

- Code development 26/09/21
- Incompressible, cylindrical, axisymmetric jet, resolvent analysis
- Conclusion:
  - resolvent analysis and modal analysis results qualitatively in good agreement with Garnaud's reference results.

Base flow (obtained by Newton)

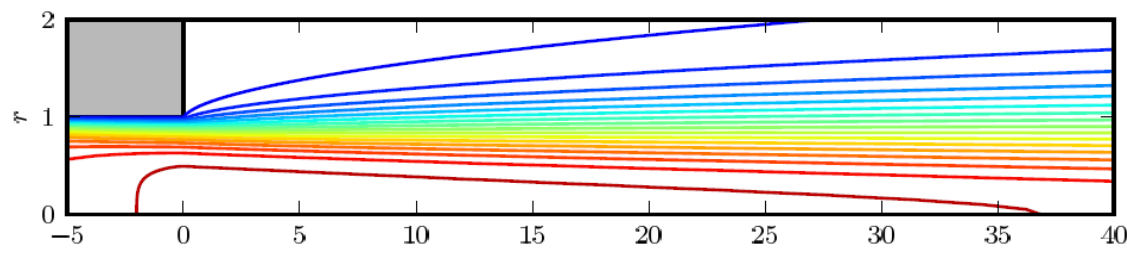
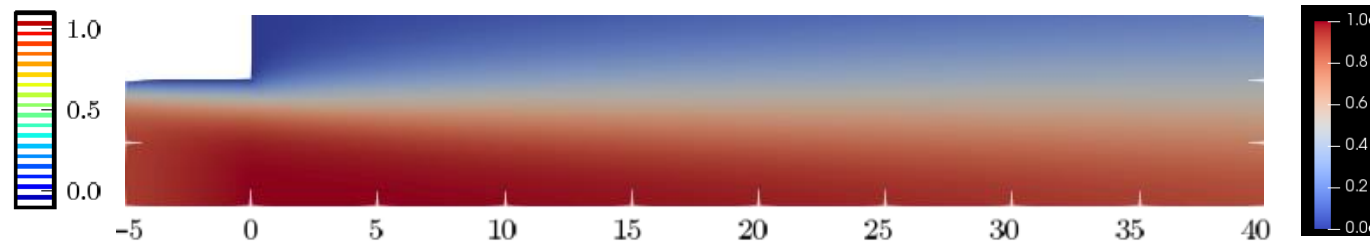
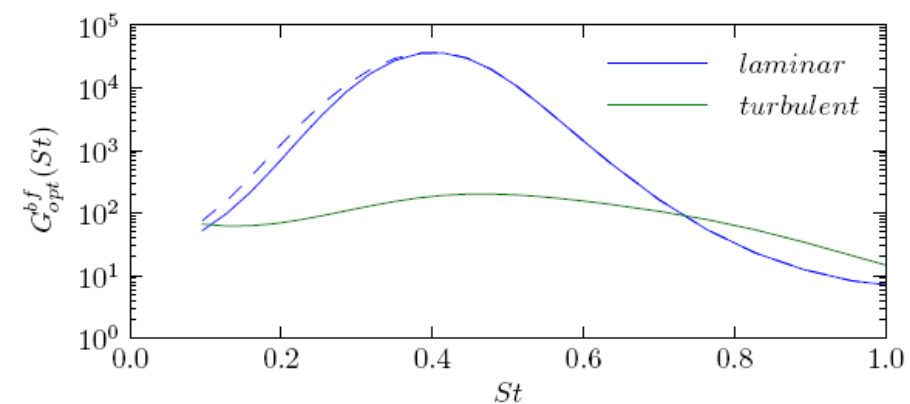


