# Advanced Physical Chemistry II

# HW

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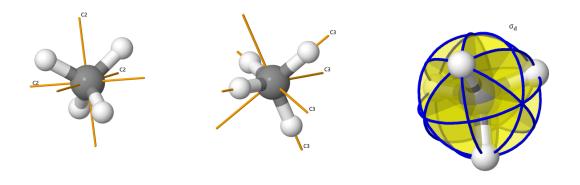
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#### Group Theory: the Exploitation of Symmetry 12

4, 12, 18, 24, 26, 31, 33, 34, 37, 40

The symmetry elements of CH<sub>4</sub> is as follows. (each  $S_4$  axis is coincident with  $C_2$  axis)<sup>1</sup>



12-12

$$C_3 \sigma_v = \sigma_v' \tag{12.1}$$

$$C_3 \sigma_v = \sigma_v'$$

$$C_3 \sigma_v'' = \sigma_v$$

$$(12.1)$$

$$(12.2)$$

12-18 For the irreducible representation of  $C_{3v}$ ,

$$\hat{E} \begin{pmatrix} u_x \\ u_y \end{pmatrix} = \begin{pmatrix} u_x \\ u_y \end{pmatrix} \tag{12.3}$$

Set the z-axis as  $C_3$  axis, we have

$$\hat{C}_{3} \begin{pmatrix} u_{x} \\ u_{y} \end{pmatrix} = \begin{pmatrix} \cos \frac{-\pi}{3} & \sin \frac{-\pi}{3} \\ -\sin \frac{-\pi}{3} & \cos \frac{-\pi}{3} \end{pmatrix} \begin{pmatrix} u_{x} \\ u_{y} \end{pmatrix} = \begin{pmatrix} -\frac{1}{2} & -\frac{\sqrt{3}}{2} \\ \frac{\sqrt{3}}{2} & -\frac{1}{2} \end{pmatrix} \begin{pmatrix} u_{x} \\ u_{y} \end{pmatrix}$$
(12.4)

$$\hat{C}_{3}^{2} \begin{pmatrix} u_{x} \\ u_{y} \end{pmatrix} = \begin{pmatrix} \cos \frac{-2\pi}{3} & \sin \frac{-2\pi}{3} \\ -\sin \frac{-2\pi}{3} & \cos \frac{-2\pi}{3} \end{pmatrix} \begin{pmatrix} u_{x} \\ u_{y} \end{pmatrix} = \begin{pmatrix} -\frac{1}{2} & \frac{\sqrt{3}}{2} \\ -\frac{\sqrt{3}}{2} & -\frac{1}{2} \end{pmatrix} \begin{pmatrix} u_{x} \\ u_{y} \end{pmatrix}$$
(12.5)

Set the x-axis on  $\sigma_v$ , we have

$$\begin{cases} \hat{\sigma}_v u_x &= u_x \\ \hat{\sigma}_v u_y &= -u_y \end{cases} \Rightarrow \hat{\sigma}_v \begin{pmatrix} u_x \\ u_y \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} u_x \\ u_y \end{pmatrix}$$
 (12.6)

Similarly, we get

$$\hat{\sigma}_v' \begin{pmatrix} u_x \\ u_y \end{pmatrix} = \begin{pmatrix} -\frac{1}{2} & \frac{\sqrt{3}}{2} \\ \frac{\sqrt{3}}{2} & \frac{1}{2} \end{pmatrix} \begin{pmatrix} u_x \\ u_y \end{pmatrix} \tag{12.7}$$

$$\hat{\sigma}_v'' \begin{pmatrix} u_x \\ u_y \end{pmatrix} = \begin{pmatrix} -\frac{1}{2} & -\frac{\sqrt{3}}{2} \\ -\frac{\sqrt{3}}{2} & \frac{1}{2} \end{pmatrix} \begin{pmatrix} u_x \\ u_y \end{pmatrix}$$
(12.8)

Thus,  $(u_x, u_y)$  forms a joint basis for the irreducible representation of  $C_{3v}$ .

