

0.1 Conformal mapping

0.1.1 Joukowski Transformation

The question is how can we transform our cylinder into something that looks like an airfoil? We achieve this by using conformal mapping which maps each part of our coordinate space to a new one.

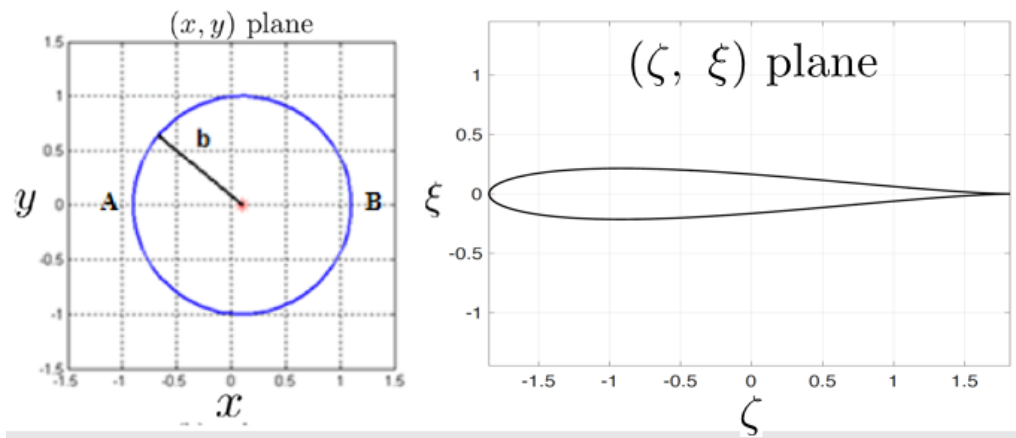
$$x = r \cos \theta \quad (1)$$

$$y = r \sin \theta \quad (2)$$

$$w = G(z) = z + \frac{\lambda^2}{z} \quad (3)$$

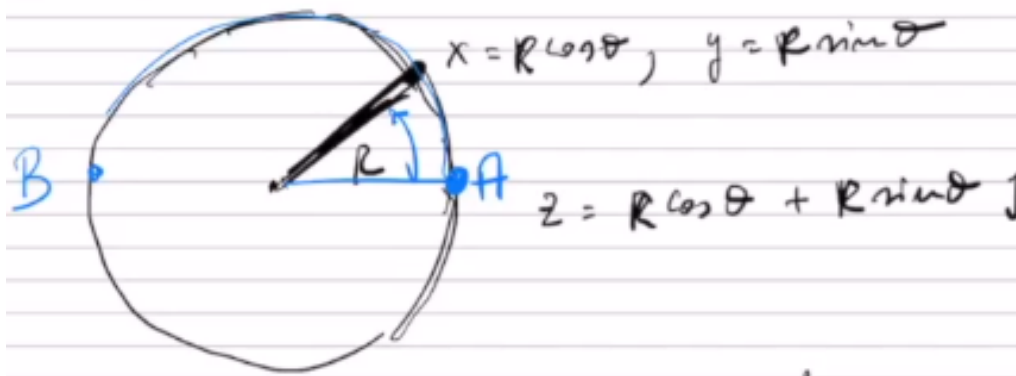
$$z = re^{i\theta} = x + yi \quad (4)$$

$$w = \zeta + \xi i \quad (5)$$



0.1.2 Flat plate and Ellipse

Lets consider a cylinder:



