A Computational Toolbox for Teaching Quantum Mechanics and Quantum Optics

Uma Ferramenta Computational para o Ensino da Mecânica Quântica e da Óptica Quântica

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We describe a collection of codes that may be useful in teaching of quantum mechanics and quantum optics courses. The toolbox has been developed using Matlab programming language, but we expect that a translation to an open source mathematics software could be easily done. These codes can be used to help visualize quantum states of light using Wigner function, how these states may be affected by loss, among other things.

Keywords: Quantum states of light, Wigner function, purity, fidelity.

Descrevemos um conjunto de códigos computacionais que podem ser úteis no ensino da mecânica quântica e da óptica quântica. Esse ferramental foi desenvolvido usando a linguagem de programação Matlab, mas esperamos que uma tradução desses códigos para uma linguagem de software aberto possa ser feito facilmente. Os códigos aqui descritos ajudam a visualizar estados quânticos da luz usando a função de Wigner, como estes estados podem ser afetados por perdas, entre outras coisas.

Palavras-chave: Estados quânticos da luz, função de Wigner, pureza, fidelidade.

1. Introduction

Quantum mechanics is one of the most crucial, and yet, challenging, topics of modern physics [1, 2]. With its application to many and important technological problems, the challenges in teaching quantum mechanics are now part not only of the physical science undergraduate curricula, but also of engineering education, especially in the area of communications [3, 4]. Some of the challenges are also faced by professors when teaching quantum optics in a graduate (or undergraduate) level. Any visualization tools in these areas of knowledge can be of great assistance towards an appropriated grasp of concepts by students [5, 6, 7, 8, 9].

We are going to describe here a set of codes developed using Matlab programming language and their use to give insight into the nature of quantum mechanics and quantum optics. In the context of our work, we use a formulation of quantum mechanics that is equivalent to the standard approach: the Wigner function [10].

If the wave function $\psi(x)$ of a system is known, we may easily determine the probability density $|\psi(x)|^2$

in position space x. On the other hand, the momen-

tum distribution, $|\phi(p)|^2$, is difficult to visualize if we

have only $\psi(x)$. That is the standard formulation of

quantum mechanics. We may alternatively approach

quantum mechanics through a function that displays

the probability distribution simultaneously in the x

and p variables. This function was introduced in 1932

by Eugene Wigner [11] and is defined as:

space x and momentum space p can be obtained from the marginals of the Wigner function, since

$$\int W(x,p)dp = \psi^*(x)\psi(x) = |\psi(x)|^2,$$

$$\int W(x,p)dx = \phi^*(p)\phi(p) = |\phi(p)|^2.$$
 (2)

The Wigner function W(x, p) is known as a quasi probability function because, even though it is real and normalized to one, it is not necessarily non-negative.

 $W(x,p) = \frac{1}{h} \int e^{-ipy/\hbar} \psi(x+y/2) \psi^*(x-y/2) dy.$ (1) Quantum mechanical probability densities in position

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This paper is organized as follows: in Section 2 we describe the use of the computational toolbox to study the harmonic oscillator. In Section 3 we show the use of the visualization toolbox for other states of the light, and in Section 4 we make some concluding remarks.

2. Example: simple harmonic oscillator

We can use the wave function for each energy state of the harmonic oscillator and Eq. (1) to find the corresponding Wigner function for each of these states. Alternatively, we may use the Algebraic Method described in [12], continue the discussion to introduce the number operator, $\hat{n} = \hat{a}^{\dagger} \hat{a}$, and the ket notation for the harmonic oscillator's eigenstates $|n\rangle$, as done in [13].