<sup>&</sup>lt;sup>1</sup>The pictures are taken from chemtube3d.

12-24 No, because  $2p_{zN}$  and  $1s_{H_A} + 1s_{H_B} + 1s_{H_C}$  both belongs to  $A_1$ .

12-26

•

$$\frac{\sqrt{12}}{6}(2\phi_4 - \phi_3) = \frac{\sqrt{12}}{6} \left[ 2\frac{1}{\sqrt{12}}(\psi_1 + 2\psi_2 + \psi_3 - \psi_4 - 2\psi_5 - \psi_6) - \frac{1}{\sqrt{12}}(2\psi_1 + \psi_2 - \psi_3 - 2\psi_4 - \psi_5 + \psi_6) \right] 
= \frac{1}{6}(0 + 3\psi_2 + 3\psi_3 + 0 - 3\psi_5 - 3\psi_6) 
= \frac{1}{2}(\psi_2 + \psi_3 - \psi_5 - \psi_6) 
= \phi'_4$$
(12.9)

•

$$\langle \phi_4' | \phi_4' \rangle = \frac{1}{3} \langle 2\phi_4 - \phi_3 | 2\phi_4 - \phi_3 \rangle$$

$$= \frac{1}{3} (4 + 1 - 4 \langle \phi_4 | \phi_3 \rangle)$$
(12.10)

while

$$\langle \phi_4 | \phi_3 \rangle = \frac{1}{12} \langle \psi_1 + 2\psi_2 + \psi_3 - \psi_4 - 2\psi_5 - \psi_6 | 2\psi_1 + \psi_2 - \psi_3 - 2\psi_4 - \psi_5 + \psi_6 \rangle$$

$$= \frac{1}{12} (2 + 2 - 1 + 2 + 2 - 1)$$

$$= \frac{1}{2}$$
(12.11)

thus

$$\langle \phi_4' \, | \, \phi_4' \rangle = \frac{1}{3} \left( 5 - 4 \times \frac{1}{2} \right) = 1 \tag{12.12}$$

 $\therefore \phi_4'$  is normalized.

•

$$S_{33}' = S_{44}' = 1 (12.13)$$

$$S_{34}' = S_{43}' = \frac{\sqrt{12}}{6} \langle \phi_3 | 2\phi_4 - \phi_3 \rangle = 0$$
 (12.14)

$$H'_{33} = \left\langle \phi_3 \middle| \widehat{\mathcal{H}} \middle| \phi_3 \right\rangle = \frac{1}{12} (12\alpha + 12\beta) = \alpha + \beta \tag{12.15}$$

$$H'_{44} = \left\langle \phi_4 \middle| \widehat{\mathcal{H}} \middle| \phi_4 \right\rangle = \frac{1}{4} (4\alpha + 4\beta) = \alpha + \beta \tag{12.16}$$

$$H'_{34} = H'_{43} = \left\langle \phi_3 \middle| \widehat{H} \middle| \phi_4 \right\rangle = \frac{1}{2\sqrt{12}}(0+0) = 0$$
 (12.17)

thus

$$\begin{vmatrix} \alpha + \beta - E & 0 \\ 0 & \alpha + \beta - E \end{vmatrix} = 0 \Rightarrow E = \alpha + \beta$$
 (12.18)

12-31

• Since  $H_{ii} = \alpha$ ,  $H_{12} = H_{23} = \beta$ ,  $S_{ij} = \delta_{ij}$ , we have

$$\begin{vmatrix} \alpha - E & \beta & 0 \\ \beta & \alpha - E & \beta \\ 0 & \beta & \alpha - E \end{vmatrix} = \beta^3 \begin{vmatrix} x & 1 & 0 \\ 1 & x & 1 \\ 0 & 1 & x \end{vmatrix} = 0$$
 (12.19)

thus

$$x^3 - 2x = 0 \Rightarrow x = 0, \pm \sqrt{2} \tag{12.20}$$

•  $C_{2v}$  group has 4 operators,  $\hat{E}, \hat{C}_2, \hat{\sigma}_v, \hat{\sigma}'_v$ .

