

Search for production of a Higgs boson and a single top quark in $\mu\mu$ final states in proton collisions at $\sqrt{s} = 13$ TeV

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Overview

- Through this project we will investigate the production of Higgs boson in association with a single top quark (tH) in proton-proton collisions with the CMS experiment of the LHC. This mechanism of production of the Higgs boson has not been observed before by any experiment.
 - Understanding the production of the Higgs boson, as well as its decays are an important part of the physical program of the CERN international laboratory experiments that try to complete the tests to verify the Standard Model, the theory of the fundamental particles



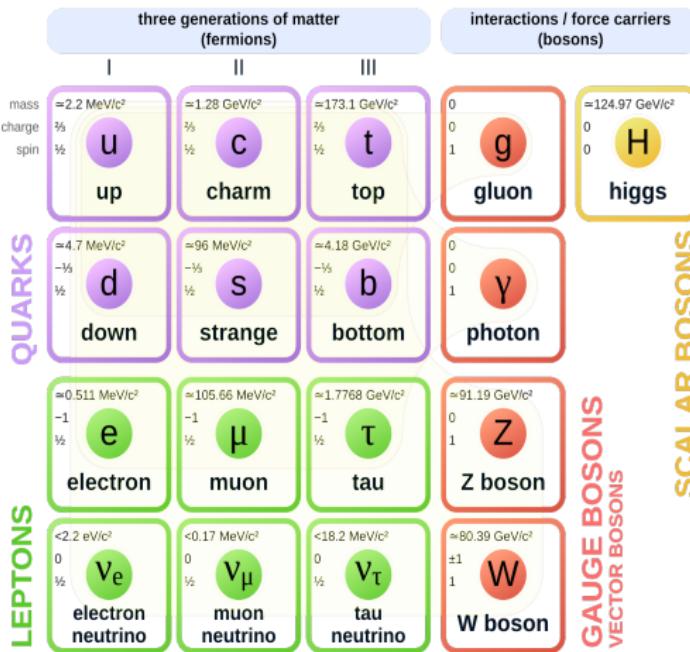
Motivation for single top Higgs (tH)

- Coupling measurement is essential to establish the nature of the Higgs
 - The exploration of Higgs production on the tH channel is subject relatively new. Measurements of CMS and ATLAS are compatible with SM predictions.
 - The tH study explores the relative sign of top-Higgs and W-Higgs couplings.
Small deviations from SM predictions could be associated with physics beyond the standard model (BSM) such as String Theory and Supersymmetry.



Standard Model

Standard Model of Elementary Particles



Electroweak SM Lagrangian

The standard model is a theory of fields with spins 0, $\frac{1}{2}$ and 1. The SM lagrangian

$$\mathcal{L} = \mathcal{L}_{\text{gauge}} + \mathcal{L}_f + \mathcal{L}_{\text{higgs}} + \mathcal{L}_{\text{yukawa}}$$

donde

- $\mathcal{L}_{\text{gauge}} = -\frac{1}{4} [F^{\mu\nu} F_{\mu\nu}] - \frac{1}{4} [G^{i\mu\nu} G^i_{\mu\nu}]$
 - $F^{\mu\nu} = \partial_\mu B_\nu - \partial_\nu B_\mu$
 - $G^{i\mu\nu} = \partial_\mu W_\nu^i - \partial_\nu W_\mu^i - g \epsilon^{ijk} W_\mu^j W_\nu^k$
 - $\mathcal{L}_f = i\bar{\Psi}_L \not{D} \Psi_L + i\bar{\psi}_R \not{D} \psi_R$ is the kinematic term for fermions
 - $D_\mu \Psi_L = (\partial_\mu + igW_\mu + ig'Y_L B_\mu) \Psi_L$
 - $D_\mu \psi_R = (\partial_\mu + ig'Y_R B_\mu) \psi_R$
 - g and g' are boson coupling constants. ¹
 - $W_\mu = \sigma^i W_\mu^i$, where σ^i are the pauli matrices.

¹L(R) refers to Left and right fields. Y is the weak hypercharge, $Y = T^3 - Q \cdot Q$ is charge of particle field and T is weak isospin. $T^3 = \pm 1/2$ for left handed doublets and $T^3 = 0$ for right handed singlets.



