

# Global HMS for Genus Two Curves

Apr. 30, 2021

Western Hemisphere Virtual Symplectic Seminar

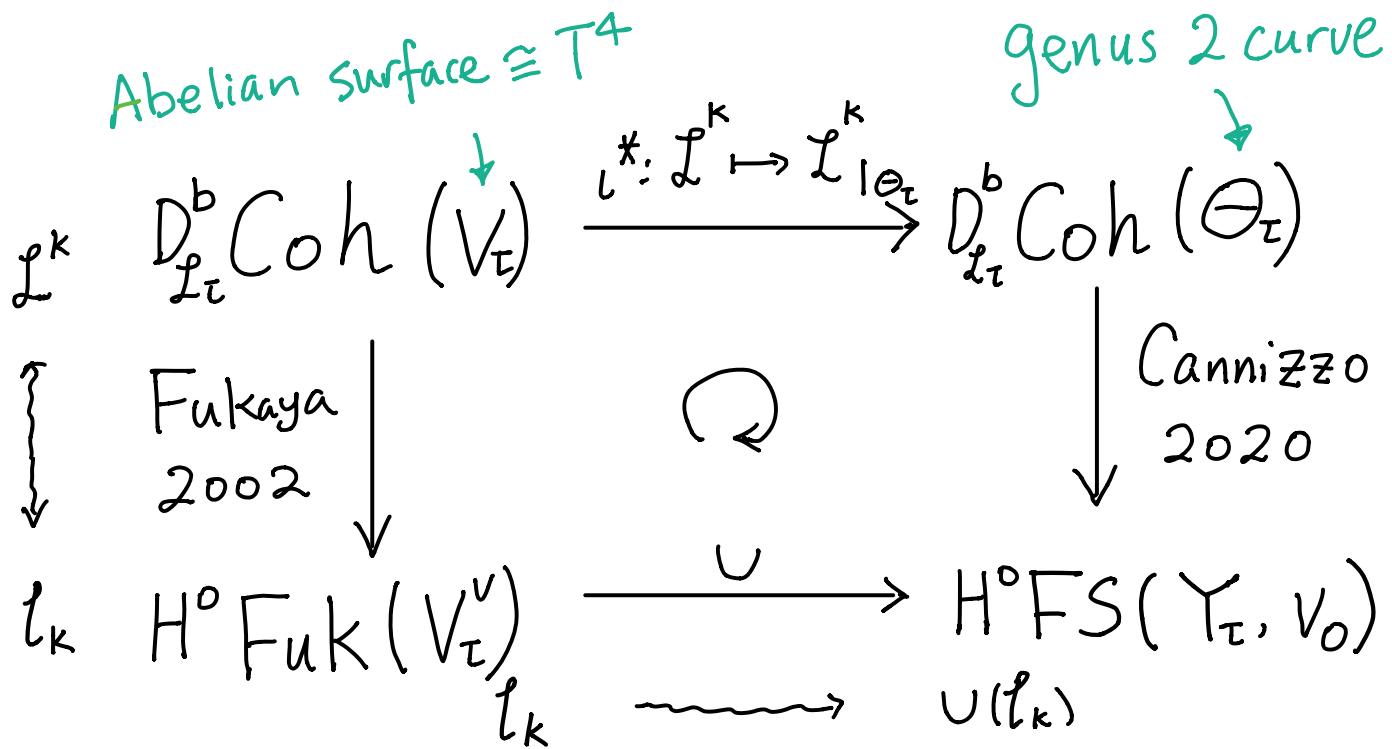
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joint work with Haniya Azam