$$\hat{E} \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} \tag{12.21}$$

$$\hat{C}_2 \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} = \begin{pmatrix} 0 & 0 & -1 \\ 0 & -1 & 0 \\ -1 & 0 & 0 \end{pmatrix} \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix}$$
(12.22)

$$\hat{\sigma}_v \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 1 & 0 \\ 1 & 0 & 0 \end{pmatrix} \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} \tag{12.23}$$

$$\hat{\sigma}_v'\begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} = \begin{pmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix} \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} \tag{12.24}$$

Collect the traces of the matrices above, we get

$$\Gamma = 3 \quad -1 \quad 1 \quad -3 \tag{12.25}$$

•

$$a_{A_1} = \frac{1}{4}[3 \times 1 + (-1) \times 1 + 1 \times 1 + (-3) \times 1] = 0$$
 (12.26)

$$a_{A_2} = \frac{1}{4}[3 \times 1 + (-1) \times 1 + 1 \times (-1) + (-3) \times (-1)] = 1$$
 (12.27)

$$a_{B_1} = \frac{1}{4}[3 \times 1 + (-1) \times (-1) + 1 \times 1 + (-3) \times (-1)] = 2$$
 (12.28)

$$a_{B_2} = \frac{1}{4} [3 \times 1 + (-1) \times (-1) + 1 \times (-1) + (-3) \times 1] = 0$$
 (12.29)

thus

$$\Gamma = A_2 + 2B_2 \tag{12.30}$$

which means the secular determinant can be written as the combination of a 1D block and a 2D block.

• Now we generate the symmetric orbitals

$$\widehat{\mathbf{P}}_{A_2} \,\psi_1 = \frac{1}{4} (\psi_1 - \psi_3 - \psi_3 + \psi_1) = \frac{1}{2} (\psi_1 - \psi_3) \tag{12.31}$$

$$\widehat{\mathbf{P}}_{B_1} \, \psi_1 = \frac{2}{4} (\psi_1 + \psi_3 + \psi_3 + \psi_1) = \psi_1 + \psi_3 \tag{12.32}$$

$$\widehat{\mathbf{P}}_{B_1} \,\psi_2 = \frac{2}{4} (\psi_2 + \psi_2 + \psi_2 + \psi_2) = 2\psi_2 \tag{12.33}$$

(12.34)

After normalization,

$$\phi_1 = \frac{1}{\sqrt{2}}(\psi_1 - \psi_3) \tag{12.35}$$

$$\phi_2 = \psi_2 \tag{12.36}$$

$$\phi_3 = \frac{1}{\sqrt{2}}(\psi_1 + \psi_3) \tag{12.37}$$

(12.38)

Thus,

$$\begin{vmatrix} \alpha - E & 0 & 0 \\ 0 & \alpha - E & \sqrt{2}\beta \\ 0 & \sqrt{2}\beta & \alpha - E \end{vmatrix} = \beta^3 \begin{vmatrix} x & 0 & 0 \\ 0 & x & \sqrt{2} \\ 0 & \sqrt{2} & x \end{vmatrix} = 0$$
 (12.39)

i.e.

$$x(x^2 - 2) = 0 (12.40)$$

٠.

$$E = 0, \pm \sqrt{2} \tag{12.41}$$

i.e.

$$E = \alpha, \alpha \pm \sqrt{2}\beta \tag{12.42}$$

12-33

1. 
$$i = j = A_1$$

$$\sum_{R} \Gamma_{A_1}(R)_{11} \Gamma_{A_1}(R)_{11} = 1 + 1 + 1 + 1 + 1 + 1 = \frac{6}{1}$$
(12.43)