Electroweak SM Lagrangian

Higgs lagrangian

$$\mathcal{L}_{\text{Higgs}} = (D_\mu \Phi)^\dagger (D^\mu \Phi) - V(\Phi^\dagger \Phi) \quad (1)$$

with

- $D_\mu \Phi = \left(\partial_\mu + (ig/2)\sigma^i W_\mu^i - i\frac{1}{2}g' B_\mu \right) \Phi$
 - $V(\Phi^\dagger \Phi) = -\mu^2 \Phi^\dagger \Phi + \frac{1}{2}\lambda (\Phi^\dagger \Phi)^2, \quad \mu^2 > 0$
 - $\Phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}$ is a SU(2) doublet.
 - σ^i are the Pauli matrices.
 - Φ is higgs field which is a complex scalar field. Φ
 - λ and μ^2 are Higgs potential parameters



SM Lagrangian

Yukawa lagrangian

$$\mathcal{L}_{yukawa} = \Gamma_{mn}^u q_{m,L} \tilde{\phi} u_{n,R} + \Gamma_{mn}^d \bar{q}_{m,L} \phi d_{n,R} + \Gamma_{mn}^e \bar{l}_{m,L} \phi e_{n,R} + h.c \quad (2)$$

where h.c is hermitian conjugate.

The matrices Γ_{mn} describe the so called Yukawa couplings between higgs doublet ϕ and the fermions.

Choosing

$$\Phi = -\frac{1}{2} \begin{pmatrix} 0 \\ v+h \end{pmatrix} \rightarrow \Phi = \frac{1}{2} \begin{pmatrix} 0 \\ v \end{pmatrix}$$

By using it on the first part of the lagrangian it is obtained

$$\mathcal{L}_{yukawa} = \frac{f_u v}{\sqrt{2}} (\bar{u}_L u_R + \bar{u}_R u_L)$$

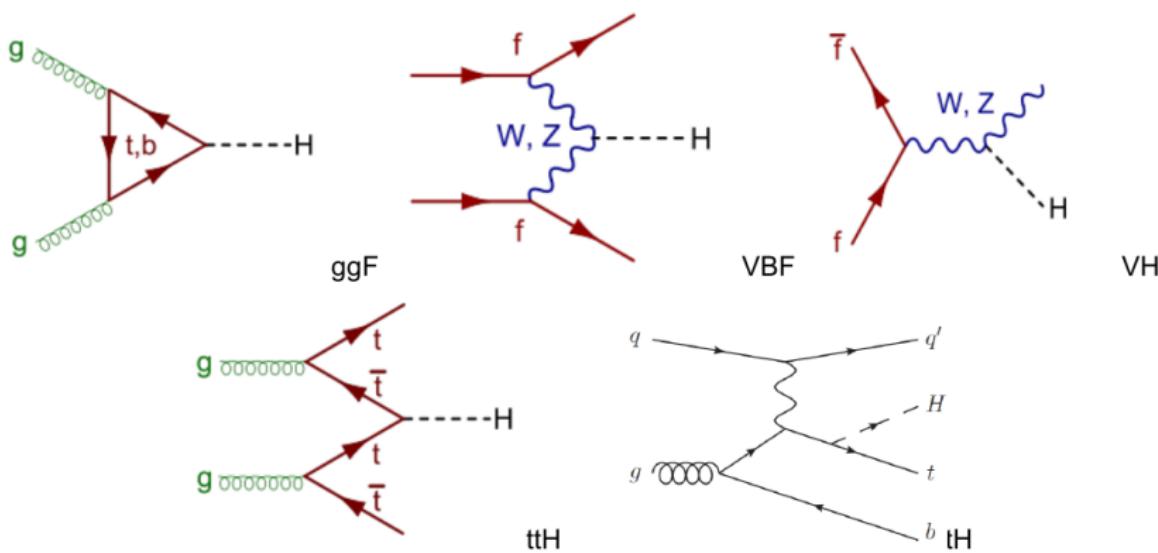
from which the masses for the fermions in question can be read off:

$$m_u = -\frac{f_u v}{\sqrt{2}}$$

which m is the mass of fermions and $f_u = \Gamma_{mn}^u$ for μ family.



