Introduction
Problem statement
Multiscale method
Benchmark
Conclusion

Multiscale Model Reduction for Neutron Diffusion Equation

Aleksandr Avvakumov ¹ Denis Spiridonov ²
Aleksandr Vasilev ²

 $^1{\rm National}$ Research Center "Kurchatov Institute", Moscow, Russia $^2{\rm North\text{-}Eastern}$ Federal University, Yakutsk, Russia

4th MHPCMP, September 9-12, 2020 Sochi



Introduction

Problem statement

We consider a bounded 2D domain Ω ($\mathbf{x} = \{x_1, x_2\} \in \Omega$) with a convex boundary $\partial \Omega$. The one-group transport equation in the diffusion approximation with one-group delayed neutron source can be written as follows

$$\frac{1}{v}\frac{\partial\phi}{\partial t} - \nabla \cdot D\nabla\phi + \Sigma_{r}\phi = \frac{1-\beta}{K_{eff}}\nu\Sigma_{f}\phi + \lambda c,
\frac{\partial c}{\partial t} + \lambda c = \frac{\beta}{K_{eff}}\nu\Sigma_{f}\phi.$$
(1)

Here $\phi(\mathbf{x},t)$ is the neutron flux at point \mathbf{x} and time t, v is the effective neutron velocity, $D(\mathbf{x})$ is the diffusion coefficient, $\Sigma_r(\mathbf{x},t)$ is the removal cross-section, $\Sigma_f(\mathbf{x},t)$ is the fission cross-section, β is the effective fraction of delayed neutrons, λ is the decay constant of the delayed neutron source.

Boundary and initial conditions

Let's determine the boundary and initial conditions for Equation (1). The albedo-type conditions are set at the boundary $\partial\Omega$:

$$D\frac{\partial \phi}{\partial n} + \gamma \phi = 0, \tag{2}$$

where n is an outer normal to the boundary $\partial\Omega$, γ is the albedo constant. Let's propose that at the initial time t=0, the reactor is in steady-state critical condition

$$\phi(\mathbf{x},0) = \phi^0(\mathbf{x}). \tag{3}$$

Time approximation

To solve the problem within the domain Ω , we approximate the system of equations (1)-(3) using the finite element method. We discretize the time derivatives using finite-difference scheme, and then bring each stationary problem to a variational formulation. For approximation in time, we use a fully implicit scheme with time step τ .

Varitional formulation

To specify the variational formulation, we multiply equations by the test function q and integrate over the domain Ω . Using the integration by parts, we obtain the following variational formulation: let's find $\phi^{n+1} \in V$ such that

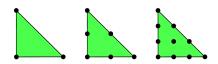
$$\int_{\Omega} \frac{1}{v} \frac{\phi^{n+1}}{\tau} q d\mathbf{x} + \int_{\Omega} D^{n+1} \nabla \phi^{n+1} \cdot \nabla q d\mathbf{x} + \int_{\Omega} \Sigma_{r}^{n+1} \phi^{n+1} q d\mathbf{x} - \int_{\Omega} \frac{1 + \lambda \tau - \beta}{K_{eff} (1 + \lambda \tau)} \nu \Sigma_{f}^{n+1} \phi^{n+1} q d\mathbf{x} + \int_{\partial \Omega} \gamma \phi^{n+1} q d\mathbf{s} =$$

$$\int_{\Omega} \frac{1}{v} \frac{\phi^{n}}{\tau} q d\mathbf{x} + \int_{\Omega} \frac{\lambda c^{n}}{1 + \lambda \tau} q d\mathbf{x}, \quad \forall q \in V,$$
(4)

where $V = H^1(\Omega)$ is the Sobolev space consisting of scalar functions v such that v^2 and $|\nabla v^2|$ have a finite integral in Ω .

Discrete problem

Further, it's necessary to pass from the continuous variational problem (4) to the discrete problem. We introduce finite-dimensional space of finite elements $V_h \subset V$ and formulate a discrete variational problem. We use standard linear basis functions as basis functions to solve the problem on the fine grid.



Time step selection algorithm

- Predictable time step: $\tilde{\tau}^{n+1} = \gamma \tau^n$ (eg $\gamma = 1.25$)
- ② Predictive solution \widetilde{y}^{n+1} : an explicit scheme, $\widetilde{t}^{n+1} = t^n + \widetilde{\tau}^{n+1}$
- **3** Estimation of approximation error: by found \widetilde{y}^{n+1} from an implicit scheme
- Step selection τ^{n+1} : $\|\psi^{n+1}\| \approx \delta$
- **Solution** on a new time layer y^{n+1} : an implicit scheme, $t^{n+1} = t^n + \tau^{n+1}$

Neutron diffusion equation

One group diffusion approximation with one group delayed neutron sources

$$\frac{1}{v}\frac{\partial \phi}{\partial t} - \nabla \cdot D\nabla \phi + \Sigma_{a}\phi = (1 - \beta)\nu\Sigma_{f}\phi + \lambda c,$$
$$\frac{\partial c}{\partial t} + \lambda c = \beta\nu\Sigma_{f}\phi.$$

