CE394M: Stress paths and invariants

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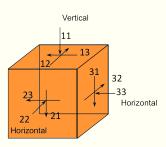
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Stresses / strains

1D consolidation / simple shear

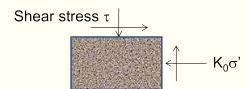
2D plane strain

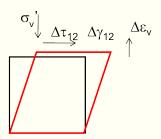


3D general (axi-symmetric as a special case)

1D simple shear

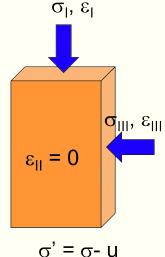
Normal effective stress σ'_{v}





2D plane strain / Mohr-Coulomb model

Stresses and strains: independent components



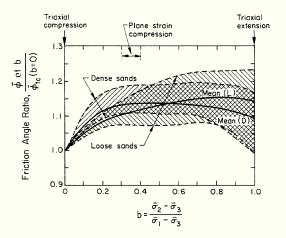
2D Mohr circle

Mohr circle to 2D stress components

Successive Mohr circles and stress path for constant σ_3 and increasing σ_1 :

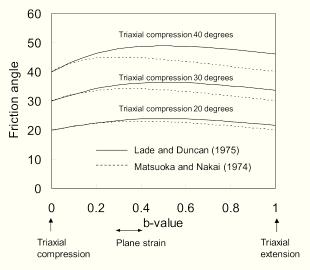
Effect of σ_{II}

Bishop (1966) defined b-value:



Effect of σ_{II} on friction angle (Kulhawy and Mayne, 1990)

Effect of σ_{II} on friction



Chapter 11., Mitchell and Soga, 2005

Effect of σ_{II} on friction

Various models fit the experimental data showing the intermediate stress effect. Lade and Duncan (1975)

$$\frac{I_1^3}{I_3} = const$$

Matsuoka and Nakai (1974)

$$\frac{I_1I_2}{I_3} = const$$

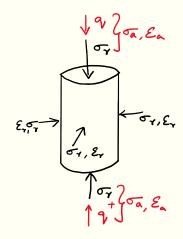
$$I_1 = \sigma_1 + \sigma_2 + \sigma_3$$

$$I_2 = \sigma_1 \sigma_2 + \sigma_2 \sigma_3 + \sigma_3 \sigma_1$$

$$I_3 = \sigma_1 \sigma_2 \sigma_3$$

Matsuoka and Nakai's model gives the same friction angle for compression and extension, whereas Lade and Duncan's model gives the ratio of the TXE friction angle (ϕ'_{TXE}) to the TXC friction angle (ϕ'_{TXC}) to be 1.08 at $\phi'_{TXC}=20^{\circ}$ to 1.15 at $\phi'_{TXC}=40^{\circ}$.

Triaxial stresses and strains: independent components



Triaxial deviatoric strain

- Deviatoric / shear strain: $\varepsilon_s = 2/3(\varepsilon_a \varepsilon_r)$
- Why 2/3?:

Triaxial stress - strain relationship

Relationship in terms of axial and radial direction:

$$\begin{bmatrix} \varepsilon_{\mathsf{a}} \\ \varepsilon_{\mathsf{r}} \end{bmatrix} = \frac{1}{E} \begin{bmatrix} 1 & -2\nu \\ -\nu & (1-\nu) \end{bmatrix} \begin{bmatrix} \sigma_{\mathsf{a}} \\ \sigma_{\mathsf{r}} \end{bmatrix} \tag{1}$$

Relationship in terms of 'p' and 'q':

$$\begin{bmatrix} p \\ q \end{bmatrix} = \begin{bmatrix} 1/3 & 2/3 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} \sigma_a \\ \sigma_r \end{bmatrix} \quad \text{invert} \quad \rightarrow \begin{bmatrix} \sigma_a \\ \sigma_r \end{bmatrix} \begin{bmatrix} 1 & 2/3 \\ 1 & -1/3 \end{bmatrix} \begin{bmatrix} p \\ q \end{bmatrix} \tag{2}$$

