

## ABSTRACT

Testing of New Front-End Electronics for the Hadron Calorimeter at CERN

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The Hadron Calorimeter (HCAL), a major component of the Compact Muon Solenoid (CMS) detector at CERN, is designed to detect hadrons produced in proton-proton collisions supplied by the Large Hadron Collider (LHC). The HCAL is currently being upgraded with new readout modules that contain Silicon Photomultipliers (SiPMs), photodetectors that accept light from scintillator tiles and convert them to charge signals. Using a test beam capable of producing the same types of particles measured in the CMS detector, the new readout modules were tested in conditions like those of the CMS detector to fully understand their functionality before installation. Since the SiPMs in the new readout modules process the signals very differently than the hybrid photodiodes in the old readout modules, the analysis of test beam data provides insight into the features of the SiPMs. A virtual SiPM simulation was also created to assist in the analysis.

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TESTING OF NEW FRONT-END ELECTRONICS FOR THE HADRON  
CALORIMETER AT CERN

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## TABLE OF CONTENTS

1	Basic Particle Physics and CERN . . . . .	1
2	Large Hadron Collider and Compact Muon Solenoid . . . . .	7
2.1	The Large Hadron Collider . . . . .	7
2.2	The CMS Detector . . . . .	8
2.2.1	Hadron Calorimeter . . . . .	9
2.2.2	Silicon Photomultipliers . . . . .	12
3	Test Beam . . . . .	14
3.1	Test Beam Setup . . . . .	14
3.2	Test Beam Analysis . . . . .	17
4	SiPM Simulation . . . . .	29
4.1	SiPM Simulation . . . . .	29
4.2	Simulation Analysis . . . . .	30

## CHAPTER ONE

### Basic Particle Physics and CERN

Ever since ancient Greece people have been trying to discover what makes up everything. Back then they said that everything was comprised of four elements: earth, fire, water, and air. Although they use a few more than four elements, scientists today are still trying to categorize the most fundamental element that makes up everything. Democritus, an ancient Greek scientist, proposed that everything if divided enough times could be broken down into atoms their most fundamental object. Although some people were a little ambitious and named something the atom that was not the most fundamental substance scientists are still searching for the atom that Democritus theorized.

In the 19th century John Dalton believed he found the most fundamental particle and named it the atom, but as time passed there was very strong evidence that even atoms had internal components and structure. This was proven when J. J. Thompson in the late 19th century found the electron, which ended up being a component of the atom. This was not only the first fundamental particle discovered, but also the beginning of particle physics as a field of study. While it was an insightful discovery that led to many inventions, there were still many unknown things about the atom. About a decade after Thompson, Robert Millikan finely measured the charge and mass of the electron, but this raised many questions about the atom. It was not until Ernest Rutherford in the early 20th century showed that the atom is mostly empty space with a very dense core called the nucleus which has an opposite charge to that of the electron. He did this by shooting alpha particles, made up of two protons and two neutrons, at gold foil and observing their in angle. In this experiment most of the alpha particles passed through unperturbed while a handful

were majorly deflected. From this Rutherford made the conclusion most of the atom was empty space while having a highly dense core called the nucleus.

These discoveries showed there were more things to explore past the atom but in parallel there were equally important theories being made. Quantum mechanics is essential to the study of sub-atomic particles as they no longer behave according to classical mechanics. During this time Max Planck was doing his work in the quantization of energy. This set the ground work for quantum mechanics and also served as a basis to Einstein's work. In 1905 Einstein formulated his theory of special relativity, which had several major impacts on the field of particle physics. One the major ones was quantization of light into particles called photons. Although light had long been thought of as a wave Einstein showed many of its particle like properties leading to the wave-particle duality of light. Eventually De Broglie extended this duality to all particles theorizing that all particles including electrons and other sub-atomic particles could be described using waves or particles.

Einstein's theory of relativity had another major impact on particle physics. Probably the most famous of Einstein's equation,  $E = mc^2$  shows the mass energy relationship. One of the implications of quantum mechanics is that if something can happen it will happen under a certain probability. This means that given the right conditions and enough energy one can simply create mass. There are other conservation laws that need to be followed that are still relevant to particle physics, but conservation of energy is one of the biggest limiters as it says to create more massive particles more energy is needed.

As these theories became more complete, they allowed for the further description of sub-atomic particles, like the states of the electrons in the atoms and the protons and neutrons that make up the nucleus of the atom. Scientist also began to find evidence of fundamental particles not in atoms. In 1932 Carl Anderson found evidence of a Muon from cosmic rays. In addition, as scientists began to probe higher

energies they found evidence of internal structure within protons and neutrons. In 1968 at the Stanford Linear Accelerator Center they found evidence of particles within protons and neutrons, which were later called quarks.

The discovery of more and more fundamental particles led to the formulation of the standard model, which is like the periodic table of elements for elementary particles. The two main groups of particles on the standard model are fermions, which have half integer spin and bosons, which have integer spin. The main physical difference this causes is from the Pauli exclusion principle, which allows fermions to build into larger arrangements. Because of this fermions are the building blocks of the macroscopic world as they makeup protons neutrons, and atoms.

As shown in table 1.1 there are several different groups of particles in the standard model. The first group is a sub category of fermions called quarks. The main thing that distinguishes quarks from other fundamental fermions is that they interact via the strong force. While they also interact via the electromagnetic and weak force the strong force is the more dominate. Quarks can exist in isolation only for a short time, but rather exist in groups of two three or possibly more. When three quarks form a particle it is called a baryon. A proton which is made up of two up and one down quarks and a neutron which is made up of two down and one up quarks are baryons. Each of the quarks inside a baryon has a different color red, green, or blue which serves as the "charge" for the strong force. Like electric charge strong force color attracts other colors and repels similar, but since the baryon has one of each it has a net color of white or nothing. In addition to the six flavors of quark listed on table 1.1 there is also a corresponding anti-quark with the same properties except for opposite charge and color. Quarks can also form in pairs called mesons, which have a quark and anti-quark. If one tries to separate the individual quarks from the respective baryons or mesons the energy required to separate them is more than the energy required to simply create more quarks. As such if something like a proton was

split into its constituent quarks, the input energy would simply form into new quarks so none of them would be alone.

The other category of fermionic particles are the leptons. Unlike the quarks, leptons are not affected by the strong force so it is more common to find them in isolation. The electron is the most well-known lepton but there are also two other flavors which are heavier copies of the electron. As with all of the particles on the standard model there are corresponding anti-particles for the each lepton with the same mass but opposite charge. There are also three neutrino flavors corresponding to each other lepton flavor. Neutrinos are very near mass-less and do not have any charge. Since the only force that affects them is the weak force they have a very low probability of interaction. This means that they essentially pass through everything without interaction. Although they do interact through processes like beta decay most neutrinos simply just go through everything.

Finally, there are the bosons. Bosons are different from both leptons and quarks in that they have integer spin, but the individual bosons differ in which fundamental force affects them. While leptons and quarks combine together to form complex arrangements such as atoms and all of matter, bosons serve as the force carriers. The photon serves as the mediator for the electromagnetic force. This means that whenever two electrons repel each other there is a photon to mediate that interaction. The strong force interaction is mediated by the gluon. The weak force is mediated by the W boson, for those interactions with changes in charge and Z bosons for those interaction with no changes in charge.

The standard model essentially serves as a list of known fundamental particles but the road to completing it was not an easy one. The electron is the only particle on the standard model that is found in isolation at energies used on earth. To find the next particle, the muon, scientist had to look at cosmic rays, which provided the energy needed to create these particles. While neutrinos are plentiful they interact



Table 1.1. Particles on the Standard Model

Particle Type	Name	Charge	Mass	Spin
Quarks	Up	$+2/3$	$2.2 \text{ MeV}/c^2$	$\pm 1/2$
	Down	$-1/3$	$4.7 \text{ MeV}/c^2$	$\pm 1/2$
	Charm	$+2/3$	$1.28 \text{ GeV}/c^2$	$\pm 1/2$
	Strange	$-1/3$	$96 \text{ MeV}/c^2$	$\pm 1/2$
	Top	$+2/3$	$173.1 \text{ GeV}/c^2$	$\pm 1/2$
	Bottom	$-1/3$	$4.18 \text{ GeV}/c^2$	$\pm 1/2$
Leptons	Electron	$-1$	$0.511 \text{ MeV}/c^2$	$\pm 1/2$
	Muon	$-1$	$105.7 \text{ MeV}/c^2$	$\pm 1/2$
	Tau	$-1$	$1.776 \text{ GeV}/c^2$	$\pm 1/2$
	Electron Neutrino	$0$	$\approx 0$	$\pm 1/2$
	Muon Neutrino	$0$	$\approx 0$	$\pm 1/2$
	Tau Neutrino	$0$	$\approx 0$	$\pm 1/2$
Gauge Bosons	Photon	$0$	$0$	$\pm 1$
	Gluon	$0$	$0$	$\pm 1$
	W boson	$\pm 1$	$80.4 \text{ GeV}/c^2$	$\pm 1$
	Z boson	$0$	$91.19 \text{ GeV}/c^2$	$\pm 1$
Scalar Boson	Higgs Boson	$0$	$125.1 \text{ GeV}/c^2$	$0$

so infrequently scientist are still struggling to create an efficient detector for them. To discover the remaining particles on the standard model they had to use particle colliders. Throughout the last few decades, scientists have built colliders to probe higher and higher energies, to make them capable of discovering more fundamental particles. In 2012, using the Large Hadron Collider at CERN scientists were able to find evidence of the Higgs boson, the final particle on the standard model. However, there are still phenomenon in the universe that cannot be explained using the standard, such as gravity and dark matter. There are several theories beyond the standard model that offer explanations for these, but they also say there are more fundamental particles to be discovered at higher energies. This is one of the main goals of the CERN collaboration.

