Reconstruction of the Acoustic Refractive Index from the Interior Transmission Eigenvalue Problem (ITEP)

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March 2024

Content

Introduction

➤ Inverse Problem? Inverse Scattering? Interior Transmission Eigenvalue Problem? Inverse Spectral Problem?

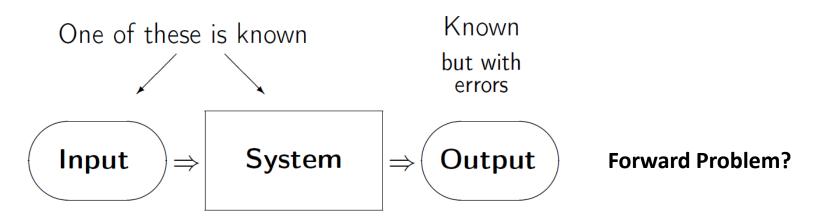
Methods

- Linear sampling method? Regularization? Finding the smallest ITE? Estimation of *n*?
- Results
- Conclusion

Introduction

Inverse Problems

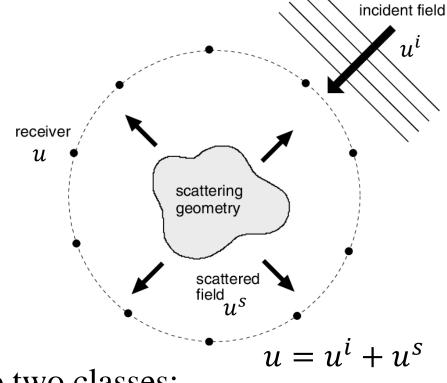
Inverse Problem



- What are the applications? Geophysical imaging, Computational Photography, Medical imaging, *Inverse Scattering*, ...
- Why are inverse problems difficult? Ill-posedness (Existence, Uniqueness, and Stability) => Need to be Regularized

Inverse Scattering

• Scattering theory is concerned with the effect of an inhomogeneous medium has on an *incident particle* or *wave*.



- Inverse scattering problem for **acoustic** and electromagnetic waves can broadly be divided into two classes:
 - ➤ Inverse obstacle problem (IOP): Determine the shape (boundary) of an impenetrable obstacle from a scattered field.
 - ➤ Inverse medium problem (IMP): Determine medium property (sound speed or permittivity)

Inverse Medium Problem (IMP)

$$\Delta u + k^2 n(x) u = 0$$
 in D ,
 $\Delta u + k^2 u = 0$ in $\mathbb{R}^3 \backslash \overline{D}$,
 $u(x) = e^{ikx \cdot d} + u^s(x)$,

Sommerfeld radiation condition $\lim_{r\to\infty} r(\frac{\partial u^s}{\partial r} - iku^s) = 0.$

where $u^i(x)=e^{ikx\cdot d}$ is the time-harmonic acoustic plane wave, $k=\omega/c_0$ the wave number, ω the angular frequency, c_0 the speed of sound in the homogeneous host medium, d the direction of propagation, r=|x|, $n=c_0^2/c^2$ is the refractive index, c^2 the speed of sound in the in the inhomogeneous medium, D the inhomogeneity.

Inverse Obstacle Problem (IOP)

$$\Delta u + k^2 u = 0 \quad \text{in } \mathbb{R}^3 \backslash \overline{D},$$

$$u(x) = e^{ikx \cdot d} + u^s(x),$$

$$u = 0 \quad \text{on } \partial D,$$

$$\lim_{r \to \infty} r(\frac{\partial u^s}{\partial r} - iku^s) = 0.$$

where *D* is an impenetrable obstacle.

• Dirichlet boundary condition corresponds to a *sound-soft* obstacle.

Far-field Pattern

• Far-field pattern, or scattering amplitude, of the scattered acoustic wave is of interest. u^s has the asymptotic behavior

$$u^{s}(x) = \frac{e^{ikr}}{r} u_{\infty}(\hat{x}, d) + \mathcal{O}\left(\frac{1}{r^{2}}\right), \qquad r = |x| \to \infty$$

where $\hat{x} \coloneqq x/|x|$.

• Far-field equation:

$$(Fg)(\hat{x}) \coloneqq \Phi_{\infty}(\hat{x}, z)$$

where F is the far-field operator, g is a density function, and, $\Phi_{\infty}(\hat{x}, z)$ is the far field pattern corresponding to the fundamental solution $\Phi(x, z)$ oh the Helmholtz equation. (Regularization methods!).