2. 
$$i = j = A_2$$
 
$$\sum_{n=0}^{\infty} \Gamma_{A_2}(R)_{11} \Gamma_{A_2}(R)_{11} = 1 + 1 + 1 + 1 + 1 + 1 + 1 = \frac{6}{1}$$
 (12.44)

3. i = j = E

$$\sum_{R} \Gamma_{E}(R)_{11} \Gamma_{E}(R)_{11} = 1 + \frac{1}{4} + \frac{1}{4} + 1 + \frac{1}{4} + \frac{1}{4} = \frac{6}{2}$$
(12.45)

$$\sum_{R} \Gamma_{E}(R)_{22} \Gamma_{E}(R)_{22} = 1 + \frac{1}{4} + \frac{1}{4} + 1 + \frac{1}{4} + \frac{1}{4} = \frac{6}{2}$$
(12.46)

$$\sum_{R} \Gamma_{E}(R)_{12} \Gamma_{E}(R)_{12} = 0 + \frac{3}{4} + \frac{3}{4} + 0 + \frac{3}{4} + \frac{3}{4} = \frac{6}{2}$$
(12.47)

$$\sum_{R} \Gamma_{E}(R)_{11} \Gamma_{E}(R)_{12} = 0 + \frac{\sqrt{3}}{4} - \frac{\sqrt{3}}{4} + 0 - \frac{\sqrt{3}}{4} + \frac{\sqrt{3}}{4} = 0$$
 (12.48)

$$\sum_{R} \Gamma_{E}(R)_{11} \Gamma_{E}(R)_{22} = 1 + \frac{1}{4} + \frac{1}{4} - 1 - \frac{1}{4} - \frac{1}{4} = 0$$
(12.49)

$$\sum_{R} \Gamma_E(R)_{12} \Gamma_E(R)_{22} = 0 + \frac{\sqrt{3}}{4} - \frac{\sqrt{3}}{4} + 0 + \frac{\sqrt{3}}{4} - \frac{\sqrt{3}}{4} = 0$$
 (12.50)

4.  $i = A_1, j = A_2$ 

$$\sum_{R} \Gamma_{A_1}(R)_{11} \Gamma_{A_2}(R)_{11} = 1 + 1 + 1 - 1 - 1 - 1 = 0$$
(12.51)

5.  $i = A_1, j = E$ 

$$\sum_{R} \Gamma_{A_1}(R)_{11} \Gamma_E(R)_{11} = 1 - \frac{1}{2} - \frac{1}{2} + 1 - \frac{1}{2} - \frac{1}{2} = 0$$
 (12.52)

$$\sum_{R} \Gamma_{A_1}(R)_{11} \Gamma_{A_2}(R)_{12} = 0 - \frac{\sqrt{3}}{2} + \frac{\sqrt{3}}{2} + 0 + \frac{\sqrt{3}}{2} - \frac{\sqrt{3}}{2} = 0$$
 (12.53)

$$\sum_{R} \Gamma_{A_1}(R)_{11} \Gamma_{A_2}(R)_{22} = 1 - \frac{1}{2} - \frac{1}{2} - 1 + \frac{1}{2} + \frac{1}{2} = 0$$
(12.54)

6.  $i = A_2, j = E$ 

$$\sum_{R} \Gamma_{A_1}(R)_{11} \Gamma_{A_2}(R)_{11} = 1 - \frac{1}{2} - \frac{1}{2} - 1 + \frac{1}{2} + \frac{1}{2} = 0$$
(12.55)

$$\sum_{R} \Gamma_{A_1}(R)_{11} \Gamma_{A_2}(R)_{12} = 0 - \frac{\sqrt{3}}{2} + \frac{\sqrt{3}}{2} + 0 - \frac{\sqrt{3}}{2} + \frac{\sqrt{3}}{2} = 0$$
 (12.56)

$$\sum_{R} \Gamma_{A_1}(R)_{11} \Gamma_{A_2}(R)_{22} = 1 - \frac{1}{2} - \frac{1}{2} + 1 - \frac{1}{2} - \frac{1}{2} = 0$$
 (12.57)