Fermions and quarks are the building blocks of all of matter. All of the atoms on the periodic table are made up of electronics, and up and down quarks, which

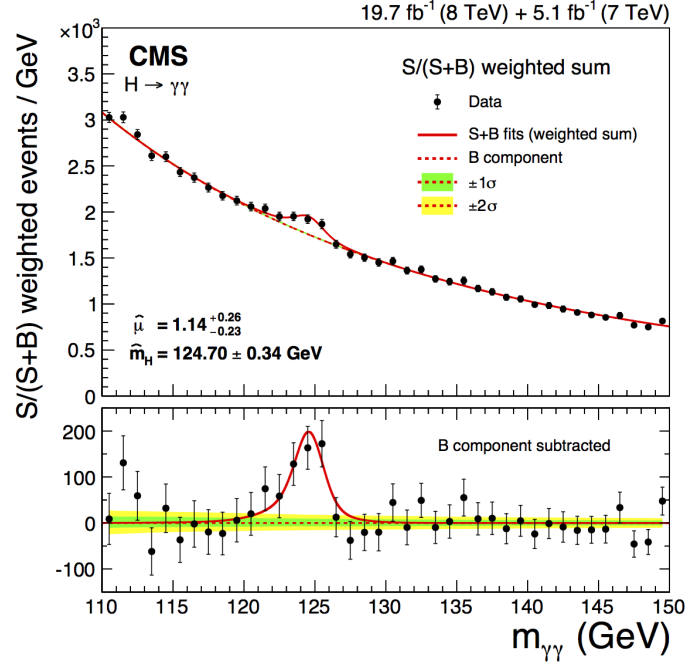


Figure 1.1. Data showing evidence of the Higgs Boson from a two photon decay.

make up protons and neutrons. Aside from the neutrinos, the particles in higher generations are extremely similar to the respective particle in the first generation except for being a substantially higher mass. Because of this these particles have very short lifetimes and are only found in high energy interactions. For instance, when cosmic rays hit the earth's atmosphere they provide the energy to create muons. To find the particles like the top quark, which is over 1000 times more massive than the muon, the colliders need to collide particles together at even higher energies. The Large Hadron Collider at CERN, is the highest energy collider humans have built to date and has started searching for particles at energies previously not achieved. One of the detectors on the LHC is the Compact Muon solenoid, which is the detector I worked on during my undergraduate career at Baylor. Chapter 2 will discuss the details of the LHC and CMS detector as more specific background to my work.

## CHAPTER TWO

### Large Hadron Collider and Compact Muon Solenoid

#### 2.1 *The Large Hadron Collider*

The process the Large Hadron Collider (LHC) uses sounds simple when put into common terms. It accelerated protons to speeds very close to the speed of light and collides them in the center of a detector to see what comes out of these collisions. However actually doing this is not simple at all. The theory of special relativity says that a particle can only approach the speed of light. As the protons in the LHC are traveling close to the speed of light it is more useful to use their energy rather than their speed. In addition, mass is often put into units of energy of the speed of light squared making the mass to energy conversion simple. To start this entire process electrons are stripped off of hydrogen gas supplying the protons for the LHC. After this the protons are put into a series of different accelerator each designed to accelerator to protons to higher and higher energies. The first accelerator is the Linac2 linear accelerator, which gets the protons up to 50 MeV. Then they are sent to the Proton Synchrotron Booster which can push them to 1.4 GeV, after which the Proton Synchrotron ring accelerates the protons to 25 GeV. At this point the protons are sorted to control the frequency at which the collisions will occur. The protons are sorted into bunches such that given a fixed point on the LHC a proton bunch passes by every 25ns. Each bunch has about 100 billion protons in it. After sorted into these bunches the protons are sent to the Super Proton Synchrotron where they achieve an energy of 450 GeV. Finally, the protons are fed into the LHC where they will be accelerated to their highest energy.

When the LHC has accelerated its particles to its peak energy it collides the particles in the middle of four different detectors. LHCb, ALICE, ATLAS, and CMS

CERN's Accelerator Complex

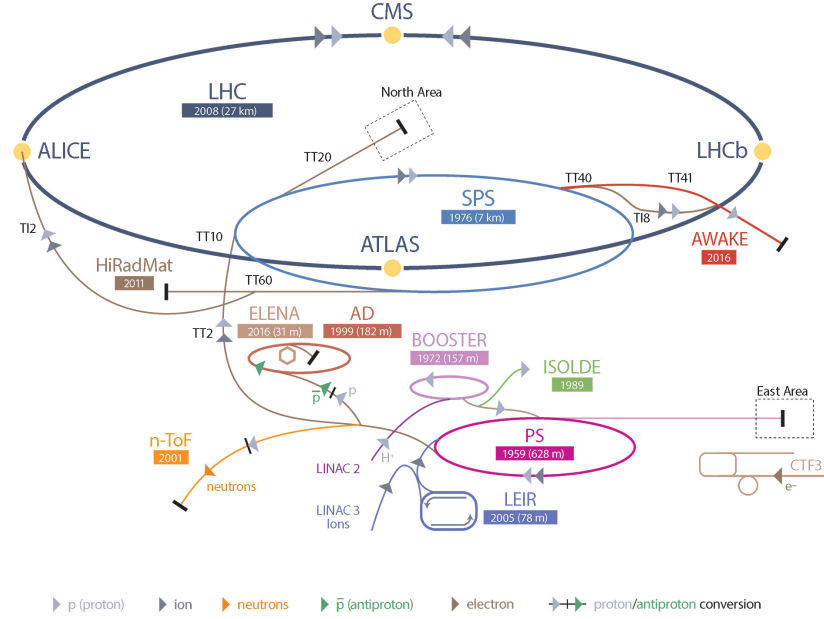


Figure 2.1. Chain of accelerators feeding into the LHC

are the names of the four detectors that were built to study collisions supplied by the LHC. The CMS detector is the one I worked on during my undergraduate career, and is discussed in detail below.

## 2.2 The CMS Detector

The CMS detector is designed to detect particles coming out of the proton-proton collisions supplied by the LHC. As many different types of particles come out of these collisions the CMS detector is split up into several different sub-detectors. Going from the closest to the collision point outward, first is the silicon tracker. This is responsible for measuring the total energy of the particles in the detector which is important for reconstructing the events using conservation of energy. Next is the Electromagnetic Calorimeter which is responsible for detecting a myriad of charged particles such as photon, and electrons. Then there is the Hadron Calorimeter which is responsible for detector hadrons, which are particles made up of quarks, such

as protons neutrons and pions. These calorimeters work by covering the possible directions out of the collision with scintillator tiles. When the particles hit these tiles they lose some of their energy depositing light proportional to the energy of the particle. The particle will eventually stop having lost all of its energy giving us a reading on its total energy. Finally, are the muon chambers. Muons are interesting particles because while they are heavy particles that will decay they have long enough lifetimes to make it through the detector. In addition, they tend to not be stopped by the detector like the two prior calorimeters almost always stop the particles they detect. Since there is a magnetic field encompassing the detector and the muons have a charge they will curve as they pass through the muon chambers. While not completely stopping the muons will give a reading in the chambers showing its path of travel. By measuring the curve of this path of travel we can measure the momentum. By combining the the data from all the different sub detectors we can get a picture of almost all of the particles that came out of the collision. There are exceptions such as neutrinos which tend to pass through anything leaving no trace. Nevertheless we are able to use this to recreate what exactly came out of the collision, which allows us to find evidence of particles like the top quark which decays so quickly it does not even make it to the detector.

### *2.2.1 Hadron Calorimeter*

One of the main jobs of the calorimeter is to measure the energy of the of the incident particle. When the particle hits the scintillator tile and loses energy it emits photons into the tile proportional to the amount of energy it lost. To measure the energy of the particle we just have to count up the photons it deposited. To do this we put a fiber around the edges of the tile so that any photons deposited will run into this fiber. The fiber is designed so that once it enters the fiber it will have total internal reflection on the edges meaning it will stay inside the fiber. The fiber then runs to a silicon photomultiplier (SiPM). The SiPM is designed to take light signals