$$z = R \cos \theta + Ri \sin \theta \quad (6)$$

$$w = G(z) = z + \frac{\lambda^2}{z} \quad (7)$$

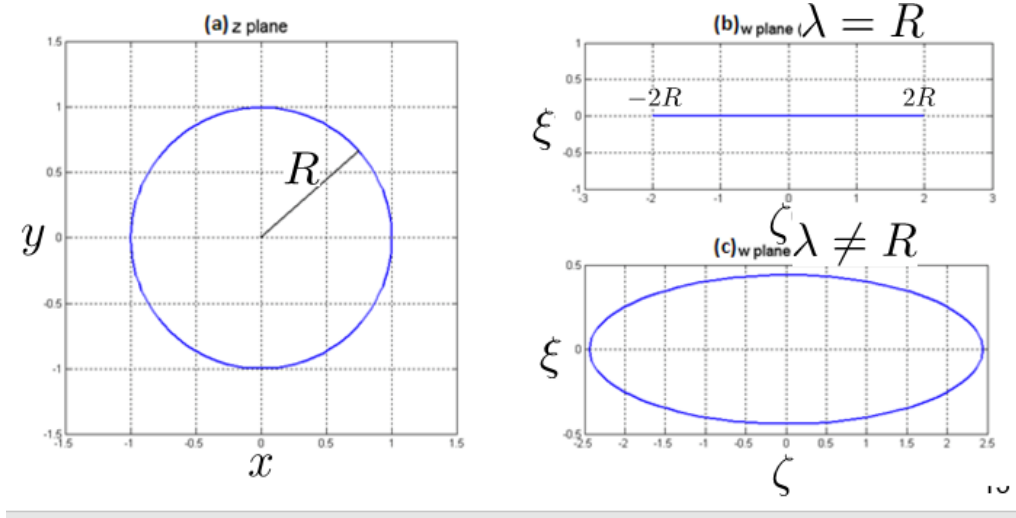
$$= R \cos \theta + Ri \sin \theta + \frac{\lambda^2}{R \cos \theta + Ri \sin \theta} \quad (8)$$

$$= R \cos \theta + Ri \sin \theta + \frac{\lambda^2(R \cos \theta - Ri \sin \theta)}{(R \cos \theta + Ri \sin \theta)(R \cos \theta - Ri \sin \theta)} \quad (9)$$

$$= R \cos \theta + Ri \sin \theta + \frac{\lambda^2(R \cos \theta - Ri \sin \theta)}{R^2} \quad (10)$$

$$w = R \cos \theta \left(1 + \frac{\lambda^2}{R}\right) + Ri \sin \theta \left(1 - \frac{\lambda^2}{R}\right) \quad (11)$$

When $\lambda = R$, we achieve a flat plate transformation as our imaginary term is always 0. When $\lambda \neq R$, we achieve an elliptical shape.



0.1.3 Aerofoils

We can further refine our conformal mapping by adding two elements to our equation, x_c and y_c . These two terms effect the center of the cylinder in x and y .

$$w = G(z) = z + \frac{\lambda^2}{z} \quad (12)$$

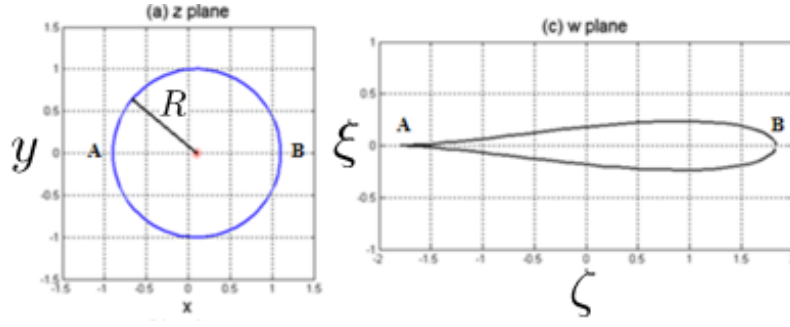
$$\lambda = R - \sqrt{x_c^2 + y_c^2} \quad (13)$$

$$z = x + x_c + (y + y_c)i \rightarrow w = \zeta + \xi i \quad (14)$$

Mapping these into their respective planes:

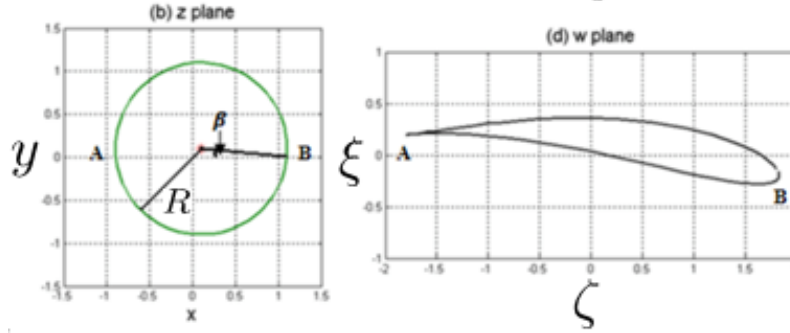
$$x_c \neq 0 \quad (15)$$

$$y_c = 0 \quad (16)$$



$$x_c \neq 0 \quad (17)$$

$$y_c \neq 0 \quad (18)$$



0.1.4 Cambered airfoil

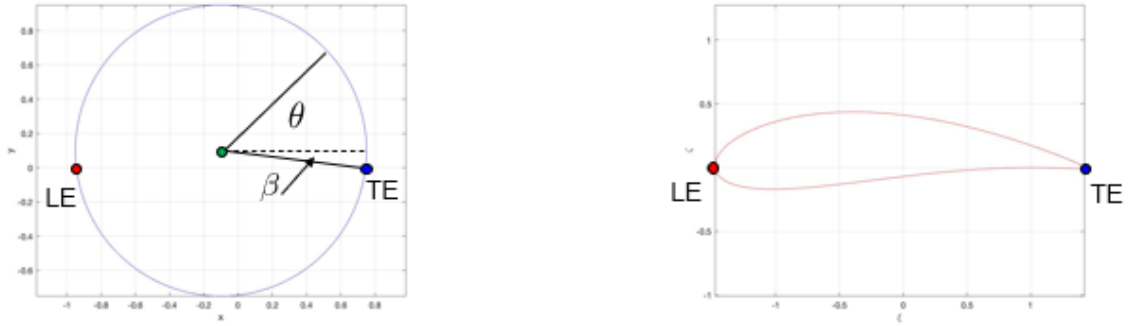


Figure 1: $x_c = -0.1$, $y_c = 0.1$, $\sin\beta = \frac{y_c}{R}$

0.2 Uniform stream + circulation

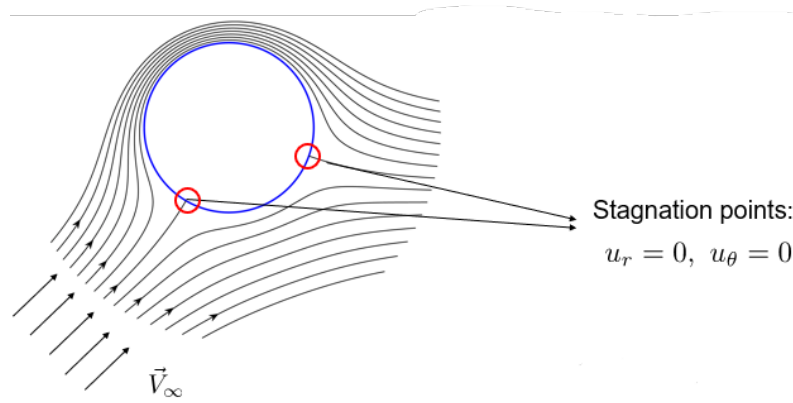
What is the correct relationship between the circulation and the angle of incidence?

$$\Gamma = f(\alpha) \quad (19)$$

$$\text{Stagnation points: } u_r = 0, \quad u_\theta = 0 \quad (20)$$

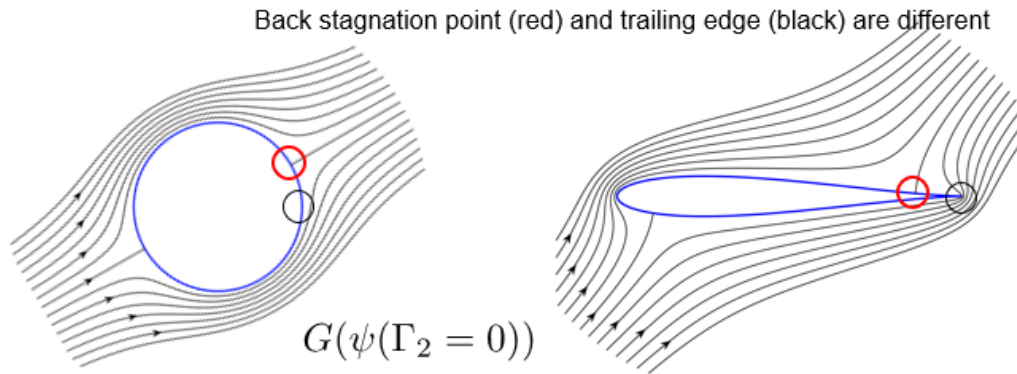
$$\phi = V_\infty r \cos(\theta - \alpha) \left(1 + \frac{R^2}{r^2} \right) - \frac{\Gamma}{2\pi} \theta \quad (21)$$

$$\psi = V_\infty r \sin(\theta - \alpha) \left(1 - \frac{R^2}{r^2} \right) + \frac{\Gamma}{2\pi} \log\left(\frac{r}{R}\right) \quad (22)$$

