# Notes of Advanced Physical Chemistry II

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## Introduction

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## 12 Group Theory: the Exploitation of Symmetry

### Matrices

 $det(\mathbf{A}) = 0 \implies \mathbf{A}$  is a singular matrix.

- 12.1 The Exploitation of the Symm of a Mol Can Be Used to Significantly Simplify Numerical Calculations
- 12.2 The Symm of Mols Can Be Described by a Set of Symm Elements

E	
$C_n$	Rotation by $360^{\circ}/n$
$\sigma$	
i	
$S_n$	

Table 1: Symmetry elements and operators

Identity

Rotation

$\sigma_h$	horizontal
$\sigma_v$	vertical
$\sigma_d$	diagonal (vertical and bisects the angle between $C_2$ axis)

Table 2

Reflection

Inversion

Rotation Reflection

$$\hat{S}_n = \hat{\sigma}_h \times \hat{C}_n \tag{12.1}$$

### 12.2.1 Point Groups of Interest to Chemists

$C_{nv}$	
$C_{nh}$	Rotation by $360^{\circ}/n$
$D_{nh}$	
$D_{nv}$	
$D_{nd}$	
$T_d$	

Table 3: Symmetry elements and operators

### The Symm Operators of a Mol Form a Group

A set of operators form a group if they satisfy:

- 1. closed under multiplication 乘法封闭
- 2. associative multiplication 乘法结合律
- 3. only one identity operator 单位元
- 4. everyone has only one inverse 逆元

#### 12.3.1 Point Group for Some Mols

No Symm Axis

 $C_1$  – nothing  $C_s$  –  $\sigma$ 

 $C_i - i$ 

 $C_n$ 

 $S_n$ 

 $C_{nv}$  –  $C_n$  and  $n\sigma_v$ 

 $C_{nh} - C_n$  and  $\sigma_h$ 

 $D_n - C_n$  and  $nC_2 \perp C_n$  e.g. 一点点交错的  $C_3H_6, C_2$  在 3 个角平分线处

 $D_{nd} - C_n(\text{also } S_{2n}) \text{ and } nC_2 \perp C_n \text{ and } n\sigma_d$ 

 $D_{nh}$  –  $C_n$  and  $nC_2 \perp C_n$  and  $\sigma_h$ 

 $T_d$  主轴是  $S_4$ 

 $O_h$ 

 $I_h$ 

- 12.4 Symm Operators Can Be Represented by Matrices
- 12.5 The  $C_{3v}$  Point Group Has a 2-D Irreducible Representation
- 12.6 The Most Important Summary of the Properties of a Point Group Is Its Character Table

basis

class same characters - in a class.
# of class = # of irred represtn.

#### notations

- 1. A:, B:, E:2D, T:3D
- 2.  $A_1$ : symm wrt  $C_2/\sigma_v$ ,  $A_2$ : antisymm wrt that.
- 3. A': symm wrt  $\sigma_h$ , A'': antisymm wrt that.
- 4.  $A_g$ :,  $A_u$ :

# 12.7 Several Mathematical Relations Involve the Characters of Irreducible Representation

notations

XU G.X.	McQuarrie	
$D^{(\nu)}(R)$		
$\chi^{(\nu)}(R)$	$\chi_j(R)$	
$n_{ u}$	$d_{j}$	dimension of repr matrix
$a_{ u}$	$a_{j}$	
$\underline{}$	h	

Table 4

order

$$\sum_{\nu} n_{\nu}^2 = g \tag{12.2}$$

character

$$\sum_{R} D_{il}^{(\nu)} D_{jm}^{*(\mu)} = \frac{g}{n_{\nu}} \delta_{\mu\nu} \delta_{ij} \delta_{lm}$$
 (12.3)

$$\sum_{R} \chi^{(\nu)}(R) \chi^{*(\mu)}(R) = g \delta_{\mu\nu}$$
 (12.4)

$$\sum_{R} \chi^{(\nu)}(R) = 0 \quad (\nu \neq A_1)$$
(12.5)

reduce a given reducible repr  $\Gamma$   $\operatorname{Suppose}$ 

$$\chi(R) = \sum_{\nu} a_{\nu} \chi^{(\nu)}(R)$$
 (12.6)

thus

$$a_{\nu} = \frac{1}{g} \sum_{R} \chi(R) \chi^{(\nu)}(R)$$
 (12.7)

- 12.8 Use Symm Arguments to Predict Which Elements in a Secular Det Equals 0
- 12.9 Generating Operators Are Used to Find LCAOs That Are Bases for IrRepr

$$\widehat{\mathbf{P}}_{j} = \frac{d_{j}}{h} \sum_{\widehat{\mathbf{R}}} \chi_{j}(\widehat{\mathbf{R}}) \, \widehat{\mathbf{R}}$$
(12.8)