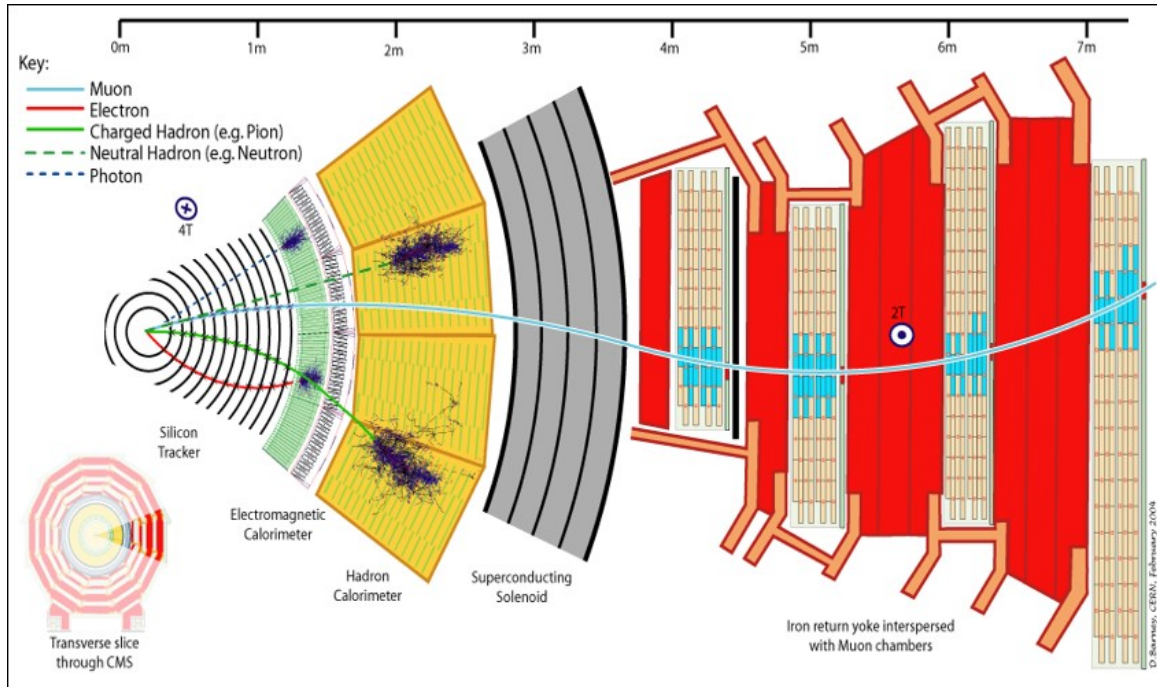


Figure 2.2. A slice of the CMS detector highlighting the different sub-detectors and showing different particles and where they are stopped

and convert them to charge signals. The light from the scintillator tiles is shined on what is called the pixel face of the SiPM. This is just a circular array of really small pixels. There are about 33000 pixels on a 3.3 mm diameter pixels face. When an individual photon hits one of the pixels on the pixel face it causes a capacitor to fire off a particular amount of charge. The total output charge of the SiPM is the sum of all of the pixels output charge. Theoretically one just needs to sum up the total amount of output charge of the SiPM and one can count the number of incident photons since each pixel will fire off a set amount of charge. The counting of the number of incident photons will give a measurement of the incident particles energy. There are some things that make the SiPMs reading not so simple. SiPMs are non-linear devices meaning the total output charge of the SiPM does not necessarily increase linearly with the number of incident photons. In fact, experiments have shown that graphing the number of incident photons vs the SiPM output charge it looks more like a square

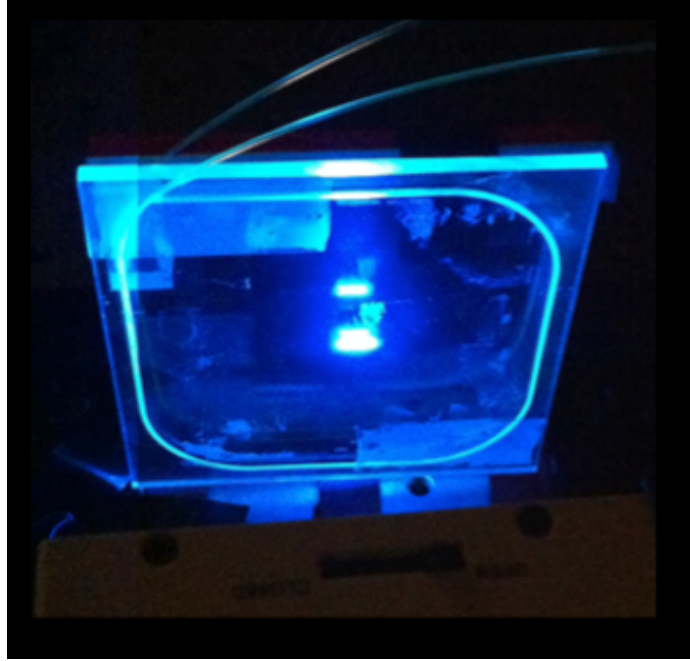


Figure 2.3. A picture of a scintillator tile with the optical fiber around the edges.

root function rather than a straight line. At low incident photon count on the order of 1000 photons the SiPM is very close to a linear device but as the incident number of photons increases the output charge does not increase at the same rate.

There are three main section to the HCAL. There is the Hadron Barrel (HB), which surrounds the beam-line like the edge of a cylinder, the Hadron End-cap (HE), which caps off the cylinder made by HB. These two parts make a cylinder which has the beam-line going through the center and the collision point in the center. Lastly there is the Hadron Front-cap (HF) which is shaped like HE but is located much further away from the collision point closer to the Muon chamber distance. To describe the geometry of the detector we use a coordinate system related to spherical coordinates with the origin being the collision point and polar angle of 0 being along the beam-line. Phi is the azimuthal angle and Eta is related to the polar angle by  $\eta = -\ln(\tan(\theta/2))$  which gives a value of 0 perpendicular to the beam-line and  $\pm\infty$  along the beam-line. Usually these coordinates are arranged into discrete integers

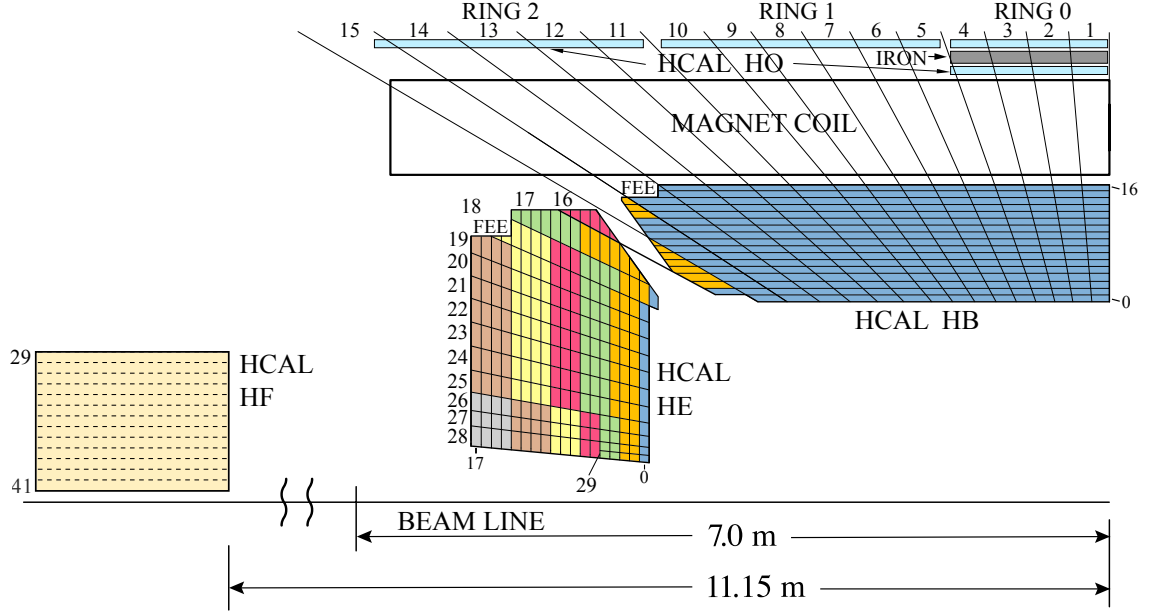


Figure 2.4. Layout of the HCAL showing a single iphi slice of HE HB and HF. Each box represents a scintillator tile

values denoted iphi and ieta, which are arranged such that a scintillator tile covers one ieta and one iphi as shown in figure 2.4 but there are exceptions to this. For the distance from the beam-line depth is used which is also a discrete value but there are often more than one scintillator tiles in a single depth.

### 2.2.2 Silicon Photomultipliers

There are several factors that contribute to SiPM non-linearity but the two main ones are cross talk and saturation. Cross talk is when a pixel is activated by a photon there is a chance that this activated pixel could activate some of its neighboring pixels artificially increase the output charge of the SiPM. When the photon hits a pixel it excites an electron causing an electron cascade and charge to flow. This excited electron has a chance to emit an IR photon which could go in any direction. If it hits one of the neighboring pixels it could activate it just like an incident photon. Saturation is when a pixel is hit in rapid succession by multiple photons. When a pixel is hit by a photon it discharges its capacitor which is the source of its output



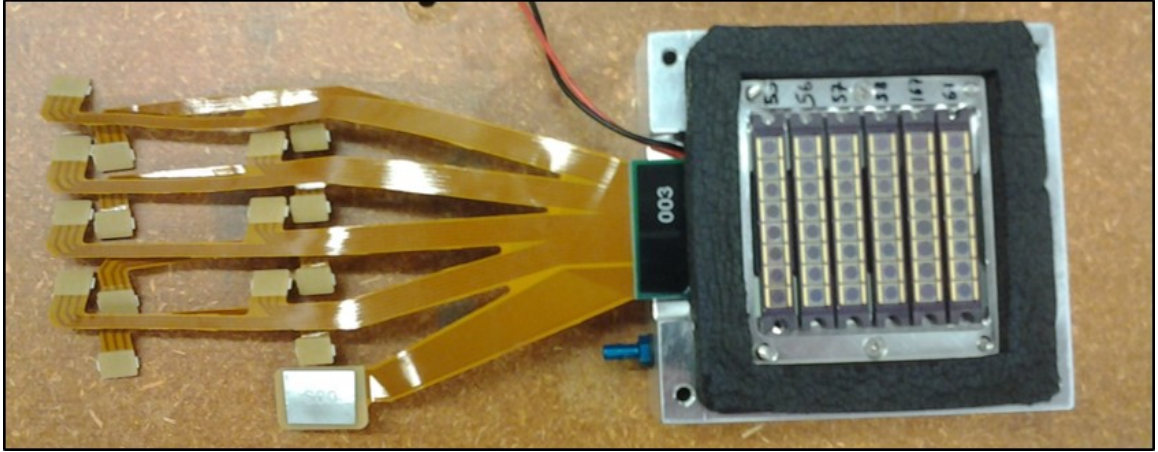


Figure 2.5. A picture of a Silicon Photomultiplier

charge. After it fires it needs some time on the order of 10 nanoseconds to recharge its capacitor. If the pixel is hit it simply fires off whatever charge it has in its capacitor at the time which could be reduced from its normal charge. Given that the SiPM has about 33000 pixels on its pixel face when the number of incident is on the order of 1000 this effect is insignificant but particles can produce much more photons where this effect can be much more significant. The non-linear properties of the SiPM makes it a bit harder to get accurate measurements, but there are ways to study them which will allow us to compensate for these affects.