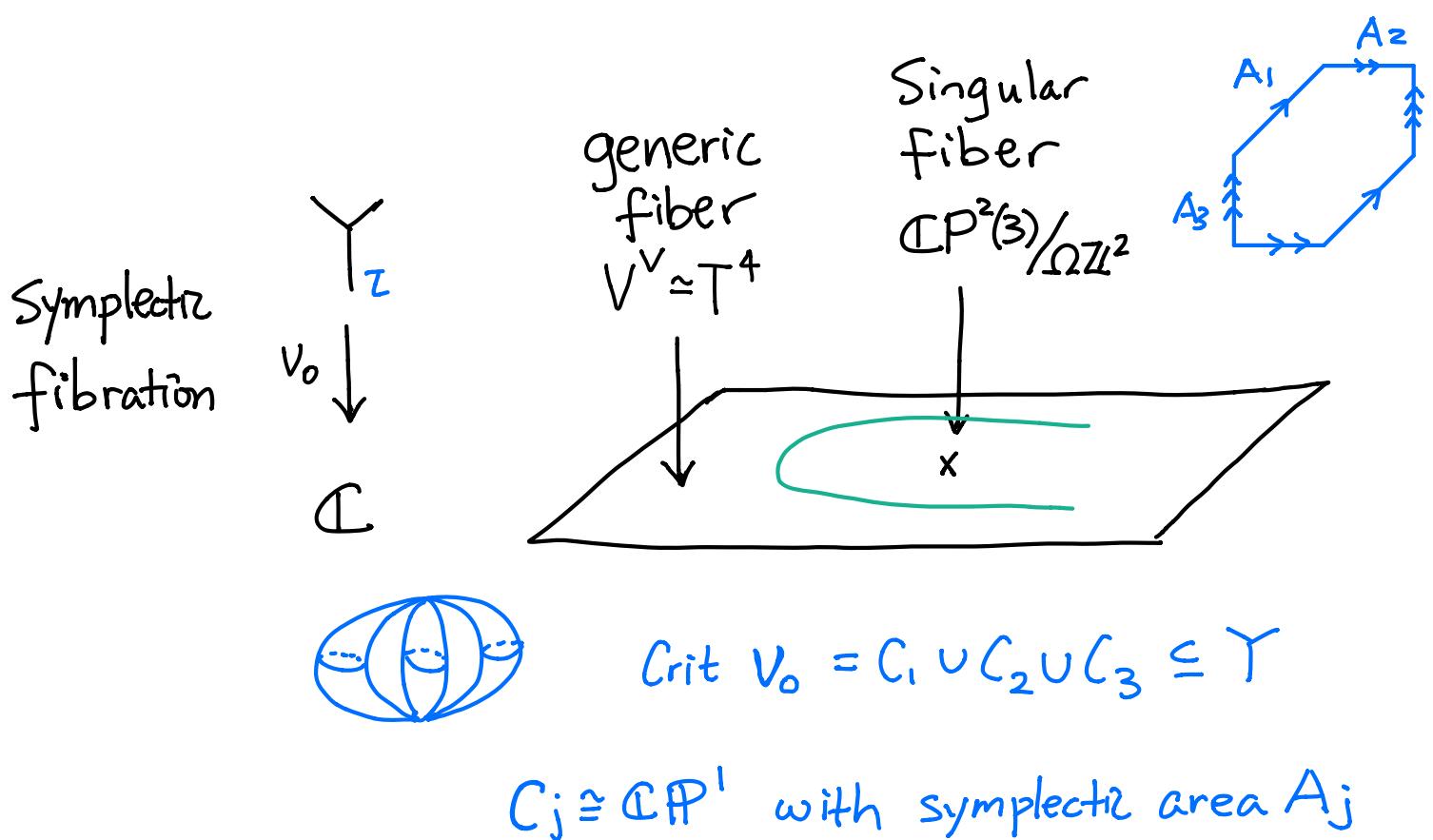
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# Catherine Cannizzo's Thesis 2019



Mirror construction based on SYZ, AAK  $\text{Bl}_{\Theta_T \times \{0\}} V_T \times \mathbb{C}$



## Genus two curve for any complex parameter

$$\tau = \begin{pmatrix} \tau_{11} & \tau_{12} \\ \tau_{21} & \tau_{22} \end{pmatrix} = B + i\Omega, \quad \text{symmetric}, \quad \det \Omega > 0$$