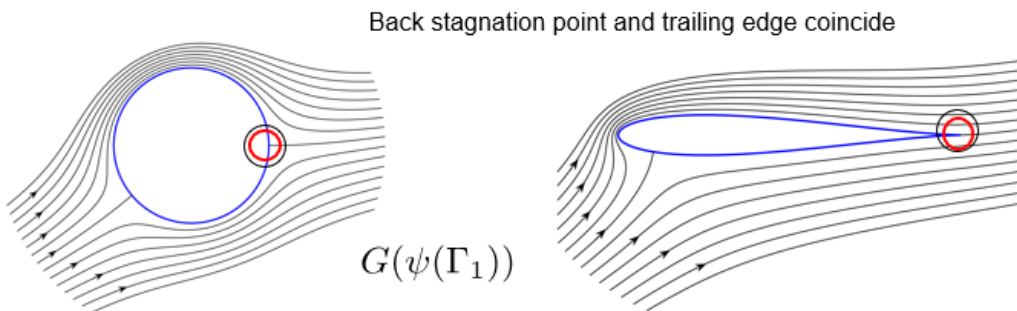


0.3 Kutta condition

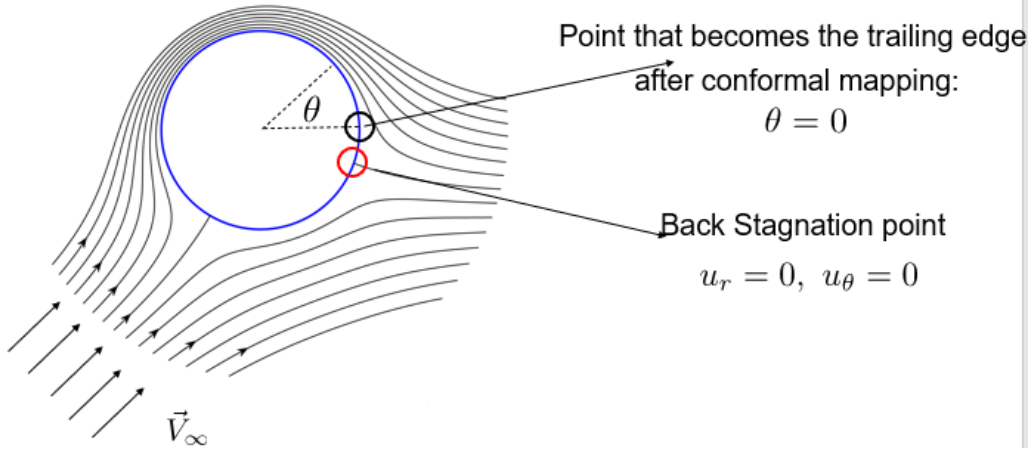
I can apply the Joukowski Transformation to the streamlines around the cylinder and obtain the streamlines around the airfoil. However I need to carefully select the circulation, Γ .



We can see that this streamline is unrealistic. This is because to curl around the trailing edge cusp, the speed goes to infinity. This is clearly impossible and flow separation would occur instead. This is the wrong circulation value for the given angle of incidence.



This streamline is more realistic as the velocity is parallel to the trailing edge cusp and its component in the vertical direction is zero. This circulation can be used to predict the lift F_L but still predicts no drag ($F_D = 0$). The next question is how do we achieve this mapping?



We can see here that for a given angle of coincidence, the circulation has to be selected by imposing that the back stagnation point (red circle) and the point corresponding to the airfoil trailing edge (black circle) coincide.

$$\phi = V_{\infty} r \cos(\theta - \alpha) \left(1 + \frac{R^2}{r^2} \right) - \frac{\Gamma}{2\pi} \theta \quad (23)$$

$$\psi = V_{\infty} r \sin(\theta - \alpha) \left(1 - \frac{R^2}{r^2} \right) + \frac{\Gamma}{2\pi} \log \left(\frac{r}{R} \right) \quad (24)$$

