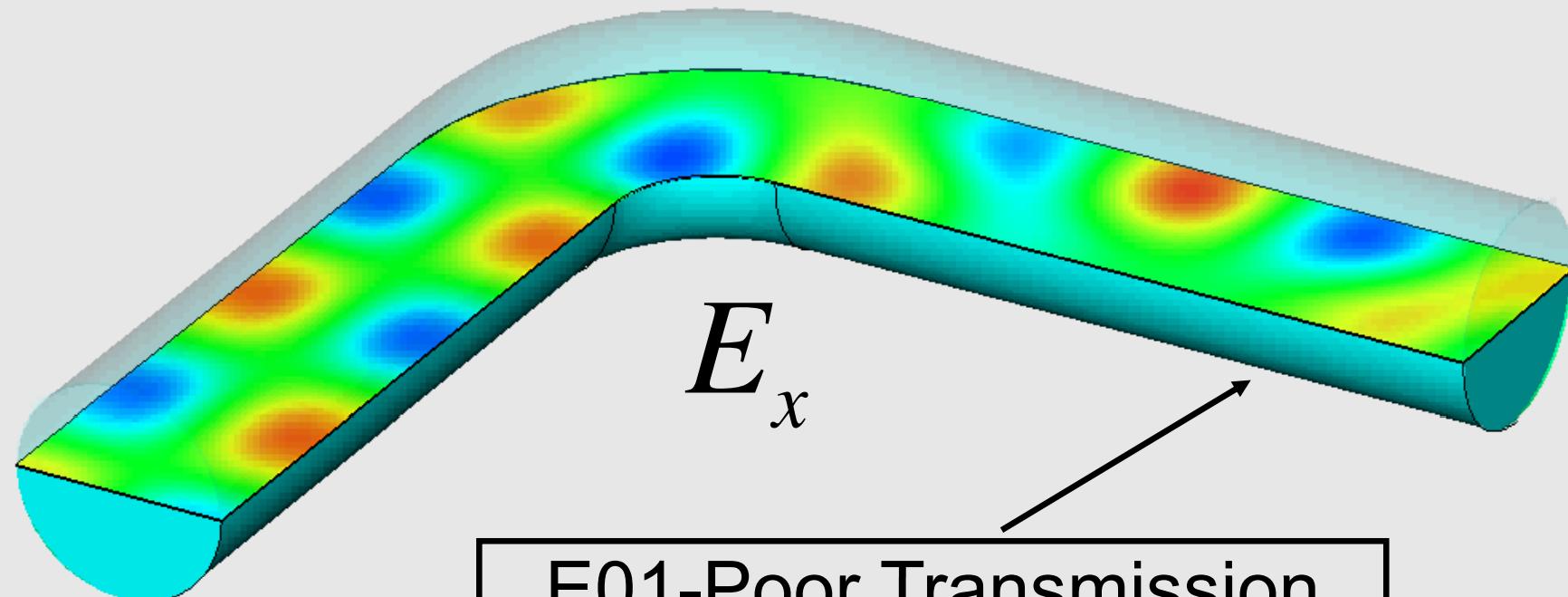


Type	= E-Field
Monitor	= e55
Maximum	= 7.77962 V/m
Plane at y	= 0
Frequency	= 5.5e+007
Phase	= 0 degrees



Perpendicular Polarization  
Better Transmission  
above cut-off of next mode

Type	= E-Field
Monitor	= e55
Maximum	= 6.49038 V/m
Plane at y = 0	
Frequency	= 5.5e+007
Phase	= 0 degrees



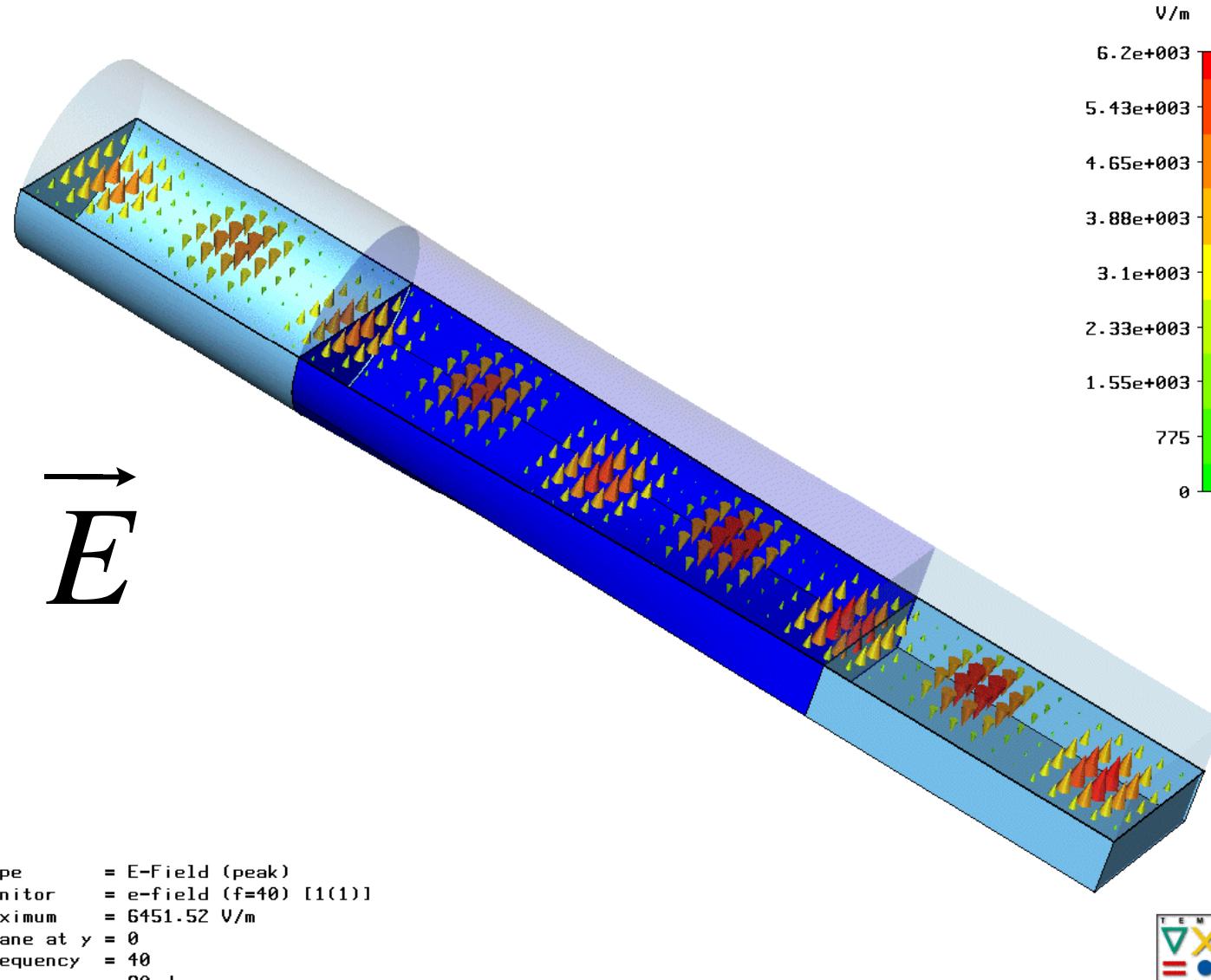
E01-Poor Transmission  
above cut-off of first mode

Type = H-Field  
Monitor = h45  
Component = y  
Plane at y = 0  
Frequency = 4.5e+007  
Phase = 0 degrees