In Cannizzo's thesis ,  $\tau = \frac{i}{2\pi} \log t \begin{pmatrix} 2 & 1 \\ 1 & 2 \end{pmatrix}$ ,  $t$  large.

$$\Theta_\tau = \left\{ \theta(x_1, x_2; \tau) = \sum_{n=(n_1, n_2) \in \mathbb{Z}^2} x_1^{n_1} x_2^{n_2} e^{\pi i n^\top \tau n} = 0 \right\} / \mathbb{Z}^2$$

¶

$$V_\tau = (\mathbb{C}^*)^2 / \tau \mathbb{Z}^2 = \{(x_1, x_2) \in (\mathbb{C}^*)^2\} /_{(x_1, x_2) \sim (\tau n) \cdot (x_1, x_2)}$$

$$\begin{aligned} \uparrow \quad & \tau \begin{pmatrix} x_1 = e^{2\pi i v_1} \\ x_2 = e^{2\pi i v_2} \end{pmatrix} & \tau(n_1, n_2) \cdot (x_1, x_2) &= \left( \frac{\tau_{11} n_1 + \tau_{12} n_2}{\tau_{21} n_1 + \tau_{22} n_2} \right) \cdot (x_1, x_2) \\ & & &= \left( e^{2\pi i (\tau_{11} n_1 + \tau_{12} n_2)} x_1, e^{2\pi i (\tau_{21} n_1 + \tau_{22} n_2)} x_2 \right) \end{aligned}$$

$$V_\tau^+ = \mathbb{C}^2 / \mathbb{Z}^2 + \tau \mathbb{Z}^2 = \{(v_1, v_2) \in \mathbb{C}^2\} / \mathbb{Z}^2 + \tau \mathbb{Z}^2$$

## Mirror to abelian surface $V_\tau$

$$\begin{aligned} V_\tau^V \cong \mathbb{R}^4 / \mathbb{Z}^4, \quad \omega_{V^V}^C &= \underbrace{\left( \sum_{j,k=1}^2 B_{jk} df_j \wedge d\theta_k \right)}_{(\rho, \theta)} + i \underbrace{\left( \sum_{j,k=1}^2 \Omega_{jk} df_j \wedge d\theta_k \right)}_{= \sum_{k=1}^2 d\xi_k \wedge d\theta_k \text{ (for } \xi = \Omega \beta\text{)}} \\ &= \sum_{k=1}^2 d\xi_k \wedge d\theta_k \end{aligned}$$

$$\ell_k = \{(\rho, \theta) \in V_\tau^V \mid \theta = -k\rho\}$$

# Mirror Landau-Ginzburg model $(Y, v_0)$

$$(x_1, x_2) \mapsto \xi = (\xi_1, \xi_2) = (\log|x_1|, \log|x_2|)$$

$$\begin{array}{ccccc} \widetilde{\Theta}_\tau & \hookrightarrow & (\mathbb{C}^*)^2 & \longrightarrow & \mathbb{R}^2 \\ \downarrow & & \downarrow & & \downarrow \\ \Theta_\tau & \hookrightarrow & V_\tau = (\mathbb{C}^*)^2 / \tau \mathbb{Z}^2 & \longrightarrow & \mathbb{R}^2 / \Omega \mathbb{Z}^2 \end{array}$$

[Abouzaid-Auroux-Katzarkov]

$$\begin{aligned} \Delta_{\widetilde{Y}_\tau} &= \left\{ (\xi, \eta) \in \mathbb{R}^3 \mid \eta \geq \psi(\xi) = \max_{n \in \mathbb{Z}^2} \{ \langle \xi, n \rangle + k(n) \} \right\} \\ &= \bigcap_{n \in \mathbb{Z}^2} \left\{ (\xi, \eta) \in \mathbb{R}^3 \mid L_n(\xi, \eta) = -n_1 \xi_1 - n_2 \xi_2 + \eta - k(n) \geq 0 \right\} \end{aligned}$$

