CE394M: Tresca and Mohr-Coulomb

Krishna Kumar

University of Texas at Austin

krishnak@utexas.edu

April 20, 2019

Overview

Constitutive modeling

2 Tresca model

Mohr-Coulomb model

Stress invariants

- The magnitudes of the component of the stress vector depend on the chosen direction of the coordinate axes (in 3D: 6 variables).
- Principal stresses always act on the same planes and have the same magnitude (invariant to the coordinate axes), but still need to define the corresponding orientations (in 3D: 6 variables).
- For isotropic materials, it is very convenient to work with alternative invariant quantities which are combinations of principal stresses.

Stress invariants

- Mean effective stress $p = \frac{1}{3}(\sigma_I + \sigma_I I + \sigma_I II)$
- Deviatoric stress: $J = \frac{1}{\sqrt{6}} \sqrt{(\sigma_I \sigma_{II})^2 + (\sigma_{II} \sigma_{III})^2 + (\sigma_{III} \sigma_I)^2}$
- Lode's angle $\theta = \tan^{-1}\left[\frac{1}{\sqrt{3}}\left(2\frac{(\sigma_{II}-\sigma_{III})}{\sigma_{I}-\sigma_{III}}-1\right)\right]$

Principal stresses can be expressed in terms of invariants:

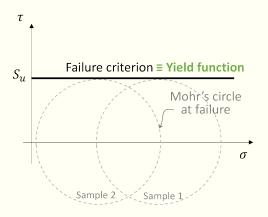
$$\begin{bmatrix} \sigma_{I} \\ \sigma_{II} \\ \sigma_{III} \end{bmatrix} = p \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix} + \frac{2}{\sqrt{3}} J \begin{bmatrix} \sin\left(\theta + \frac{2\pi}{3}\right) \\ \sin\theta \\ \sin\left(\theta - \frac{2\pi}{3}\right) \end{bmatrix}$$

Tresca model

Simulation of undrained behavior of saturated clay Failure criteria:

$$\tau_f = s_u$$

where τ_f is the shear stress at failure. s_u is the undrained strength.

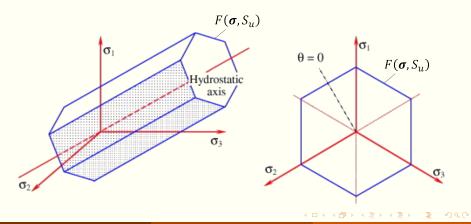


Tresca model

Yield function

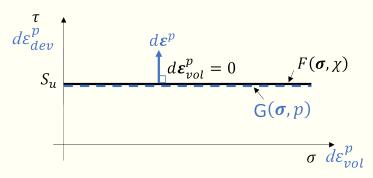
$$F(\sigma, W_p) = \sigma_I - \sigma_{III} - 2s_u = 0$$

= $J \cos \theta - s_u = 0$



Tresca model

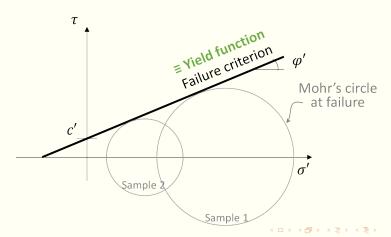
- The model is intended to predict the undrained behavior of saturated clay, hence it should predict zero volumetric strains.
- This fact constrains the plastic potential, and an associated flow rule should be considered.
- In general, this is a perfectly plastic model (no hard/soft law)
- Parameters required (3): s_u, E_u, ν_u .



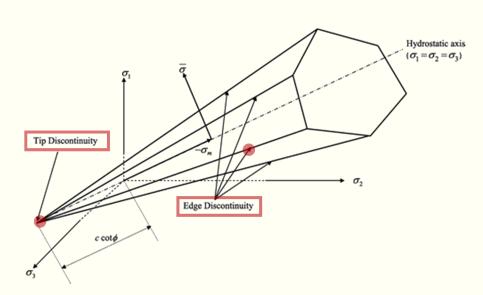
Mohr-Coulomb model

Failure criteria: $\tau_f = c' = \sigma_f' \tan \phi'$.

$$F(\sigma', c', \phi') = \frac{1}{2}(\sigma'_I - \sigma'_{III}) - \frac{1}{2}(\sigma'_I + \sigma'_{III})\sin\phi' - c'\cos\phi' = 0.$$

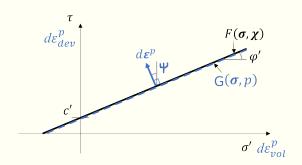


Mohr-Coulomb model



Mohr-Coulomb: flow rule?