12-34

a.

$$\sum_{R} \Gamma_i(R)_{nn} \Gamma_i(R)_{n'n'} = \frac{h}{d_i} \delta_{nn'}$$
(12.58)

Since

$$\sum_{n} \sum_{n'} \sum_{R} \Gamma_{i}(R)_{nn} \Gamma_{i}(R)_{n'n'} = \sum_{R} \left( \sum_{n} \Gamma_{i}(R)_{nn} \right) \left( \sum_{n'} \Gamma_{i}(R)_{n'n'} \right) = \sum_{R} [\chi_{i}(R)]^{2}$$
(12.59)

and

$$\sum_{n} \sum_{n'} \frac{h}{d_i} \delta_{nn'} = \frac{h}{d_i} d_i = h \tag{12.60}$$

thus

$$\sum_{R} [\chi_i(R)]^2 = h \tag{12.61}$$

b.

$$\sum_{R} \Gamma_i(R)_{nn} \Gamma_j(R)_{n'n'} = \frac{h}{d_i} \delta_{ij} \delta_{nn'}$$
(12.62)

Since

$$\sum_{n} \sum_{n'} \sum_{R} \Gamma_i(R)_{nn} \Gamma_j(R)_{n'n'} = \sum_{R} \left( \sum_{n} \Gamma_i(R)_{nn} \right) \left( \sum_{n'} \Gamma_j(R)_{n'n'} \right) = \sum_{R} \chi_i(R) \chi_j(R)$$
(12.63)

and

$$\sum_{n} \sum_{n'} \frac{h}{d_i} \delta_{ij} \delta_{nn'} = \frac{h}{d_i} \delta_{ij} d_i = h \delta_{ij}$$
(12.64)

thus

$$\sum_{R} \chi_i(R)\chi_j(R) = h\delta_{ij}$$
(12.65)

i.e.

$$\sum_{R} \chi_i(R)\chi_j(R) = 0 \quad (i \neq j)$$
(12.66)

- c. That has been obtained in 12-34.b.
- 12-37 The point group is  $D_{3h}$  and  $\Gamma = 4$  1 -2 -4 -1 2.

$$a_{A_1'} = \frac{1}{12}(4+2-6-4-2+6) = 0 (12.67)$$

$$a_{A_2'} = \frac{1}{12}(4+2+6-4-2-6) = 0$$
 (12.68)

$$a_{E'} = \frac{1}{12}(8 - 2 + 0 - 8 + 2 + 0) = 0$$
(12.69)

$$a_{A_1''} = \frac{1}{12}(4+2-6+4+2-6) = 0 (12.70)$$

$$a_{A_2''} = \frac{1}{12}(4+2+6+4+2+6) = 2$$
 (12.71)

$$a_{E''} = \frac{1}{12}(8 - 2 + 0 + 8 - 2 + 0) = 1$$
 (12.72)

thus

$$\Gamma = 2A_2'' + E'' \tag{12.73}$$

$$\widehat{\mathbf{P}}_{A_2''} \,\psi_1 = \frac{1}{12} (\psi_1 + 2\psi_1 + 3\psi_1 + \psi_1 + 2\psi_1 + 3\psi_1) = \psi_1 \tag{12.74}$$

$$\widehat{\mathbf{P}}_{A_2''} \psi_2 = \frac{1}{12} [\psi_2 + (\psi_3 + \psi_4) + (\psi_2 + \psi_3 + \psi_4) + \psi_2 + (\psi_3 + \psi_4) + (\psi_2 + \psi_3 + \psi_4)]$$

$$= \frac{1}{3} (\psi_2 + \psi_3 + \psi_4)$$
(12.75)