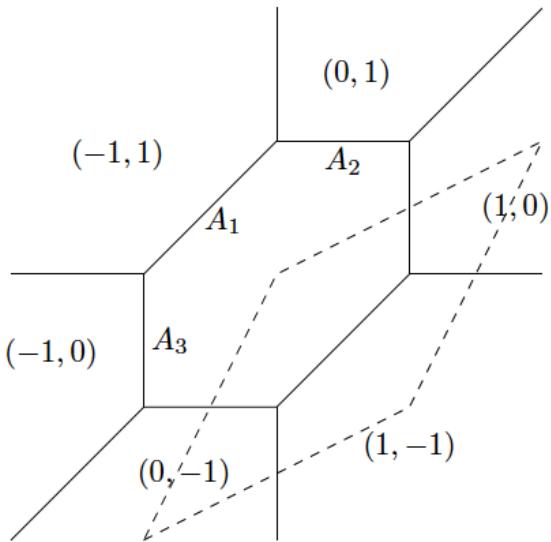
$$v_n = \begin{pmatrix} -n_1 \\ -n_2 \\ 1 \end{pmatrix} \quad \text{normal to facet } \{L_n(\xi, \eta) = 0\}$$

$$\Delta_{Y_\tau} = \Delta_{\widetilde{Y}_\tau} / \Omega \mathbb{Z}^2 \quad \begin{matrix} \text{(quotient near singular fiber} \\ \text{of } v_0 : \widetilde{Y} \rightarrow \mathbb{C} \end{matrix}$$

$$(\Omega m) \cdot (\xi_1, \xi_2, \eta) = (\xi_1 + (\Omega m)_1, \xi_2 + (\Omega m)_2, \eta - k(m) + \langle \xi, m \rangle)$$

Case 1 :

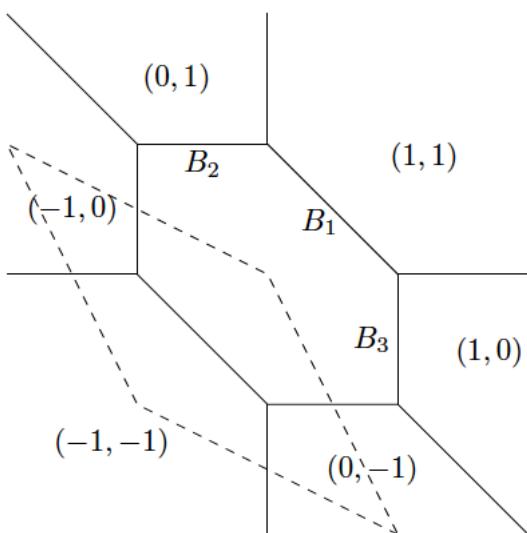
$$\begin{cases} \Omega_{12} = \Omega_{21} > 0 \\ \Omega_{11} - \Omega_{12} > 0 \\ \Omega_{22} - \Omega_{12} > 0 \end{cases}$$



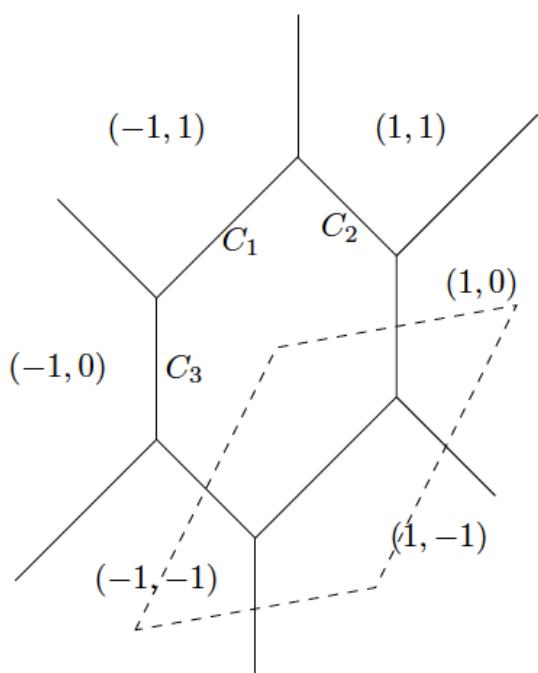
$$\begin{pmatrix} \Omega_{11} & \Omega_{12} \\ \Omega_{21} & \Omega_{22} \end{pmatrix} = \begin{pmatrix} A_1 + A_2 & A_1 \\ A_1 & A_1 + A_3 \end{pmatrix}$$

