

Notes of 141A

Solid State Physics

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7 Energy Bands

7.1 Nearly Free Elec Model

$$\begin{aligned}
 U(x) &= U_0 \left[\cos^4 \frac{\pi x}{a} - \frac{3}{8} \right] \\
 &= \dots \\
 &= U_0 \left(\frac{1}{2} \cos \frac{2\pi x}{a} + \frac{1}{8} \cos \frac{4\pi x}{a} \right)
 \end{aligned} \tag{7.1}$$

$$(\lambda_k - \epsilon)C_k + \sum_G U_G C_{k-G} = 0 \tag{7.2}$$

solve reduced problem near degenerate points

$$\begin{pmatrix} \lambda_k - \epsilon & U_g \\ U_g^* & \lambda_{k-g} - \epsilon \end{pmatrix} \begin{pmatrix} C_k \\ C_{k-g} \end{pmatrix} = 0 \quad (g = \frac{2\pi}{a}) \tag{7.3}$$

$$(\lambda_k - \epsilon)^2 - |U_g|^2 = 0 \tag{7.4}$$

$$\epsilon = \frac{\hbar^2 k^2}{2m} \pm \frac{U_0}{4} \tag{7.5}$$

or

$$\epsilon = \frac{\hbar^2 (k + G_1)^2}{2m} \pm \frac{U_0}{16} \tag{7.6}$$

7.2 Bloch Oscillators

Semi-classical eq. of motion

$$\hbar \frac{d\mathbf{k}}{dt} = \mathbf{F} = -e\mathbf{E} \tag{7.7}$$

$$v_{grp} = \frac{1}{\hbar} \frac{d\varepsilon}{dk} \tag{7.8}$$

consider

$$\varepsilon(k) = \varepsilon_0 (1 - \cos ka) \tag{7.9}$$

$$v = \frac{a\varepsilon_0}{\hbar} \sin ka \tag{7.10}$$

$$\begin{aligned}
 x &= \int \frac{a\varepsilon_0}{\hbar} \sin ka dt \\
 &= \frac{a\varepsilon_0}{\hbar} \int \sin ka dk \frac{dt}{dk} \\
 &= \dots \\
 &= -\frac{\varepsilon_0}{eE} \left(\cos \frac{-eEa}{\hbar} - 1 \right)
 \end{aligned} \tag{7.11}$$

$$T = \frac{\Delta k}{dk/dt} = \frac{2\pi/a}{eE/\hbar} = \dots \quad (7.12)$$

scattering time $\tau \gg T$.

suppose $\mathbf{F} = m^* \mathbf{a}$

$$m^* = \hbar \frac{dk/dt}{dv_g/dt} = \hbar \left(\frac{dv_g}{dk} \right)^{-1} = \hbar^2 \left(\frac{\partial^2 \varepsilon}{\partial k^2} \right)^{-1} \quad (7.13)$$

7.3 Consequence of Band Structure

even # of elec's per unit cell – insulator – C, Si, Ge

odd # of elec's per unit cell – metal – Cu, Ag, Au

Ins – $E_g > 2 \text{ eV}$

Semiconductor – $E_g < 2 \text{ eV}$

7.3.1 Density of States

Van Hove Singularities of DoS

8 Semiconductors

8.1 Photoconductivity

$$\mathbf{J}_n = \frac{1}{V} \int_{unocc} e D(\mathbf{k}_e) \mathbf{v}_e(\mathbf{k}_e) d^3 \mathbf{k}_e \quad (8.1)$$

$$\frac{d\mathbf{k}_e}{dt} = \frac{-e}{\hbar} (\mathbf{E} + \frac{1}{c} \mathbf{v}_e \times \mathbf{B}) \quad (8.2)$$

$$\mathbf{v}_e(\mathbf{k}_e) = \frac{1}{\hbar} \frac{\partial E_v(\mathbf{k}_e)}{\partial \mathbf{k}_e} \quad (8.3)$$

Holes: Def

$$\mathbf{k}_n = -\mathbf{k}_e \quad E_n = -E_v \quad (8.4)$$

$$\mathbf{v}_n(\mathbf{k}_n) = \frac{1}{\hbar} \frac{\partial E_n(\mathbf{k}_n)}{\partial \mathbf{k}_n} = \frac{1}{\hbar} \frac{\partial (-E_v(-\mathbf{k}_e))}{\partial (-\mathbf{k}_e)} = \mathbf{v}_e(\mathbf{k}_e) \quad (8.5)$$