Can we consider an associative flow rule? Failure criteria: $F(\sigma', c', \phi') = G(\sigma', c', \phi')$



 ε^p is inclined at an angle of ϕ' , and indicates negative (tensile) plastic strains. This results in a dilatant volumetric plastic strain $(d\varepsilon_{vol}^p)$, in which the angle of dilation ψ is equal to the effective friction angle (ϕ') .

Mohr-Coulomb: Drawbacks (associative)

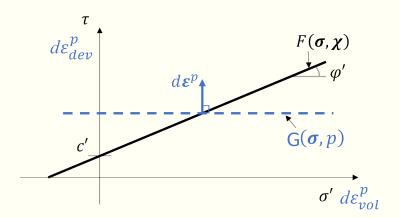
- **①** The magnitude of the plastic volumetric strains $(d\varepsilon)$ is much larger than that observed in real soils.
- Once the soil yields it will dilate for ever! (note: real soil, which may dilate initially on meeting the failure surface, will often reach a constant volume condition at large strains).

Solutions:

- \bullet Adopt a non-associated flow rule with a more realistic angle of dilation ψ
- @ Allow angle of dilation vary with plastic strain (hardening/softening law) $\psi = f(\varepsilon)$

Mohr-Coulomb: non-associative model

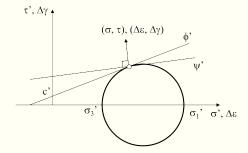
Consider a non-associated model with $\psi=0$. In which case $d\varepsilon_{vol}^P=0$. Parameters (5): c', ϕ', ψ, E, ν .



Mohr-Coulomb: Yield and Potential fn

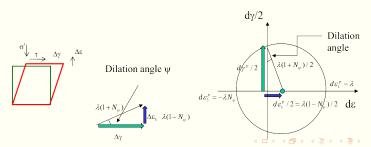
Yield fn:
$$F = \sigma_1' - \sigma_3' N_\phi + 2c\sqrt{N_\phi} = 0$$
., where $N_\phi = \frac{1+\sin\phi'}{1-\sin\phi'}$
Potential fn: $G = \sigma_1' - \sigma_3' N_\psi + constant = 0$., where $N_\psi = \frac{1+\sin\psi'}{1-\sin\psi'}$





Mohr-Coulomb: plastic strains

$$\begin{split} d\varepsilon_1^p &= \lambda \frac{\partial G}{\partial \sigma_1} = \lambda \cdot 1 = \lambda \\ d\varepsilon_2^p &= \lambda \frac{\partial G}{\partial \sigma_2} = \lambda \cdot 0 = 0 \quad \text{Zero strain increment in 2-dir} \\ d\varepsilon_3^p &= \lambda \frac{\partial G}{\partial \sigma_3} = \lambda \cdot (-N_\psi) = -\lambda N_\psi \end{split}$$

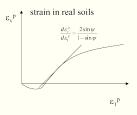


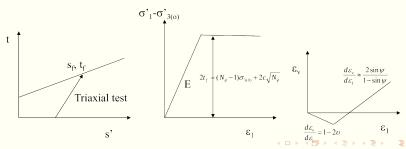
Interpretation of drained TXC

Plastic volumetric strain:

$$\begin{split} d\varepsilon_{v}^{p} &= d\varepsilon_{1}^{p} + d\varepsilon_{2}^{p} + d\varepsilon_{3}^{p} = \lambda(1 - N_{\psi}) \\ &= \lambda \frac{2\sin\psi}{1 - \sin\psi} \\ \frac{d\varepsilon_{v}^{p}}{d\varepsilon_{1}^{p}} &= \frac{\frac{2\sin\psi}{1 - \sin\psi}}{\lambda} = \frac{2\sin\psi}{1 - \sin\psi} \end{split}$$

Note that ψ changes with plastic strain in real soils





Mohr-Coulomb: Tension cut-off

- For c = 0, the standard Mohr-Coulomb criterion allows for tension. Allowable stress increase with cohesion.
- In reality, soil can sustain none or only very small tensile stresses.
- This behavior is usually included by specifying a tension cut-off, which is equivalent to include an additional yield surface: $F(\sigma, \sigma_t) = \sigma_i \sigma_t = 0$.