## CHAPTER THREE

### Test Beam

#### *3.1 Test Beam Setup*

SiPMs are one of the main features of new electronics that are being installed in the CMS detector. These new readout modules main function is to take the signals from the detector, the light signals from incident particles and convert them to digitized charge signals which can be later stored and analyzed. There are several new features with these new readout modules such as new charge integrator and encoder (QIE) chips, but the main subject of my work was the SiPMs. The main way the SiPM non-linearity and other things about the new readout modules are being study is with a test beam. The test beam is a linear particle beam that comes from the SPS accelerator when it does not feed into the LHC. The particles from this beam are not at the energies of the protons in the LHC but they are much easier to make and control. There are several different experiments that use this beam but the HCAL has a mock test stand that can be moved into the path of the beam when experiments need to be run. The CMS detector is 100 meters underground and is in a cavern that due to radiation and tight enclosures makes the detector difficult to access. Though modifications and maintenance can be during shutdowns this is not the ideal way to run tests on new electronics. The test stand at the test beam location called H2 can be easily accessed observed and is on ground level. In addition the energy and particle type of the beam can be easily controlled. The test stand is designed to be similar to a portion of the HCAL on CMS. As such when we shoot particles from the test beam at the test stand with the new electronics on it we can measure its response to experiment like conditions. Using this test beam we can shoot particles like pions at the test stand which has the new readout module installed and

look at the response. We can vary the energy of the pions which should increase the number of incident photons on the SiPMs in the readout modules and measure their response to this. In this way we can find the non-linearity of the SiPMs and create a correction curve based on this.

Because we are looking for statistical in the data we need to take a lot of data. This means we take runs with several thousand events or particle hitting the test stand and take several of these runs to ensure high statistics. On the other hand this means we want to minimize the amount of information stored per event so we can store a lot of them. One of the key systems that helps this is the trigger system. In the path of the beam before the test stand we put scintillator tiles not designed to stop the particle entirely but rather to send a generic signal to the control room. A total of four tiles were placed in the path of the beam. When we saw signals from these tiles we know that a particle is coming from the beam. This serves as our trigger signal. The trigger signal then goes to the back-end electronics that stores the data from the detector. In this way we only store data coming during the short amount of time on the order of 25 nanoseconds when we know there is relevant information coming. There is also a ready signal that goes from the back-end electronics to the trigger. Since a particle may come before the back-end electronics are finished storing the relevant data a trigger signal that comes without the ready signal will simply be ignored.

Although built to be similar to the HCAL on the CMS detector the test stand is smaller and has some differences. When looking at the data from the test beam it is important to be sure that the e-map, which traces the channels that receive the data to the actual scintillator tiles on the detector, is correct. To account for the difference between the test stand and the HCAL a new e-map needed to be created. With an accurate e-map we can create an accurate picture of the incident particles hitting the test stand. The test beam allows us to control the energy and type of



Figure 3.1. A picture of a the test stand at H2. The scintillator tiles are covered in the black tedlar and the new readout modules are connected in the back. The stand can be shifted so the beam can be aimed at different sections.

particles incident on the test stand. This allows us to study things like SiPM non-linearity as we can look at the output of the SiPMs under something like 50 GeV pions then increase to energy to something like 150 GeV and compare the different responses. However the particle does not deposit its energy in a single scintillator tile. For something like a pion it will lose a part of its energy with each scintillator tile it hits and eventually be stopped by the detector but its total energy will be spread throughout the detector. Since each particle in the beam has roughly the same energy and we can isolate the event of a single particle hitting the detector we are to find energy lost in each scintillator tile but looking at the complete picture of the detector.

SiPM non-linearity is not the only thing that is being researched using the test beam. In the CMS detector particles come so fast that the signals of individual particles blend together. In the test beam the particles are more spread out so it is easier to isolate the signal of individual particles. This makes it easier to extract the

pulse shape of the readout modules which is the shape of the output charge vs time. It is also important to see if the pulse shape changes significantly under different circumstances such as higher or lower energies. Using the pulse shape of the readout modules we can then extract the individual signals in the CMS detector.

### 3.2 *Test Beam Analysis*

With a proper e-map we can get a picture of the incident particle and observe its path and where it deposited its energy. Since the particle does not deposit all of its energy in a single scintillator tile this is critical when trying to recreate the particle energy. It is also important that there are often a variable amount of scintillator tiles in a single depth channel as shown in figure 3.2 which shows 4 scintillator tiles in depth 5 and 3 scintillator tiles in depth 4. Looking at the results in the plot to the left it explains why there is more energy deposited in depth 5 rather than depth 4.

Muons are useful for many things but when trying to recreate the energy of the particle it is easier to use pions. The reason for this is that muons tend to go through the detector depositing a small portion of the energy but not being entirely stopped. This can be seen in figure 3.2 which shows the muons depositing a consistent amount of energy in each of the depths even the far ones. On the other hand figure 3.3 shows a similar plots with a pion run. This shows that the pions hit the scintillator tiles in depth 2 depositing a large portion of their energy and depositing less and less in the next depths with basically nothing in the last depth. There is also more significant spread compared to the muon run which shows the muons going straight through while the pions leave some energy in the neighboring channels.

To recreate the energy of the incident particle we can look at plots such as figure 3.4 which shows the mean output charge in a 50GeV pion run. By looking at this output charge and the output charge of the other channels we can find what portion of the particles energy was deposited in this channel and then going by the

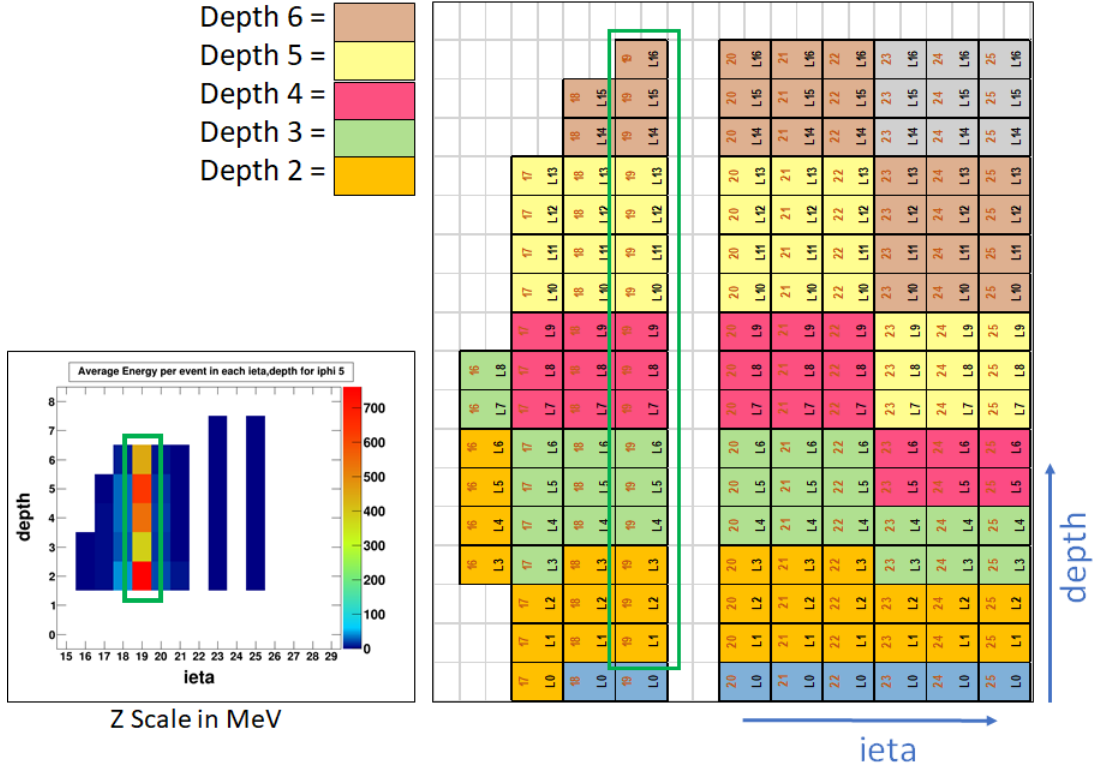


Figure 3.2. This shows the relationship between the emap on the left showing which contains a geometric layout of the scintillator tiles in a single iphi slice and the data plots in a run with 150 GeV muons.