Strain:

$$\begin{bmatrix} \varepsilon_{\mathsf{v}} \\ \varepsilon_{\mathsf{s}} \end{bmatrix} = \begin{bmatrix} 1 & 2 \\ 2/3 & -2/3 \end{bmatrix} \begin{bmatrix} \varepsilon_{\mathsf{a}} \\ \varepsilon_{\mathsf{r}} \end{bmatrix} \tag{3}$$

Volume - shear coupling

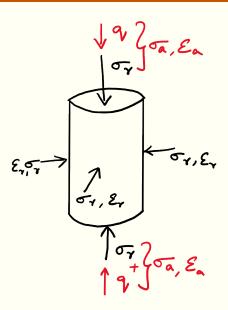
Equation (1) and eq. (2) in eq. (3) to relate $\varepsilon_v, \varepsilon_s$ to p, q:

$$\begin{bmatrix} \varepsilon_{v} \\ \varepsilon_{s} \end{bmatrix} = \begin{bmatrix} 3(1-2\nu)/E & 0 \\ 0 & 2(1+\nu)/3E \end{bmatrix} \begin{bmatrix} p \\ q \end{bmatrix}$$

Volume shear coupling

Usually volume change and shear deformation are decoupled (i.e. bulk modulus and shear modulus), but in soils they are coupled:

Triaxial test



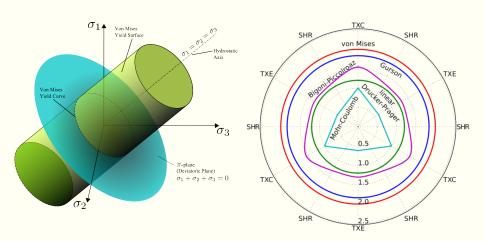
Mohr circle to 3D stress components

Successive Mohr circles and stress path for constant σ_3 and increasing σ_1 :

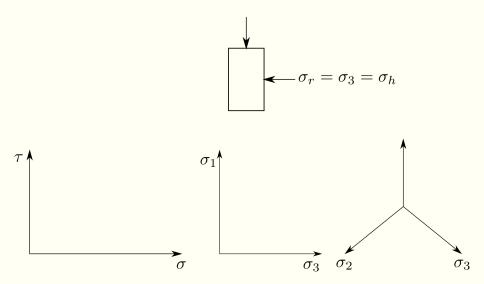
Stress paths p-q

Different stress paths for initially hydrostatic stress conditions:

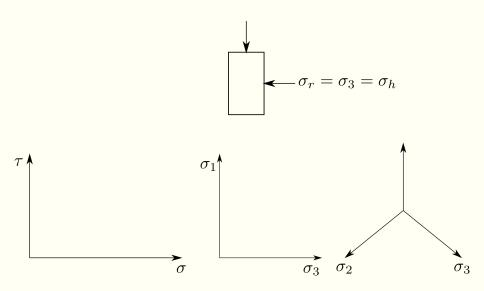
Pi-plane



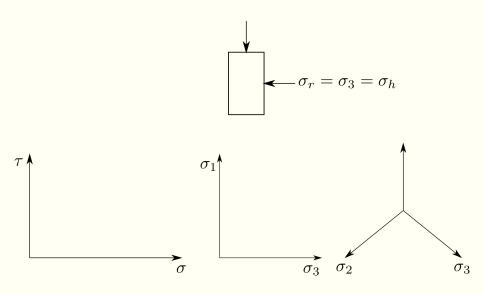
π plane: Triaxial compression



π plane: Triaxial extension



π plane: Random stress paths



Stress and strain invariants in 3D

6 stresses and strains

- Mean pressure:
- ullet Deviator stress: $q=\sqrt{3/2}\sqrt{s_{11}^2+s_{22}^2+s_{33}^2+2 au_{12}^2+2 au_{23}^2+2 au_{31}^2}$
- ullet where $s_{11}=\sigma_{11}'-p', s_{22}=\sigma_{22}'-p', s_{33}=\sigma_{33}'-p'$
- Volumetric strain:
- Deviatoric strain:

$$\varepsilon_s = \sqrt{2/3}\sqrt{e_{11}^2 + e_{22}^2 + e_{33}^2 + \gamma_{12}^2/2 + \gamma_{23}^2/2 + \gamma_{31}^2/2}$$