Case 2 :

$$\begin{cases} \Omega_{12} = \Omega_{21} < 0 \\ \Omega_{11} - \Omega_{12} > 0 \\ \Omega_{22} - \Omega_{12} > 0 \end{cases}$$



$$\begin{pmatrix} \Omega_{11} & \Omega_{12} \\ \Omega_{21} & \Omega_{22} \end{pmatrix} = \begin{pmatrix} B_1 + B_2 & -B_1 \\ -B_1 & B_1 + B_3 \end{pmatrix}$$



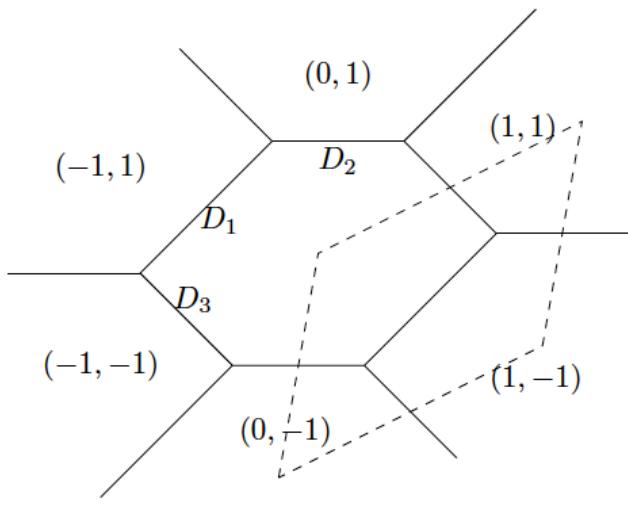
Case 3 :

$$\begin{cases} \Omega_{12} = \Omega_{21} > 0 \\ \Omega_{11} - \Omega_{12} < 0 \\ \Omega_{22} - \Omega_{12} > 0 \end{cases}$$

$$\begin{pmatrix} \Omega_{11} & \Omega_{12} \\ \Omega_{21} & \Omega_{22} \end{pmatrix} = \begin{pmatrix} C_1 + C_2 & \frac{C_1 - C_2}{2} \\ \frac{C_1 - C_2}{2} & C_3 \end{pmatrix}$$

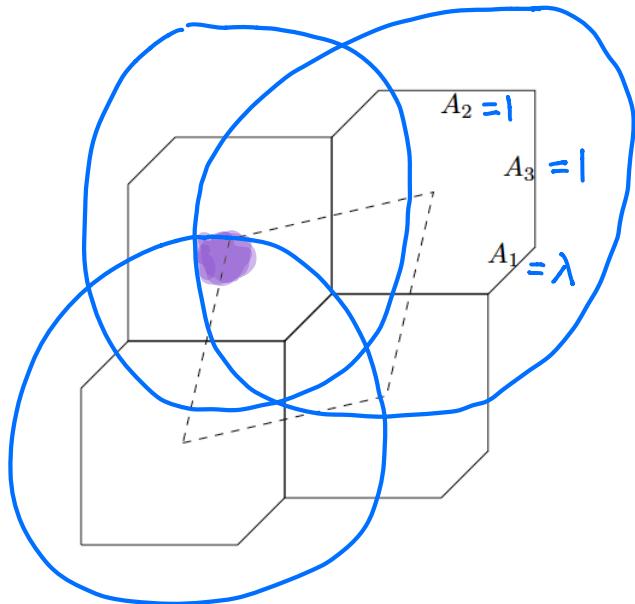
Case 4:

$$\begin{cases} \Omega_{12} = \Omega_{21} > 0 \\ \Omega_{11} - \Omega_{12} > 0 \\ \Omega_{22} - \Omega_{12} < 0 \end{cases}$$



$$\begin{pmatrix} \Omega_{11} & \Omega_{12} \\ \Omega_{21} & \Omega_{22} \end{pmatrix} = \begin{pmatrix} D_2 & \frac{D_1 - D_3}{2} \\ \frac{D_1 - D_3}{2} & D_1 + D_3 \end{pmatrix}$$