Higgs production mechanisms



tH production mechanisms

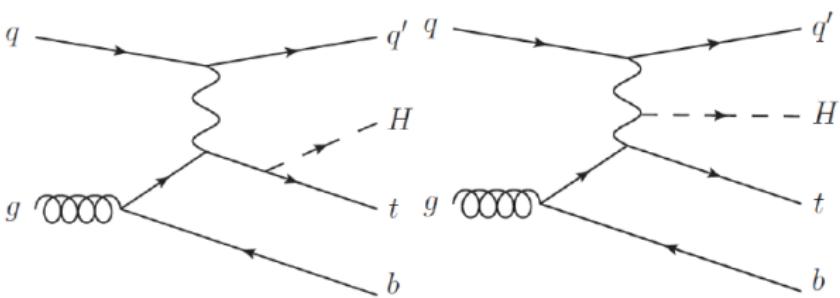


Figure: tH mechanism. Higgs radiated from a top quark (left). Higgs radiated from a W boson (right)

tH production mechanisms

- In a proton proton collision, the production cross section of the single top plus Higgs boson (tHq) process is driven by a destructive interference of two main diagrams (see Fig. 1.), where the Higgs couples to either the W boson or the top quark.
 - A second process, where the Higgs and top quark are accompanied by a W boson (tHW) has similar behavior, albeit with a weaker interference pattern.
 - However, in the presence of new physics, there may be relative opposite signs between the tH and WH couplings which lead to constructive interference and enhance the cross sections by an order of magnitude or more.

Cross section

- Classical definition: When two particles interact, cross section is the area transverse to their relative motion within which they must meet in order to scatter from each other.
 - Quantum definition: Cross section describes the likelihood of two particles interacting under certain conditions[1][8]
 - Experimentally

$$d\sigma = \frac{\text{number of particles scattered into solid angle } \Delta\Omega}{(\text{number of particles incident})(\text{scattering centers/area})} \quad (3)$$

- Cross sections are expressed in barns , where $1 \text{ barn} = 10^{-34} \text{ cm}^{-2}$

Cross section

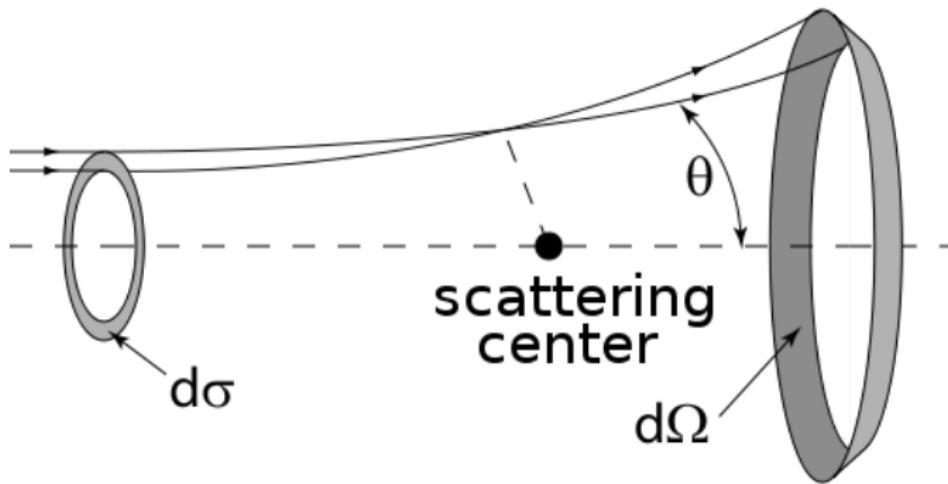


Figure: Drawing of an idealized scattering process showing the differential solid angle $\Delta\Omega$ and the scattering angle θ

Cross section

The reaction rate N_R is determined by the total cross section σ and the incident flux L .

L is called luminosity and it is measured in $\text{cm}^{-2}\text{s}^{-1}$. [8]

$$N_R = \sigma L \quad (4)$$

Higgs production Cross section

Higgs boson production cross sections in pp collisions for $\sqrt{s} = 13\text{TeV}$ (in pico barn). Integrated luminosity of 35.9 fb^{-1} for Run 2²

Production mechanism	σ (picobarns pb)	Number of events
ggF	48.93	1756587
VBF	3.78	135702
WH	1.35	48465
ZH	0.88	31592
t <bar>tH</bar>	0.50	18255
tH (only)	0.015	560.39

¹Data taken from The cern collaborarion "Higgs Physics the HL-LHC and HE-LHC" 2019, CERN-LPCC-2018-04