• where $e_{11}=\varepsilon_{11}-\varepsilon_{v}/3, e_{22}=\varepsilon_{22}-\varepsilon_{v}/3, e_{33}=\varepsilon_{33}-\varepsilon_{v}/3$

Stress and strain invariants in 3D

3 Principal stresses and strains

- Mean pressure:
- Deviator stress: $q=\sqrt{1/2}\sqrt{(\sigma_I'-\sigma_{II}')^2+(\sigma_{II}'-\sigma_{III}')^2+(\sigma_{III}'-\sigma_{I}')^2}$
- Volumetric strain:
- Deviatoric strain: $\varepsilon_s = \frac{\sqrt{2}}{3} \sqrt{(\varepsilon_I' \varepsilon_{II}')^2 + (\varepsilon_{II}' \varepsilon_{III}')^2 + (\varepsilon_{III}' \varepsilon_I')^2}$

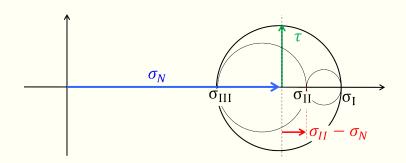
In triaxial condition (principal stresses/strains)

The missing stress invariant

General stresses:

• Principal stresses:

The Lode parameter

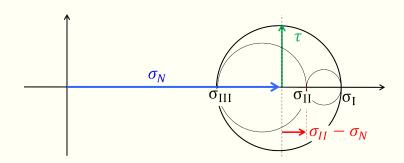


Maximum shear stress:

Normal stress:

Position of intermediate principal stress:

The Lode parameter



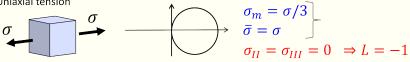
Maximum shear stress:

Normal stress:

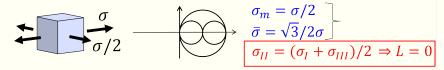
Position of intermediate principal stress:

Which states have the same Lode parameter?

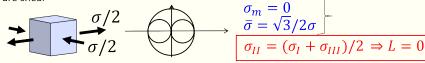
Uniaxial tension



Plane strain tension



Pure shear

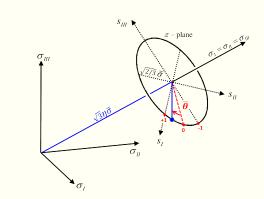


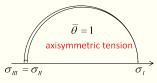
Lode angle parameter

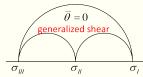
Lode angle parameter (Lode, 1926)

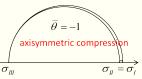
$$L = \frac{2\sigma_{II} - \sigma_{I} - \sigma_{III}}{\sigma_{I} - \sigma_{III}}$$

Lode angle $\bar{\theta} \approx -L$









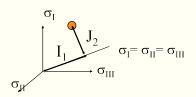
Stress invariant

Principal stresses (3 components) ⇔ Invariants (3 components)

•
$$I_1 =$$

•
$$J_2 =$$

•
$$J_2 =$$



Lode angle:

•
$$\sin 3\theta = \frac{-3\sqrt{3}}{2} \frac{J_3}{J_2^{3/2}}$$
.

$$\bullet \ \theta = \tan^{-1} \left[\frac{1}{\sqrt{3}} \left(2 \frac{(\sigma_2 - \sigma_3)}{(\sigma_1 - \sigma_3)} - 1 \right) \right]$$

- 0
- •