Interior Transmission Eigenvalue Problem (ITEP)

- Are there any incident wave u^i such that scattered field u^s is identically zero?
 - > The answer to this question leads to the interior transmission eigenvalue problem (ITEP).
- ITEP: find k > 0 and $v, w \in C^2(D)$ such that:

$$\Delta w + k^2 n(x)w = 0 \quad \text{in } D$$

$$\Delta v + k^2 v = 0 \quad \text{in } D$$

$$w = v \quad \text{on } \partial D$$

$$\frac{\partial w}{\partial x} = \frac{\partial v}{\partial x} \quad \text{on } \partial D$$

The value of k for which there exist a non-trivial solution (v, w) for above equations are the interior transmission eigenvalues (ITE).

Inverse Spectral Problem (ISP)

- **Spectral Inversion**: Can one hear the shape of a drum? [Mark Kac 1966]
- Assume the shape of the scatterer is known => Determine the unknown refractive index n(x) from transmission eigenvalues.
- ITEs is required => <u>Inverse Spectral Problem for Transmission Eigenvalues</u>
- Note:
 - ✓ Real transmission eigenvalue can be measured by scattering data
 - ✓ ITEs provide information about the material properties of scatterer (n(x)).

Project Overview

Main Steps of the project:

- 1) Extract the ITEs from far-field data
- 2) Estimate the acoustic refractive index (n(x)) from the first ITE.

Methods

Linear sampling method (LSM)

• Consider the far-field operator $F: L^2(\Omega) \to L^2(\Omega)$ as:

$$(Fg)(\hat{x}) = \int_{\Omega} u_{\infty}(\hat{x}, d)g(d) \, \mathrm{d}s(d).$$

where $u_{\infty}(\hat{x}, d)$ is the far-field pattern corresponding to incident direction d and observation direction \hat{x} .

• The far-field equation can be written as a $\underbrace{\text{linear inverse problem}}_{\text{of determining }g}$ from knowledge of

$$f = \frac{e^{i\pi/4}}{\sqrt{8\pi k}} e^{-ik\hat{x}\cdot y}$$

where g and f are related by

$$Fg = f$$
.

• If k is **not** a transmission eigenvalue, considering Ω as a unit circle/sphere:

 $||g||_{L^2(\Omega)} \to \infty$ when y approaches the boundary of the scatterer from the inside.

LSM: Regularization

- Fg = f is not in general solvable since f is **not** in the range (column space) of F (measurement noise, ...).
- Tikhonov regularization: find approximate solution g_{α} as solution to:

$$\min_{g} \left\{ \|F_{\delta} g_{\alpha} - f\|_{L^{2}(\Omega)}^{2} + \alpha \|g_{\alpha}\|_{L^{2}(\Omega)}^{2} \right\}$$

where F_{δ} is the far-field operator affected by measurement noise of the order of magnitude δ and α is the regularization parameter.

• By solving above equation, one can obtain:

$$g_{\alpha} = (F_{\delta}^{T} F_{\delta} + \alpha I)^{-1} F_{\delta}^{T} f$$

• Choosing the regularization parameter: <u>Morozov's discrepancy principle</u>, L-Curve method, Generalized Cross Validation, ...

LSM: Obtaining ITE

- Assume the shape of *D* found by LSM.
- LSM can be expected to fail when k is a ITE and in particular the norm of the (regularized) solution to

$$(Fg)(\hat{x}) = 1$$

should be large for such values of k.

So, determining ITE from the far-field data

• ITEP (slide 11) can be transformed into a *fourth-order equation* for u: = $w - v \in H_0^2(D)$

$$(\Delta + k^2 n) \frac{1}{n-1} (\Delta + k^2) u = 0, \quad \text{in } D$$

$$u = 0 \quad \text{and} \quad \frac{\partial u}{\partial v} = 0, \quad \text{on } \partial D$$

• Weak form:

$$\int_{D} \frac{1}{n-1} (\Delta u + k^2 u) (\Delta \overline{\phi} + k^2 n \overline{\phi}) \, \mathrm{d}x = 0, \quad \forall \phi \in H_0^2(D)$$

 ϕ : test function

Estimation of n from the first real ITE

- Use the Galerkin method to discretize the weak form!
- Assume $\{\varphi_i\}_{i=1}^{\infty}$ is a set of eigenfunctions for the problem (clamped plate problem):

$$L\varphi_i = \mu_i \varphi_i \quad \text{in } D$$

$$\phi_i = 0, \qquad \frac{\partial \phi_i}{\partial \psi} = 0 \quad \text{on } \partial D$$

where L is a forth order elliptic operator (Bilaplacian operator).