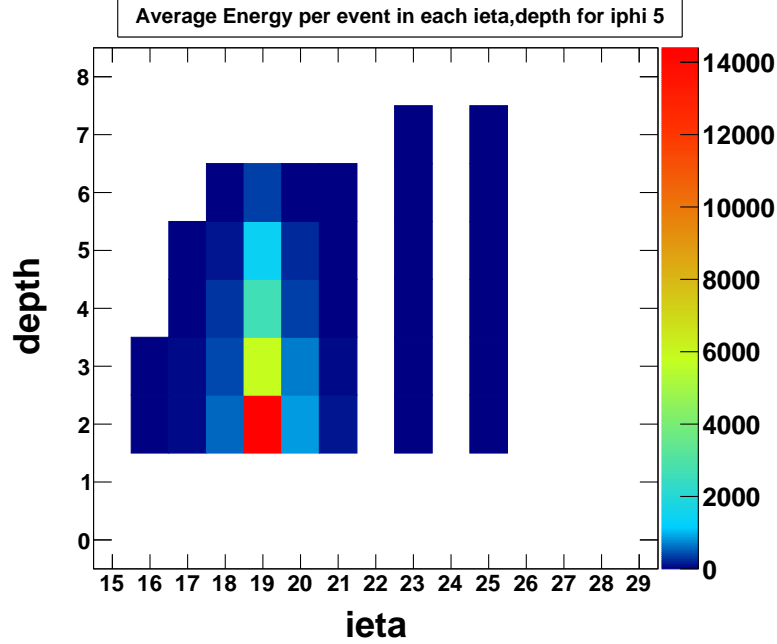


Figure 3.3. This shows the recreation of an event with 50GeV pions showing the pions aimed at iphi 5 ieta 19 and the energy deposited in each channel z scale in MeV.

known energy of the particle how much energy was deposited to produce this output charge.

One of the key things in doing a pulse shape analysis using test beam data is the timing information. In the test beam we take two different timing information. The first comes from the QIE chips in the readout modules. The readout modules measure the output charge of the SiPM over 250ns and the information is binned into 10 time samples each 25ns long. The output pulse of the SiPM is about 75ns long so it usually stays confined to 3 or 4 time samples. In addition we also take what is called a Time to Digital Converter (TDC) value. This value stores at what time in the 250ns span the output charge of the SiPM crossed a threshold with a 0.5ns resolution. Since the output charge starts out below this threshold, the TDC value should give the the start of the output pulse of the SiPM. For instance, a TDC value of 20 in time sample 4 means the pulse crossed the threshold charge value 10

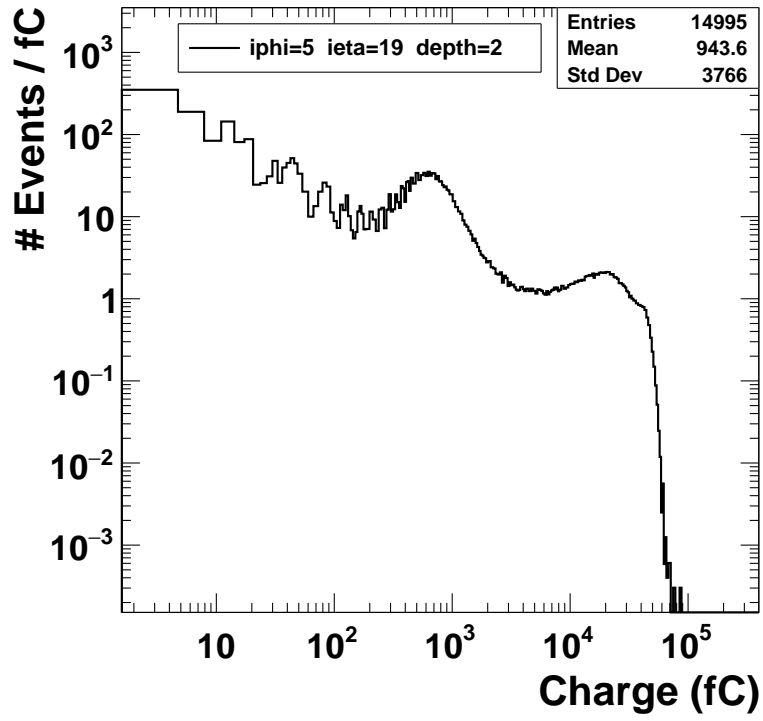


Figure 3.4. This shows the number of events that produces output charge in the SiPM. The peak in the middle shows the average output charge in response to the incident particles at this energy.



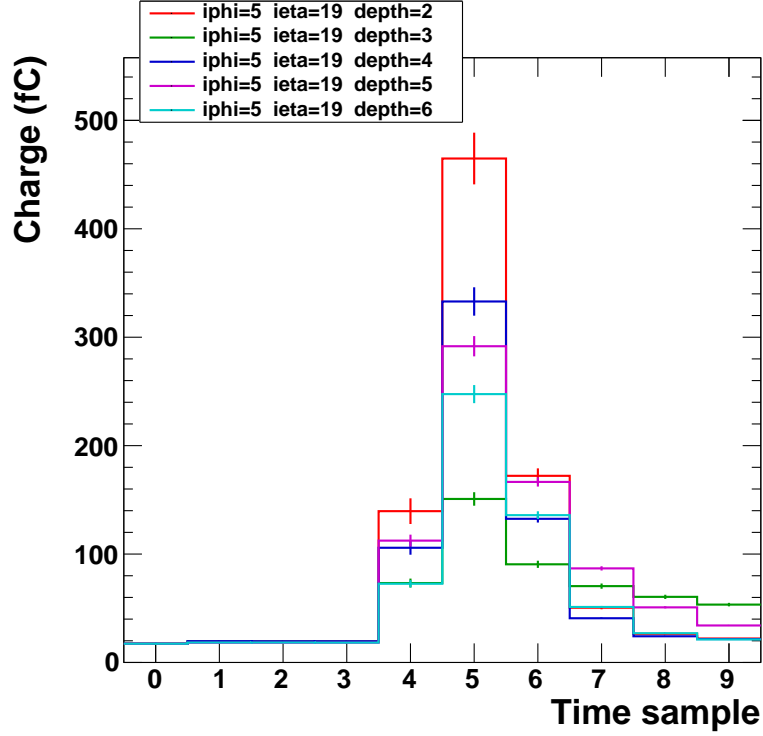


Figure 3.5. A histogram of the output pulse of the readout modules binned in 25ns bins. This is from a run with 150GeV muons aimed at iphi5 and ieta 19. This plot shows the output pulse of the different depths at that location.

nanoseconds after the start of time sample 4 or 110 nanoseconds after the start of data taking. One of the problems with the TDC value is that it could be affected by the amplitude of the output pulse as a high amplitude will cross this threshold earlier. To extract the pulse shape we can do something called a phase scan.

When we take a run using the test beam we do not just take data from one single particle but several thousand stored in their own individual events. For instance, when we took a run with 250 GeV pions that run has several thousand events recording the data when the particles all at about 250 GeV hit our test stand. Because of the nature of the test beam and also just quantum mechanics there is some randomness to each event. To do a phase scan we take advantage of this randomness by looking at pulses with the same total charge but different TDC values. Theoretically,

two pulses with the same charge but different TDC values should be exactly the same except for a time shift. These pulses could look different in data though as different parts of the pulse will be put into different time samples. A phase scan takes the difference in binned pulse value from several different pulses all with similar overall charge but a range of TDC values. An example of this is taking the difference of the charge value from time sample 4 in a pulse with a TDC value of 20 also in time sample 4 and the charge value from time sample 4 in a pulse with a TDC value of 21 also in time sample 4 gives the pulse height at nanosecond 110 in the 250 nanosecond long data set. Doing this process over all the time samples with more pulses with TDC values ranging from 0-49 all in time sample 4 will give a more analog version of the pulse shape as shown in figure 3.6.

There is also the trigger information. This timing information comes from the test beam equipment that signals the back end electronics to store the data from the readout modules. The trigger information theoretically does not have the same discrepancy as TDC information but most people still use the TDC from the QIE chips. The relationship between the two signals can be shown in figure 3.7. Theoretically this timing information could be used for a phase to obtain an analog pulse shape like that in figure 3.6.

To do a phase scan properly we need to use events that have very similar output charge. There are a variety of reasons for this but one of them is that there is evidence that the pulse shape may change with output charge. To study this we simply have to do a phase scan using event with different output charge and see if there are any significant changes. Although the particles will have about the energy the test beam is set to it is common to have several outliers. To minimize this disturbance we take in all the events from several different runs energies ranging all the way from 30 GeV to 300 GeV and only use the events that fall within a certain charge range. Figures 3.8 - 3.12 show pulse shapes obtained using several different charge ranges.

**Fit Testbeam Data**

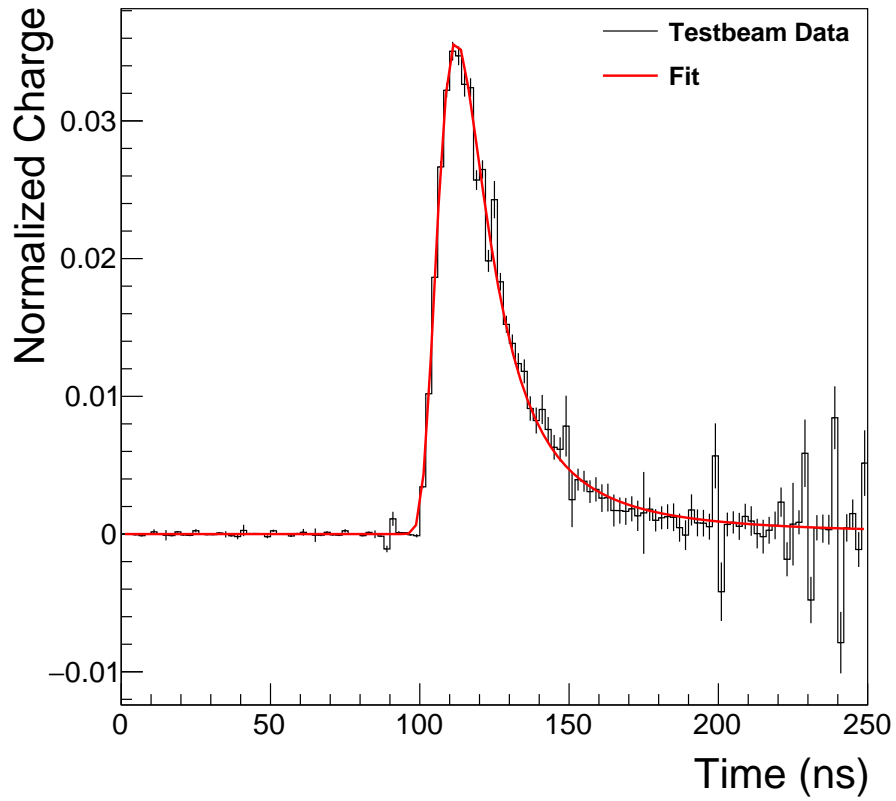


Figure 3.6. A plot showing a more analog version of the SiPM pulse shape obtained using a phase scan on test beam data. It is also fitted by a Landau Gaussian function which is what we expected from theory. The total output charge of the runs used in this range from 10,000-29,000 femtocoulombs.