In the toolbox, there is a routine for generating the pure state vector for photon number eigenstate n in the photon number basis, as show bellow. When calling this routine, the students need to specify two variables: the number of photons, n, and the photon number at which the Hilbert space is truncated (since the routines where developed to study Quantum State Tomography, we need to specify where we truncate the Hilbert space).

```
function psi = generate_fock_vector(n, maxPhoton)
% Creates photon number state
% psi = generate_fock_vector(n, maxPhoton) returns the pure state vector
% for photon number eigenstate n in the photon number basis. maxPhoton is
% the photon number at which the Hilbert space is truncated.
psi = zeros(maxPhoton+1,1);
psi(n+1,1) = 1;
```

The following lines of code plot the Wigner function of a vacuum state (n = 0). Notice that to plot the Wigner function of any n photon state, all that we need to do is change the number of photons n in the routine. The resulting plot for the vacuum state can be seen in Figure 1.

```
% The infinite dimensional state space for the harmonic oscillator will be % represented in the photon number basis.

We will truncate the Hilbert space at maxPhotonNumber photons.

maxPhotonNumber = 10;

% First, pre-compute a lot of numbers, such as coefficients for Hermite
% polynomials, factorials, binomial coefficients.

S = init_tables(maxPhotonNumber);

% Make state vector for Fock state.

n = 0; % number of photons

psi = generate fock_vector(n, maxPhotonNumber)
% The Fock state may suffer from some loss by passing through a
% medium with etaState efficiency.
etaState = 1;
rho = apply_loss(psi,etaState,S);
% Now it must be represented by a density matrix, rho.

wignerStepSize = 0.1;
[x,p] = meshgrid(-4:wignerStepSize:4,-4:wignerStepSize:4);
wigner2 = wigner(rhol, x,p);
mesh(x,p,wigner2); xlabel('x'); ylabel('p'); zlabel('W(x,p)');
```

Figure 2 shows the plot of W(x, p) of a n = 4 number state. In here it can be seen the negative part of the

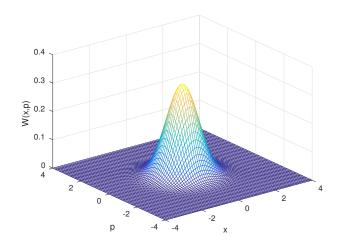


Figura 1: Wigner function of a vacuum state n = 0.

Wigner function, usually related to the nonclassicality of the state.

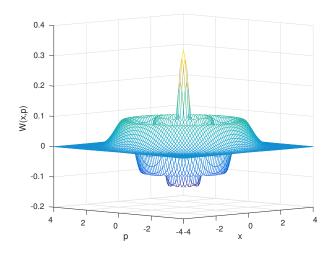


Figura 2: Wigner function of a number state n = 4.

3. Example: other states of light

In the context of a quantum optics course, it is natural to study other states of light, such as coherent and squeezed states. Moreover, these states may also be introduced in an upperlevel quantum mechanics course. A coherent state is an eigenstate of the annihilation operator, $\hat{a}|\alpha\rangle = \alpha|\alpha\rangle$, where α is a complex number. Coherent states are minimum uncertainty states.

A squeezed state is a state of light where fluctuations are reduced below the standard quantum limit in one quadrature at the expense of increased fluctuations in the conjugate quadrature. The toolbox provide routines to generate a coherent and a squeezed vacuum state, as shown bellow.

To plot the Wigner function of a coherent state and a squeezed vacuum state, we use the lines of code showed. We may plot the Wigner function of coherent states of different amplitude by changing the value for α in the code. When plotting the Wigner function of a squeezed vacuum state, we need to specify if we are using the state's variance (as in the example showed) or the ratio of the state's variance with respect to vacuum's variance (that is 1/2). Plotting different squeezed vacuum states is done by changing the variance or the ratio.

In Figure 3 we have the Wigner function of a coherent state with $\alpha=1$. When students compare Figures 1 and 3, they see that a coherent state is a displaced vacuum state. Figure 4 shows the Wigner function of a squeezed vacuum state with variance equals to 1/6 of vacuum variance.