• By the Galerkin method, u_k is approximated as:

$$u_k^N = \sum_{i=1}^{(N)} c_i \phi_i$$

Estimation of n from the first real ITE

• Put u_k^N in the weak form:

$$\left[A^{(N)} - (k^{(N)})^2 B^{(N)} + (k^{(N)})^4 C^{(N)}\right] c = 0$$

a quadratic eigenvalue problem

where

$$A^{(N)} := \int_{D} \frac{1}{n(x) - 1} \Delta \phi_i \Delta \bar{\phi}_j \, \mathrm{d}x$$

$$B^{(N)} := -\left(\int_{D} \frac{n(x)}{n(x) - 1} \Delta \phi_{i} \overline{\phi}_{j} dx + \int_{D} \frac{1}{n(x) - 1} \phi_{i} \Delta \overline{\phi}_{j} dx\right)$$

$$C^{(N)} := \int_{D} \frac{n(x)}{n(x) - 1} \phi_i \bar{\phi}_j \, \mathrm{d}x$$

are $N \times N$ matrices and $\boldsymbol{c} = [c_1, c_2, ..., c_N]^T$, i, j = 1, ..., N.

Estimation of n from the first real ITE - Pseudocode

```
Given k_0; %analytically or computationally
RI = [n_{min}: step: n_{max}]; %refractive indices
     for idx = 1: length(RI)
          n = RI(idx)
          Compute A^{(N)}, B^{(N)}, C^{(N)}
          Solve \left[A^{(N)} - (k^{(N)})^2 B^{(N)} + (k^{(N)})^4 C^{(N)}\right] c = 0
          Find k_0^{(N)}
          Compute Err(idx) = norm(k_0 - k_0^n)
     end
     [:,idx_{opt}] = \min(Err)
     n_{opt} = RI(idx_{opt})
```

Results

LSM for Shape Reconstruction – 2D

- Shape: a unit circle centered at the origin.
- $k = 2\pi$
- Far-field patterns calculated:
 - A. Analytically [Colton and Kirsch 1996]

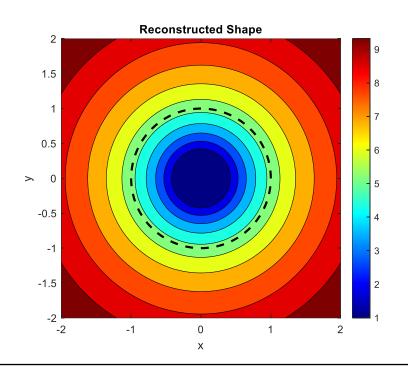
$$u_{\infty}(\phi;\theta) = -e^{-i\pi/4} \sqrt{2/\pi k} \sum_{n=-\infty}^{\infty} \frac{J_n(ka)}{H_n^1(ka)} e^{in(\phi-\theta)}$$

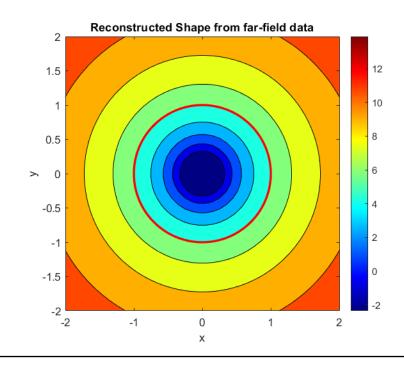
- B. Numerically by μ -diff toolbox [Thierry et al. 2015] in MATLAB
- **Tikhonov** Regularization

LSM for Shape Reconstruction – 2D

Number of incident angles = 180

Number of scattered angles = 180





Impenetrable unit circle centered at the origin- **Analytical** far-field patterns.

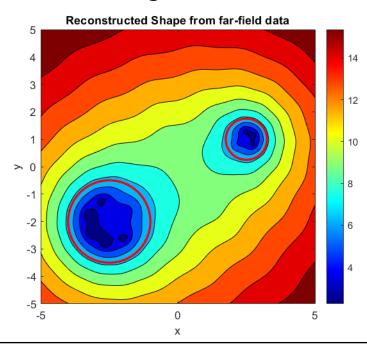
Impenetrable unit circle centered at the origin ($\alpha = 10^{-5}$).