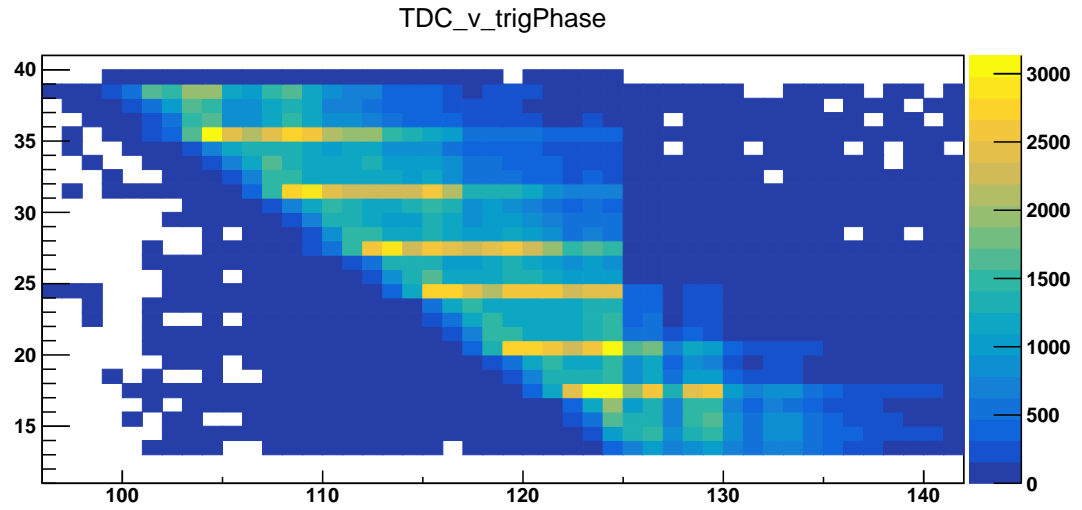


Figure 3.7. A histogram showing the statistical relationship between the trigger timing information and the QIE TDC information.

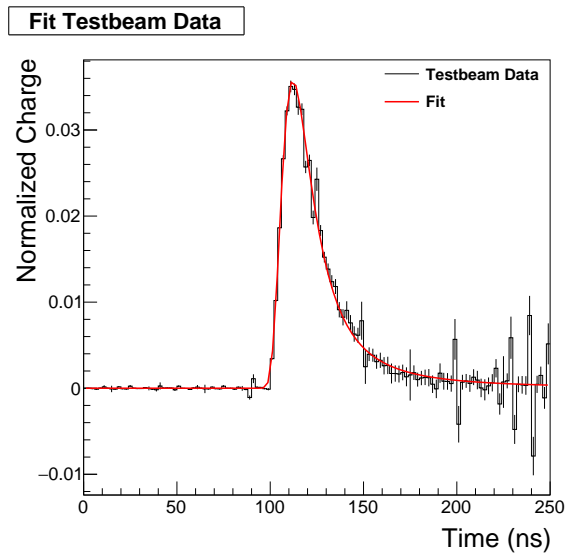


Figure 3.8. Fitted pulse shape 10,000-29,000 fC

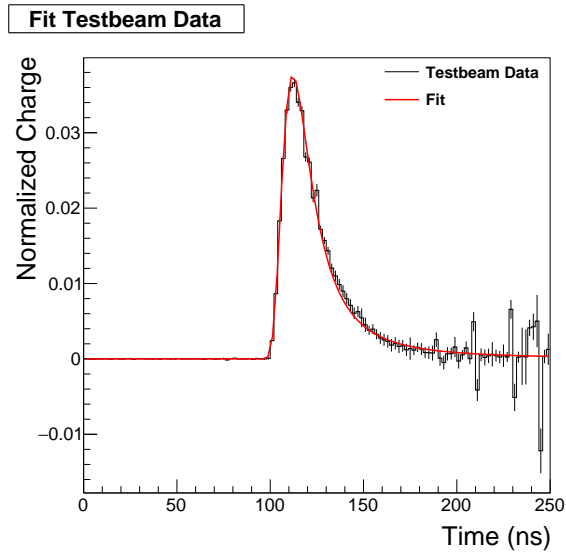


Figure 3.9. Fitted pulse shape 29,000-50,000 fC

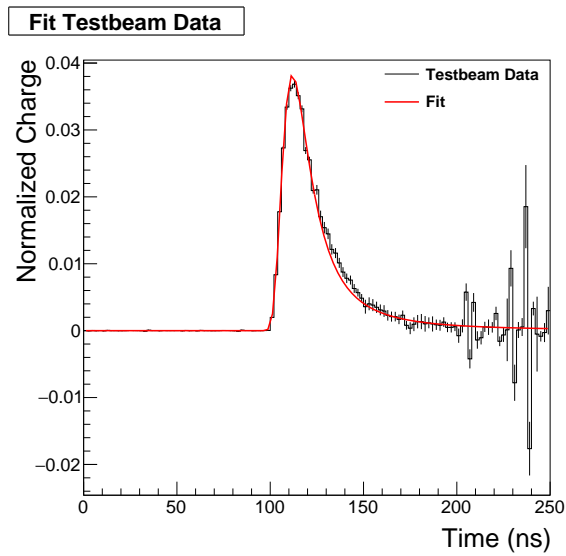


Figure 3.10. Fitted pulse shape 50,000-80,000 fC

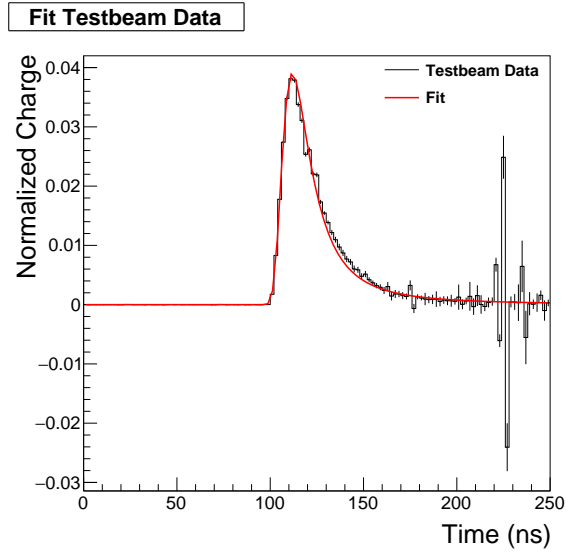


Figure 3.11. Fitted pulse shape 80,000-125,000 fC

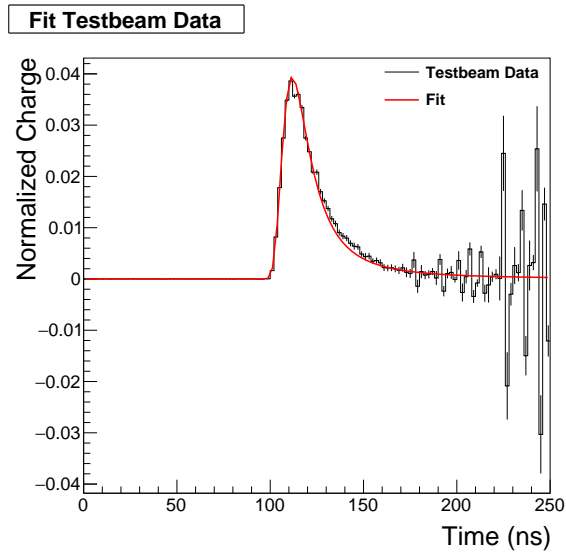


Figure 3.12. Fitted pulse shape 125,000-168,000 fC

The output charges of the different events can be shown in figure 3.13 which shows the number of events that output a certain charge. This figure can also be used to start the analysis on the SiPM non-linearity. One key thing about the non-linearity is that at low charge output and correspondingly low particle energy the SiPM is very near linear, and it is only at high output charge does this effect become significant. This means to study non-linearity one needs to use particles at sufficiently high energies. After the test beam was over and the analysis of the data was started it was discovered that the test beam was not really capable of producing particles at energies in the non-linear range of the SiPM. To illustrate this looking at figure 3.13 it is apparent that in the 300 GeV run, an energy close to the maximum the test beam is capable of producing, there are minimally few events that output charge above 300,000 fC. The SiPM outputs approximately 40 fC for every incident photon electron which means the highest energy events create about 6000 incident photo electrons. When using a simulation of the SiPM it is apparent that this is still very much in the linear range of the SiPM. One thing to note is that not all of a single particle's energy goes to a single SiPM. In figure 3.4 is shows that while the majority of the particle's energy is deposited in the first group of scintillator tiles it hits there is still a significant portion that is in neighboring channels. Summing over all the channels should give the total energy of the particle but this means that 300,000 fC does not directly translate to 300 GeV. As the energy of the particle increases its energy will spread out more which means a 600 GeV pion will not necessarily deposit 12000 photo electrons in the first channel it hits.

