Notes of JU Guoxing TD&SP

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23 Photon

23.1 Classic Thermodynamics of Electromagnetic Radiation

Energy density

$$u = \frac{U}{V} = n\hbar\omega \tag{23.1.1}$$

pressure

$$pV = \frac{1}{3}Nm\langle v^2 \rangle \tag{23.1.2}$$

for photon

$$p = \frac{1}{3}nmc^2 (23.1.3)$$

24 Phonon

- 24.1 Einstein Model
- 24.2 Debye Model

25 Relativistic Gas

25.1

25.2 Ultrarelativistic Gas

$$E = pc = \hbar kc \tag{25.2.1}$$

single-particle partition function

$$Z_1 = \int_0^\infty e^{-\beta\hbar kc} g(k) dk \qquad (25.2.2)$$

where

$$g(k)dk = \frac{Vk^2dk}{2\pi^2}$$
 (25.2.3)

Let $x = \beta \hbar kc$

$$Z_{1} = \frac{V}{2\pi^{2}} \left(\frac{1}{\beta \hbar c}\right)^{3} \int_{0}^{\infty} e^{-x} x^{2} dx = \frac{V}{\pi^{2}} \left(\frac{1}{\beta \hbar c}\right)^{3}$$
 (25.2.4)

or
$$Z_1 = \frac{V}{\Lambda^3}$$

N-particle partition fxn

$$Z_N = \frac{Z_1^N}{N!} (25.2.5)$$

25.3 Adiabatic Expansion

$$VT^3 = Cons. (25.3.1)$$

26 Real Gases

26.1 vdW Gas

$$\left(p + \frac{a}{V_m^2}\right)(V_m - b) = RT$$
(26.1.1)

critical point (临界点)

$$\left(\frac{\partial p}{\partial V}\right) = -\frac{RT}{(V-b)^2} + \frac{2a}{V^3} = 0 \tag{26.1.2}$$

$$\left(\frac{\partial^2 p}{\partial V^2}\right) = \dots \tag{26.1.3}$$

. . .

critical vol

$$V_c = 3b \tag{26.1.4}$$

critical temp

$$T_c = \frac{8a}{27Rb} {26.1.5}$$

critical pres

$$p_c = \frac{a}{27b^2} \tag{26.1.6}$$

thus

$$\frac{p_c V_c}{RT_c} = \frac{3}{8} {26.1.7}$$

26.2 Dieterici Equation

vdW Eq.

$$p = \frac{RT}{V - b} - \frac{a}{V_m^2}$$

$$\equiv p_{rep} + p_{attr}$$
(26.2.1)

Dieterici Eq.

$$p = p_{rep} e^{-\frac{a}{RTV}}$$
 (26.2.2)

26.3 Virial Expansion

$$\frac{pV_m}{RT} = 1 + \frac{B}{V_m} + \frac{C}{V_m^2} + \cdots {26.3.1}$$

B, C, etc., are called virial coefficients.

Try to drive B:

$$U = \sum_{i=1}^{N} \frac{p_i^2}{2m} + \sum_{i < j} V(|\mathbf{r}_i - \mathbf{r}_j|)$$
 (26.3.2)

$$Z = \frac{1}{N!h^{3N}} \tag{26.3.3}$$

...

$$p = \dots (26.3.4)$$

Using hard sphere model

$$V(r) = \begin{cases} \infty & r < r_0 \\ -\varepsilon \left(\frac{r_0}{r}\right)^6 & r > r_0 \end{cases}$$
 (26.3.5)

we have

$$B = b - \frac{a\beta}{N} \tag{26.3.6}$$

where

$$a = \dots \quad b = \dots$$
 (26.3.7)

26.4 The Law of Corresponding States

Def

$$\tilde{p} = \frac{p}{p_c}, \quad \tilde{V} = \frac{V}{V_c}, \quad \tilde{T} = \frac{T}{T_c}$$
 (26.4.1)

thus

$$\left(\tilde{p} + \frac{3}{\tilde{V}^2}\right) = \frac{8\tilde{T}}{3\tilde{V} - 1} \tag{26.4.2}$$

27 Cooling Real Gas

27.1 The Joule Expansion

U unchanged. Irreversible.

Joule coefficient

$$\mu_J = \left(\frac{\partial T}{\partial V}\right)_{IJ} \tag{27.1.1}$$

thus

$$\mu_{J} = -\left(\frac{\partial T}{\partial U}\right)_{V} \left(\frac{\partial U}{\partial V}\right)_{T} = -\frac{1}{C_{V}} \left[T\left(\frac{\partial p}{\partial T}\right)_{V} - p\right]$$
(27.1.2)

For ideal gas

$$\mu_J = 0 (27.1.3)$$

For vdW gas

$$\mu_J = -\frac{a}{C_V V^2} \tag{27.1.4}$$

$$\Delta T = \dots \tag{27.1.5}$$

27.2 Isothermal Expansion

$$\left(\frac{\partial U}{\partial V}\right)_T = T \left(\frac{\partial p}{\partial T}\right)_V - p \tag{27.2.1}$$

$$\Delta U = \int_{V_1}^{V_2} \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right] dV$$
 (27.2.2)

For ideal gas

$$\Delta U = 0 \tag{27.2.3}$$

For vdW gas

27.3 Joule-Kelvin Expansion

aka Joule-Thomson expansion.

28 Phase Transitions

28.1 Latent Heat

28.2 Chemical Potential and Phase Changes

Phase equilibrium

$$\mu_1 = \mu_2 \tag{28.2.1}$$

28.3 Clausius-Clapeyron Equation

$$\mu_1(p,T) = \mu_2(p,T) \tag{28.3.1}$$

$$\mathrm{d}\mu_1 = \mathrm{d}\mu_2 \tag{28.3.2}$$

with Gibbs-Duhem Eq.:

$$d\mu = -sdT + vdp \tag{28.3.3}$$

where $s = \frac{S}{N}, \ v = \frac{V}{N}$ thus

$$-s_1 dT + v_1 dp = -s_2 dT + v_2 dp$$
 (28.3.4)

$$\frac{\mathrm{d}p}{\mathrm{d}T} = \frac{s_2 - s_1}{v_2 - v_1} \tag{28.3.5}$$

Def latent per particle

$$\ell = T\Delta s \tag{28.3.6}$$

thus

$$\frac{\mathrm{d}p}{\mathrm{d}T} = \frac{\ell}{T(v_2 - v_1)} \tag{28.3.7}$$

or

$$\frac{\mathrm{d}p}{\mathrm{d}T} = \frac{L}{T(V_2 - V_1)} \tag{28.3.8}$$