Boundary condition

$$D\frac{\partial \phi}{\partial n} + \gamma_a \phi = 0.$$

Initial conditions

$$\phi(0) = \phi^0, c(0) = c^0.$$

Calculated formulas

Denote vectors and matrix $\varphi = \{\varphi, s\}$, $\psi = \{\psi_1, \psi_2\}$,

$$A = \begin{pmatrix} -\nabla \cdot D\nabla + \Sigma_{a} - (1-\beta)\nu\Sigma_{f} - \lambda & 0\\ 0 & \lambda - \beta\nu\Sigma_{f} \end{pmatrix}.$$

The approximation error

$$\begin{split} \widetilde{\psi}^{n+1} &= (A^{n+1} - A^n)\varphi^n + A^{n+1}(\widetilde{\varphi}^{n+1} - \varphi^n) \\ &= \widetilde{\tau}^{n+1} \left(\frac{A^{n+1} - A^n}{\widetilde{\tau}^{n+1}} \varphi^n + A^{n+1} \frac{\widetilde{\varphi}^{n+1} - \varphi^n}{\widetilde{\tau}^{n+1}} \right). \end{split}$$

We match error $\widetilde{\psi}^{n+1}$ with step $\widetilde{\tau}^{n+1}$, and ψ^{n+1} with step τ^{n+1} :

$$\bar{\tau}^{n+1} = \gamma_{n+1} \tau^n, \quad \gamma_{n+1} = \frac{\delta}{\|\widetilde{\psi}^{n+1}\|} \gamma.$$

The needed time step

$$\tau^{n+1} \leq \overline{\tau}^{n+1}, \quad \tau^{n+1} \leq \widetilde{\tau}^{n+1}, \quad \tau^{n+1} = \max \big\{ \tau^0, \min \big\{ \gamma_{n+1}, \gamma \big\} \tau^n \big\}.$$

The approximation error has the first order in time

$$\widetilde{\psi}^{n+1} = \mathcal{O}(\widetilde{\tau}_{n+1}).$$

In view of this, we set

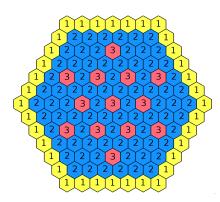
$$\|\widetilde{\psi}^{n+1}\| \leq \|(A^{n+1}-A^n)\varphi^n + A^{n+1}(\widetilde{\varphi}^{n+1}-\varphi^n)\|.$$

Calculated formula for time step

$$\gamma_{n+1} = \frac{\delta}{\|(A^{n+1} - A^n)\varphi^n + A^{n+1}(\widetilde{\varphi}^{n+1} - \varphi^n)\|} \gamma.$$



IAEA-2D benchmark



- Without reflector
- One group of *instantaneous* and *delayed* neutrons
- Modeling effect of immersion or extraction of control rods

Scenario

Define the scenario of process:

- The spectral problem is solved (initial condition);
- Calculation for the non-stationary model in the range 0 to 0.5 sec;
- **②** At a moment of 0.1 sec the value Σ_a for the zone 3 changes to \pm 0.000625.

Software

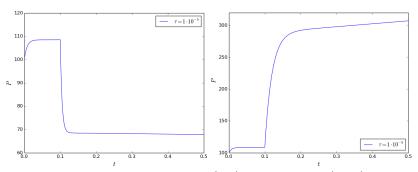








Nuclear power

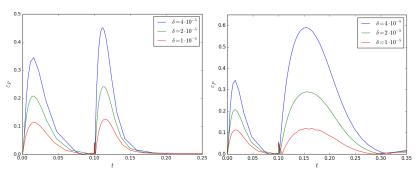


Nuclear power for immersion (left) and extraction (right).

$$P(t) = a \sum_{g=1}^{G} \int_{\Omega} \Sigma_{fg} \phi_g dx,$$

where a – normalization factor.

Error

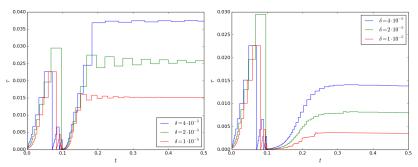


Error for immersion (left) and extraction (right).

$$\epsilon_P(t) = |P_{ref} - P|,$$

where P_{ref} – reference solution.

Time step



Time steps for immersion (left) and extraction (right).

Counting time and number of steps

	immersion			extraction		
δ	$\max(\epsilon_P)$	n	t, sec	$\max(\epsilon_P)$	n	t, sec
$4 \cdot 10^{-5}$	0.450	136	16	0.590	241	35
$2\cdot 10^{-5}$	0.241	159	20	0.290	373	62
$1\cdot 10^{-5}$	0.125	270	37	0.120	773	145

Reference solution: fixed time step 10^{-5} , number of steps – 50000, counting time – 2130 sec.

Acknowledgements

This work was supported by the Russian Foundation for Basic Research (grants 16-08-01215 and 18-31-00315) and by the grant of the Russian Federation Government 14.Y26.31.0013.

Conclusion

- An algorithm for automatic time step selection for numerical solution of neutron diffusion problems has been developed.
- The solution is obtained using guaranteed stable implicit schemes, and the step choice is performed with the use of the solution obtained by an explicit scheme.
- Calculation results obtained for a neutron diffusion problems demonstrate reliability of the proposed algorithm for time step choice.

Thank you for your attention!