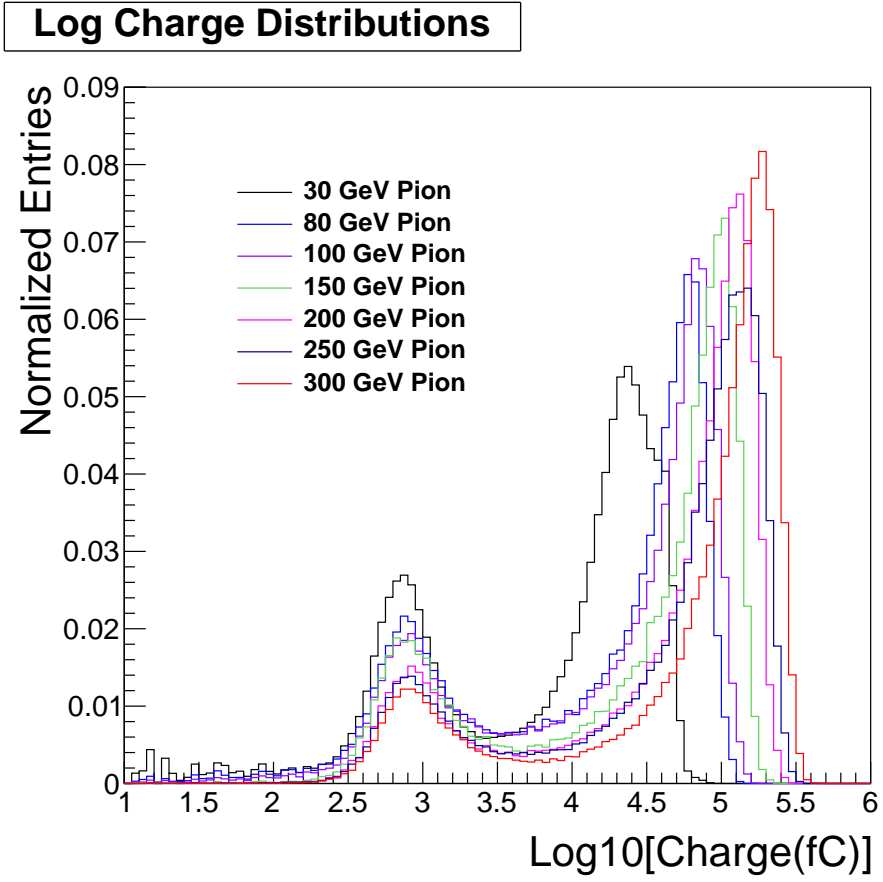


Figure 3.13. A plot showing the charge output of the different events over the different runs used to find the pulse shape.



## CHAPTER FOUR

### SiPM Simulation

#### 4.1 *SiPM Simulation*

In addition to studying the new electronics with the test beam, we also did a SiPM simulation which could use theoretical signal from the detector and give the SiPMs response to them. When the light from the detector goes to the SiPM a Y11 light mixer is used so that the light is spread across the pixel face. The light mixer has a unique shape in which it shines the light onto the pixel face. Using a mathematical representation of this pulse shape we can simulate how light is shined on the pixel face. We can then create a virtual pixel face that has accurate geometric representation of the pixel face. Each of these virtual pixels could be activated by the incoming Y11 light pulse and would respond by firing off a set amount of charge just like an actual SiPM would. When these pixels are activated they also have a certain probability to activate a neighboring pixel which will simulate the cross talk of the SiPM. Each pixel also has a recharge rate so if it is hit in rapid succession it will fire with a reduced amount of charge emulating the saturation effect. Using this simulation we easily adjust the number of input photons and get very good details of the SiPM response. We can also easily turn different things off to see how significant individual effects are. This simulation is capable of producing more fine tuned data and in greater quantities than test beam data. In addition, it is able to simulate situations that are not possible in the test beam. On the other hand the accuracy of this simulation is heavily dependent on the data we input into it.

Since the program can simulate the non-linearity effects of the SiPM, we can use theoretical inputs to the SiPM to study its effects. The light a particle deposits in a scintillator tile can vary in time the Y11 light mixer ensures that the input light

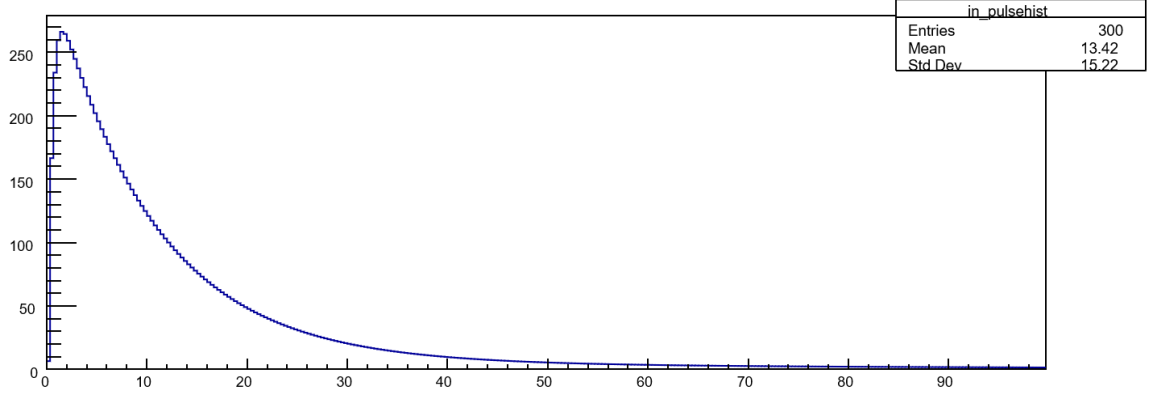


Figure 4.1. A graph of the Y11 pulse shape with nanoseconds on the x axis and number of incident photons on the y axis

pulse shape remains constant. This means to change energy of the incident particle we just have to increase the number of photons in the input pulse while keeping the same shape. The output of the virtual SiPM is normalized so if the simulation was a completely linear device its output "charge" should exactly equal the number of input photons.

## 4.2 Simulation Analysis

To do the analysis using the SiPM simulation, we simply measure the total output of the SiPM over a light pulse and compare it to the number of input photons. Since there is a degree of randomness we will use the average pulse over several inputs. Compared to using test beam data the analysis with the simulation is simple. In test beam the energy of the particle is spread over several SiPMs while in the simulation we have a very fine control of the input energy to the single SiPM. With a plot of input number of photons vs SiPM output charge we then create a correction curve that is necessary to make the graph linear. This is process for doing a non-linearity analysis using test beam data so we can compare the two correction curves to see if there is agreement. With other analyses it is common to simply divide out the cross talk. Given that there is a 15.4% chance of crosstalk with a large signal of several

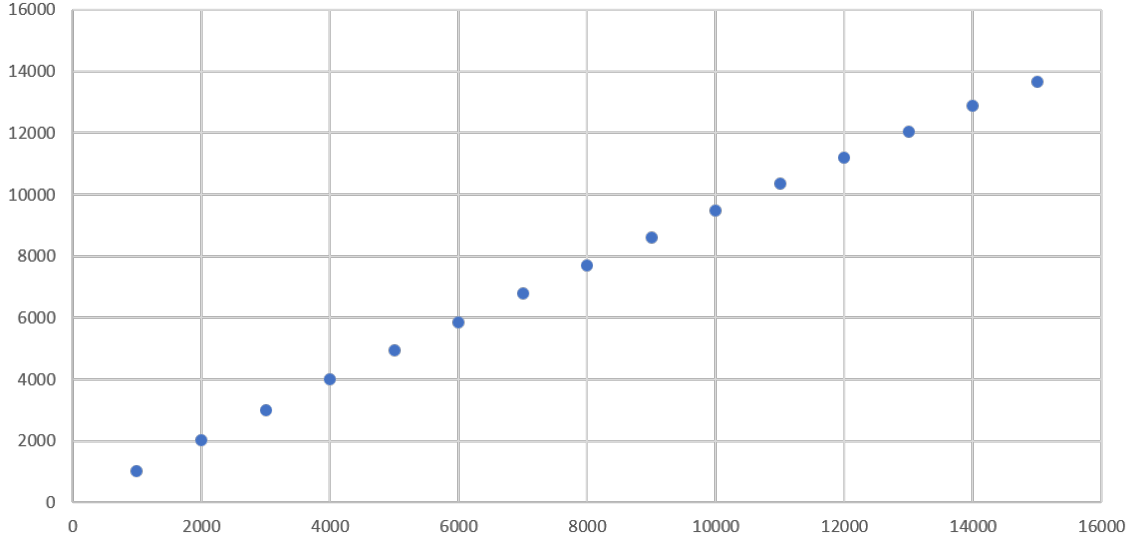


Figure 4.2. Graph of number of incident photons vs output charge of the SiPM simulation. Given the units of the simulation a perfectly linear device would have all data points falling on the  $y=x$  line but as shown as the number of incident photons is increased the data points fall short of this line.

thousand photons it is assumed that 15.4% of the signal from the SiPM is due to cross talk so that number is simply divided out. This means the majority of the effects that the correction curve is compensating for is saturation.