Figure 4.2 in Garnaud Ph. D thesis



new code result

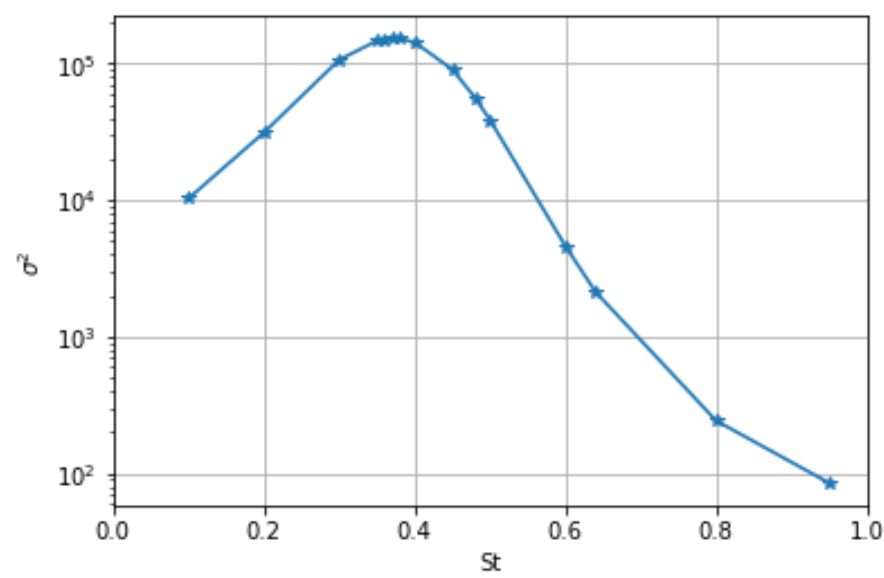
Resolvent gain curve



We work with laminar  $x_\infty=40$

Figure 6.6: Optimal amplification of periodic body forcing for the laminar and turbulent mean flows. The dashed line displays the results obtained for domain of length  $x_\infty = 60$  ( $x_\infty = 40$  is used otherwise).

Figure 6.6 in Garnaud Ph. D thesis



Note: no restriction of forcing domain inside the nozzle, as Garnaud did in his calculation. This can lead to quantitative differences.

