

Notes of JU Guoxing TD&SP

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29 Bose-Einstein and Fermi-Dirac Distributions

29.1

29.2 Wave Function of Identical Particles

29.3 The Statistics of Identical Particles

suppose energy of each particle is E

$$\mathcal{Z} = \sum_N \sum_{\alpha} e^{\beta(\mu N - E_{\alpha})} = \sum_N e^{N\beta(\mu - E)} \quad (29.3.1)$$

thus

$$\langle n \rangle = \frac{\sum_N N e^{N\beta(\mu - E)}}{\sum_N e^{N\beta(\mu - E)}} = -\frac{1}{\beta \mathcal{Z}} \frac{\partial \mathcal{Z}}{\partial E} = -\frac{1}{\beta} \frac{\partial \ln \mathcal{Z}}{\partial E} \quad (29.3.2)$$

for fermions

$$\mathcal{Z} = \sum_{N=0}^1 e^{N\beta(\mu - E)} = 1 + e^{\beta(\mu - E)} \quad (29.3.3)$$

Fermi-Dirac distribution fxn

$$f_D(E) \equiv \langle n \rangle = \frac{e^{\beta(\mu - E)}}{1 + e^{\beta(\mu - E)}} = \frac{1}{e^{\beta(E - \mu)} + 1} \quad (29.3.4)$$

for bosons

$$\mathcal{Z} = \sum_{N=0}^{\infty} e^{N\beta(\mu - E)} = \frac{1}{1 - e^{\beta(\mu - E)}} \quad (29.3.5)$$

Bose-Einstein distribution fxn

$$f_B(E) \equiv \langle n \rangle = \frac{e^{\beta(\mu - E)}}{1 - e^{\beta(\mu - E)}} = \frac{1}{e^{\beta(E - \mu)} - 1} \quad (29.3.6)$$

30 Quantum Gases and Condensates

30.1 non-interacting q fluid

$$\mathcal{Z} = \prod_k \mathcal{Z}_k^{2S+1} \quad (30.1.1)$$

where \mathcal{Z}_k refers to (29.3.3)(29.3.5)

i.e.

$$\mathcal{Z}_k = (1 \pm e^{(\mu-E)})^{\pm 1} \quad (30.1.2)$$

grand potential

$$\begin{aligned} \Phi &= -\frac{1}{\beta} \ln \mathcal{Z} \\ &= \mp \frac{2S+1}{\beta} \sum_k \ln(1 \pm e^{(\mu-E)}) \end{aligned} \quad (30.1.3)$$

$$\begin{aligned} &= \mp \frac{1}{\beta} \int_0^\infty \ln(1 \pm e^{(\mu-E)}) g(E) dE \\ g(k) dk &= \frac{4\pi k^2 dk}{(2\pi/L)^3} (2S+1) = \frac{(2S+1)V k^2 dk}{2\pi^2} \end{aligned} \quad (30.1.4)$$

$$g(E) dE = \frac{(2S+1)V E^{1/2} dE}{4\pi^2} \left(\frac{2m}{\hbar^2} \right)^{3/2} \quad (30.1.5)$$

thus

$$\begin{aligned} \Phi &= \mp \frac{(2S+1)V}{4\pi^2 \beta} \left(\frac{2m}{\hbar^2} \right)^{3/2} \int_0^\infty \ln(1 \pm e^{\beta(\mu-E)}) E^{1/2} dE \\ &= \mp \frac{(2S+1)V}{4\pi^2 \beta} \left(\frac{2m}{\hbar^2} \right)^{3/2} \int_0^\infty \ln(1 \pm e^{\beta(\mu-E)}) \frac{2}{3} d(E^{3/2}) \\ &= \mp \frac{2}{3} \frac{(2S+1)V}{4\pi^2 \beta} \left(\frac{2m}{\hbar^2} \right)^{3/2} \left\{ \left[\ln(1 \pm e^{\beta(\mu-E)}) E^{3/2} \right] \Big|_0^\infty - \int_0^\infty E^{3/2} \frac{\mp \beta e^{\beta(\mu-E)}}{1 \pm e^{\beta(\mu-E)}} dE \right\} \\ &= \mp \frac{2}{3} \frac{(2S+1)V}{4\pi^2 \beta} \left(\frac{2m}{\hbar^2} \right)^{3/2} \left\{ 0 \pm \beta \int_0^\infty E^{3/2} \frac{1}{e^{\beta(E-\mu)} \pm 1} dE \right\} \\ &= \frac{2}{3} \frac{(2S+1)V}{4\pi^2} \left(\frac{2m}{\hbar^2} \right)^{3/2} \int_0^\infty \frac{E^{3/2}}{e^{\beta(E-\mu)} \pm 1} dE \end{aligned} \quad (30.1.6)$$