To be sure the simulation is designed correctly it is useful to compare things like non-linearity to the actual SiPM. Figure 4.3 shows the non-linearity of the SiPM from shining a laser directly on the SiPM. Figure 4.4 shows a similar plot constructed using the simulation. There is significant disagreement between these two sets of data.

Using the simulation we can also look at the pulse shape of the SiPM. This pulse shape should theoretically be the same obtained from test beam data. One of the main things that is necessary to look at is does the pulse shape change significantly with a change in input energy. As shown in figure 4.5 does not significantly change with different energies.

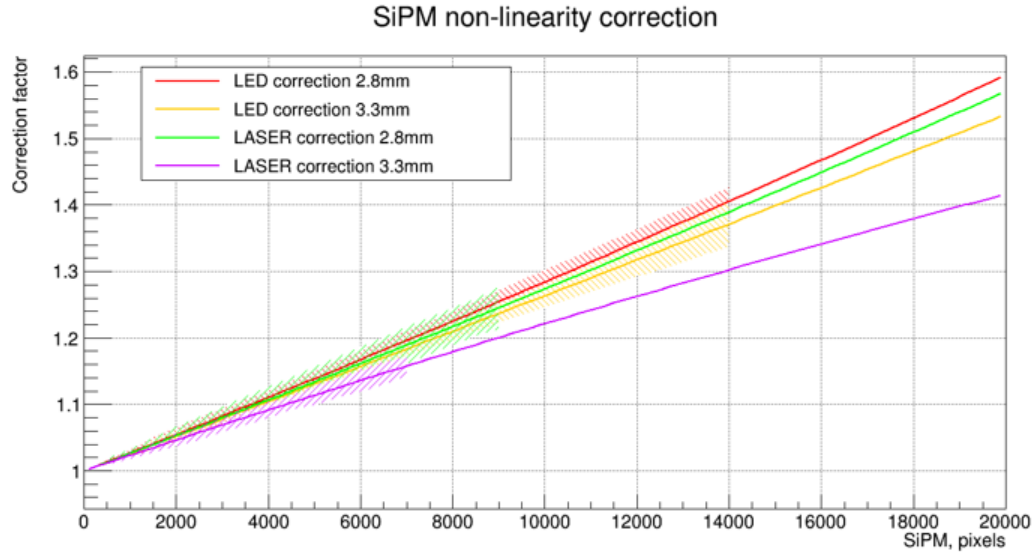


Figure 4.3. Graph of the correction factor vs. number of pixels fired. The correction factor is the number the output needs to be multiplied by in order to obtain the linear output value. This means a perfectly linear device would always have a correction factor of 1. This data was obtained from shining a laser directly on the SiPM.

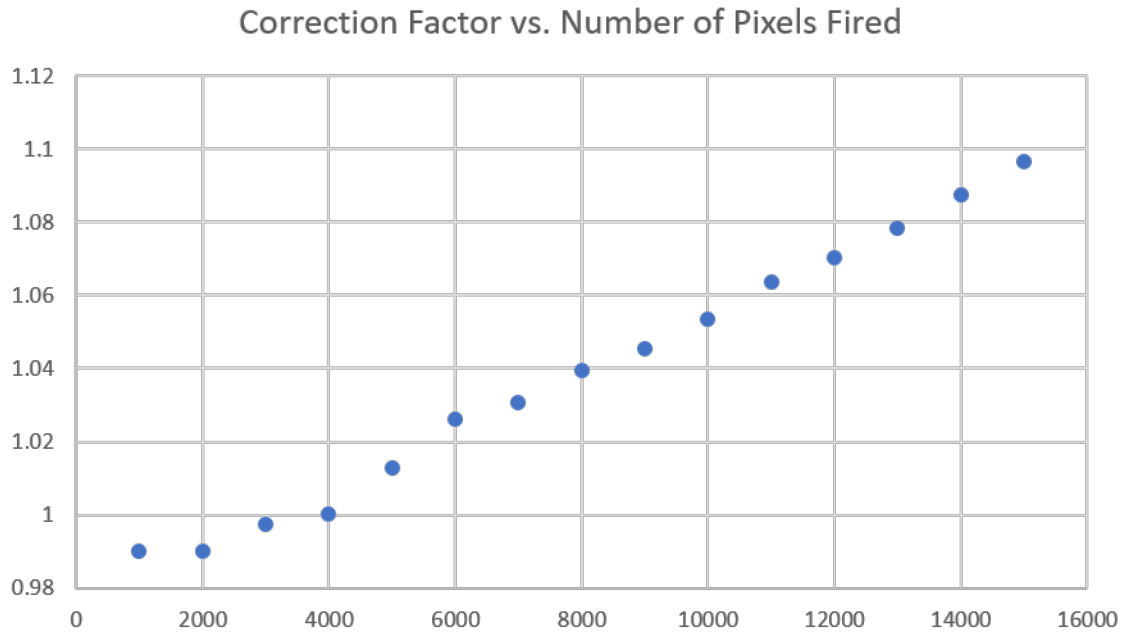


Figure 4.4. Graph of the correction factor vs. number of pixels fired. The correction factor is the number the output needs to be multiplied by in order to obtain the linear output value. This means a perfectly linear device would always have a correction factor of 1. This data was obtained from the SiPM simulation.

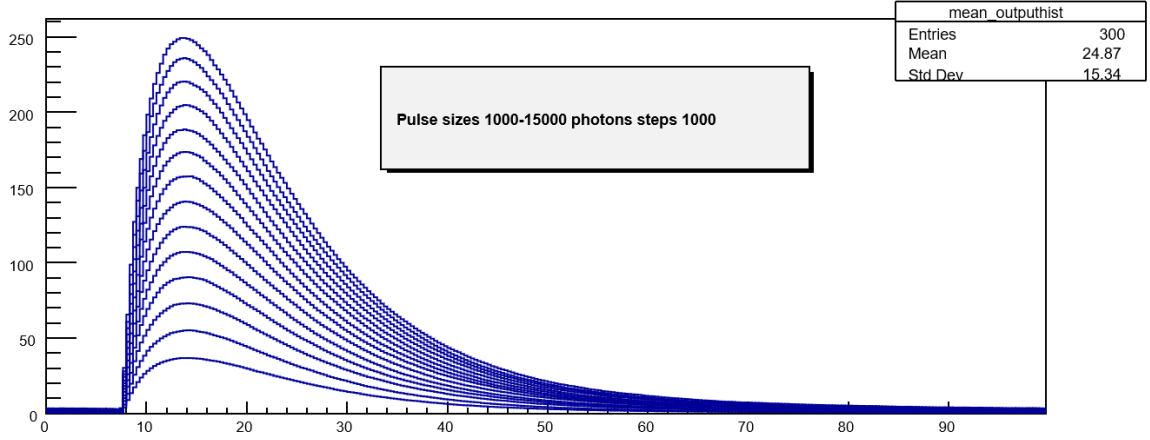


Figure 4.5. Pulse shapes from the SiPM simulation. They are stacked on top of each other to highlight any differences from an increase input photon count.

One effect of the saturation that is less understood and harder to implement is that the pixels can sometimes recharge faster. This effect is due to the QIE chips in the readout modules which in addition to taking and processing the charge from the SiPMs also supply the charge that recharges the capacitors for each individual pixel. The capacitors for the pixels recharge exponentially just like normal RC circuits. Normally the time constant for the capacitors is about 9ns which means the capacitor is basically completely charges after 20ns but this is when the SiPM is outputting a very low amount of charge to the QIE chips. When the SiPM is outputting more charge the QIE chips then supply more charge to the SiPM recharging the pixels faster. When the pixels recharge fast they will fire off closer to their maximum charge even when hit in rapid succession. This means as saturation becomes more prevalent with a high number of incident photons this reduces the effects of saturation. The effect of different recharge time for the pixels can be see in figure 4.6.

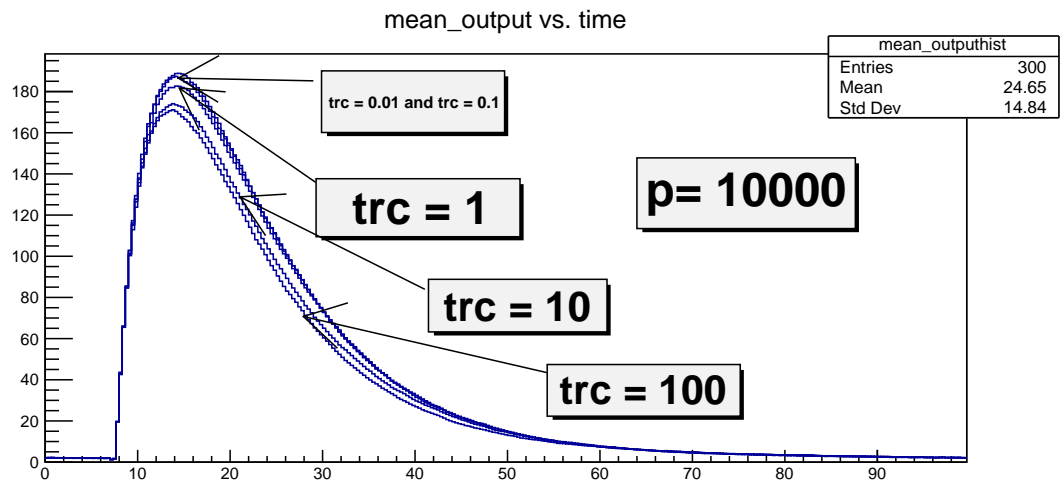


Figure 4.6. Output pulses of the SiPM simulation with 10000 incidents photons comparing the effect of changing the recharge time constant on the pixels. The TRC value is the recharge time constant.