# Resolvent modes

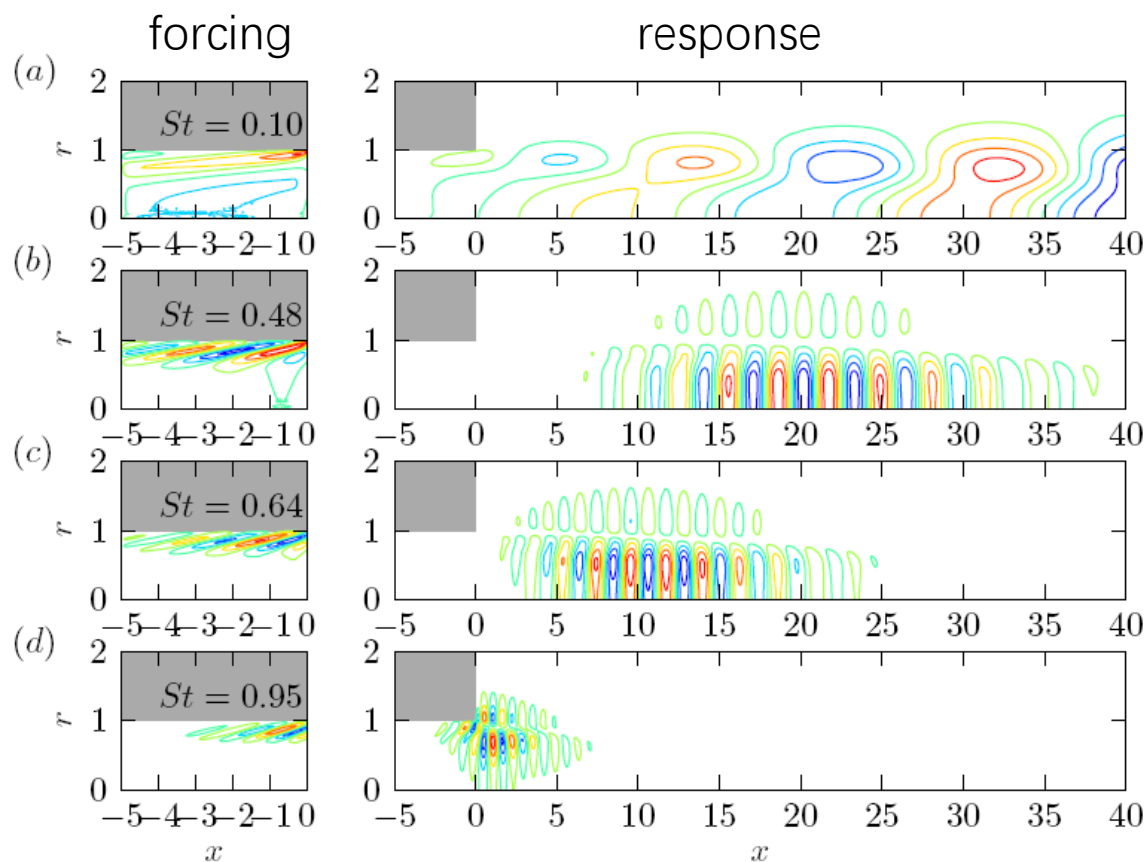