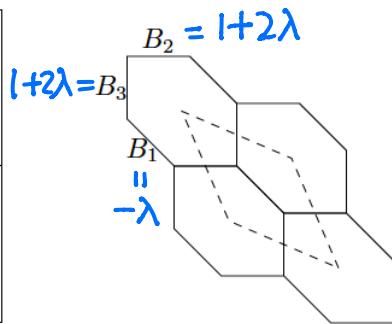
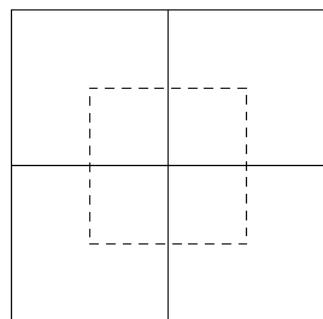
Example  $\Omega = \begin{pmatrix} 1+\lambda & \lambda \\ \lambda & 1+\lambda \end{pmatrix}$



$$\lambda > 0$$

$$\lambda = 0$$

$$\lambda < 0$$



# Symplectic structure on $\tilde{Y}$

of Guillemin, Kanazawa-Lau

On  $\tilde{Y}$ :

$$\tilde{G}(\xi, \eta) = \sum_{n \in \mathbb{Z}^2} \chi_n(\xi, \eta) L_n(\xi, \eta) \log L_n(\xi, \eta)$$

$$\Psi^{jk} = \frac{\partial^2 G}{\partial \xi_j \partial \xi_k} = \frac{\partial p_j}{\partial \xi_k}, \quad j, k = 1, 2, 3 := \eta, \quad \xi_\eta := \eta$$

$$p_j = \frac{\partial \tilde{G}}{\partial \xi_j}$$

$$\xi_j = \frac{\partial \tilde{F}}{\partial p_j}$$

$$\Psi_{jk} = (\Psi^{jk})^{-1} = \frac{\partial \xi_j}{\partial p_k} = \frac{\partial^2 \tilde{F}}{\partial p_j \partial p_k}$$

$$\omega = \sum_{j,k=1}^3 \Psi_{jk} dp_j \wedge d\theta_k = \frac{i}{2} \sum_{j,k=1}^3 \Psi_{jk} du_j \wedge du_k = \sum_{k=1}^3 d\xi_k \wedge d\theta_k$$

Complex toric coordinates

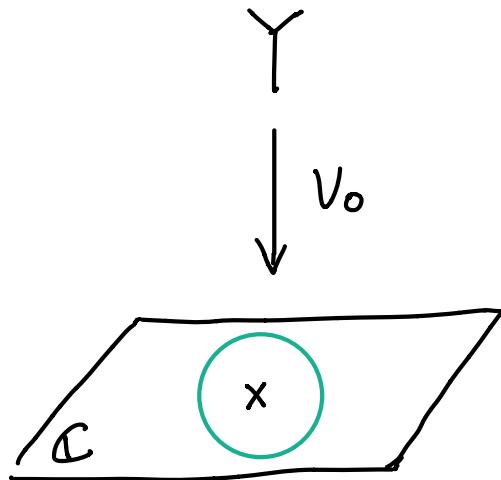
$$(t_1 = e^{u_1}, t_2 = e^{u_2}, t_3 = e^{u_3}) \in (\mathbb{C}^*)^3$$