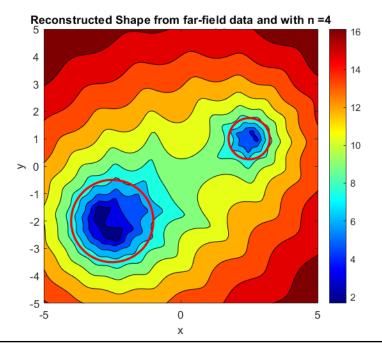
LSM for Shape Reconstruction – 2D

Shape: circles centered at (-2.5, 2.5) and (-2, 1) with a radius of 1.5 and 0.75, respectively

Far-field by μ -diff toolbox. **Tikhonov** Regularization

Number of incident angles = 180 Number of scattered angles = 180



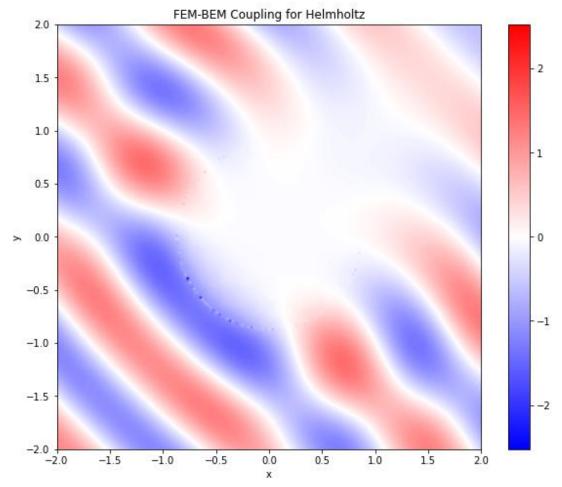


Impenetrable circles centered at (-2.5, 2.5) and (-2, 1) with a radius of 1.5 and 0.75, respectively. ($\alpha = 10^{-8}$).

Penetrable circles centered at (-2.5, 2.5) and (-2, 1) with a radius of 1.5 and 0.75, respectively and with n = 4 ($\alpha = 10^{-8}$).

LSM for Shape Reconstruction – 3D

- **Shape:** a unit sphere centered at the origin
- $k = 2\pi$
- **Far-field** patterns calculated:
 - By coupling the finite-element method
 (FEM) and the boundary-element method
 (BEM) in Python.
 - FEM was implemented by FEniCS [https://fenicsproject.org/].
 - BEM was implemented by Bempp [https://bempp.com/]



Total (near) field scattered from a unit sphere centered at the origin, cross section view at z = 0, n = 1/4

LSM for Shape Reconstruction – 3D

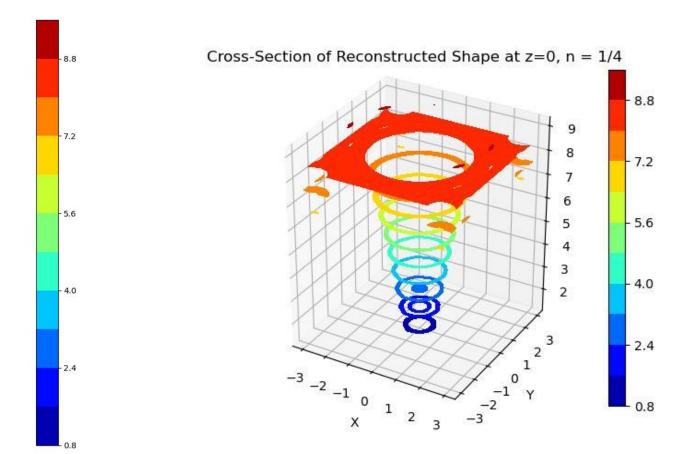
Shape: a unit sphere located at the origin with n = 1/4.

Far-field by FEM-BEM. **Tikhonov** Regularization ($\alpha = 10^{-10}$)

Number of incident angles = 81 Number of scattered angles = 81

Cross-Section of Reconstructed Shape at z=0, n=1/4

The figure shows a slice at z = 0.



LSM for Shape Reconstruction – 3D

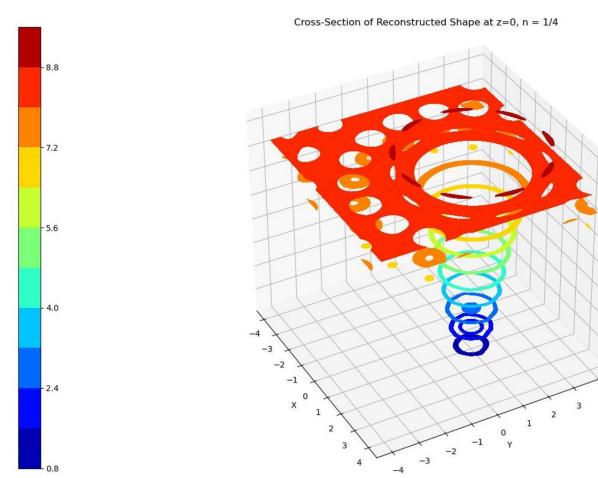
Shape: A unit sphere located at (1,1) with n = 1/4.

