

Notes of JU Guoxing TD&SP

hebrewsnabla

2018年5月17日

19 Equipartition of Energy

19.1 Equipartition Theorem

20 Partition Function

Energy Fluctuation:

$$\langle (E - \langle E \rangle)^2 \rangle$$

Relative fluctuation:

$$\frac{\langle (E - \langle E \rangle)^2 \rangle}{\langle E \rangle^2} \propto \frac{1}{N} \quad (20.0.1)$$

20.1

20.2 Obtain functions of State

20.2.1 Internal E

$$U = \frac{\sum_i E_i e^{-\beta E_i}}{\sum_i e^{-\beta E_i}} = \frac{\sum_i E_i e^{-\beta E_i}}{Z} \quad (20.2.1)$$

while

$$-\frac{dZ}{d\beta} = \sum_i E_i e^{-\beta E_i} \quad (20.2.2)$$

thus

$$U = -\frac{1}{Z} \frac{dZ}{d\beta} = -\frac{d \ln Z}{d\beta} \quad (20.2.3)$$

20.2.2 Entropy

E.g. 20.3

a) 2-level system. Energy level $-\Delta/2, \Delta/2$.

$$Z = \quad (20.2.4)$$

Find U, F, S

Discussion:

Def: characteristic temperature (特征温度) $k_B T_{ch} = E$

1) High temperature limit: $\beta\Delta = \frac{\Delta}{k_B T} \ll 1$

i.e. $T \gg T_{ch}$

2) Low temperature limit: $\beta\Delta = \frac{\Delta}{k_B T} \gg 1$

$$U = -\frac{\Delta}{2} \quad (20.2.5)$$

ground state occupied.

特征温度附近，热容量有极大值，称为Schottky反常。

b) simple harmonic oscillator

$$Z = \quad (20.2.6)$$

Discussion:

Def: Einstein characteristic Temp $k_B \theta_E = \hbar\omega$

1) $T \gg \theta_E$

20.3

20.4 Combining Partition Functions

21

21.3 very good

22 Chemical Potential

22.1 Definition

$$dU = TdS - pdV + \mu dN \quad (22.1.1)$$

$$\mu = \left(\frac{\partial U}{\partial N} \right)_{S,V} \quad (22.1.2)$$

22.2 Meaning of CP

$$dS = \left(\frac{\partial S}{\partial U} \right)_{N,V} dU + \left(\frac{\partial S}{\partial V} \right)_{N,U} dV + \left(\frac{\partial S}{\partial N} \right)_{U,V} dN \quad (22.2.1)$$

22.3 Grand Partition Function (巨配分函数)

system ϵ, N, V

reservoir $U \gg \epsilon, \mathcal{N} \gg N$

thus, entropy of reservoir

$$S(U - \epsilon, \mathcal{N} - N) = S(U, \mathcal{N}) - \frac{1}{T}\epsilon + \frac{\mu}{T}N \quad (22.3.1)$$

$$\Omega = e^{S/k_B} = e^{S(U, \mathcal{N})} \quad (22.3.2)$$

...

grand partition function

$$\Xi = \sum_{N=0}^{\infty} \sum_s e^{\beta(\mu N - E_s)} \quad (22.3.3)$$

or

$$\Xi = \sum_{N=0}^{\infty} \sum_s e^{\alpha N - \beta E_s} \quad (\alpha = \beta\mu) \quad (22.3.4)$$

$$\begin{aligned} \langle N \rangle &= \frac{1}{\Xi} \sum_N \sum_s N e^{\alpha N - \beta E_s} \\ &= \frac{1}{\Xi} \sum_N \sum_s \left(-\frac{\partial}{\partial \alpha} \right) e^{\alpha N - \beta E_s} = -\frac{\partial}{\partial \alpha} \ln \Xi \end{aligned} \quad (22.3.5)$$

$$U = \langle E \rangle = \frac{1}{\Xi} \sum_N \sum_s E_s e^{\alpha N - \beta E_s} \quad (22.3.6)$$

$$= \frac{1}{\Xi} \sum_N \sum_s \left(-\frac{\partial}{\partial \beta} \right) e^{\alpha N - \beta E_s} = -\frac{\partial}{\partial \beta} \ln \Xi \quad (22.3.7)$$

$$X =$$

22.4 Grand Potential

$$\Phi_G = -k_B T \ln \Xi \quad (22.4.1)$$

Since

$$S = k_B \left(\ln \Xi - \alpha \frac{\partial}{\partial \alpha} \ln \Xi - \beta \frac{\partial}{\partial \beta} \ln \Xi \right) \quad (22.4.2)$$

$$\langle N \rangle = -\frac{\partial}{\partial \alpha} \ln \Xi \quad (22.4.3)$$

$$U = -\frac{\partial}{\partial \beta} \ln \Xi \quad (22.4.4)$$

thus

$$S = k_B(\ln \Xi + \alpha \langle N \rangle + \beta U) = k_B \ln \Xi - \frac{\mu}{T} \langle N \rangle + \frac{1}{T} U \quad (22.4.5)$$

$$\Phi_G = U - TS - \mu \langle N \rangle = F - \mu \langle N \rangle \quad (22.4.6)$$

$$d\Phi_G = -SdT - pdV - \langle N \rangle d\mu \quad (22.4.7)$$

22.5

$$S = \frac{1}{T}U + \frac{p}{T}V - \frac{\mu}{T}N \quad (22.5.1)$$

i.e.

$$U - TS + pV = \mu N \quad (22.5.2)$$

thus

$$G = \mu N \quad (22.5.3)$$

(μ - single particle Gibbs function)

$$\Phi_G = F - \mu N = F - G = -pV \quad (22.5.4)$$

Differentiate (22.5.3)

$$dG = \mu dN + N d\mu \quad (22.5.5)$$

while

$$dG = -SdT + Vdp + \mu dN \quad (22.5.6)$$

thus

$$SdT - Vdp + N d\mu = 0 \quad (22.5.7)$$

which is Gibbs-Duhem Equation.