Velocity components:

$$u_r = \frac{1}{r} \frac{\partial \psi}{\partial \theta} = V_{\infty} r \cos(\theta - \alpha) \left(1 - \frac{R^2}{r^2} \right) \quad (25)$$

$$u_{\theta} = -\frac{\partial \psi}{\partial r} = -V_{\infty} \sin(\theta - \alpha) \left(1 + \frac{R^2}{r^2} \right) - \frac{\Gamma}{2\pi r} \quad (26)$$

On the cylinder $r = R$:

$$u_r = 0 \quad (27)$$

$$u_{\theta} = -2V_{\infty} \sin(\theta - \alpha) - \frac{\Gamma}{2\pi R} \quad (28)$$

Kutta condition:

$$u_{\theta}(\theta = 0) = 0 \quad (29)$$

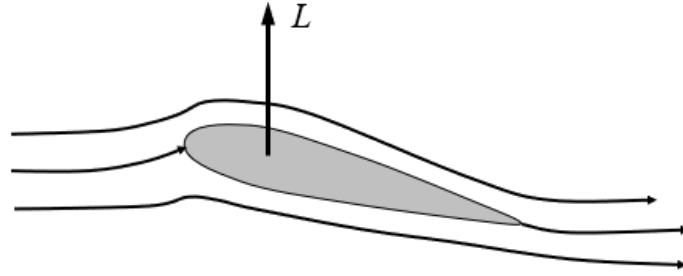
$$\Gamma = -4\pi V_{\infty} R \sin(-\alpha) = 4\pi V_{\infty} R \sin \alpha \quad (30)$$

For a cambered aerofoil:

$$u_{\theta}(\theta = -\beta) = 0 \quad (31)$$

$$\Gamma = -4\pi V_{\infty} R \sin(-\beta - \alpha) = 4\pi V_{\infty} R \sin(\beta + \alpha) \quad (32)$$

0.4 Aerofoil lift



The lift is given by:

$$L = \rho V \Gamma = 4\pi V_{\infty}^2 R \sin \alpha \quad (33)$$

The lift coefficient is:

$$c_L = \frac{L}{\frac{1}{2}\rho V_{\infty}^2 c} = 9\pi \frac{R}{c} \sin \alpha \quad (34)$$

Flow is inviscid so zero drag.

0.5 Flow past an Aerofoil

0.5.1 How was the circulation created?

When an airfoil is still on the ground there is no flow and total circulation is zero. As soon as the airfoil takes off a negative circulation is created at the back of the airfoil and a positive one is stored within the airfoil boundary layer.

0.5.2 Stalling

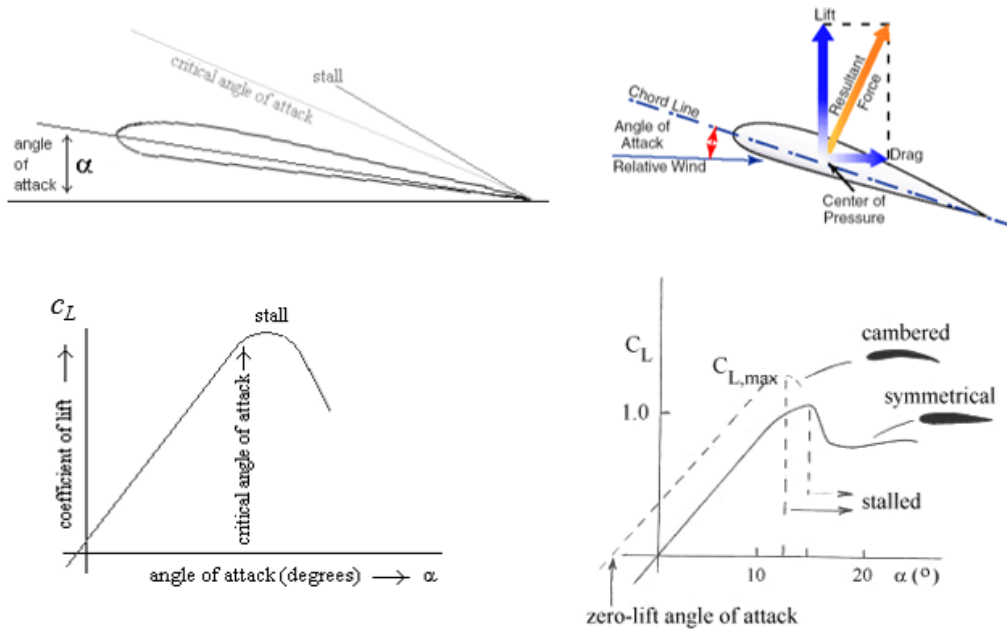
If the angle of attack becomes too great and boundary layer separation occurs on the top of the aerofoil the pressure pattern will change dramatically. In such a case, the lift component is insufficient to overcome the weight of the aircraft and disaster is imminent. This phenomenon is known as stalling. When stalling occurs, all, or most, of the 'suction' pressure is lost and the plane will suddenly drop from the sky. The solution to this is to put the plane into a dive to regain the boundary layer. A transverse lift force is then exerted on the wing which gives the pilot some control and allows the plane to be pulled out of the dive.