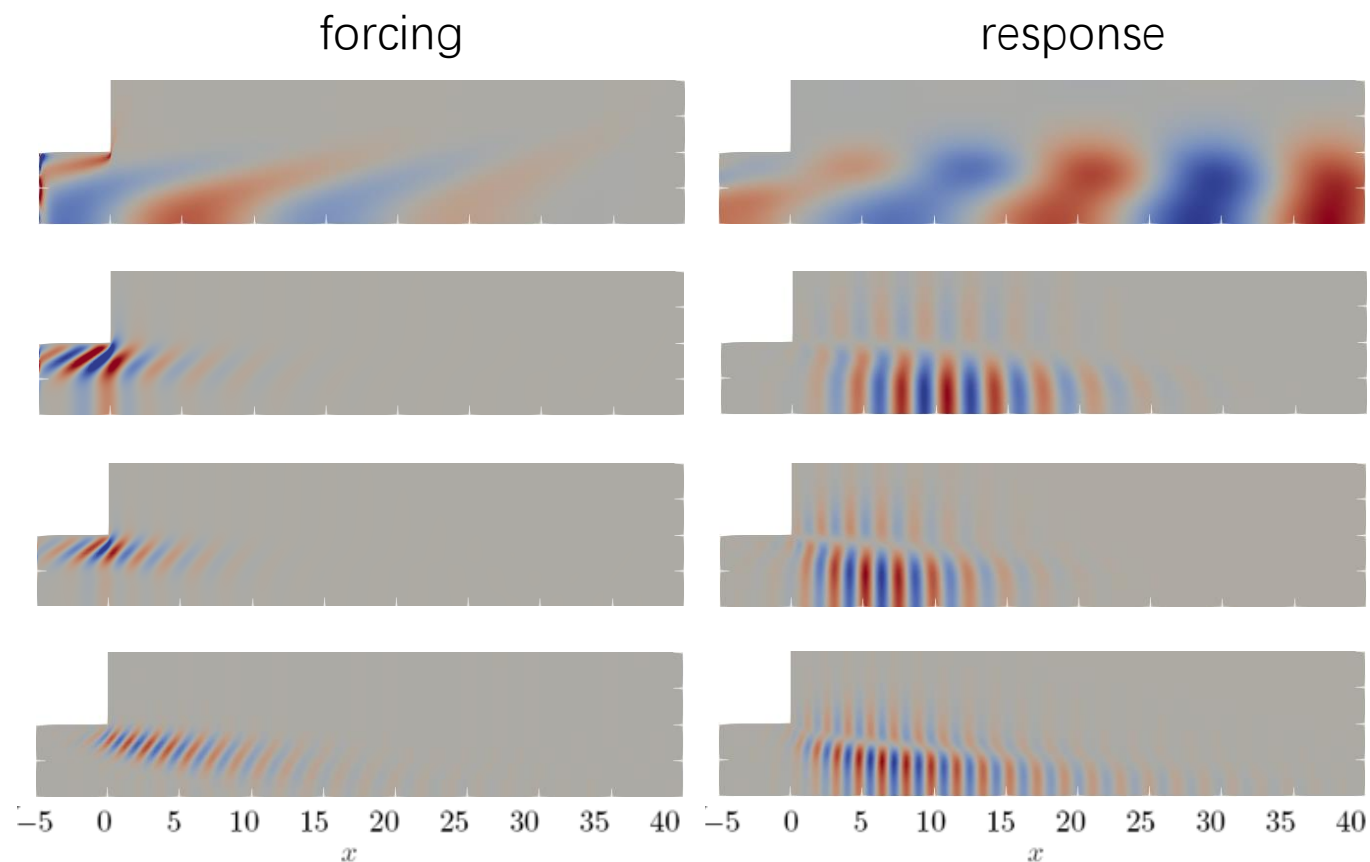