Higgs production cross section

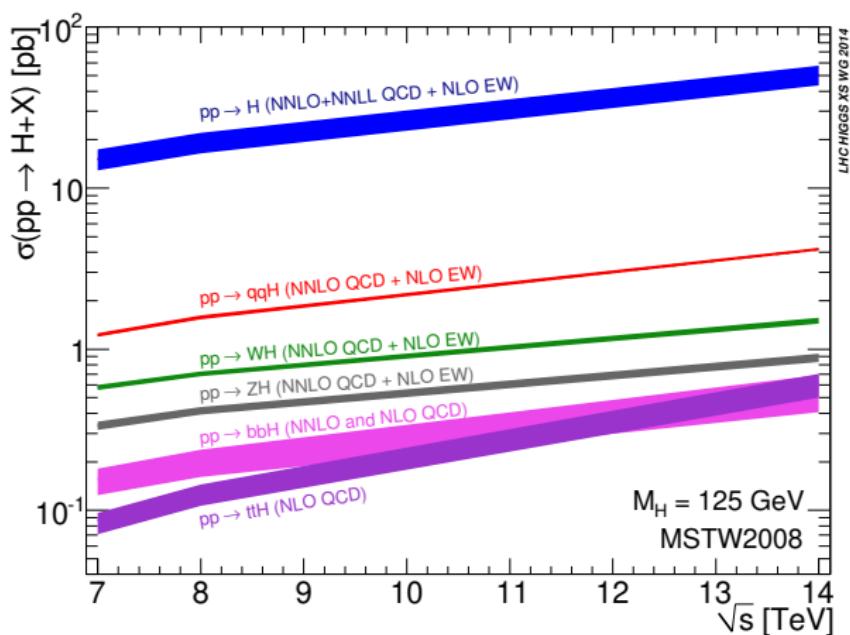


Figure: Higgs boson production cross sections as a function of the centre-of-mass-energies

Branching ratio

In particle physics, the branching ratio for a decay process is the ratio of the number of particles which decay via a specific decay mode with respect to the total number of particles which decay via all decay modes.

$$\text{BR} = \frac{\Gamma_i}{\sum_i \Gamma_i} \quad (5)$$

Where $\Gamma = \sum_i \Gamma_i$ is the total decay width (sum of all partial widths) of the particle and is related to lifetime of the particle: $\Gamma = 1/\tau$. Since the dimension of Γ is the inverse of time, in our system of natural units, it is measured in inverse seconds.⁰

⁰Cleaves H.J. (2011) Branching Ratio. In: Gargaud M. et al. (eds) Encyclopedia of Astrobiology. Springer, Berlin, Heidelberg

Higgs Branching ratio

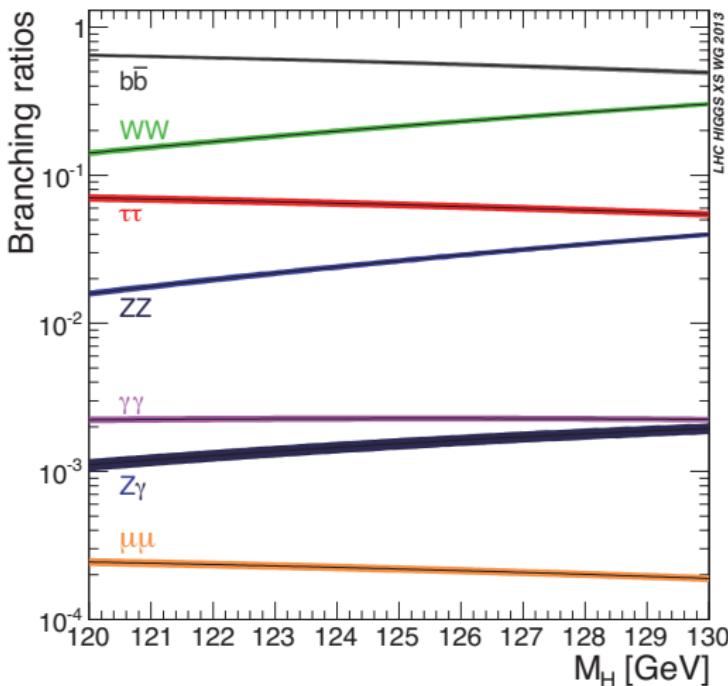


Figure: Standard Model Higgs boson decay branching ratios

Higgs Branching ratios per channel

SM Higgs boson branching ratios for $M_H = 125$ GeV

Higgs decay	Branching ratio (BR)
$H \rightarrow b\bar{b}$	50.82%
$H \rightarrow W^+W^-$	21.5%
$H \rightarrow \tau^+\tau^-$	6.27%
$H \rightarrow ZZ$	2.61%
$H \rightarrow \gamma\gamma$	0.227%
$H \rightarrow Z\gamma$	0.153%
$H \rightarrow \mu^+\mu^-$	0.0217%