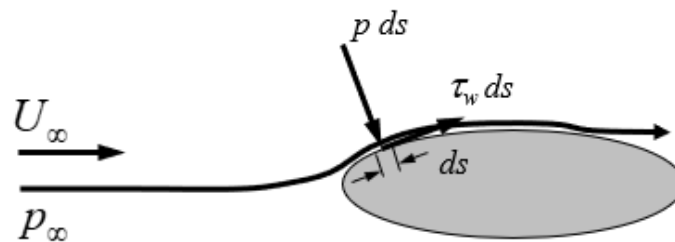
Boundary layer separation occurs on the upper surface (where the adverse pressure gradients are large) and produces a different c_p distribution from the unseparated case. As skin friction is predominant for the unstalled wing, the profile drag is sensitive to any increase in form drag.

- At a small angle of attack α : separation is close to trailing edge, wake is thin, low form of drag
- As the angle of attack α increases: separation moves along the top of the aerofoil, wake width increases, form drag increases
- The critical angle of attack α_{crit} : value of α form maximum c_L , stall angle
- If the angle of attack $\alpha > \alpha_{crit}$: separation from most of the upper surface, wider wake with turbulence, considerable form drag

The exact value of the critical angle of attack depends on the type and shape of the aerofoil.



0.6 Form drag and skin friction



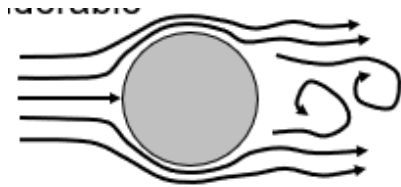
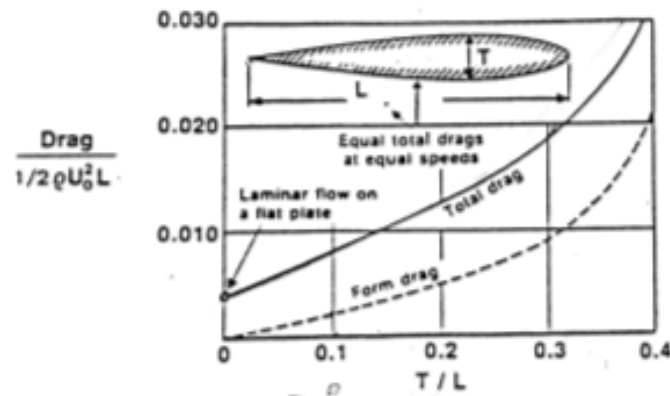
Force due to pressure p on element of surface area ds is $p ds$. The force $p ds$ has a component in the stream direction and if this component is integrated over the whole body surface, it gives the drag due to pressure distribution or the form drag. Form drag is the resultant of the forces normal to the surface (pressure distribution: normal stresses). The same element of area ds experiences a shear stress τ_w due to the velocity of the gradient normal to the surface and the associated shear force is $\tau_w ds$. This also has a component in the stream direction which when integrated over the body surface gives the drag due to skin friction. Skin friction drag is resultant of forces tangential to the surface (shear stresses).

0.7 Total drag

Total drag (also known as profile drag) is given by:

$$\text{Total drag} = \text{Form drag} + \text{Skin friction} \quad (35)$$

Separation is what leads to large form drag. The benefits of streamlining can be considerable.



Bluff Body

Form drag dominant
Skin friction insignificant
 $c_D \approx 1.0$ at high Re



Aerofoil

Skin friction dominant
Low form drag
 $c_D \approx 0.01$ at high Re