Far-field by FEM-BEM. **Tikhonov** Regularization ($\alpha = 10^{-10}$)

Number of incident angles = 81 Number of scattered angles = 81

Cross-Section of Reconstructed Shape at z=0, n=1/4The figure shows a slice at z=0.

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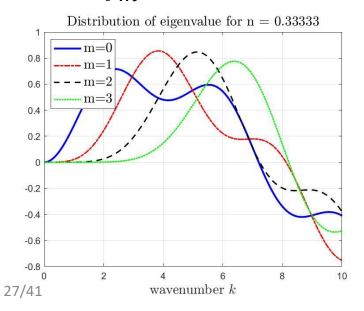


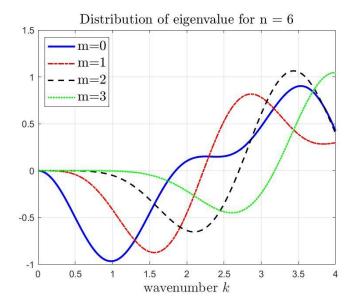
Extraction of the Smallest ITE from Far-field - 2D

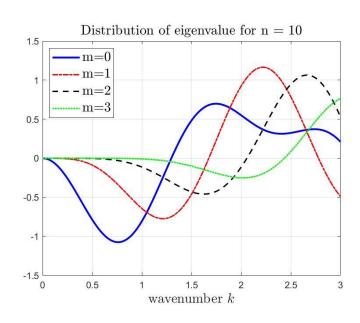
- **Shape:** unit circle centered at the origin
- **Analytical solution:** for a circle with constant n(x) and radius r, the first transmission eigenvalue is the lowest positive for which [Cakoni et al. 2007]:

$$\det\begin{pmatrix} J_m(kr) & J_m(k\sqrt{n}r) \\ J'_m(kr) & J'_m(k\sqrt{n}r) \end{pmatrix} = 0, \qquad m = 0, 1, \dots$$

where J_m are the Bessel functions of the first kind.

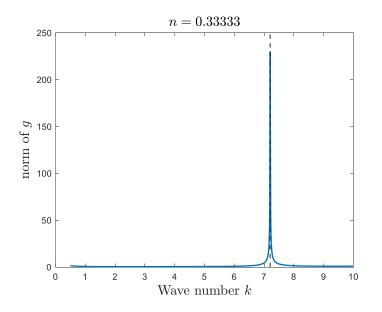


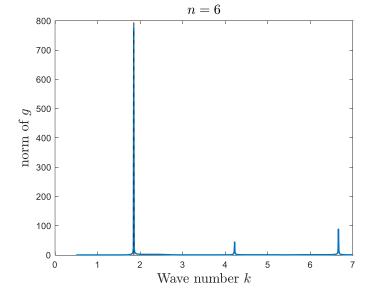


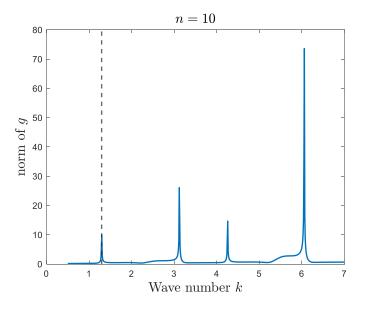


Extraction of the Smallest ITE from Far-field - 2D

- Using LSM (slide 17)
- **Far-field** by μ -diff toolbox. **Tikhonov** Regularization ($\alpha = 10^{-10}$)
- Number of incident angles = 61; Number of scattered angles = 61



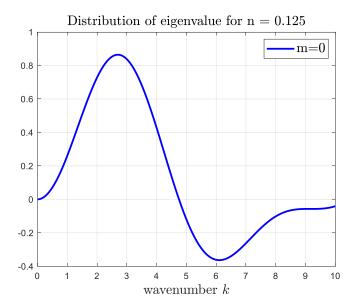


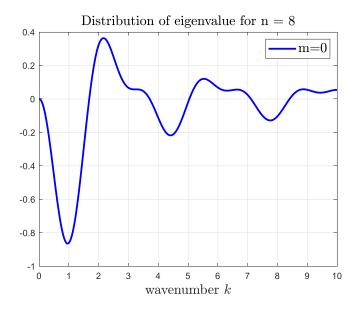


Extraction of the Smallest ITE from Far-field - 3D

- **Shape:** a unit sphere centered at the origin
- Analytical solution: for a sphere with constant n(x) and radius r, the first transmission eigenvalue is obtained from [Leung and Colton 2012]:

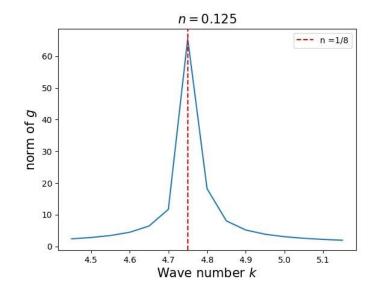
$$d(k) \coloneqq \frac{\sqrt{n}\sin(k)\cos(\sqrt{n}k) - \cos(k)\sin(\sqrt{n}k)}{\sqrt{n}k} = 0$$

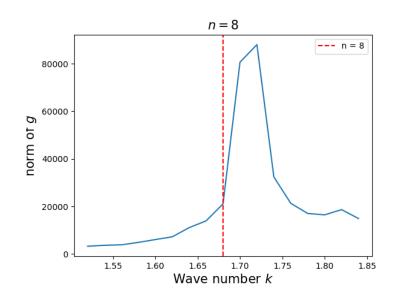




Extraction of the Smallest ITE from Far-field - 3D

- Using LSM (slide 17)
- **Far-field** by FEM-BEM. **Tikhonov** Regularization ($\lambda = 10^{-12}$)
- Number of incident angles = 81; Number of scattered angles = 81





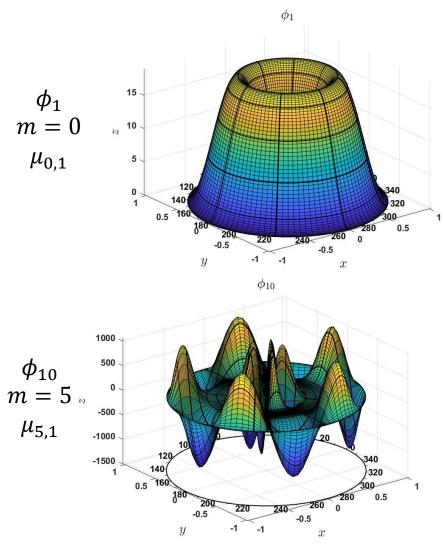
• For unite circle, now we construct the basis $\{\varphi_i\}_{i=1}^N$ for the clamped-plate problem (slide 19) with $L = \Delta \Delta$ (bilaplacian operator). In polar coordinates [Chakraverty 2008]:

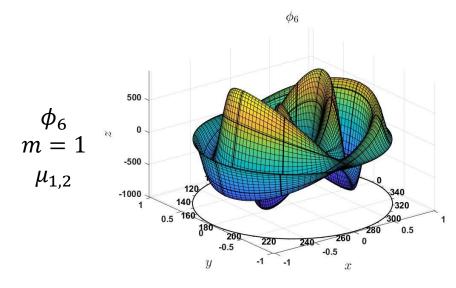
$$\varphi_i(r,\theta) = \left[a_i J_m(\mu_{m,j}r) + I_m(\mu_{m,j}r) \right] \cos(i\theta)$$

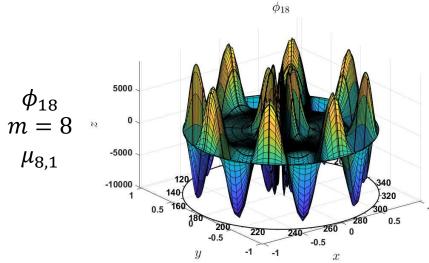
where I_m is modified (hyperbolic) Bessel function of first kind and eigenvalues $\mu_{m,j}$ can be computed from the relation:

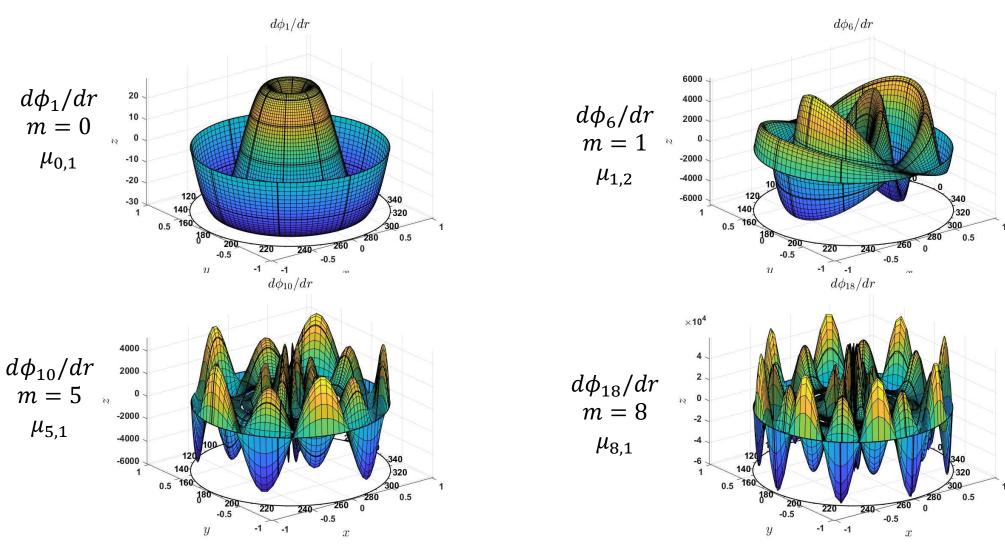
$$\det \begin{pmatrix} J_m(\mu_{m,j}r) & J'_m(\mu_{m,j}r) \\ I_m(\mu_{m,j}r) & I'_m(\mu_{m,j}r) \end{pmatrix} = 0, \qquad m = 0, 1, \dots$$

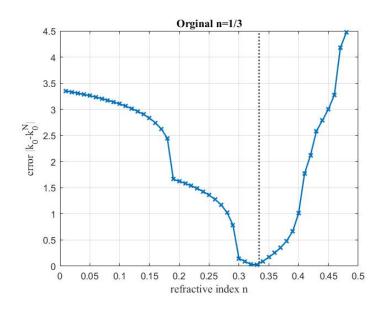
• **Note:** $\mu_{m,j}$ should be ordered and use the appropriate Bessel functions. j denotes j^{th} zero of above relation for each m.

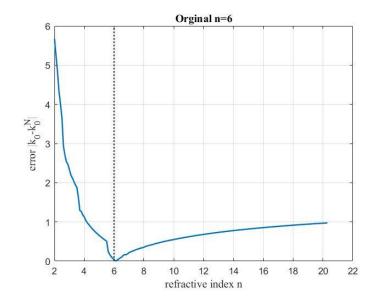


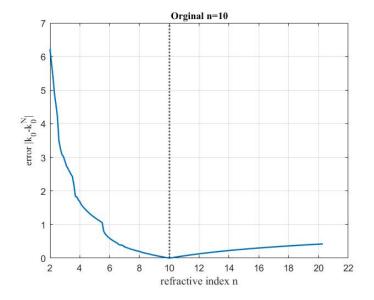












• For unite sphere, now we construct the basis $\{\varphi_i\}_{i=1}^N$ for the bilaplacian operator (slide 19) with $L = \Delta \Delta$. In spherical coordinates (assuming azimuthal symmetry):

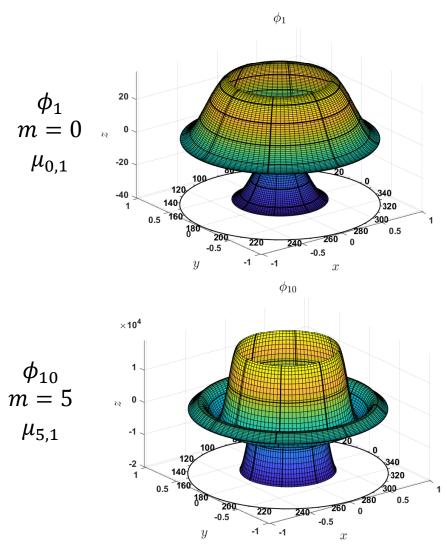
$$\varphi_i(r,\theta) = \left[a_i j_m(\mu_{m,j}r) + i_m(\mu_{m,j}r)\right] P_m(\cos(\theta))$$

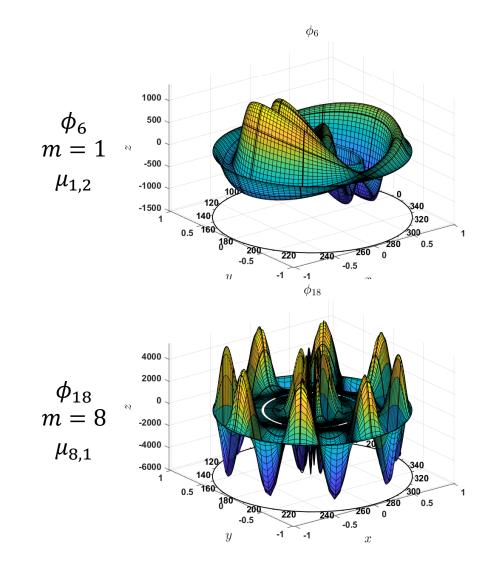
where j_m , i_m , and P_m are the spherical Bessel functions of the first kind, the Modified spherical Bessel functions of the first kind, and Legendre polynomials.