Holes: + charge, evolves like particle, $m^* > 0$

8.2 Intrinsic Mobility

$$\mathbf{J}_{tot} = \mathbf{J}_e + \mathbf{J}_h \quad (8.6)$$

$$\sigma_{tot} = \sigma_e + \sigma_h \quad (8.7)$$

$$\sigma_e = \frac{n_e e^2 \tau_e}{|m_e^*|} \quad (8.8)$$

$$\sigma_h = \frac{n_h e^2 \tau_h}{|m_h^*|} \quad (8.9)$$

def: mobility

$$\mu = \frac{e\tau}{m^*} \quad (8.10)$$

$$\sigma_{tot} = n_e e \mu_e + n_h e \mu_h \quad (8.11)$$

8.3 Impurity Conductivity / Doping

$$n_e(T) \neq n_h(T) \Rightarrow \text{doping} \quad (8.12)$$

Impurity atoms that can give up an electron are called donors.

8.3.1 Donor States

the donated elec moves in the coulomb potential $-e/\epsilon r$

$$N_d = \frac{\#donors}{Vol.} \quad N_a = \frac{\#acceptors}{Vol.} \quad (8.13)$$

$$n_e(T) = 2 \left(\frac{m_e k_B T}{2\pi \hbar^2} \right)^{3/2} e^{-(E_c - \mu)/k_B T} \quad (8.14)$$

$$n_p(T) = 2 \left(\frac{m_h k_B T}{2\pi \hbar^2} \right)^{3/2} e^{-(\mu - E_v)/k_B T} \quad (8.15)$$

Intrinsic:

$$n_i(T) = \sqrt{n_e(T) n_p(T)} \quad (8.16)$$

$$\mu_i(T) = E_v + \frac{1}{2} E_g + \frac{3}{4} k_B T \ln \frac{m_h}{m_e} \quad (8.17)$$

$$\frac{N_d - N_a}{n_i(T)} = 2 \sinh \frac{\mu - \mu_i}{k_B T} \quad (8.18)$$

8.4 Hall Effect w/ 2 Carrier Types

Kittel 8.3

mobilities:

$$\mu_e = \frac{e\tau_e}{m_e} \quad \mu_h = \frac{e\tau_h}{m_h} \quad (8.19)$$

(m_e is eff mass)

Recall conductivity tensor w/ B-field

$$\sigma = \frac{\sigma_0}{(1 + \omega_c^2 \tau^2)^2} \begin{pmatrix} 1 & -\omega_c \tau \\ \omega_c \tau & 1 \end{pmatrix} \quad (8.20)$$

where $\sigma_0 = \frac{ne^2\tau_e}{m_e}$, $\omega_c = \frac{eB}{m_e c}$.
transverse current ($\omega_c \tau \ll 1$)

$$j_y(e) = \sigma_0(\omega_c \tau E_x + E_y) = ne\mu_e \left(\frac{\mu_e B}{c} E_x + E_y \right) \quad (8.21)$$

$$j_y(h) = pe\mu_h \left(\frac{-\mu_h B}{c} E_x + E_y \right) \quad (8.22)$$

longitudinal curr

$$j_x = j_x(e) + j_x(h) = (ne\mu_e + pe\mu_h)E_x \quad (8.23)$$

tot transverse curr = 0, thus

$$(n\mu_e^2 - p\mu_h^2) \frac{eB}{c} E_x + (n\mu_e + p\mu_h) e E_y = 0 \quad (8.24)$$

$$E_y = -E_x \frac{B}{c} \frac{n\mu_e^2 - p\mu_h^2}{n\mu_e + p\mu_h} \quad (8.25)$$

$$R_H = \frac{E_y}{j_x B} = -\frac{1}{ec} \frac{n\mu_e^2 - p\mu_h^2}{(n\mu_e + p\mu_h)^2} \quad (8.26)$$

8.5 Tight-binding

start w/ AOs (1-D)

$$\psi_k(r) = \sum_j c_j \phi(r - r_j) \quad (8.27)$$

(ϕ is s orb)

Bloch Th.

$$\psi_k(r) = e^{ikr} u(r) \quad (8.28)$$

if $c_j = \frac{1}{\sqrt{N}} e^{ikr}$

$$\langle \psi_k | \hat{\mathbf{H}} | \psi_k \rangle = \frac{1}{N} \sum_j \sum_m e^{ik(r_j - r_m)} \langle \phi_m | \hat{\mathbf{H}} | \phi_j \rangle \quad (8.29)$$

let $\rho_m = r_m - r_j$, suppose only nearest neighbor interacts

$$\langle \psi_k | \hat{\mathbf{H}} | \psi_k \rangle = \frac{1}{N} \sum_m e^{ik\rho_m} \int dV \phi^*(r - \rho_m) \hat{\mathbf{H}} \phi(r) \quad (8.30)$$