$$\widehat{\mathbf{P}}_{E''}\,\psi_1 = \frac{1}{12}(2\psi_1 - 2\psi_1 + 0 + 2\psi_1 - 2\psi_1 + 0) = 0 \tag{12.76}$$

$$\widehat{\mathbf{P}}_{E''}\,\psi_2 = \frac{1}{12}[2\psi_2 - (\psi_3 + \psi_4) + 0 + 2\psi_2 - (\psi_3 + \psi_4) + 0] = \frac{1}{6}(2\psi_2 - \psi_3 - \psi_4) \tag{12.77}$$

$$\widehat{\mathbf{P}}_{E''} \,\psi_3 = \frac{1}{12} [2\psi_3 - (\psi_2 + \psi_4) + 0 + 2\psi_3 - (\psi_2 + \psi_4) + 0] = \frac{1}{6} (2\psi_3 - \psi_2 - \psi_4) \tag{12.78}$$

$$\widehat{\mathbf{P}}_{E''} \,\psi_4 = \frac{1}{12} [2\psi_4 - (\psi_3 + \psi_2) + 0 + 2\psi_4 - (\psi_3 + \psi_2) + 0] = \frac{1}{6} (2\psi_4 - \psi_2 - \psi_3) \tag{12.79}$$

Remove linear dependence and do normalization, we get

$$\phi_1 = \psi_1 \tag{12.80}$$

$$\phi_2 = \frac{1}{\sqrt{3}}(\psi_2 + \psi_3 + \psi_4) \tag{12.81}$$

$$\phi_3 = \frac{1}{\sqrt{6}} (2\psi_2 - \psi_3 - \psi_4) \tag{12.82}$$

$$\phi_4 = \frac{1}{\sqrt{6}}(2\psi_3 - \psi_2 - \psi_4) \tag{12.83}$$

thus

$$\begin{vmatrix} \alpha - E & \sqrt{3}\beta & 0 & 0 \\ \sqrt{3}\beta & \alpha - E & 0 & 0 \\ 0 & 0 & \alpha - E & -\alpha/2 - E/2 \\ 0 & 0 & -\alpha/2 - E/2 & \alpha - E \end{vmatrix} = \beta^4 \begin{vmatrix} x & \sqrt{3} & 0 & 0 \\ \sqrt{3} & x & 0 & 0 \\ 0 & 0 & x & -x/2 \\ 0 & 0 & -x/2 & x \end{vmatrix} = 0$$
 (12.84)

i.e.

$$(x^2 - 3)x^2 = 0 (12.85)$$

thus

$$x = 0, \pm \sqrt{3} \tag{12.86}$$

i.e.

$$E = \alpha, \alpha \pm \sqrt{3}\beta \tag{12.87}$$

Since there are  $4 \pi$  electrons,

$$E_{\pi} = 2(\alpha + \sqrt{3}\beta) + 2\alpha = 4\alpha + 2\sqrt{3}\beta$$
 (12.88)

12-40

$$a_{A_1'} = \frac{1}{12}(3+0+3+3+0+3) = 1$$
 (12.89)

$$a_{A_2'} = \frac{1}{12}(3+0-3+3+0-3) = 0 (12.90)$$

$$a_{E'} = \frac{1}{12}(6+0+0+6+0+0) = 1$$
 (12.91)

$$a_{A_1''} = \frac{1}{12}(3+0+3-3+0-3) = 0 (12.92)$$

$$a_{A_2''} = \frac{1}{12}(3+0-3-3+0+3) = 0 (12.93)$$

$$a_{E''} = \frac{1}{12}(6+0+0-6+0+0) = 0 \tag{12.94}$$

thus

$$\Gamma = A_1' + E' \tag{12.95}$$

From the character table of  $D_{3h}$ , we can see that  $A'_1$  is corresponding to s orbital and E' is corresponding to  $p_x, p_y$  or  $d_{x^2-y^2}, d_{xy}$  orbitals.