Students from both courses can explore the fundamental superposition principle of quantum mechanics

```
% The infinite dimensional state space for the harmonic oscillator will be
% represented in the photon number basis.
% We will truncate the Hilbert space at maxPhotonNumber photons.
maxPhotonNumber = 10;

% First, pre-compute a lot of numbers, such as coefficients for Hermite
% polynomials, factorials, binomial coefficients.
S = init_tables(maxPhotonNumber);

% Make state vector for coherent state.
alpha = 1; % amplitude of coherent state
psi = generate_coherent_vector(alpha, maxPhotonNumber);
% The coherent state may swiffer from some loss by passing through a
% medium with etaState efficiency.
etaState = 1;
rho = apply_loss(psi,etaState,S);
% Now it must be represented by a density matrix, rho.
wignerStepSize = 0.1;
[x,p] = meshgrid(-4:wignerStepSize:4,-4:wignerStepSize:4);
wigner2 = wigner(rho, x,p);
mesh(x,p,wigner2); xlabel('x'); ylabel('p'); zlabel('W(x,p)');
```

```
% The infinite dimensional state space for the harmonic oscillator will be
% represented in the photon number basis.
% We will truncate the Hilbert space at maxPhotonNumber photons.
maxPhotonNumber = 10;
% First, pre-compute a lot of numbers, such as coefficients for Hermite
% polynomials, factorials, binomial coefficients.
S = init_tables(maxPhotonNumber);
% Make state vector for squeezed vacuum state.
v = 1/6; %variance
ratioSwitch = 'true variance';
psi = generate_squeezed_vacuum_vector(v, maxPhotonNumber, ratioSwitch);
% The squeezed vacuum state may suffer from some loss by passing through a
% medium with etaState efficiency.
etaState = 1;
rho = apply_loss(psi,etaState,S);
% Now it must be represented by a density matrix, rho.
wignerStepSize = 0.1;
[x,p] = meshgrid(-4:wignerStepSize:4,-4:wignerStepSize:4);
wigner2 = wigner(rho, x,p);
mesh(x,p,wigner2); xlabel('x'); ylabel('p'); zlabel('W(x,p)');
```

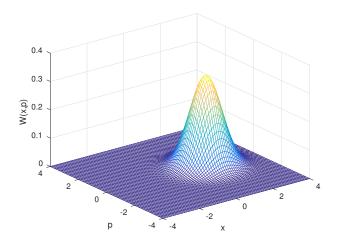


Figura 3: Wigner function of a coherent state with $\alpha = 1$.

using superpositions of coherent states, also known as Schrödinger cat states, such as:

$$|\text{cat}\rangle = \frac{1}{\sqrt{2(1 + e^{-2|\alpha|^2}\cos\theta)}} \left(e^{i\theta}|-\alpha\rangle + |\alpha\rangle\right).$$
 (3)

The routines to generate a cat state with amplitude α and to plot the corresponding Wigner function for such states are shown bellow.

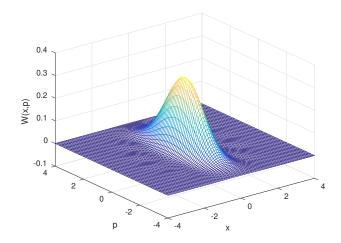


Figura 4: Wigner function of a squeezed vacuum state with variance equals to 1/6 of vacuum variance.

```
function catPsi = generate cat vector(alpha, theta, maxPhotons)
% Creates Schrodinger cat state
% catPsi = generate cat vector(alpha, theta, maxPhoton) returns the pure
% state vector for a Schrodinger cat state (namely,
% e^{i*theta}|-alpha>+|alpha>) in the photon number basis. maxPhoton is
% the photon number at which the Hilbert space is truncated or the table
if isstruct(maxPhotons)
   maxPhotons = maxPhotons.photons;
constant = (exp(1i .* theta) .*
sqrt(factorial(0:maxPhotons));
catPsi = catPsi.';
                                                              (-alpha).^(\theta:maxPhotons) \ + \ alpha.^(\theta:maxPhotons)) \ ./
normalization = exp(-abs(alpha).^
catPsi = normalization .* catPsi;
                                                                   ^2./2)./sqrt(2.*(1+exp(-2*abs(alpha)^2).*cos(theta)));
catPsi = normalize(catPsi,'check');
 % The infinite dimensional state space for the harmonic oscillator will be
% represented in the photon number basis.
% We will truncate the Hilbert space at maxPhotonNumber photons.
 maxPhotonNumber = 10;
 % First, pre-compute a lot of numbers, such as coefficients for Hermite % polynomials, factorials, binomial coefficients.
 % polynomials, factorials, binomi
S = init tables(maxPhotonNumber);
 % Make state vector for Schrodinger cat state.
alpha = 2; % amplitude of coherent states in the superposition
phase = 0; % phase between superposition
psi = generate_cat_vector(alpha, phase, S);
% The Schrodinger_cat state suffers from some loss by passing through a
  % medium with etaState efficiency
  rho = apply_loss(psi,etaState,S);
% Now it must be represented by a density matrix, rho.
 wignerStepSize = 0.1:
 wigherStep512e = 0.1;
[x,p] = meshgrid(-5:wignerStepSize:5,-5:wignerStepSize:5);
wigner2 = wigner(rho, x,p);
mesh(x,p,wigner2); xlabel('x'); ylabel('p'); zlabel('W(x,p)');
```