## 13 Molecular Spectroscopy

13.1

	micro	far IR	IR	visible & UV
$f/\mathrm{Hz}$				
$\lambda/\mathrm{m}$				
$\bar{\nu}/\mathrm{cm}^{-1}$				
$\lambda/\mathrm{m}$ $\bar{ u}/\mathrm{cm}^{-1}$ $E/\mathrm{J}\mathrm{mol}^{-1}$				
process				

Table 5

- 13.2 Rotational Transitions Accompany Vibrational Transitions
- 13.3
- 13.4
- 13.5 Overtones Are Observed in Vibrational Spectra

$$G(v) = \tilde{\nu}_e \left( v + \frac{1}{2} \right) - \tilde{x}_e \tilde{\nu}_e \left( v + \frac{1}{2} \right)^2$$
(13.1)

 $\tilde{x}_e$ : anharmonicity cons.

13.6 Electronic Spectra Contains Electronic, Vibrational and Rotational Info

$$\tilde{v}_{obs} = \tilde{T}_e + \dots \tag{13.2}$$

### 13.7 Franck-Condon Principle Predicts the Relative Intensities of Vibronic Transitions

# 13.8 The Rotational Spectrum of a Polyatomic Mols Depends Upon the Principal Moments of Inertia of the Mol

$$\begin{pmatrix}
I_{xx} & I_{xy} & I_{xz} \\
I_{xy} & I_{yy} & I_{yz} \\
I_{xz} & I_{yz} & I_{zz}
\end{pmatrix} \xrightarrow{\text{diagnalization}} \begin{pmatrix}
I_{A} & & \\ & I_{B} & \\ & & I_{C}
\end{pmatrix}$$
(13.3)

	top	requisition
$I_C = I_B > I_A = 0$		
$I_C = I_B = I_A$	sph top	$2C_n, n \geq 3$
$I_C = I_B > I_A$	prolate symm top	
$I_C > I_B = I_A$	oblate symm top	
$I_C \neq I_B \neq I_A$	asymm	

Table 6

#### 13.9 The Vibrations of Polyatomic Mols Are Represented by Normal Coordinates

# 13.10 Normal Coordinates Belong to Irreducible Representations of Mol Point Groups Contribution to $\chi(R)$ per unmoved atom

$\hat{\mathbf{R}}$	contribution per unmoved atom

Table 7

Now we get  $\Gamma_{3N}$ .

Subtract the irreducible representations corresponding to translational (x, y, z) and rotational  $(R_x, R_y, R_z)$  degrees of freedom, we get  $\Gamma_{vib}$ .

#### 13.11 Selection Rules Are Derived from TD Perturbation Theory

Consider a mol interacting w/ EM radiation. The EM field

$$\mathbf{E} = \mathbf{E}_0 \cos 2\pi \nu t \tag{13.4}$$

$$\hat{\mathbf{H}}^{(1)} = -\boldsymbol{\mu} \cdot \mathbf{E} = -\boldsymbol{\mu} \mathbf{E}_0 \left( e^{i \, 2\pi\nu t} + e^{-i \, 2\pi\nu t} \right) / 2 \tag{13.5}$$

$$\Psi(t) = a_1(t)\Psi_1(t) + a_2(t)\Psi_2(t)$$
(13.6)

$$a_1(t) \widehat{H}^{(1)} \Psi_1 + a_2(t) \widehat{H}^{(1)} \Psi_2 = i \hbar \left( \Psi_1 \frac{da_1}{dt} + \Psi_2 \frac{da_2}{dt} \right)$$
 (13.7)

$$a_1(t) \left\langle \psi_2 \left| \widehat{\mathbf{H}}^{(1)} \right| \Psi_1 \right\rangle + a_2(t) \left\langle \psi_2 \left| \widehat{\mathbf{H}}^{(1)} \right| \Psi_2 \right\rangle = \mathrm{i} \, \hbar \left( 0 + \frac{\mathrm{d} a_2}{\mathrm{d} t} \, \mathrm{e}^{-\mathrm{i} \, E t / \hbar} \right) \tag{13.8}$$

$$i\hbar \frac{\mathrm{d}a_2}{\mathrm{d}t} = e^{-i(E_1 - E_2)t/\hbar} \left\langle \psi_2 \left| \widehat{\mathbf{H}}^{(1)} \right| \psi_1 \right\rangle$$
 (13.9)

$$\frac{\mathrm{d}a_2}{\mathrm{d}t} \approx \dots \tag{13.10}$$

13.12 The Selection Rule in the Rigid-Rotator Approx Is  $\Delta J=\pm 1$ 

$$\langle J', M' | \mu_z | J, M \rangle = \int_0^{2\pi} d\phi \int_0^{\pi} Y_{J'}^{M'*} \mu_z Y_J^M \sin\theta d\theta$$
  