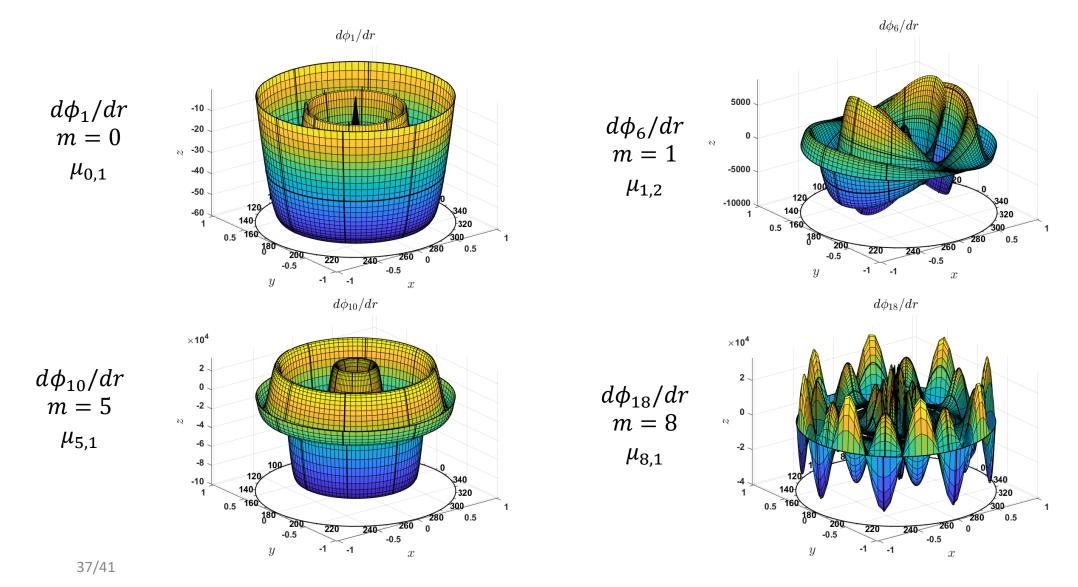
• Eigenvalues $\mu_{m,j}$ can be computed from the relation:

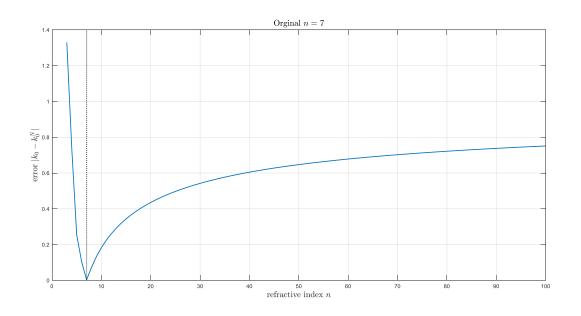
$$\det \begin{pmatrix} j_m(\mu_{m,j}r) & j'_m(\mu_{m,j}r) \\ i_m(\mu_{m,j}r) & i'_m(\mu_{m,j}r) \end{pmatrix} = 0, \qquad m = 0, 1, ...$$

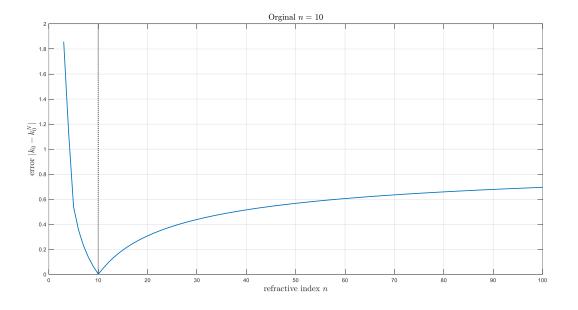
• **Note:** $\mu_{m,j}$ should be ordered and use the appropriate Bessel functions. j denotes j^{th} zero of above relation for each m.











Conclusion

Conclusion

- In this project we established the fundamental basis for the reconstruction of the acoustic refractive index from ITEP.
- The project has two main parts: A) Extraction of the ITEs from far-field data; and B) Estimation of the acoustic refractive index (n(x)) from the smallest ITE.
- 2D cases were the implementation of the previous works, **BUT** 3D cases were completely new and implemented during this project.
- For the first time, the FEM-BEM method has been used for this problem.

References

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