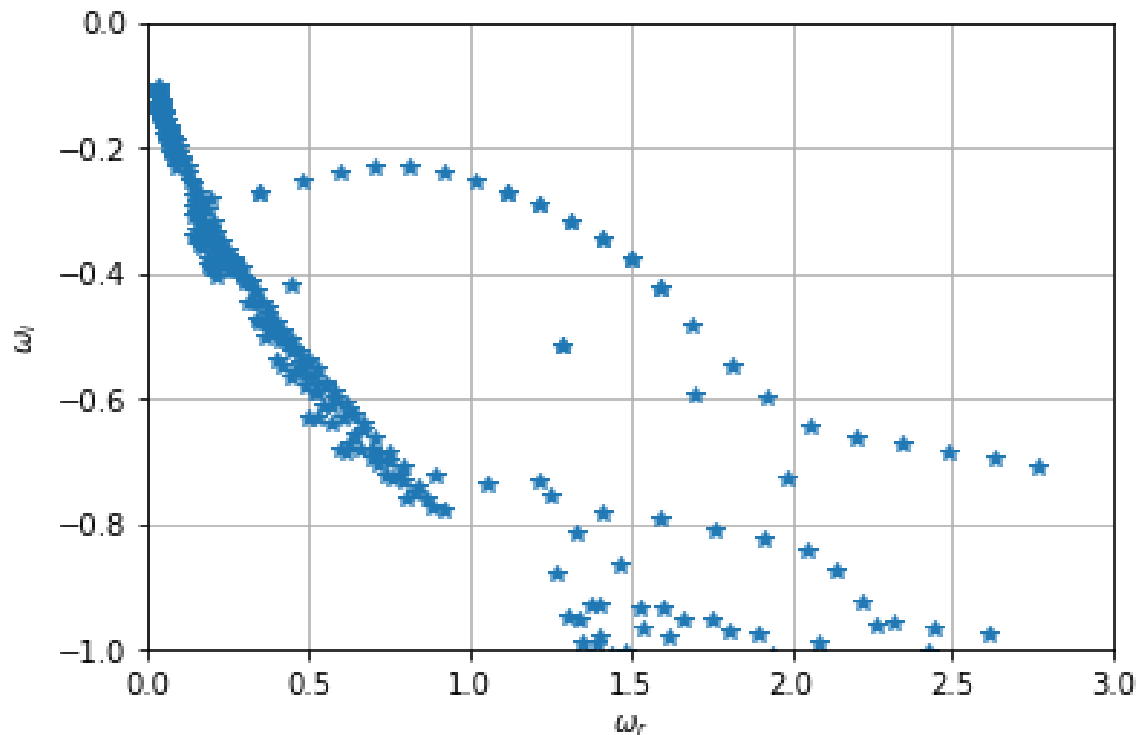
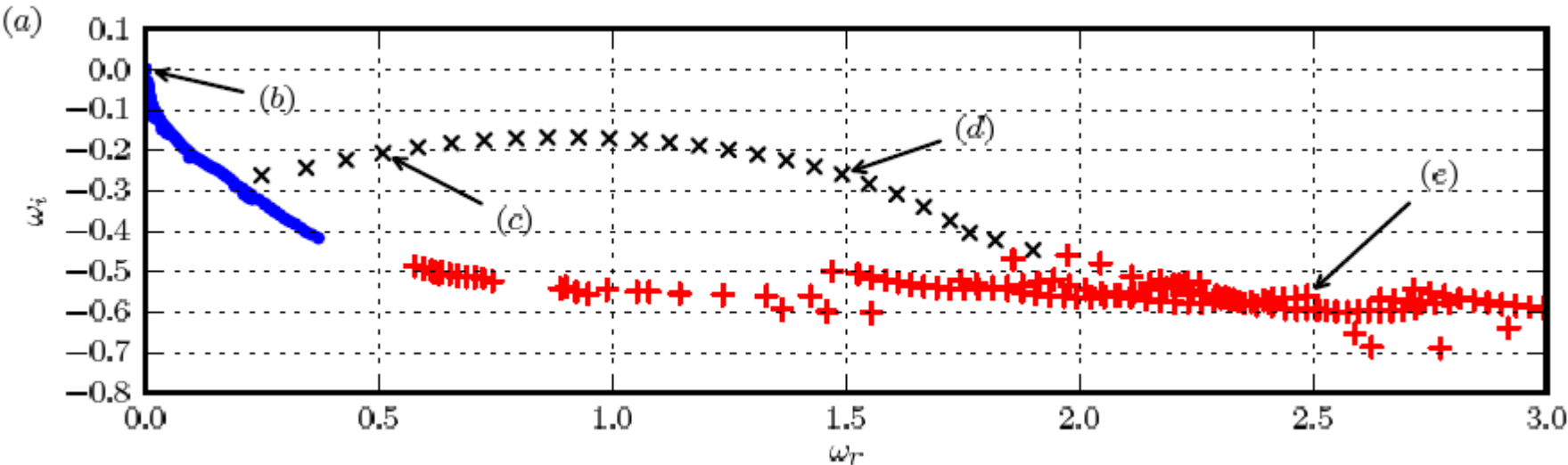
Figure 6.7 in Garnaud Ph. D. thesis