$(r - r_j \rightarrow r)$

Let

$$-\varepsilon_0 = \int dV \phi^*(r) \hat{\mathbf{H}} \phi(r) \quad -t = \int dV \phi^*(r - \rho) \hat{\mathbf{H}} \phi(r) \quad (8.31)$$

$$E = -\varepsilon_0 - t \sum_{n.n.} e^{-ik\rho_{n.n.}} \quad (8.32)$$

for cubic lattice (1-D)

$$E = -\varepsilon_0 - t(1 + 2 \cos ka) \quad (8.33)$$

consider a lattice w/ 2-atom basis

$$\psi_k(r) = \alpha_k \psi_k^A(r) + \beta_k \psi_k^B(r) \quad (8.34)$$

$$H_{AB} = \langle \psi_k^A | \hat{\mathbf{H}} | \psi_k^B \rangle = \begin{pmatrix} AHA & AHB \\ BHA & BHB \end{pmatrix} = \begin{pmatrix} -\varepsilon_1 & \sum_{n.n.} e^{-ik\rho_{n.n.}} \\ \sum_{n.n.} e^{ik\rho_{n.n.}} & -\varepsilon_2 \end{pmatrix} \quad (8.35)$$

$$H_{AB} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} = E_k \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \quad (8.36)$$

$$H = \begin{pmatrix} -\varepsilon_1 & t(1 + e^{-ika}) \\ t(1 + e^{ika}) & -\varepsilon_2 \end{pmatrix} \quad (8.37)$$

$$\det(H - E_k I) = 0 \quad (8.38)$$

$$E_k = \frac{1}{2}(\varepsilon_1 + \varepsilon_2) \pm \sqrt{\dots} \quad (8.39)$$

8.6 p-n Junctions

p-n jcn

– diode

– half-transistor

– solar cell

– LED

– laser

$$I = I_{CB} + I_{VB} \quad (8.40)$$

conduction band, valence band

$$I_{CB} = I_{nr} - I_{ng} \quad (8.41)$$

recomb, generation

$$I_{nr} \sim e^{-e(\Delta\phi-V)\beta} \quad (8.42)$$

$$I_{ng} \sim e^{-e\Delta\phi\beta} \quad (8.43)$$

$$ICB = Ing(e^{eV\beta} - 1) \quad (8.44)$$

$$IVB = I_{hr} - I_{hg} \quad (8.45)$$

$$IVB = I_{hg}(e^{eV\beta} - 1) \quad (8.46)$$

$$I = (Ing + I_{hg})(e^{eV\beta} - 1) \quad (8.47)$$

LED

Solar cell

Schottkey barrier

p-n-p jcn

9 Tight Binding

Start: Localized picture

$$\begin{aligned} \Psi_{\mathbf{k}} &= \frac{1}{\sqrt{N}} \sum_{j=1}^N C_{\mathbf{k},j} \phi(\mathbf{r} - \mathbf{r}_j) \\ &= \frac{1}{\sqrt{N}} \sum_{j=1}^N e^{i\mathbf{k} \cdot \mathbf{r}_j} \phi(\mathbf{r} - \mathbf{r}_j) \quad (\text{Bloch Th.}) \end{aligned} \quad (9.1)$$

$$\begin{aligned} \varepsilon(\mathbf{k}) &= \langle \Psi_{\mathbf{k}} | \hat{\mathbf{H}} | \Psi_{\mathbf{k}} \rangle = \frac{1}{N} \left\langle \sum_{j=1}^N e^{i\mathbf{k} \cdot \mathbf{r}_j} \phi(\mathbf{r} - \mathbf{r}_j) \left| \hat{\mathbf{H}} \right| \sum_{j=1}^N e^{i\mathbf{k} \cdot \mathbf{r}_j} \phi(\mathbf{r} - \mathbf{r}_j) \right\rangle \\ &= \frac{1}{N} \sum_j \sum_n e^{i\mathbf{k} \cdot (\mathbf{r}_j - \mathbf{r}_n)} \langle \phi(\mathbf{r} - \mathbf{r}_n) | \hat{\mathbf{H}} | \phi(\mathbf{r} - \mathbf{r}_j) \rangle \end{aligned} \quad (9.2)$$