= ... (13.11)

- 13.13 The Harmonic-Oscillator Selection Rule Is  $\Delta \nu = \pm 1$
- 14 Nuclear Magnetic Resonance Spectroscopy
- 14.1 Nuclei Have Intrinsic Spin Angular Momenta
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- 14.6 Spin-Spin Coupling

$$\hat{\mathbf{H}} = -\gamma B_0 (1 - \sigma_1) \hat{\mathbf{I}}_{z1} - \gamma B_0 (1 - \sigma_2) \hat{\mathbf{I}}_{z2} + 2\pi J_{12} \hat{\mathbf{I}}_1 \cdot \hat{\mathbf{I}}_2$$
(14.1)

$$\widehat{\mathbf{H}}^{(0)} = -\gamma B_0 (1 - \sigma_1) \hat{\mathbf{I}}_{z1} - \gamma B_0 (1 - \sigma_2) \hat{\mathbf{I}}_{z2} \quad \widehat{\mathbf{H}}^{(1)} 2\pi J_{12} \hat{\mathbf{I}}_1 \cdot \hat{\mathbf{I}}_2$$
 (14.2)

$$\psi_1^{(0)} = \alpha(1)\alpha(2) \qquad \qquad \psi_2^{(0)} = \beta(1)\alpha(2) \tag{14.3}$$

$$\psi_3^{(0)} = \alpha(1)\beta(2) \qquad \qquad \psi_4^{(0)} = \beta(1)\beta(2) \tag{14.4}$$

$$E_1^{(0)} = -\gamma B_0 \left( 1 - \frac{\sigma_1 + \sigma_2}{2} \right) \qquad E_2^{(0)} = -\gamma B_0 (\sigma_1 - \sigma_2)$$
 (14.5)

$$E_3^{(0)} = \gamma B_0(\sigma_1 - \sigma_2) \qquad \qquad E_4^{(0)} = \gamma B_0 \left( 1 - \frac{\sigma_1 + \sigma_2}{2} \right) \tag{14.6}$$

perturbed to 1st order

$$E_1 = -\gamma B_0 \left( 1 - \frac{\sigma_1 + \sigma_2}{2} \right) + \frac{2\pi J_{12}}{4} \tag{14.7}$$

$$E_2 = -\gamma B_0(\sigma_1 - \sigma_2) - \frac{2\pi J_{12}}{4} \tag{14.8}$$

$$E_3 = \gamma B_0(\sigma_1 - \sigma_2) - \frac{2\pi J_{12}}{4} \tag{14.9}$$

$$E_4 = \gamma B_0 \left( 1 - \frac{\sigma_1 + \sigma_2}{2} \right) + \frac{2\pi J_{12}}{4} \tag{14.10}$$

Since

$$\nu_0 = \frac{\gamma B_0}{2\pi} \tag{14.11}$$

$$\nu_{1 \to 2} = \nu_0 (1 - \sigma_1) - \frac{J_{12}}{2} \tag{14.12}$$

$$\nu_{1 \to 3} = \nu_0 (1 - \sigma_2) - \frac{J_{12}}{2} \tag{14.13}$$

$$\nu_{2 \to 4} = \nu_0 (1 - \sigma_2) + \frac{J_{12}}{2} \tag{14.14}$$

$$\nu_{3\to 4} = \nu_0(1-\sigma_1) + \frac{J_{12}}{2} \tag{14.15}$$

- 14.7 Spin-Spin Coupling Between Chemically Equivalent Protons
- 14.8 The n+1 Rule
- 14.9 2nd-Order Spectra

Only for the case in which

$$J << \nu_0 |\sigma_1 - \sigma_2| \tag{14.16}$$

the n+1 spectra.

## 15 Lasers, Laser Spectroscopy and Photochemistry

15.1

15.2 The Dynamics of Spectroscopic Transitions between the Electronic States absorption

$$-\frac{\mathrm{d}N_1(t)}{\mathrm{d}t} = B_{12}\rho_{\nu}(\nu_{12})N_1(t) \tag{15.1}$$

emission

- 15.3 Population Inversion, 2-Level System
- 15.4 Population Inversion, 3-Level System
- 15.5 What is Inside a Laser?
- 15.6 He-Ne Laser
- 15.7 High-Resolution Laser Spectroscopy

Hyperfine structure

## 15.8 The Dynamics of Photochemistry Process

def: quantum yield

$$\Phi = \frac{\text{\# mols undergoing reaction}}{\text{\# photons absorbed}}$$
 (15.2)