$$u_j = p_j + i\theta_j$$

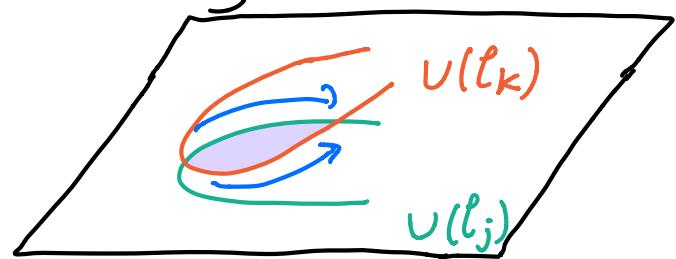
# Monodromy

Proposition:  $\varphi_{2\pi}(l_k)$  is Hamiltonian isotopic to  $l_{k+1}$

↑  
Monodromy, i.e. parallel transport along a  
circle  $re^{it}$  in the base  $\mathbb{C}$



Proposition useful for: Computing differentials



## Parallel transport

$$* \quad v_0 = t_3 = e^{\xi\eta + i\theta\eta} = e^{\frac{\partial G}{\partial \eta}(\xi, \eta)} e^{i\theta\eta}$$

\* Horizontal lift  $(\frac{\partial}{\partial \theta})^\#$  of  $\frac{\partial}{\partial \theta}$  is in the span of the Hamiltonian vector field generated by  $\log|v|$

$$d(\log|v|) = \bar{\Psi}^{\eta_1} d\xi_1 + \bar{\Psi}^{\eta_2} d\xi_2 + \bar{\Psi}^{\eta\eta} d\eta$$

$$\left(\frac{\partial}{\partial \theta}\right)^{\#} = \frac{\partial}{\partial \theta_{\eta}} + \frac{\Psi^{\eta_1}}{\Psi^{\eta\eta}} (\xi, \eta) \frac{\partial}{\partial \theta_1} + \frac{\Psi^{\eta_2}}{\Psi^{\eta\eta}} (\xi, \eta) \frac{\partial}{\partial \theta_2}$$

\* Parallel transport is

$$\Psi_t(\xi_1, \xi_2, \eta, \theta_1, \theta_2, \theta_{\eta})$$

$$= \left( \xi_1, \xi_2, \eta, \theta_1 + \frac{\Psi^{\eta_1}}{\Psi^{\eta\eta}}(\xi, \eta)t, \theta_2 + \frac{\Psi^{\eta_2}}{\Psi^{\eta\eta}}(\xi, \eta)t, \theta_{\eta} + t \right)$$

\* Proof of the proposition :

- On a fiber  $v_0^{-1}(-r)$ , has an Hamiltonian vector field

$$X = \sum_{j=1}^2 \underbrace{\left( \frac{\Psi^{\eta_j}}{\Psi^{\eta\eta}}(\xi, \eta) + \sum_{l=1}^2 \Omega^{jl} \xi_l \right)}_{=: f_j(\xi)} \frac{\partial}{\partial \theta_j}.$$

$$\text{s.t. } \phi_X^{2\pi}(l_{k+1}) = \varphi_{2\pi}(l_k)$$

- $\phi_X^t$  is well-defined on the quotient  $Y = \tilde{Y}/\Omega_{\mathbb{Z}^2}$

Reason: moment map is  $\Omega\mathbb{Z}^2$ -equivariant

can show  $f_j(\xi + \Omega m) = f_j(\xi) - m$ .

\* Extra fun fact:

Fix  $b > 0$ , consider

$$R_n(b) := \left\{ \xi \in \overset{\circ}{\Delta}_{\tilde{Y}} \mid L_m(\xi, \eta) \geq b \text{ if } m \neq n \right\}$$

As  $L_n(\xi, \eta) \rightarrow 0$  on  $R_n(b)$ , we have

$$\frac{\Psi^{\eta_j}}{\Psi^{\eta\eta}} (\xi, \eta) \rightarrow \begin{pmatrix} -n_1 \\ -n_2 \end{pmatrix} \quad \begin{matrix} \text{normal vector} \\ \text{fits intuition} \end{matrix}$$