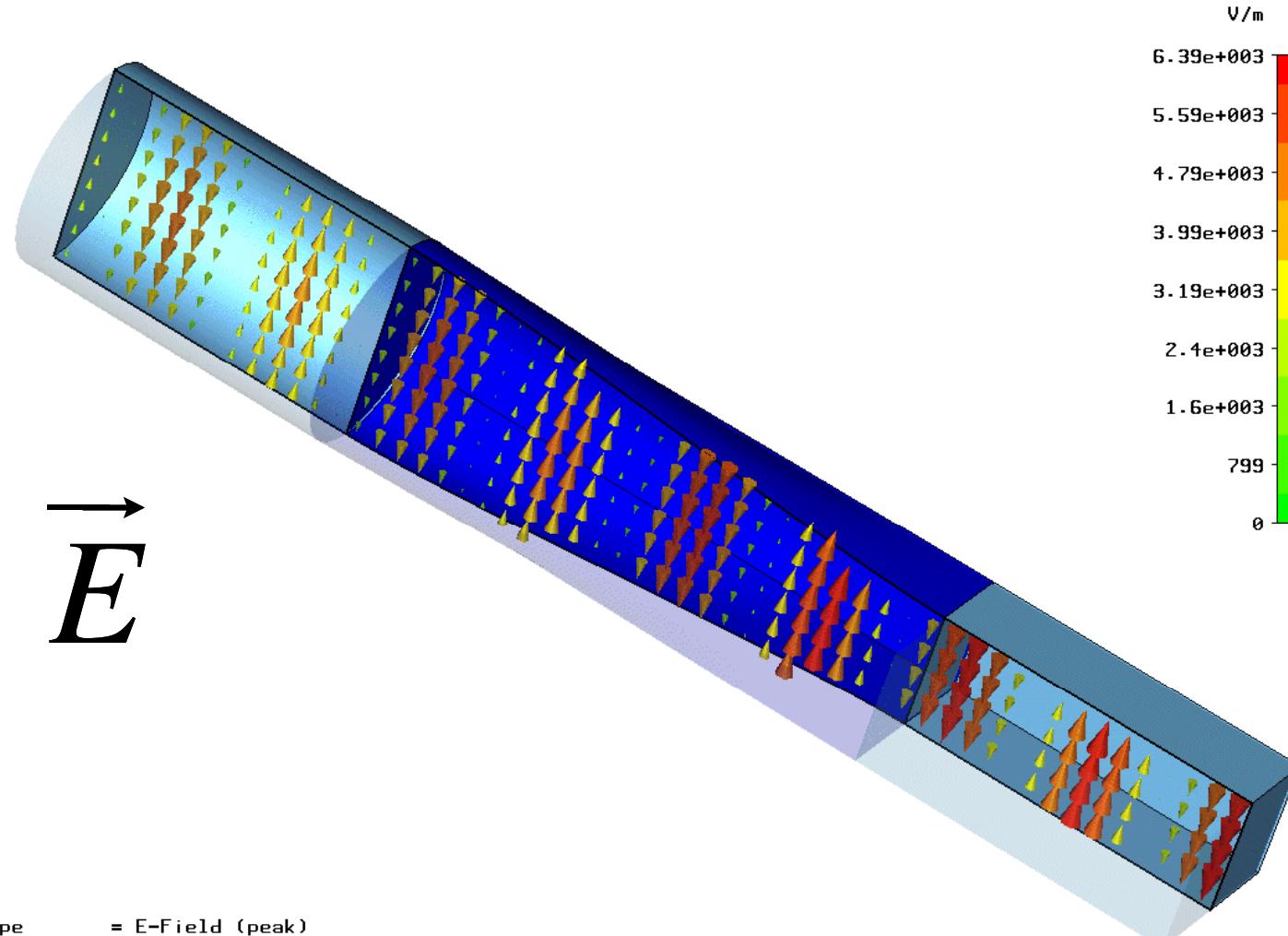


## Rectangular-Circular Transition

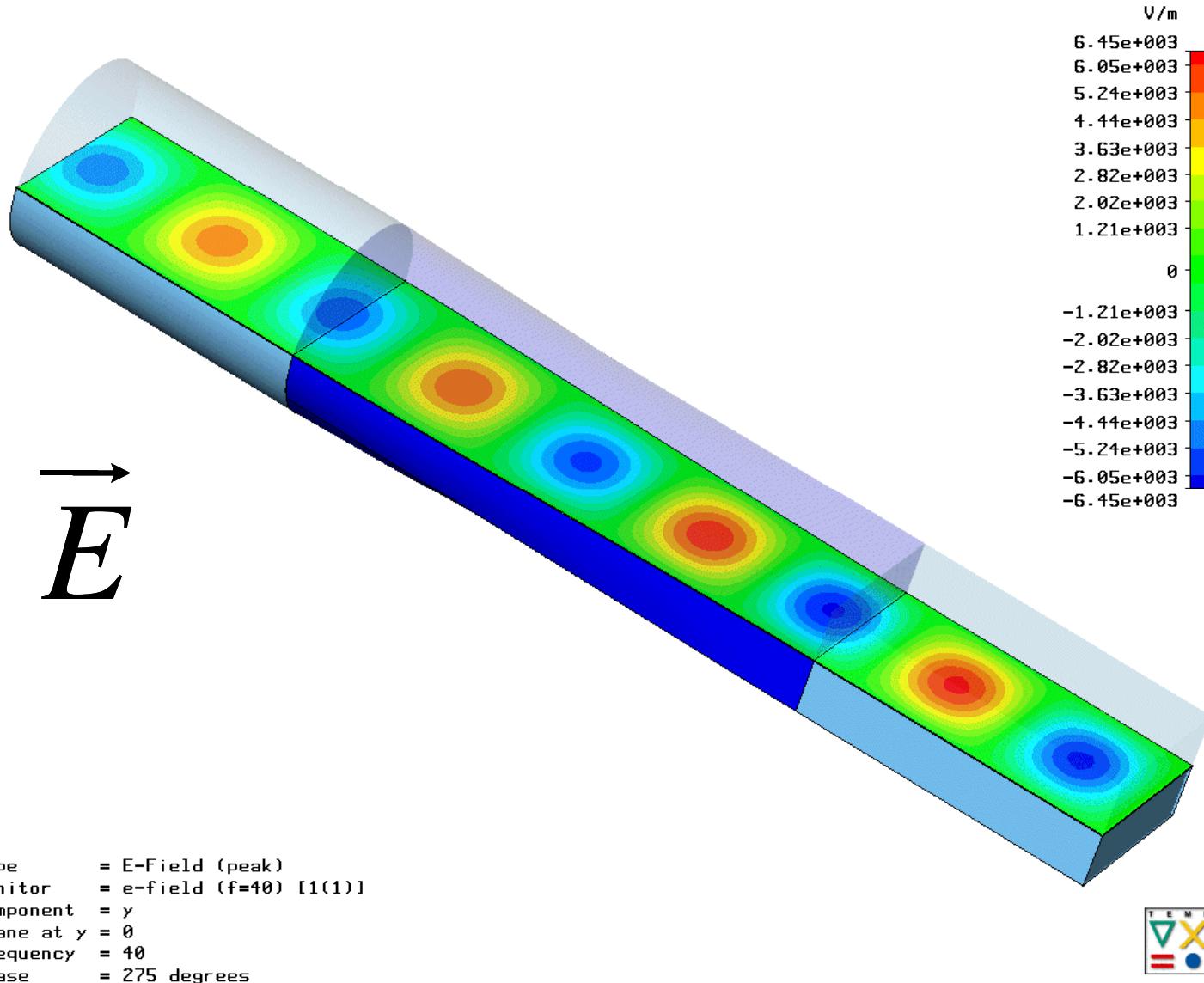
- Transition from rectangular to circular waveguides:



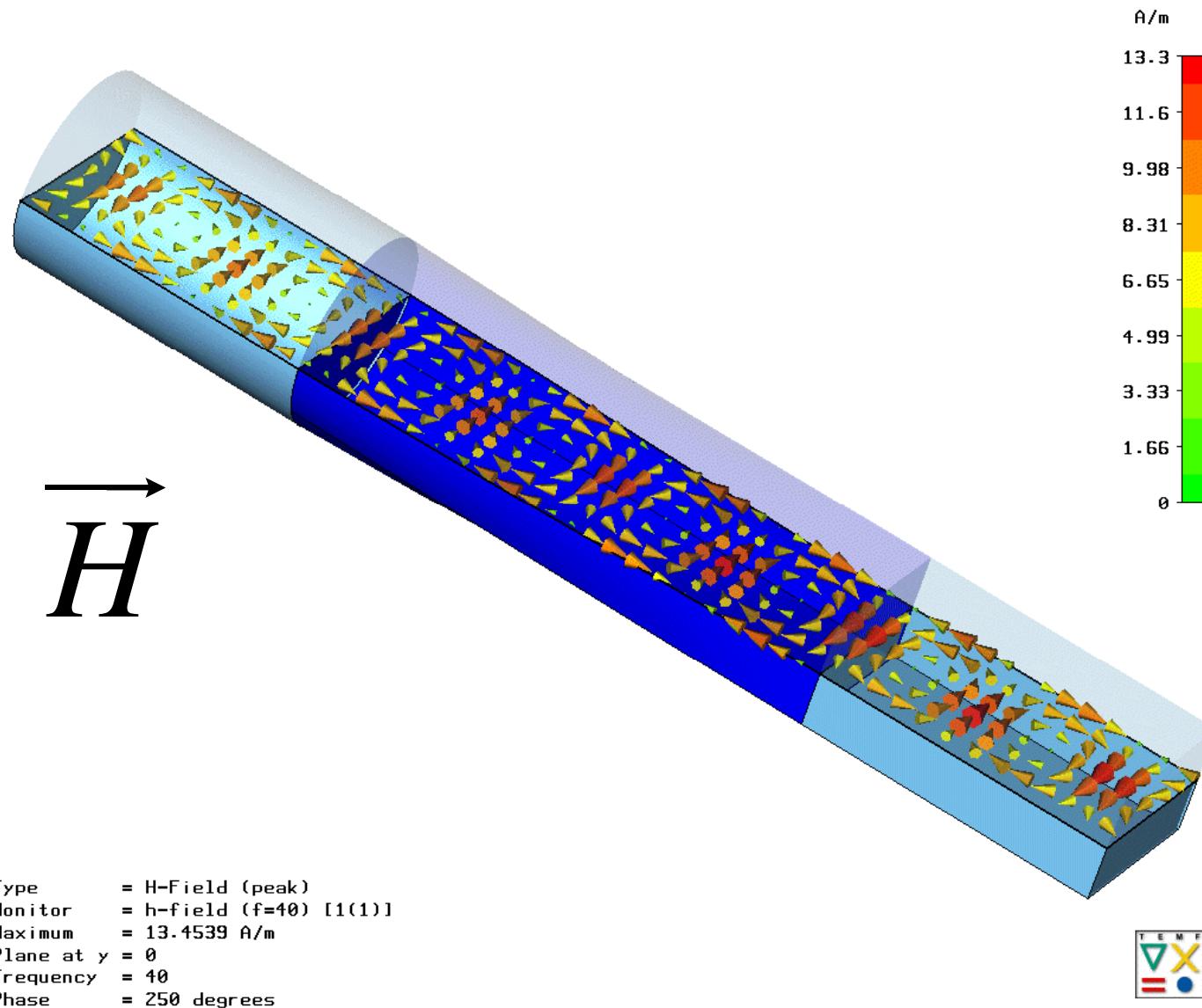
- Transition from rectangular to circular waveguides:



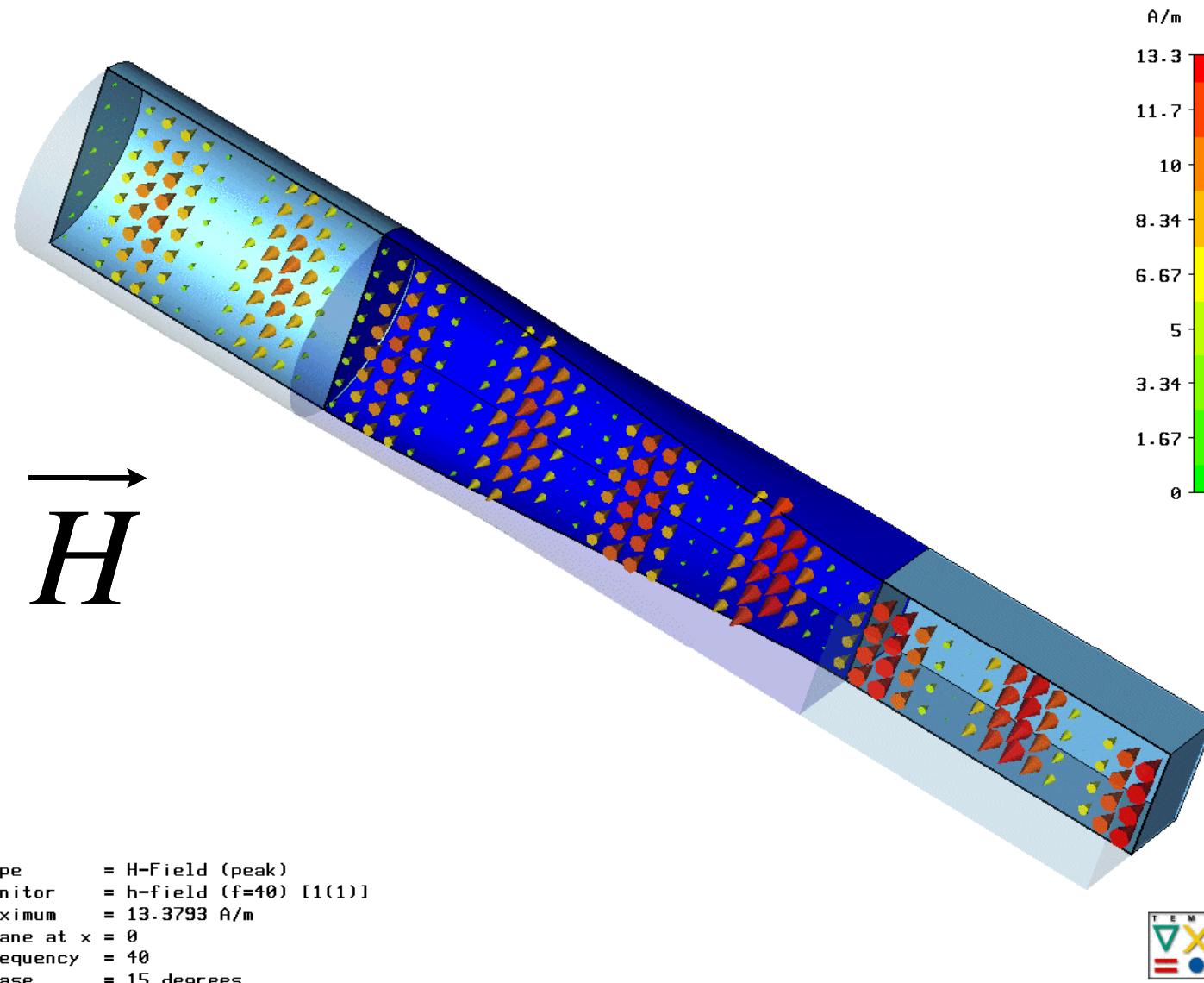
- Transition from rectangular to circular waveguides:



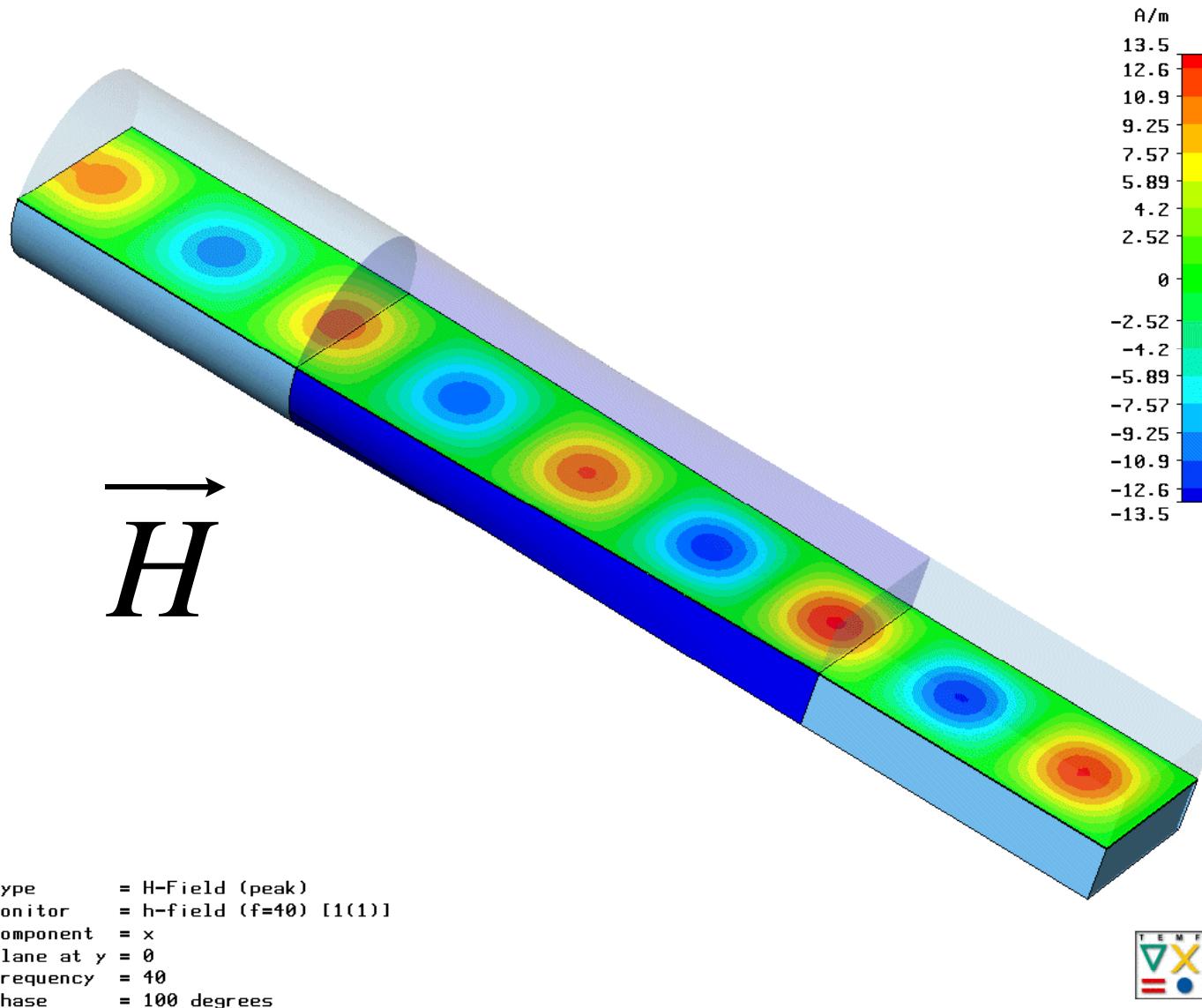
- Transition from rectangular to circular waveguides:



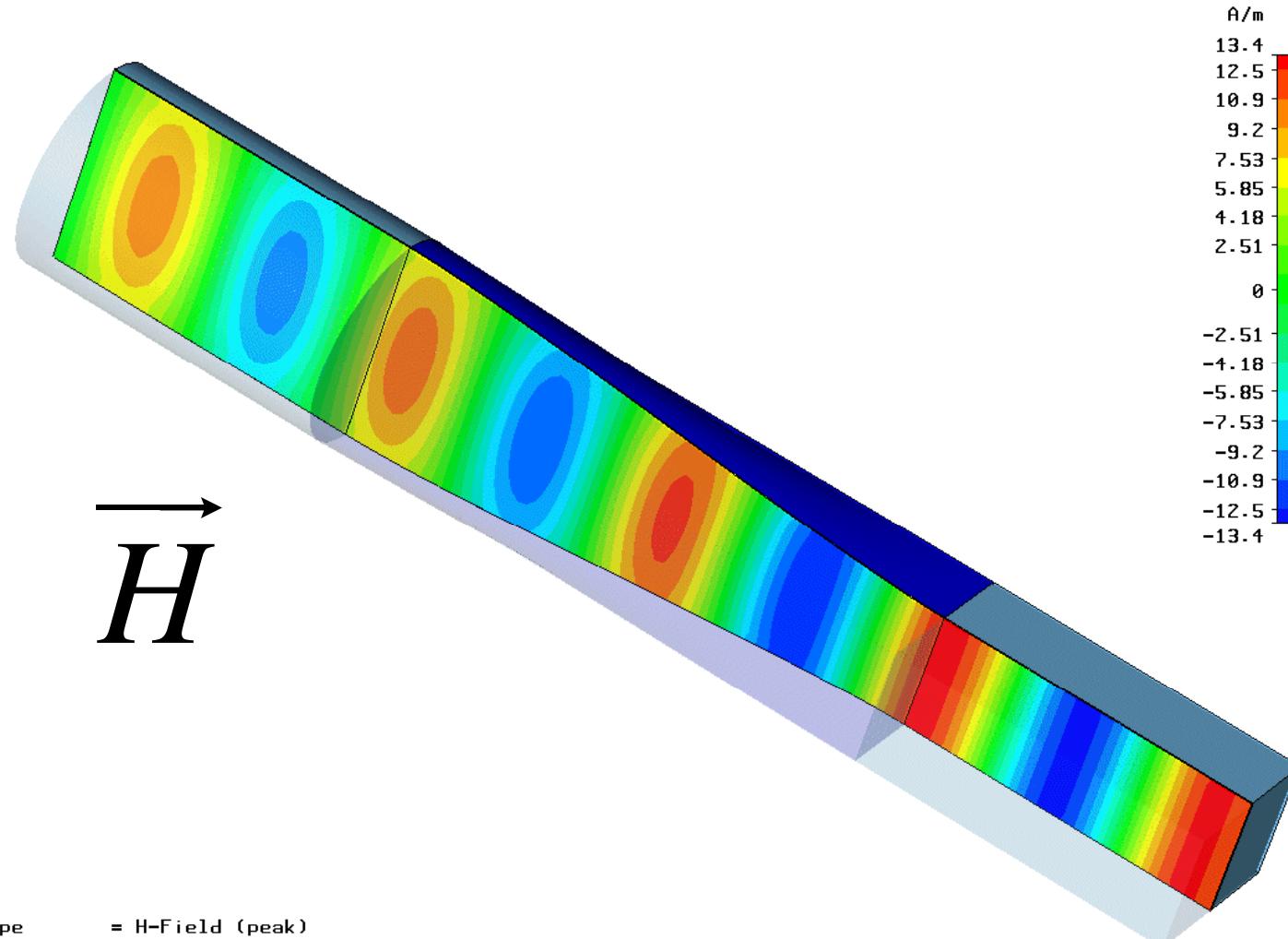
- Transition from rectangular to circular waveguides:



- Transition from rectangular to circular waveguides:



- Transition from rectangular to circular waveguides:

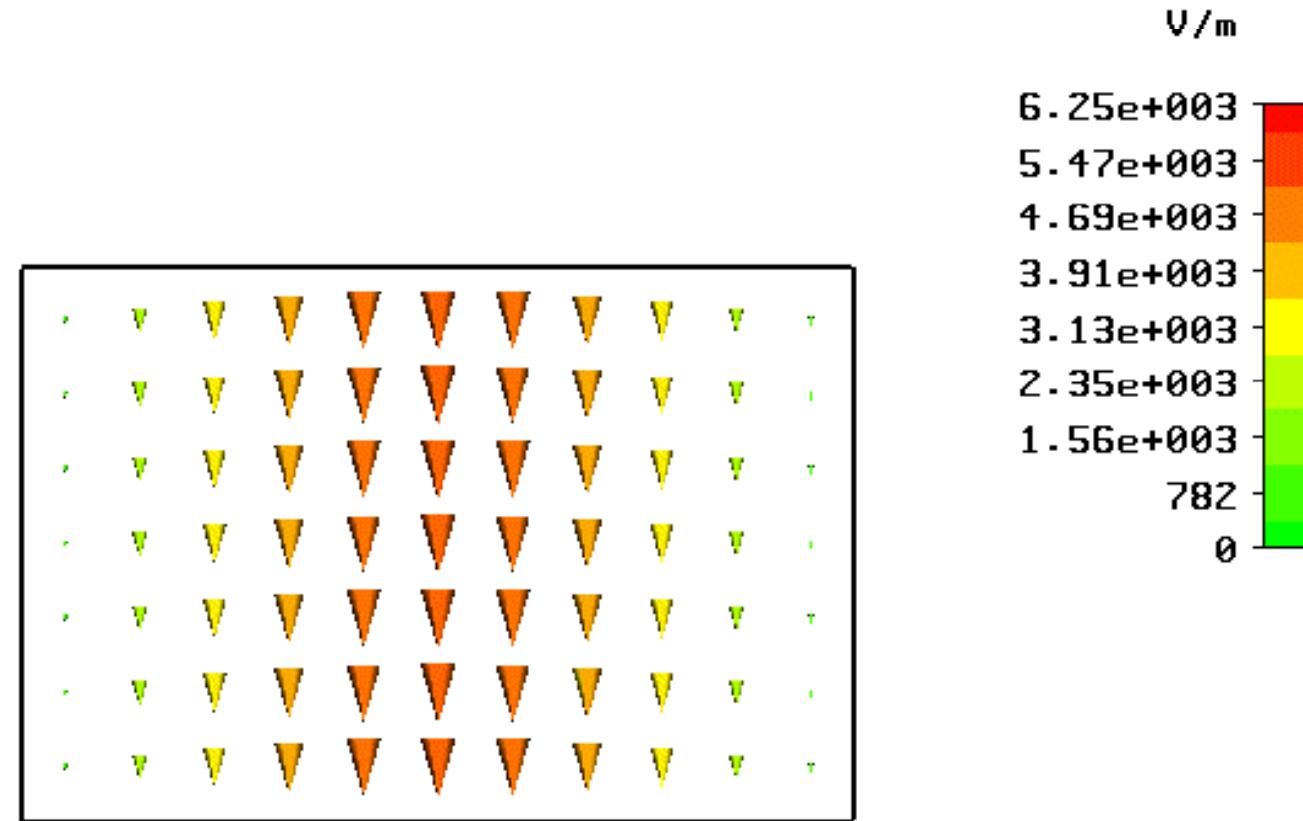


```
Type      = H-Field (peak)
Monitor   = h-field (f=40) [1(1)]
Component = x
Plane at x = 0
Frequency = 40
Phase     = 40 degrees
```



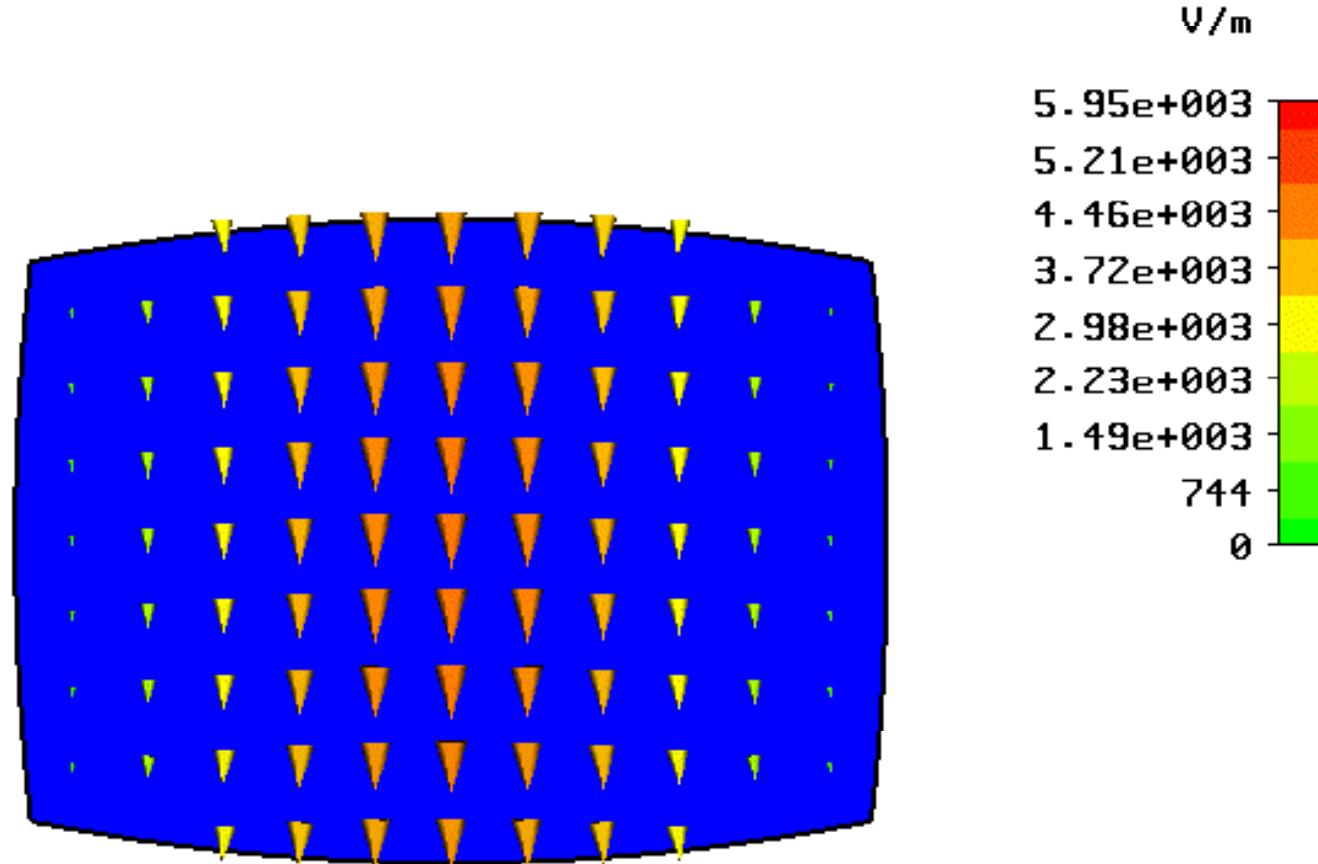
- Transition from rectangular to circular waveguides:

$\vec{E}$



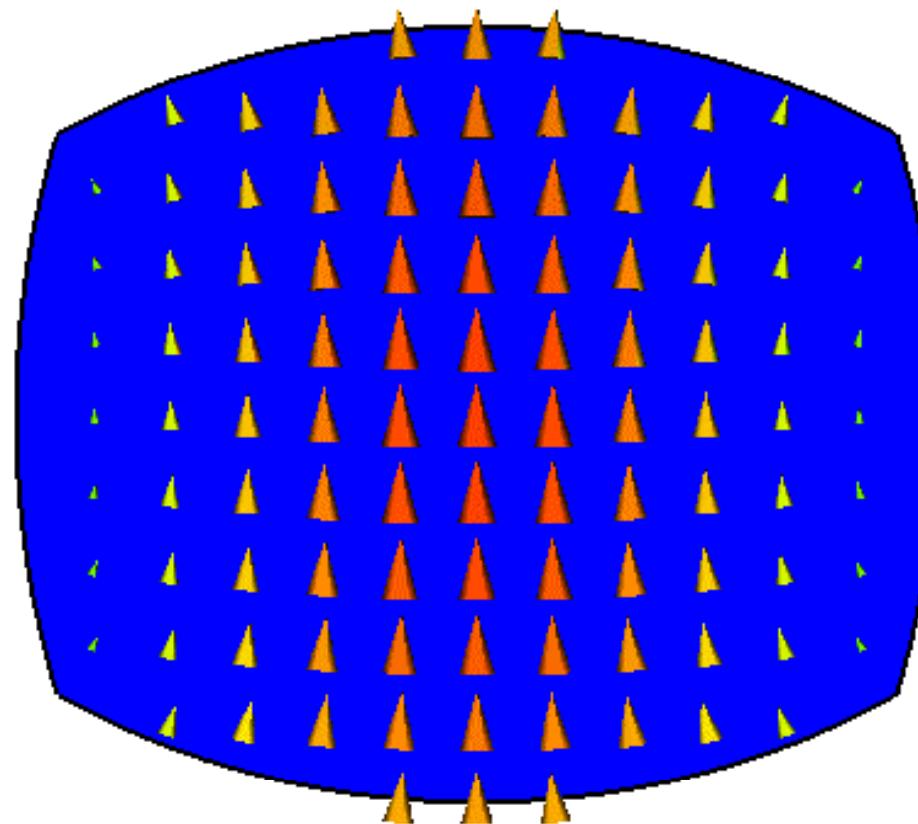
- Transition from rectangular to circular waveguides:

$\vec{E}$

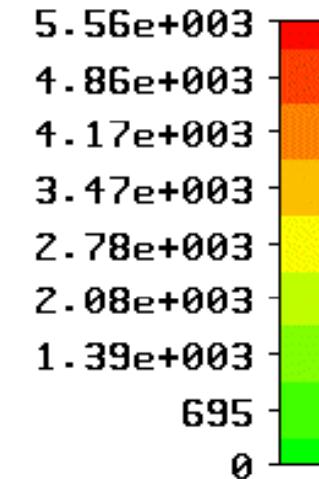


- Transition from rectangular to circular waveguides:

$$\vec{E}$$

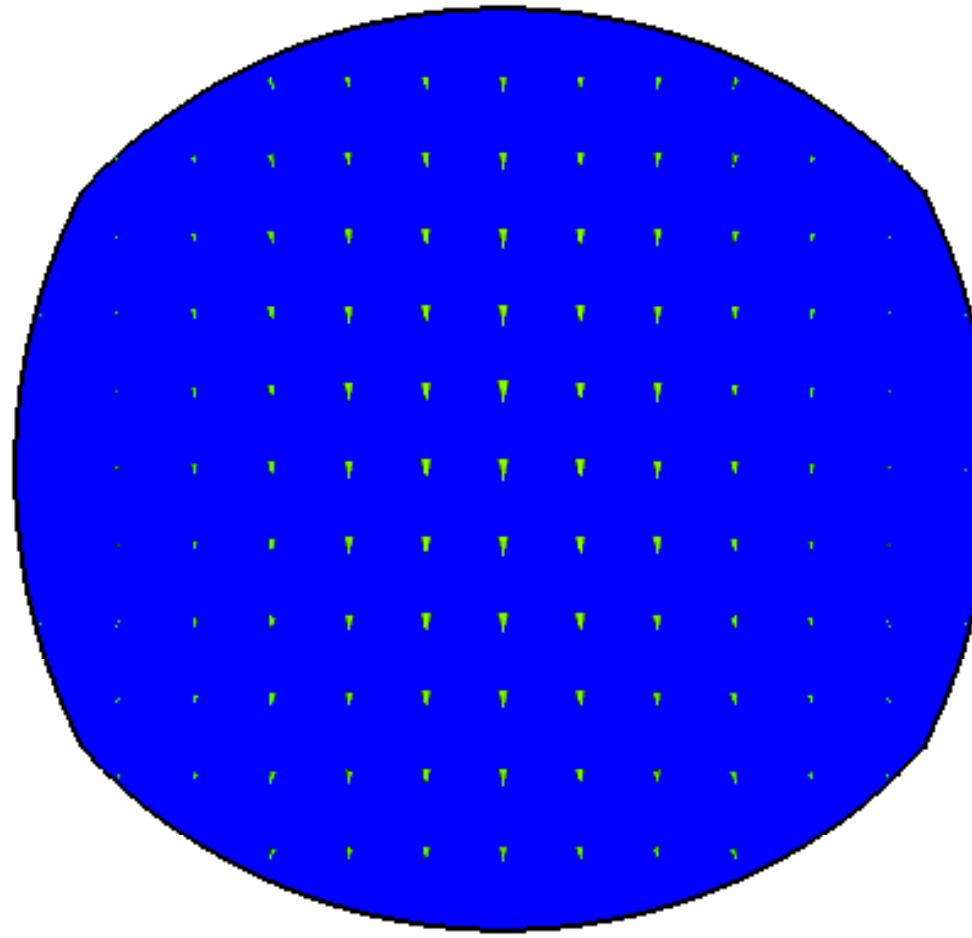


$V/m$

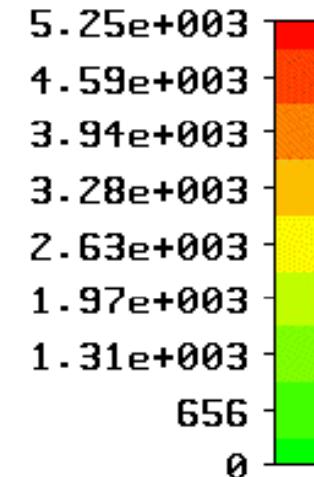


- Transition from rectangular to circular waveguides:

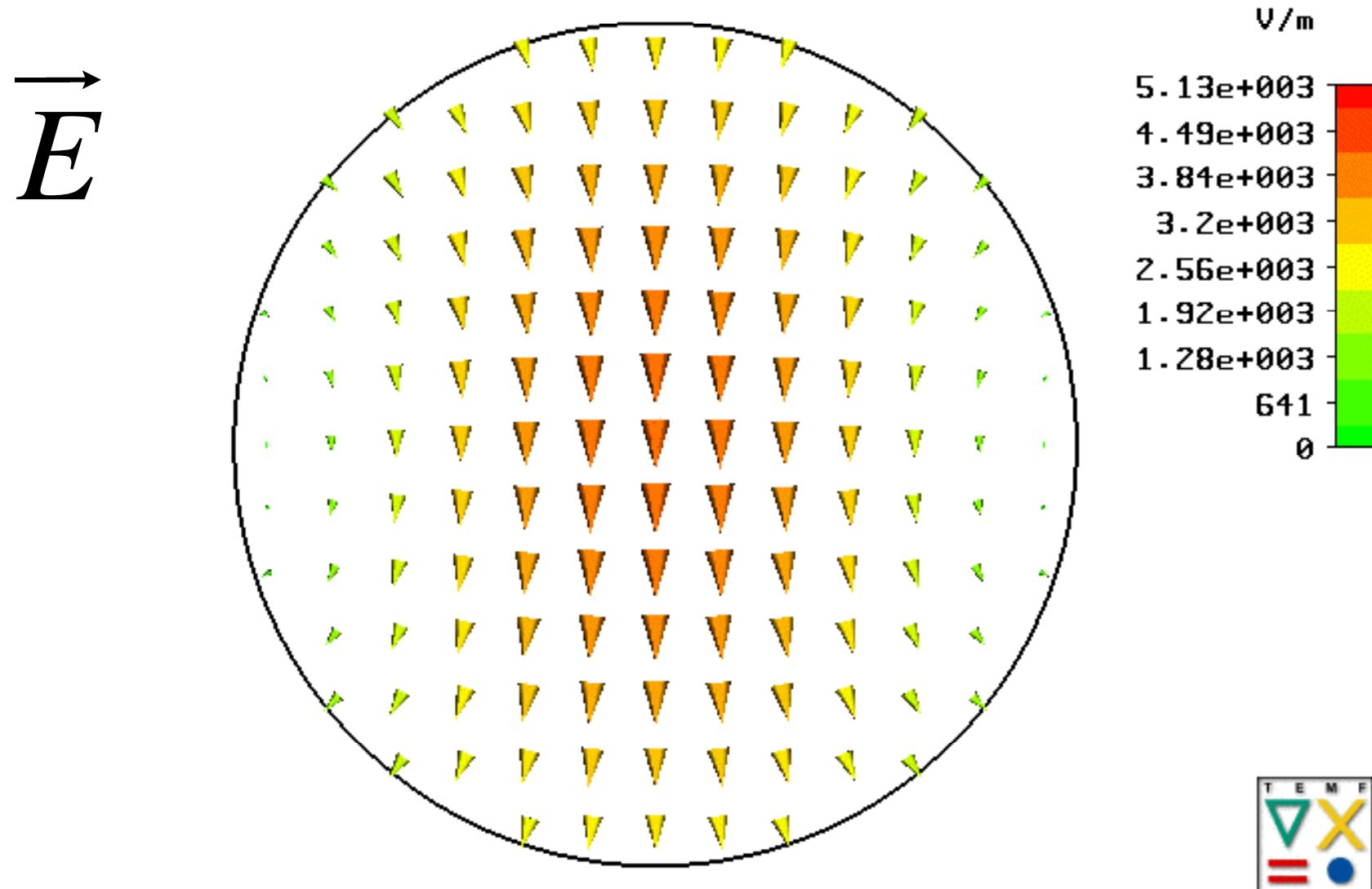
$$\vec{E}$$



V/m



- Transition from rectangular to circular waveguides:



## Coaxial Waveguides

- General solution: combination of Bessel and Neumann functions (field is 0 on the z axis, therefore the Neumann function can also be present in the solution)
- General solution:

$$Z_m(Kr) = A J_m(Kr) + B N_m(Kr).$$

- Practical application: modes with cutoff frequency 0

- **TEM waves (cutoff 0 Hz):**

$$\omega_c = K c = 0 \Rightarrow K = 0$$

☞ Additionally, no dependency on the angular coordinate

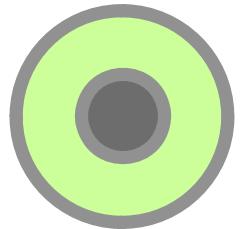
$$\mu = 0$$

☞ Equation after variable separation (see circular waveguide)

$$\frac{1}{f} \frac{\partial^2 f}{\partial r^2} + \frac{1}{fr} \frac{\partial f}{\partial r} + \frac{1}{r^2} \frac{1}{g} \frac{\partial^2 g}{\partial \varphi^2} + \frac{1}{h} \frac{\partial^2 h}{\partial z^2} + k^2 = 0$$

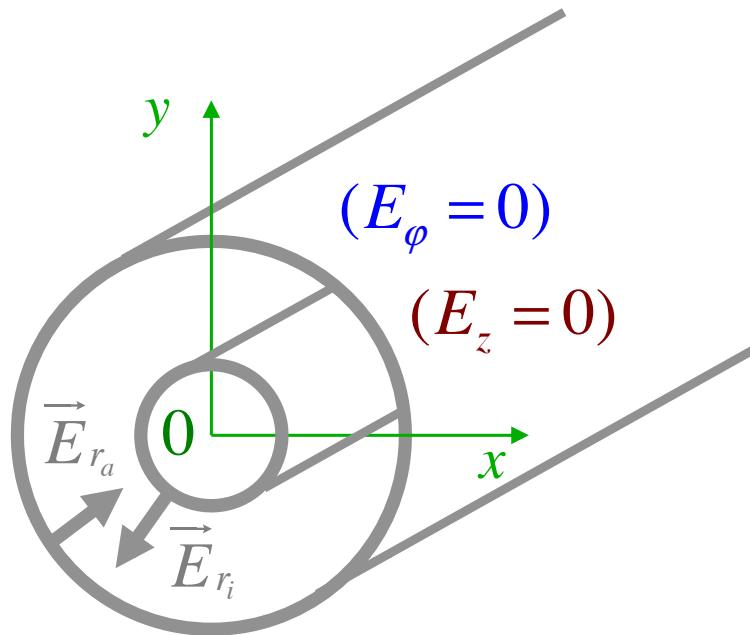
$$K^2 = \frac{1}{f} \frac{\partial^2 f}{\partial r^2} + \frac{1}{fr} \frac{\partial f}{\partial r} - \frac{1}{r^2} \mu^2 = 0$$

Coaxial WG



$$K = 0, \quad \mu = 0 \quad \Rightarrow \boxed{\frac{\partial^2 f}{\partial r^2} + \frac{1}{r} \frac{\partial f}{\partial r} = 0}$$

- TEM-solution (No component in propagation direction):



$K = 0, \mu = 0$  (Separation variables)

$$A_z = C_1 \ln \frac{r}{C_2} e^{-jkz} \quad (\text{Vector potential})$$

TEM: no Bessel-ODE

Field vector components:

Continuity:

$$E_\phi(r = r_i, \varphi, z) = 0$$

$$E_\phi(r = r_a, \varphi, z) = 0$$

$$E_z(r = r_i, \varphi, z) = 0$$

$$E_z(r = r_a, \varphi, z) = 0$$

$$\left. \begin{aligned} E_r &= \frac{1}{r} C_1 \cdot e^{-jkz} \\ H_\varphi &= \frac{1}{Zr} C_1 \cdot e^{-jkz} \end{aligned} \right\} \text{Static solution} \cdot e^{-jkz}$$

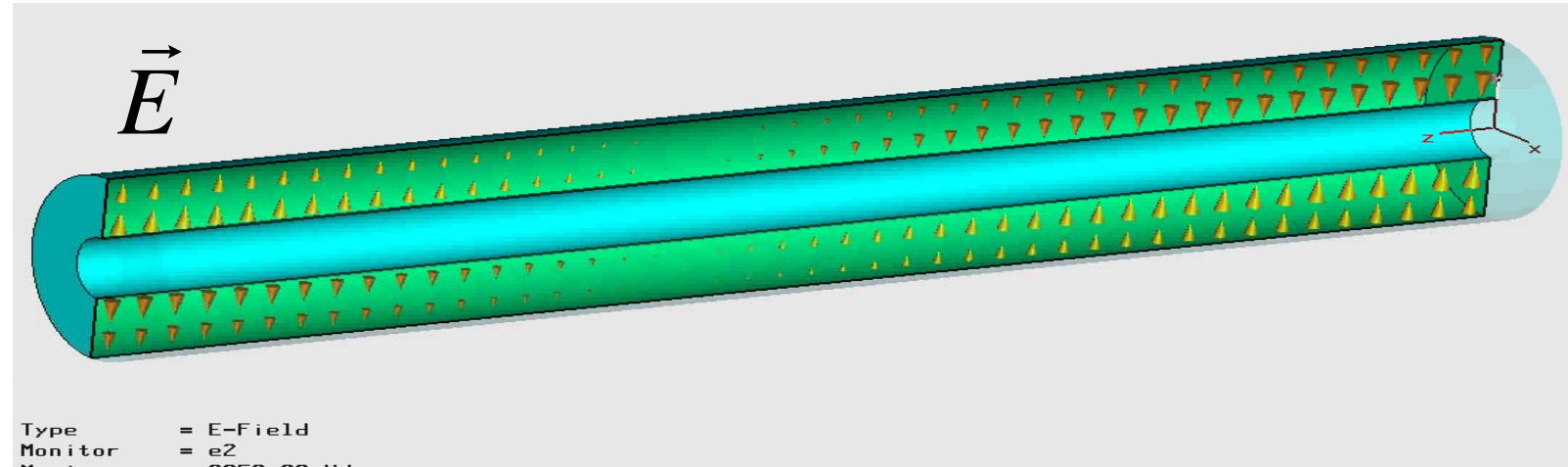
$$E_\varphi = E_z = H_r = H_z = 0$$

- TEM-Mode:

$$r_i = 1\text{ mm}, r_a = 3\text{ mm}$$

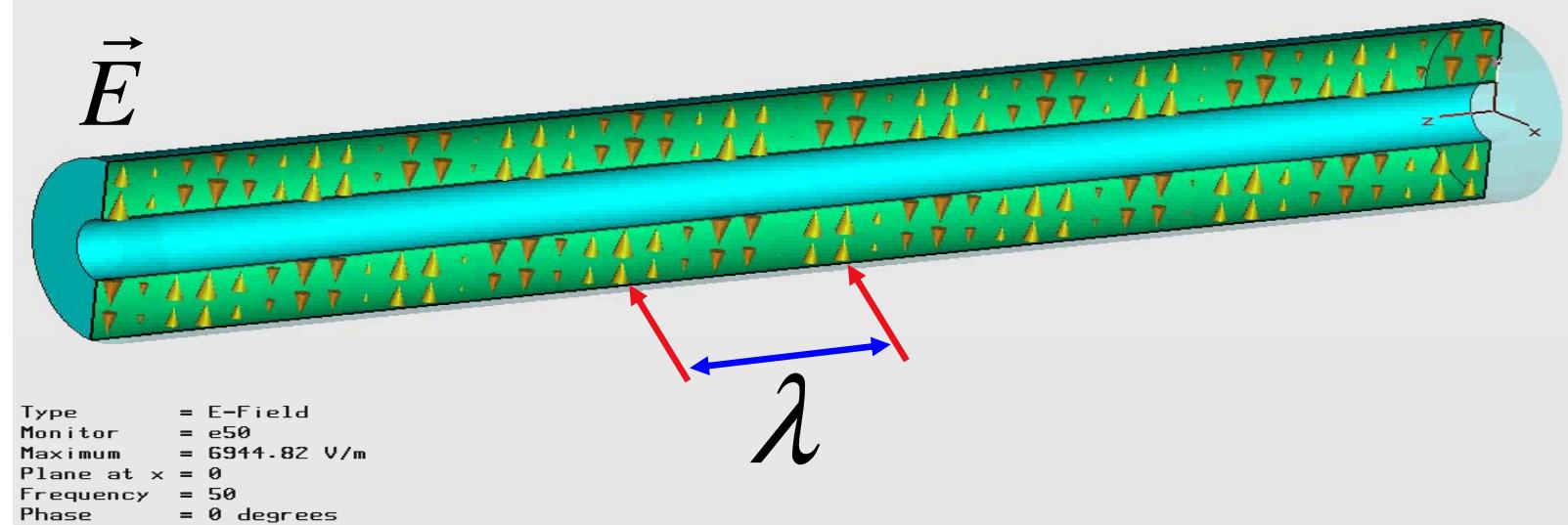
$$f = 2\text{ GHz}$$

$$\lambda = 150\text{ mm}$$



$$f = 50\text{ GHz}$$

$$\lambda = 6\text{ mm}$$

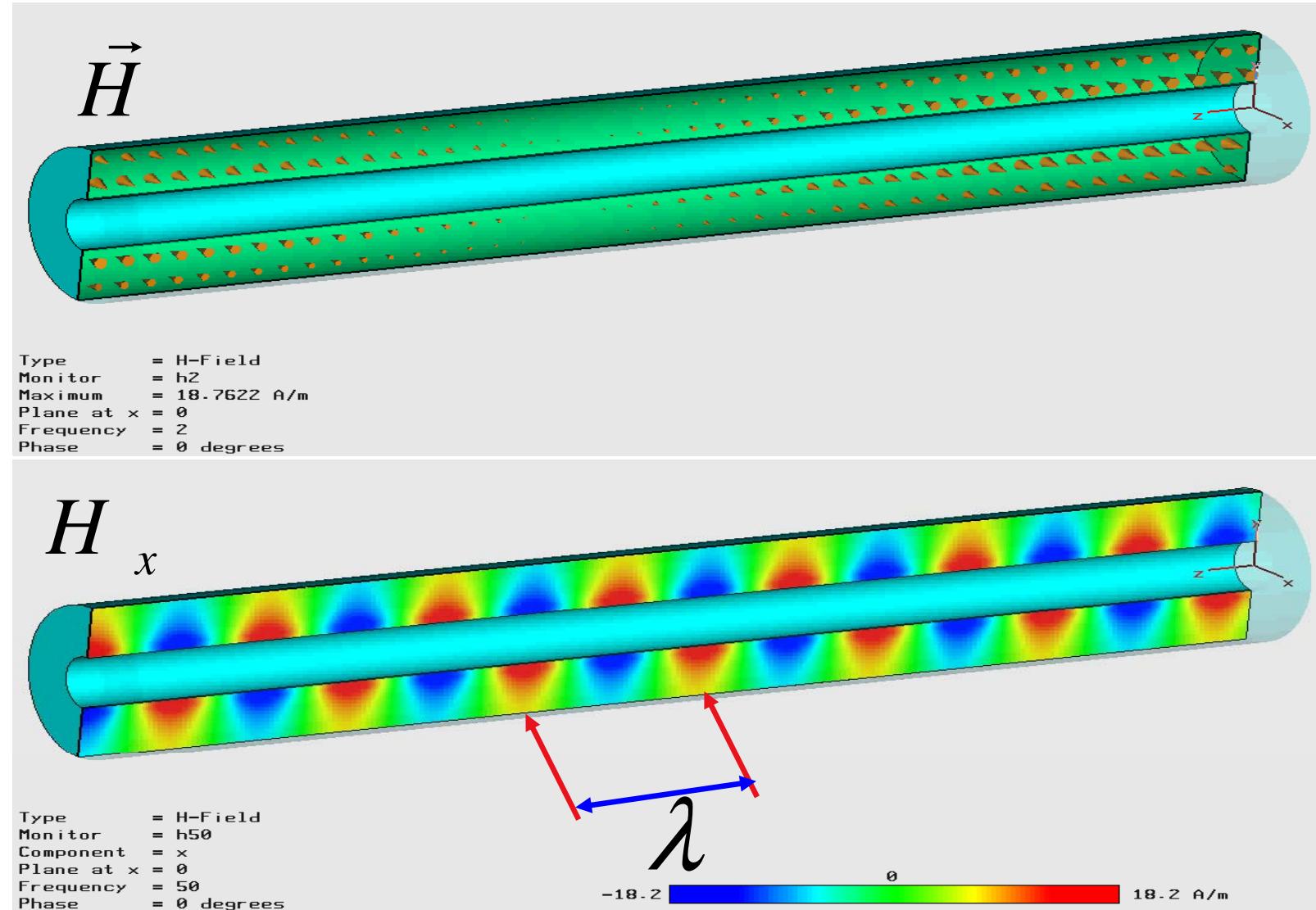


- TEM-Mode:

$$r_i = 1\text{ mm}, r_a = 3\text{ mm}$$

$$f = 2\text{ GHz}$$

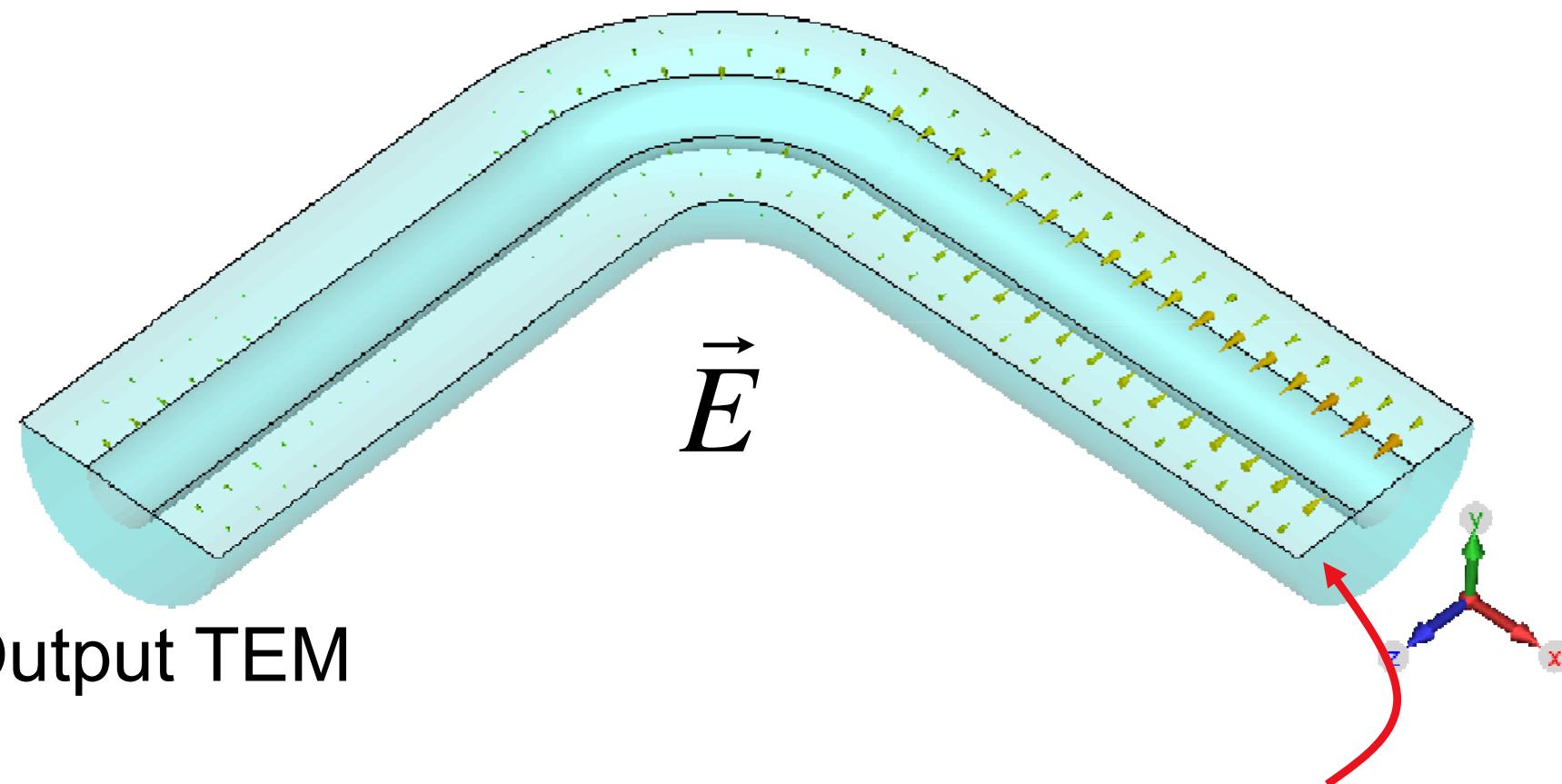
$$\lambda = 150\text{ mm}$$



$$f = 50\text{ GHz}$$

$$\lambda = 6\text{ mm}$$

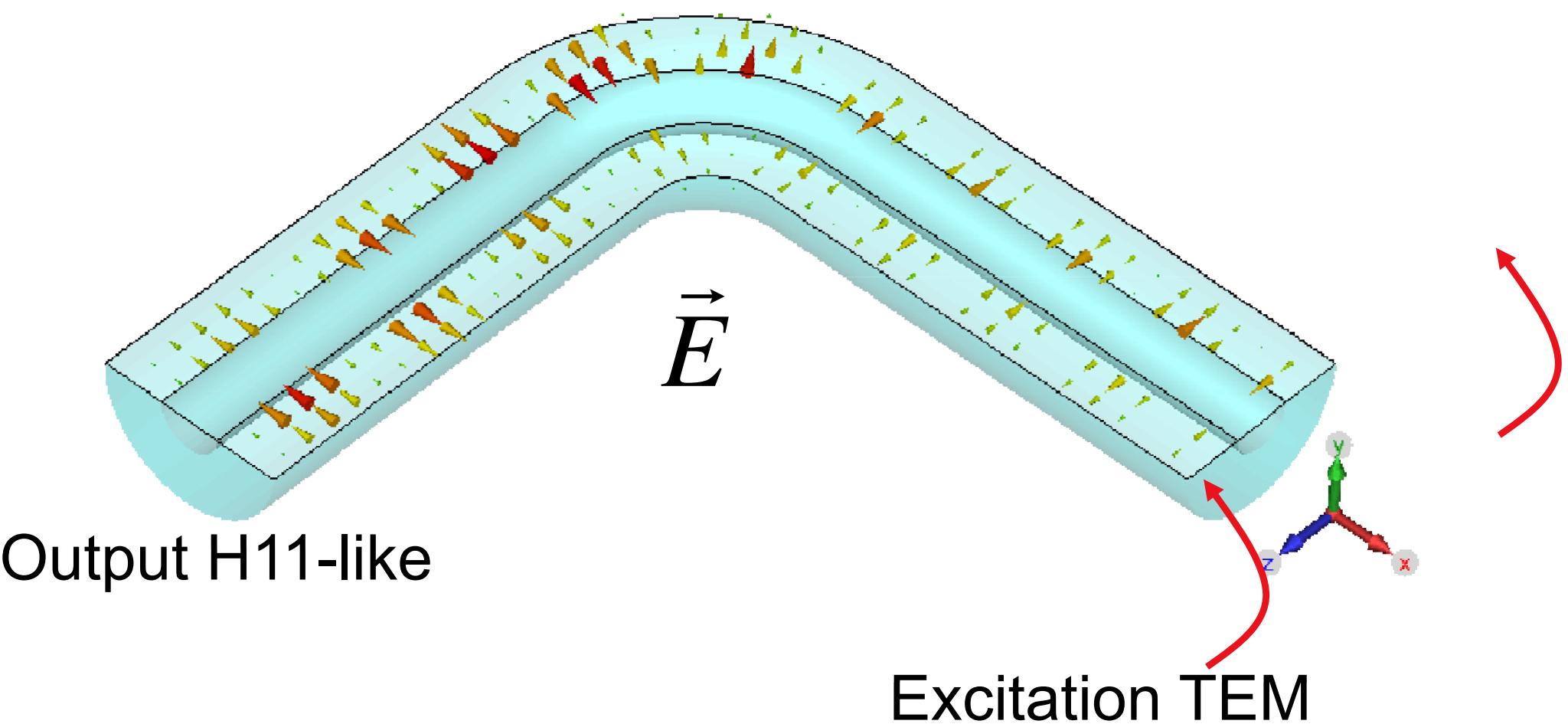
$f = 2 \text{ GHz}$



Output TEM

Excitation TEM

$f = 40 \text{ GHz}$



Excitation TEM

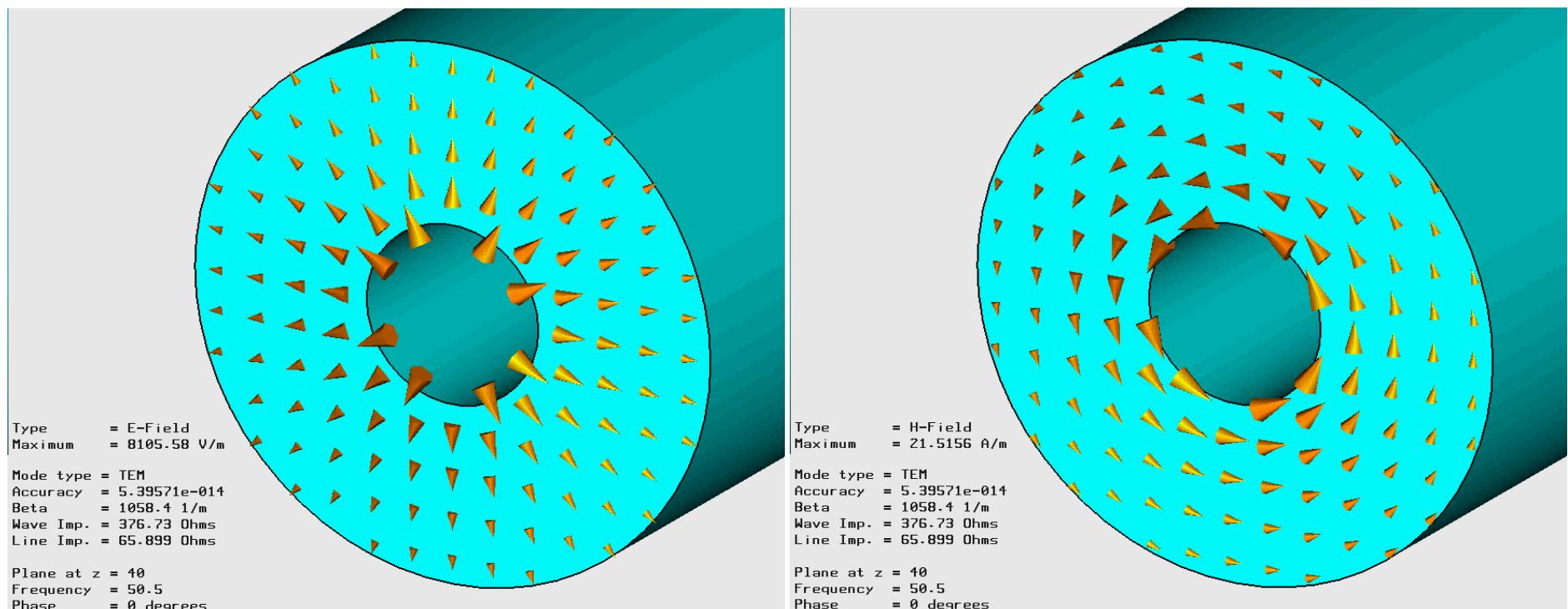
- TEM-Mode:

$$E_z = 0 \quad H_z = 0$$

$$f_{\text{cutoff}} = 0 \text{Hz}$$

$$\vec{E}$$

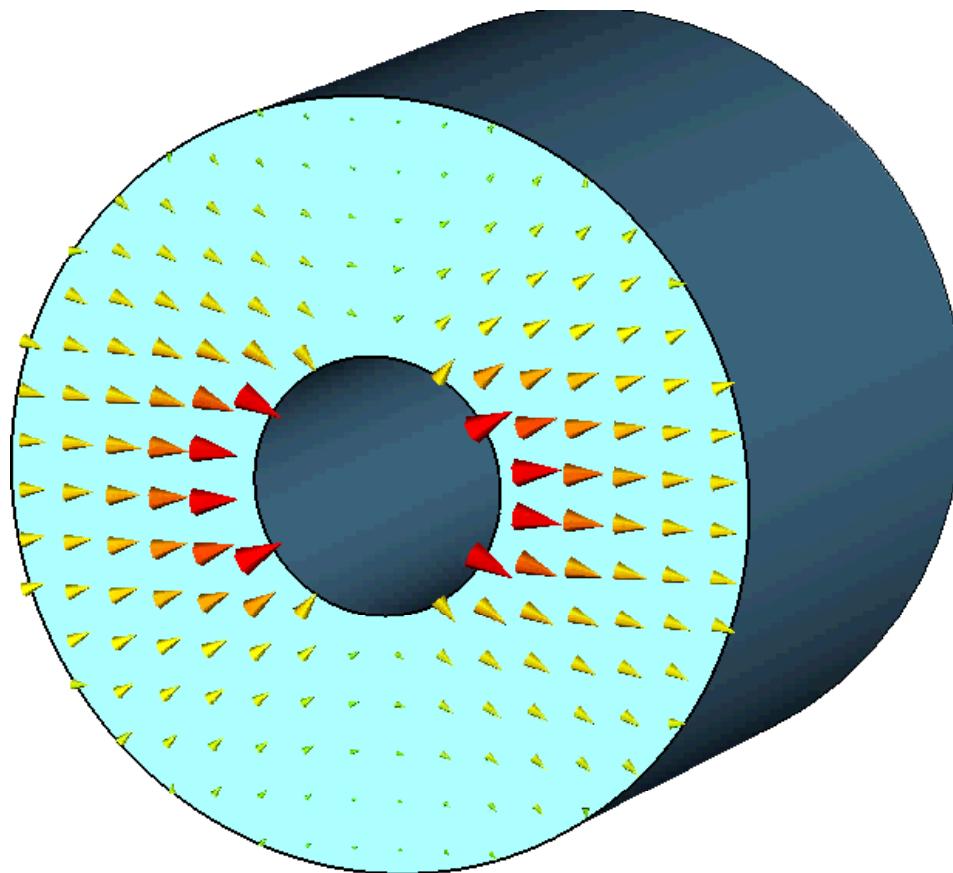
$$\vec{H}$$



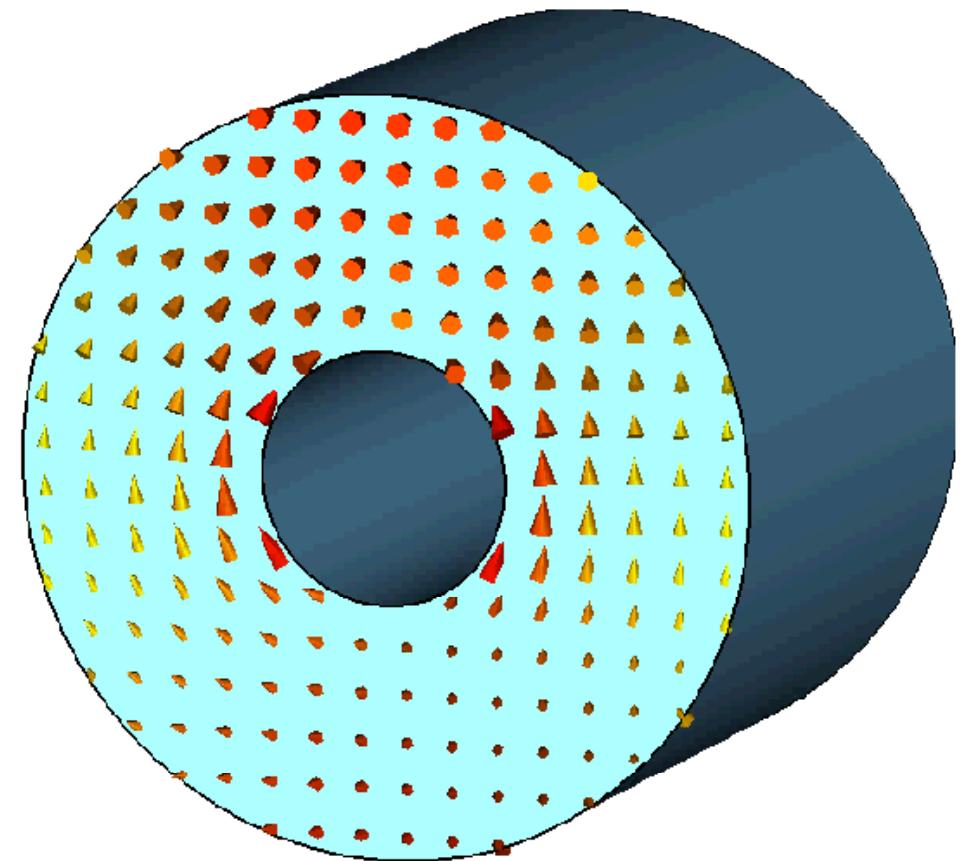
- Higher Coax-Modes:

$$r_i = 1\text{ mm}, r_a = 3\text{ mm}$$

$$f_{cutoff} = 24.5\text{ GHz}$$



$\vec{E}$  at  $\omega t = 0$

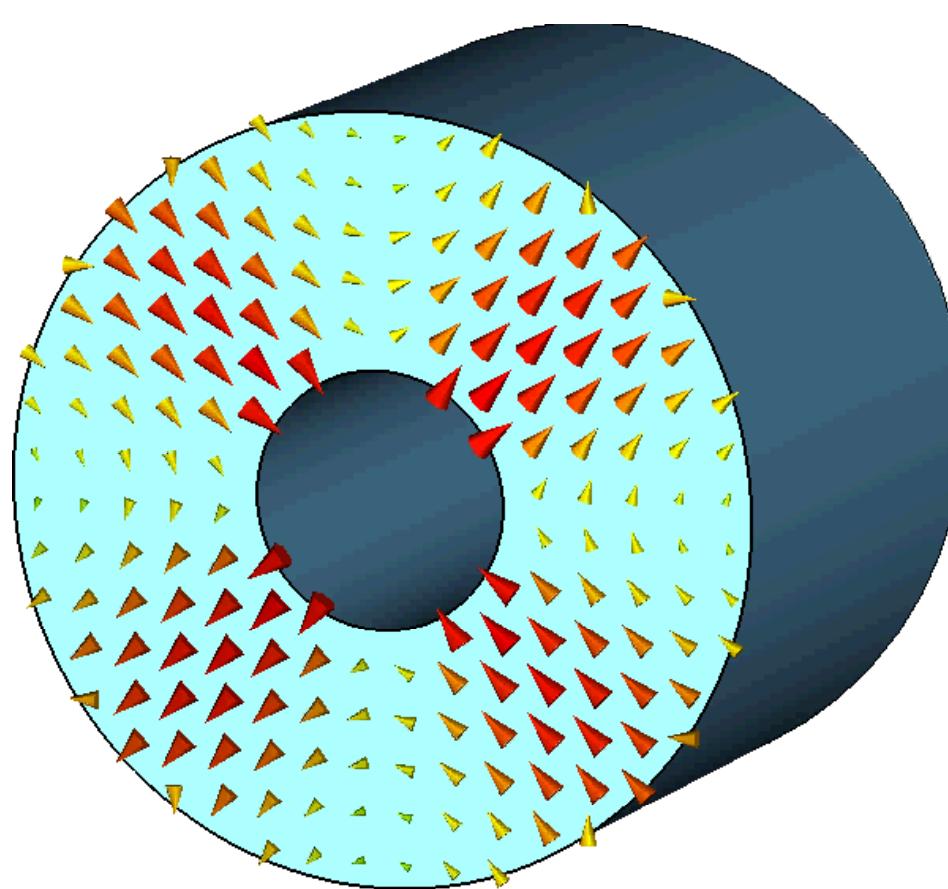


$\vec{H}$  at  $\omega t = \pi / 2$

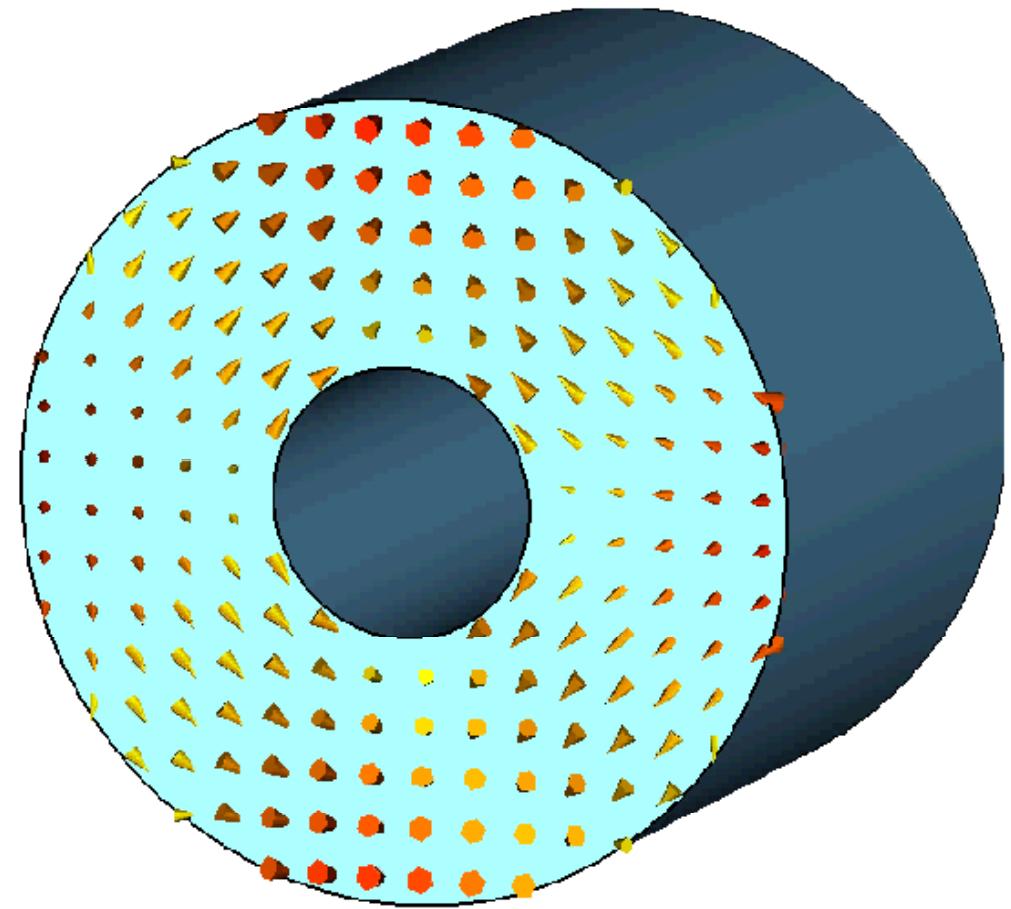
- Higher Coax-Modes:

$$r_i = 1\text{ mm}, r_a = 3\text{ mm}$$

$$f_{cutoff} = 46.3\text{ GHz}$$

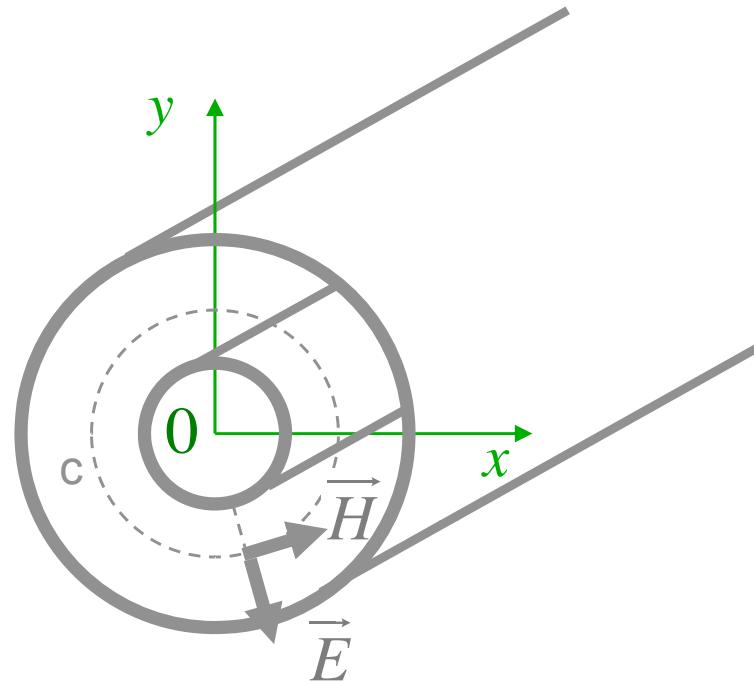


$\vec{E}$  at  $\omega t = 0$



$\vec{H}$  at  $\omega t = \pi / 2$

- Definitions of characteristic impedances:



Characteristic field impedance:

$$Z = \frac{|\vec{E}|}{|\vec{H}|}$$

Characteristic line impedance:

$$Z = \frac{U}{I}$$

$$U(z) = \int_{r_i}^{r_a} \vec{E} \cdot d\vec{s} \quad (\text{voltage})$$

$$I(z) = \oint_c \vec{H} \cdot d\vec{s} \quad (\text{current})$$

- Definitions of characteristic impedances:

Characteristic line impedance:

$$Z = \frac{U}{I}$$

$$U(z) = \int_{r_i}^{r_a} \vec{E} \cdot \vec{ds} \quad (\text{voltage})$$

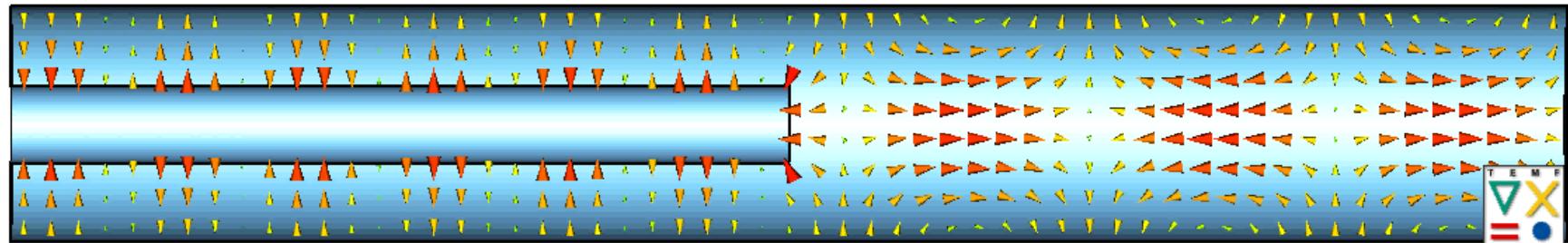
$$I(z) = \oint_c \vec{H} \cdot \vec{ds} \quad (\text{current})$$

- The characteristic line impedance only depends on the dimensions of the cross-section and on the permittivity / permeability of the dielectric

- Transition from coaxial cable to circular waveguides:

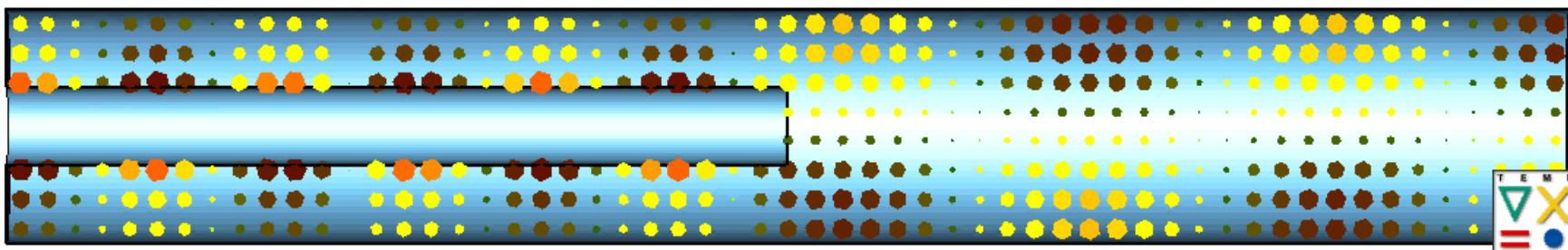
$$(r_i = 1 \text{ mm}, r_a = 3 \text{ mm})$$

☞ Electric field:  $f = 57 \text{ GHz}$



TEM-mode → TM<sub>01</sub>-Mode

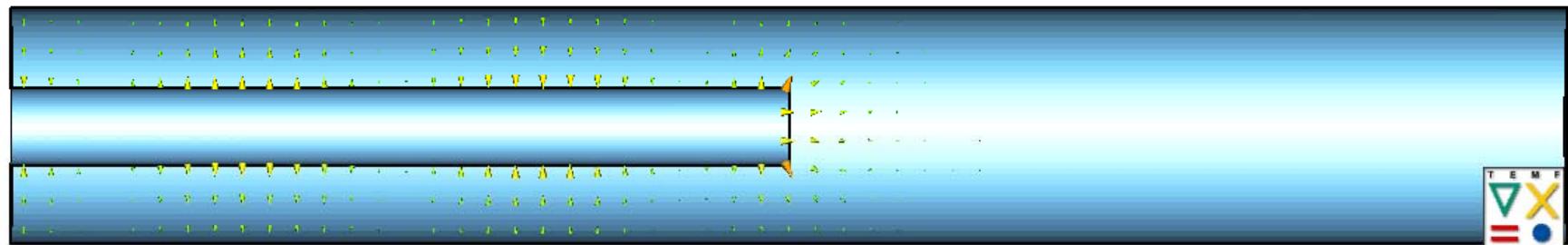
☞ Magnetic field:



- Transition from coaxial cable to circular waveguides:

- ☞ Electric field:

$$f = 25 \text{ GHz} \quad (f_c = 38.25 \text{ GHz} !)$$

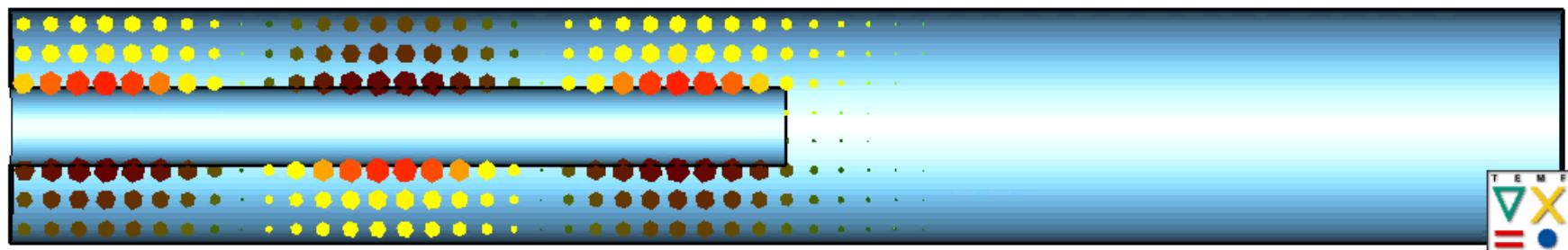


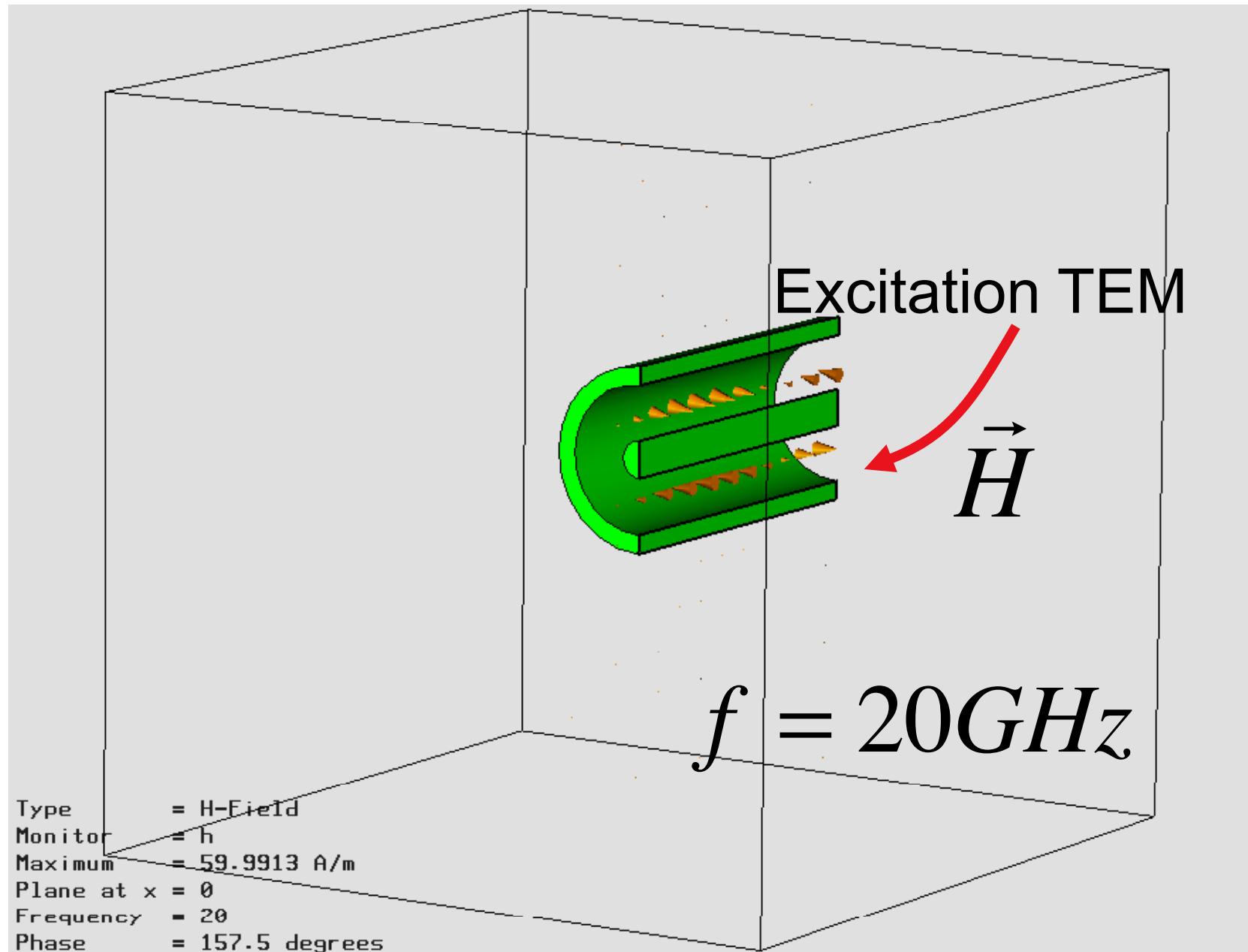
TEM-mode

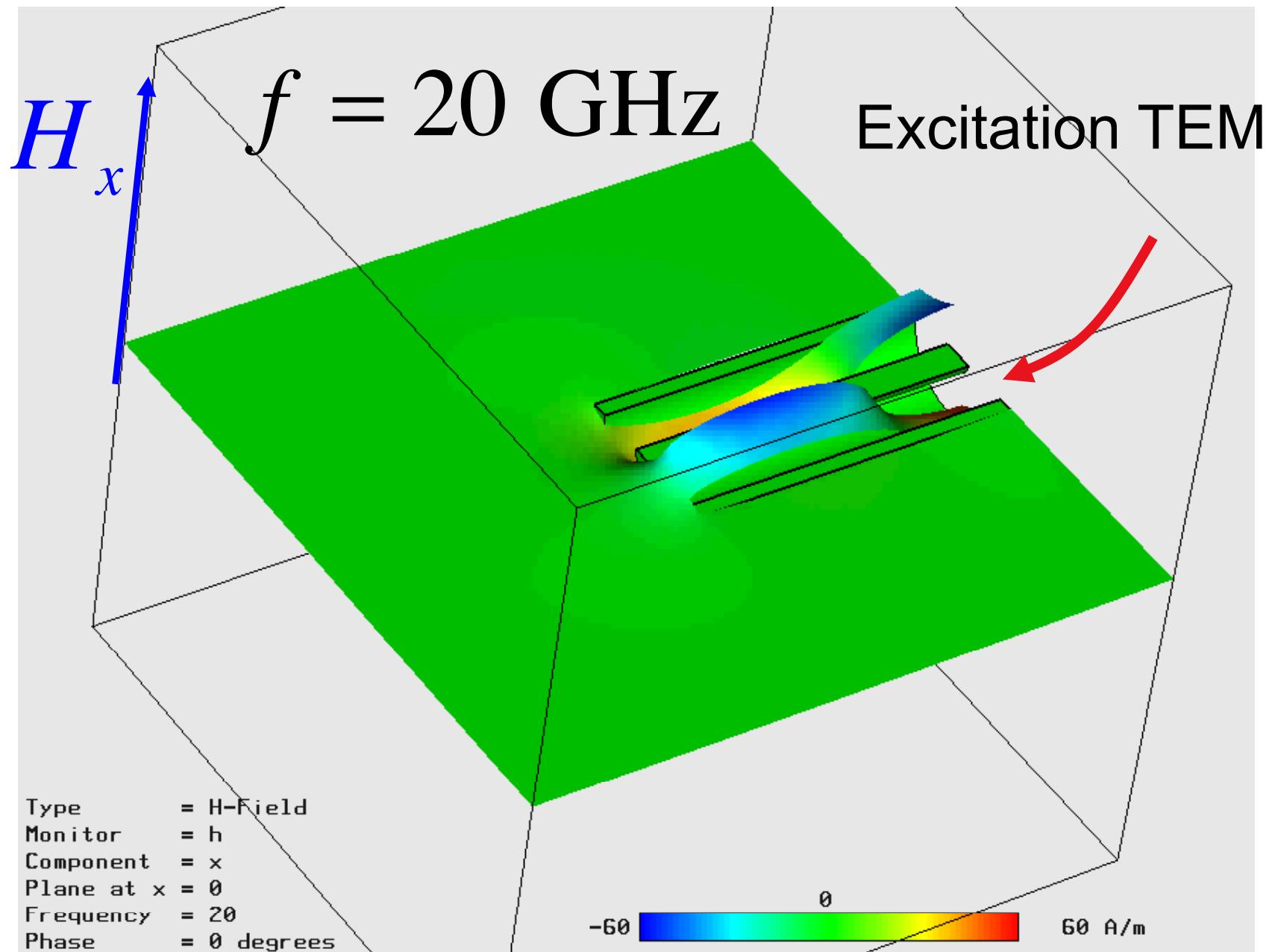


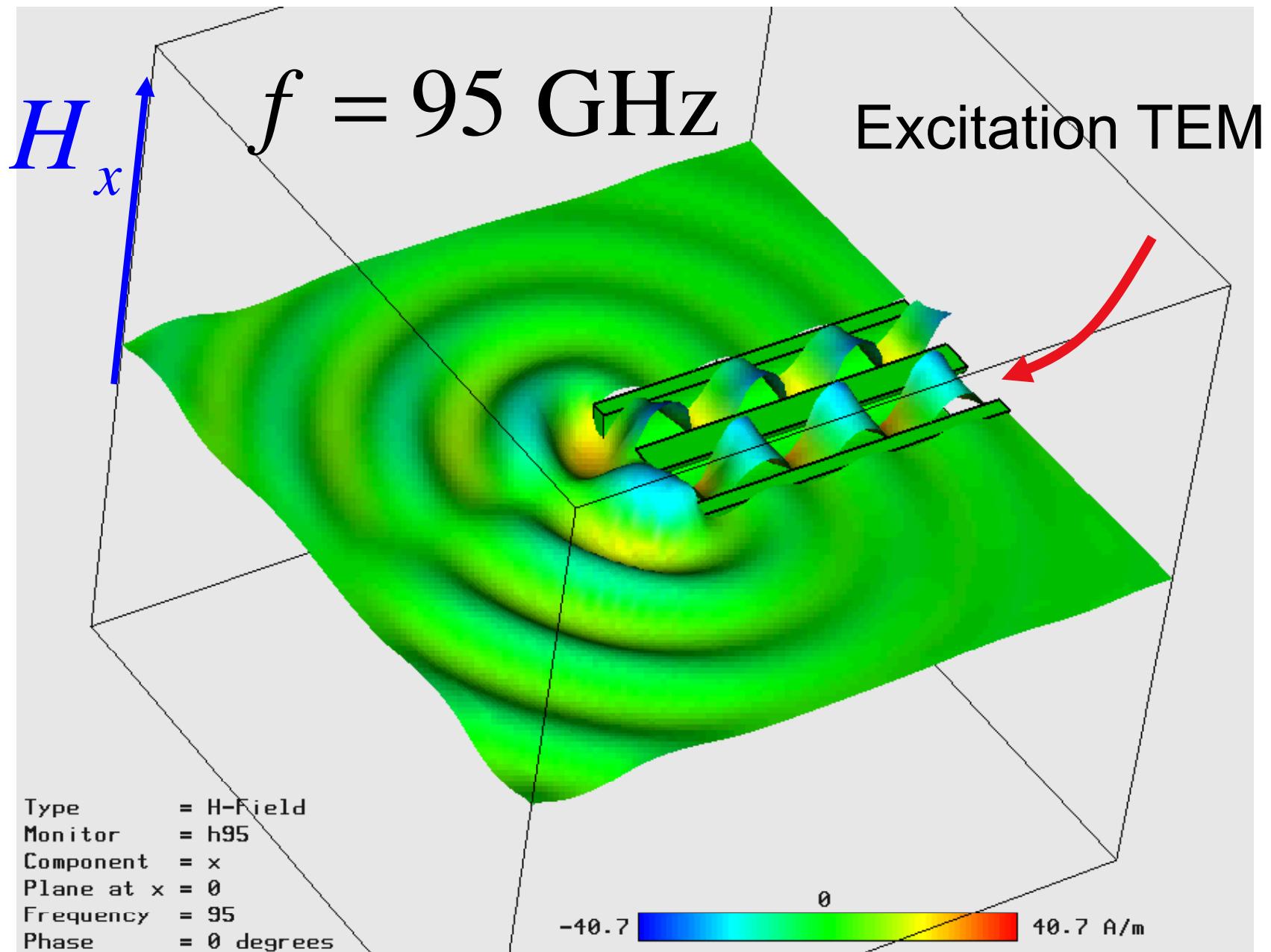
TM<sub>01</sub>-Mode below cutoff !

- ☞ Magnetic field:









End