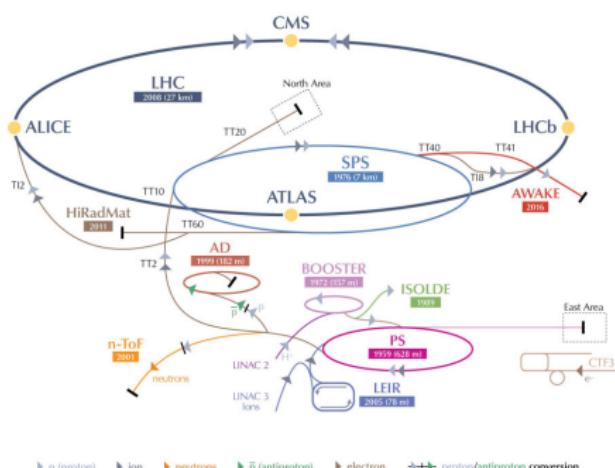
$\mu\mu$ same sign decay rate

Table of decay chains for tH. Expected number of final events assuming 560 produced tH events. l represents μ^\pm, e^\pm, τ^\pm .

Decay chain	BR	Events
$tH \rightarrow W^+ b W^+ W^- \rightarrow \mu^+ \nu_\mu b \mu^+ \nu_\mu q \bar{q}'$	2.096×10^{-3}	1.173
$tH \rightarrow W^+ b W^+ W^- \rightarrow \mu^+ \nu_\mu b \mu^+ \nu_\mu l^- \bar{\nu}_l$	3.37×10^{-4}	0.899
$tH \rightarrow W^+ b \tau^+ \tau^- \rightarrow \mu^+ \nu_\mu b \mu^+ \nu_\mu \bar{\nu}_\tau l^- \bar{\nu}_l \nu_\tau$	3.637×10^{-4}	0.203
$tH \rightarrow W^+ b W^+ W^- \rightarrow \tau^+ \bar{\nu}_\tau b \mu^+ \nu_\mu q \bar{q} \rightarrow \mu^+ \nu_\mu \bar{\nu}_\tau \bar{\nu}_\tau b \mu^+ \nu_\mu q \bar{q}$	1.890×10^{-4}	0.105
$tH \rightarrow W^+ b \tau^+ \tau^- \rightarrow \mu^+ \nu_\mu b \nu_\tau \mu^+ \nu_\mu \bar{\nu}_\tau q \bar{q}$	1.681×10^{-4}	0.094
$tH \rightarrow W^+ b W^+ W^- \rightarrow \tau^+ \bar{\nu}_\tau b \mu^+ \nu_\mu l^- \bar{\nu}_l \rightarrow \mu^+ \nu_\mu \bar{\nu}_\tau \bar{\nu}_\tau b \mu^+ \nu_\mu l^- \bar{\nu}_l$	3.045×10^{-5}	0.017
$tH \rightarrow W^+ b ZZ \rightarrow q \bar{q} b ZZ \rightarrow q \bar{q} b \mu^+ \mu^- \mu^+ \mu^-$	1.966×10^{-5}	0.011
$tH \rightarrow W^+ b \tau^+ \tau^- \rightarrow \tau^+ \bar{\nu}_\tau b \mu^+ \nu_\mu \bar{\nu}_\tau q \bar{q}' \nu_\tau \rightarrow \mu^+ \nu_\mu \bar{\nu}_\tau \bar{\nu}_\tau b \mu^+ \nu_\mu \bar{\nu}_\tau q \bar{q}' \nu_\tau$	1.549×10^{-5}	0.008

LHC

CERN's Accelerator Complex



LHC Large Hadron Collider SPS Super Proton Synchrotron PS Proton Synchrotron

AD, Antisecotor Decalocutor; **CTE2**, Clic Test Facility; **AWAKE**, Advanced Wakefield Experiment; **ESCAPE**, Electron Sausage Instability Experiment.