Figure 5 shows the Wigner function of a cat state with amplitude of coherent states in the superposition of $\alpha = 2$ and superposition phase $\theta = 0$.

The effect of loss in any of the states showed here can be considered and observed by changing the etaState parameter when using the lines of codes. Let us consider, for example, the effect of loss in the cat state of Figure 5. For that, consider that the state passed through a medium with 90% efficiency (etaState = 0.9). The Wigner function for the state after the loss is shown in figure 6.

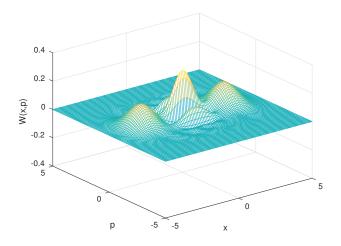


Figura 5: Wigner function of a cat state with amplitude of coherent states in the superposition $\alpha=2$ and superposition phase $\theta=0$.

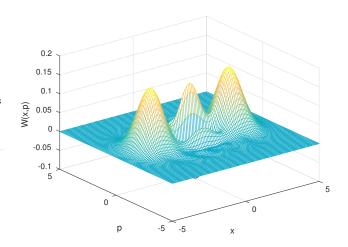


Figura 6: Wigner function of a cat state with amplitude of coherent states in the superposition $\alpha = 2$ and superposition phase $\theta = 0$ after going through a medium with 90% efficiency.

We then may use the following lines of code to calculate how "close" the state with loss is of the original state. That is done using the concept of fidelity, defined, for any two density matrices, ρ and σ , as

$$F(\rho,\sigma) = \left(Tr\left(\sqrt{\sqrt{\rho}\sigma\sqrt{\rho}}\right)\right)^2,\tag{4}$$

where the fidelity is equal to 1 if and only if the two states do coincide.

In the example given, the fidelity is F = 0.7171.

```
% The infinite dimensional state space for the harmonic oscillator will be % represented in the photon number basis.

We will truncate the Hilbert space at maxPhotonNumber photons.
maxPhotonNumber = 10;

% First, pre-compute a lot of numbers, such as coefficients for Hermite % polynomials, factorials, binomial coefficients.

S = init_tables(maxPhotonNumber);

% Make state vector for Schrodinger cat state.
alpha = 2; % amplitude of coherent states in the superposition phase = 0; % phase between superposition phase, S);

% The Schrodinger cat state suffers from some loss by passing through a % medium with etaState efficiency.
etaStatel = 1; rhol = apply_loss(psil,etaStatel,S);

% Now it must be represented by a density matrix, rho.
psi2 = generate_cat_vector(alpha, phase, S);
etaState2 = 0.9;
rho2 = apply_loss(psi2,etaState2,S);

F = fidelity(rho1, rho2)
```

4. Conclusion

The toolbox of computational routines described in this paper has been proved useful for teaching quantum mechanics and quantum optics courses. The lines of code developed for this toolbox facilitates the visualization of quantum states using the Wigner quasiprobability distibution, also called the Wigner function.

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