new code result

Note: no restriction of forcing domain inside the nozzle, as Garnaud did in his calculation. This can lead to quantitative differences.

Modal analysis --- Eigenvalues



# Code review (brief)

Anything with label “lowMach” or “lowMach\_reacting” has not been implemented. Only “incompressible” is implemented.

# Code structure

- Mesh
- Matlab: calculate eigenvalues from Matlab for modal analysis if required
- jet\_incompressible: result path folder
- command\_yaj.py: command file
- yaj.py: main script

# command\_yaj.py

- Newton
- Resolvent
- Eigenvalues (modal analysis)

```
14 import_flag=1 #1: import base flow from file #0: not import
15 flow_mode='incompressible' #currently only incompressible is implemented.
16 yo=yaj(MeshPath,flow_mode,datapath,import_flag)
```

← We need to import base flow not only when doing resolvent analysis and modal analysis, but also when doing Newton iteration from a previous calculated flow field.

If put it to 0, Newton iteration starts from the flow field given by self.InitialConditions().

self.OperatorNonlinear()

- rhs of nonlinear governing equations
- A good reference for nonlinear governing equations is [https://demichie.github.io/NS\\_cylindrical/](https://demichie.github.io/NS_cylindrical/)
- Can's code defines the differential operators in cylinder coordinates for axisymmetric case. The second-order stress tensor in grad\_cyl is reduced from (19) in ref.

```
#####  
# Can's code for cylindrical coordinates  
# Define operator for axi-symmetric polar coordinate system  
# in x,r,th it comes [Sxx Sxr Sxth ]  
#                      [Srx Srr Srth ]  
#                      [Sthx Sthr Sthth]  
  
# Gradient with x[0] is x and x[1] is r  
def grad_cyl(v):  
    return as_tensor([[v[0].dx(0), v[0].dx(1), 0],  
                      [v[1].dx(0), v[1].dx(1), 0],  
                      [0, 0, v[1]/self.r]])  
  
def div_cyl(v):  
    c=(1./self.r)*(self.r*v[1]).dx(1) + v[0].dx(0)  
    return c
```

$$\nabla \mathbf{u} = \begin{bmatrix} \frac{\partial u_r}{\partial r} & \frac{1}{r} \frac{\partial u_r}{\partial \theta} - \frac{u_\theta}{r} & \frac{\partial u_r}{\partial z} \\ \frac{\partial u_\theta}{\partial r} & \frac{1}{r} \frac{\partial u_\theta}{\partial \theta} + \frac{u_r}{r} & \frac{\partial u_\theta}{\partial z} \\ \frac{\partial u_z}{\partial r} & \frac{1}{r} \frac{\partial u_z}{\partial \theta} & \frac{\partial u_z}{\partial z} \end{bmatrix} \quad (19)$$



self.OperatorNonlinear()

- The axisymmetric equations read in (24)-(26) in ref. (\rho and g are not considered.)
- Integration by parts is used for diffusion and pressure term, with neglecting boundary parts.

```
#mass
F = div_cyl(u)*self.Test[1]*self.r*dx
#momentum
F -= inner(grad(u)*u, self.Test[0])*self.r*dx
F -= self.mu*inner(grad_cyl(u), grad_cyl(self.Test[0]))*self.r*dx
F -= -inner(p, div_cyl(self.Test[0]))*self.r*dx
return F
```

$$\frac{\partial u_r}{\partial t} + u_r \frac{\partial u_r}{\partial r} + u_z \frac{\partial u_r}{\partial z} = -\frac{1}{\rho} \frac{\partial p}{\partial r} + \frac{\mu}{\rho} \left\{ -\frac{u_r}{r^2} + \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial u_r}{\partial r} \right) + \frac{\partial^2 u_r}{\partial z^2} \right\} + g_r, \quad (24)$$

$$\frac{\partial u_z}{\partial t} + u_r \frac{\partial u_z}{\partial r} + u_z \frac{\partial u_z}{\partial z} = -\frac{1}{\rho} \frac{\partial p}{\partial z} + \frac{\mu}{\rho} \left\{ \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial u_z}{\partial r} \right) + \frac{\partial^2 u_z}{\partial z^2} \right\} + g_z, \quad (25)$$

while the continuity equation is

$$\frac{1}{r} \frac{\partial}{\partial r} (r u_r) + \frac{\partial u_z}{\partial z} = 0. \quad (26)$$

- When doing non-axisymmetric, one may simply consider using the full second-order tensor in (19).

self.Newton()

```
126 self.mu=self.mu*self.play #update viscosity
170 self.play-=self.minus_play #increment of viscosity
```

```
154 if self.label=='incompressible':
155     #write results in private_path for a given mu
156     u_r,p_r = self.ru.split()
157     ll="mu"+str(np.round(self.play, decimals=2))
158     File(self.dnspath+self.private_path+"u"+f"{ifile:03d}"+l
159     File(self.dnspath+self.private_path+"baseflow"+f"{ifile:
160     print(self.dnspath+self.private_path+"baseflow"+f"{ifile
161 elif self.label=='lowMach':
162     pass
163
164     #write result of current mu
165     File(self.dnspath+"u"+"pvd") << u_r
166     File(self.dnspath+"baseflow000"+"xml") << self.ru.vector()
167     print(self.dnspath+"baseflow000"+"xml written!")
```