Let

$$\boldsymbol{\rho}_m = \mathbf{r}_j - \mathbf{r}_n \quad \boldsymbol{\xi}_n = \mathbf{r} - \mathbf{r}_n \quad (9.3)$$

$$\begin{aligned} \varepsilon(\mathbf{k}) &= \frac{1}{N} \sum_m \sum_n e^{i\mathbf{k} \cdot \boldsymbol{\rho}_m} \langle \phi(\boldsymbol{\xi}_n) | \hat{\mathbf{H}} | \phi(\boldsymbol{\xi}_n - \boldsymbol{\rho}_m) \rangle \\ &= \sum_m e^{i\mathbf{k} \cdot \boldsymbol{\rho}_m} \langle \phi(\boldsymbol{\xi}) | \hat{\mathbf{H}} | \phi(\boldsymbol{\xi} - \boldsymbol{\rho}_m) \rangle \\ &\equiv \sum_m e^{i\mathbf{k} \cdot \boldsymbol{\rho}_m} I(\boldsymbol{\rho}_m) \end{aligned} \quad (9.4)$$

Overlap integral

$$I(\rho_m) = \left\langle \phi(\mathbf{r}) \left| \hat{\mathbf{H}} \right| \phi(\mathbf{r} - \boldsymbol{\rho}_m) \right\rangle \quad (9.5)$$

Now neglect all non-neighbor interaction.

when $\rho_m = \rho(t)$: hopping amplitude.)

$$-t = I(\rho) = \left\langle \phi(\mathbf{r}) \left| \hat{\mathbf{H}} \right| \phi(\mathbf{r} - \boldsymbol{\rho}) \right\rangle \quad (9.6)$$

when $\rho_m = 0$

$$-\alpha = \left\langle \phi(\mathbf{r}) \left| \hat{\mathbf{H}} \right| \phi(\mathbf{r}) \right\rangle \quad (9.7)$$

thus

$$\varepsilon(\mathbf{k}) = -\alpha - t \sum_m e^{i\mathbf{k} \cdot \boldsymbol{\rho}_m} \quad (9.8)$$

9.0.1 1-D Crystal

$$\begin{aligned} \varepsilon(k) &= \dots = -\alpha - 2t \cos ka \\ &= -\alpha - 2t \left(1 - \frac{k^2 a^2}{2} \right) \\ &= -\alpha - 2t + tk^2 a^2 \end{aligned} \quad (9.9)$$

let

$$m^* = \frac{\hbar^2}{2ta^2} \quad (9.10)$$

$$\varepsilon(k) = -\alpha - 2t + \frac{\hbar^2 k^2}{2m^*} \quad (9.11)$$

10 Screening

$$D(k, \omega) = \varepsilon(k, \omega) E(k, \omega) \quad (10.1)$$

10.0.1 Static Screening

$\omega = 0$

$$E(r) = \sum_k e^{ik \cdot r} E(k) \quad (10.2)$$

$$\rho(r) = \sum_k e^{ik \cdot r} \rho(k) \quad (10.3)$$

$$D(r) = \sum_k e^{ik \cdot r} D(k) \quad (10.4)$$

$$\nabla \cdot E = \nabla \cdot \left(\sum_k e^{ik \cdot r} E(k) \right) = 4\pi \sum_k e^{ik \cdot r} \rho(k) \quad (10.5)$$

$$\nabla \cdot D = \nabla \cdot \left(\sum_k e^{ik \cdot r} \varepsilon(k) E(k) \right) = 4\pi \sum_k e^{ik \cdot r} \rho_{ext}(k) \quad (10.6)$$

\therefore

$$\varepsilon(k) = \frac{\rho_{ext}(k)}{\rho(k)} = 1 - \frac{\rho_{ind}(k)}{\rho(k)} \quad (10.7)$$

In terms of potential, ...

10.0.2 Calculating ρ_{ind} Using Thomas-Fermi Theory of Screening

$$-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \psi_i(r) = E_i \psi_i(r) \quad (10.8)$$

$$V = -e\phi(r) \quad (10.9)$$

$\phi(r)$ varies slowly enough that

$$E_i = \frac{\hbar^2 k^2}{2m} - e\phi(r) \quad (10.10)$$

10.1 Graphene Tight-binding

General tight-binding for 2 atom-basis

$$\psi_k = a_k \psi_k^{(A)}(r) + b_k \psi_k^{(B)}(r) \quad (10.11)$$

$$\psi_k^{(B)}(r) = \sum_j c_j \phi^{(B)}(r - r_j) \quad (10.12)$$