- When doing Newton iterations, we sometimes need to start with a larger viscosity to get an initial field, and then gradually decrease the viscosity. Thus we add a prefactor `self.play` multiplied in front of `self.mu`. For example, when doing `lowMach` for cooled cylinder, we may need to set `self.play=10` at the beginning, and then gradually decrease `self.play` to 1.
- For this reason, the results issued from each iteration are saved twice: a folder `self.private_path='doing/'` records all the history files; the current results are saved in `self.datapath`, and they are covered each time when a new iteration is done.
- However, in the present case, we can directly set `self.play=1`, and Newton can converge.

# Boundary conditions

Get all free indices without Dirichlet condition

175  `def get_indices(self):`


- Newton

```
189     def BoundaryConditions(self):  
  
202         # define boundary conditions for Newton/timestpper  
203         if self.label=='incompressible':  
204             ux_tanh=Expression('tanh(5*(1-x[1]))', degree=2)  
205             bcs_inflow_x = DirichletBC(self.Space.sub(0).sub(0), ux_tanh, inlet)  
206             bcs_inflow_r = DirichletBC(self.Space.sub(0).sub(1), 0, inlet)  
207             bcs_wall=DirichletBC(self.Space.sub(0), (0,0), wall)  
208             bcs_symmetry=DirichletBC(self.Space.sub(0).sub(1), 0, symmetry)  
209             return [bcs_inflow_x,bcs_inflow_r,bcs_wall,bcs_symmetry]
```

- Resolvent (open ux at inlet)

```
def BoundaryConditionsPerturbations(self):  
  
227     if self.label=='incompressible':  
228         bcs_inflow = DirichletBC(self.Space.sub(0).sub(1), 0, inlet)  
229         bcs_wall=DirichletBC(self.Space.sub(0), (0,0), wall)  
230         bcs_symmetry=DirichletBC(self.Space.sub(0).sub(1), 0, symmetry)  
231         return [bcs_inflow,bcs_wall,bcs_symmetry]
```

- Modal (closed ux at inlet)

```
def BoundaryConditionsPerturbations(self):  
  
227  if self.label=='incompressible':  
228     bcs_inflow = DirichletBC(self.Space.sub(0), (0,0), inlet)  
229     bcs_wall=DirichletBC(self.Space.sub(0), (0,0), wall)  
230     bcs_symmetry=DirichletBC(self.Space.sub(0).sub(1), 0, symmetry)  
231     return [bcs_inflow,bcs_wall,bcs_symmetry]
```

self.Resolvent()

- Implementation of  $P^H Q^H R^H M_r R Q P f = \sigma^2 I^H M_f I f$
- The matrix P, Q are real, so the Hermitian is the same as the Transpose. Only R is complex. LU decomposition is used for R.
- Detailed comments are given in the code.
- Some details with different notations can be found in document by Lutz: resolvent\_doc.pdf.
- The quadrature matrix is required to compensate the quadrature in resolvent operator R, because  $R=(A-i*\omega B)^{-1}$ , and there is quadrature in A and B.
- The matrix I reshapes forcing matrix Mf (m\*m) to  $I^T * M_f * I$  (n\*n). Normally the matrix I is the same as P. But this is not the case if we use values not equal to 0 or 1 in P to damp (but not entirely eliminate) the forcing in some regions.
- The output functionalities only work for setting nb of resolvent modes as 1, presently.

self.Eigenvalues()

- We find in some cases, shift-invert method works better in Matlab than Python especially for reacting flow, and we have not exactly understood why.
- For this reason, three operational mode by setting flag\_mode  
0: save matrix as .mat with file name "savematt"; 1: load result matrix from .mat with file name "loadmatt"; 2: calculate eigenvalues in python
- One could use mode 0 first and then use get\_eigs.m in folder Matlab/. Then use mode 1 to load and output the eigenmodes.