LEIP: Low Energy Ion Beam - UNIG: Union Accelerator - T-5: Neutral-Tone Of Effects - MIP-DMT: High Production for Materials

LHC parameters

Quantity	Number
Circumference	26 659 m
Dipole operating temperature	1.9 K (-271.3°C)
Number of magnets	9593
Number of main dipoles	1232
Number of main quadrupoles	392
Number of RF cavities	8 per beam
Nominal energy, protons	6.5 TeV
Nominal energy, ions	2.56 TeV/u (energy per nucleon)
Nominal energy, protons collisions	13 TeV
No. of bunches per proton beam	2808
No. of protons per bunch (at start)	1.2×10^{11}
Number of turns per second	11245
Number of collisions per second	1 billion

Table. LHC characteristics for run 2.

LHC

Accelerator operation energies

Accelerator	Energy
Linac 2	50 MeV
PS Booster	1.4 GeV
Proton Scyncroton (PS)	25 GeV
SPS	450 GeV
LHC	6.5 TeV

CMS detector

CMS DETECTOR

Total weight : 14,000 tonnes
Overall diameter : 15.0 m
Overall length : 28.7 m
Magnetic field : 3.8 T

STEEL RETURN YOKE
12,500 tonnes

SILICON TRACKERS
Pixel ($100 \times 150 \mu\text{m}$) $\sim 16\text{m}^2$ $\sim 66\text{M}$ channels
Microstrips ($80 \times 180 \mu\text{m}$) $\sim 200\text{m}^2$ $\sim 9.6\text{M}$ channels

SUPERCONDUCTING SOLENOID
Niobium titanium coil carrying $\sim 18,000\text{A}$

MUON CHAMBERS
Barrel: 250 Drift Tube, 480 Resistive Plate Chambers
Endcaps: 468 Cathode Strip, 432 Resistive Plate Chambers

PRESHOWER
Silicon strips $\sim 16\text{m}^2$ $\sim 137,000$ channels

FORWARD CALORIMETER
Steel + Quartz fibres $\sim 2,000$ Channels

CRYSTAL
ELECTROMAGNETIC
CALORIMETER (ECAL)
 $\sim 76,000$ scintillating PbWO₄ crystals

HADRON CALORIMETER (HCAL)
Brass + Plastic scintillators $\sim 7,000$ channels

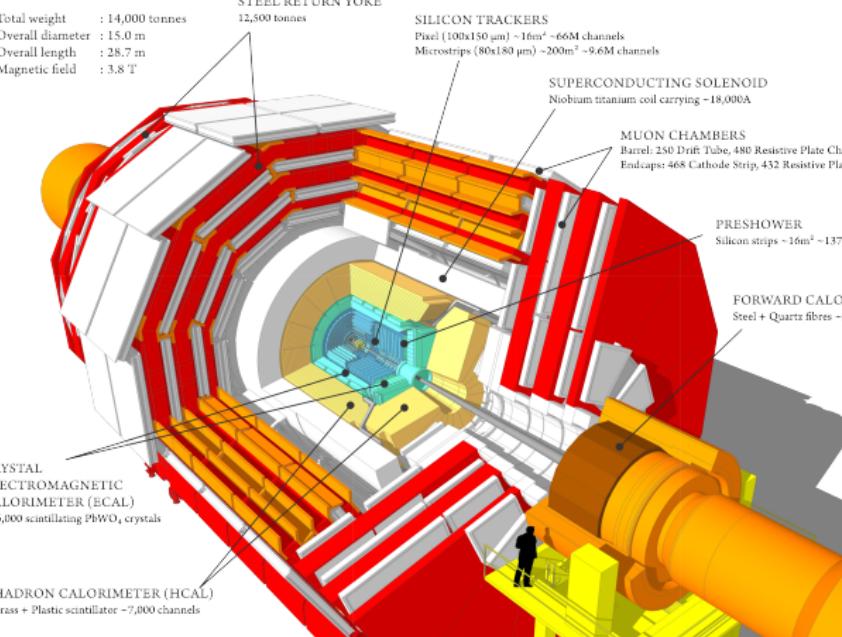


Figure: Compact muon solenoid

CMS integrated luminosity

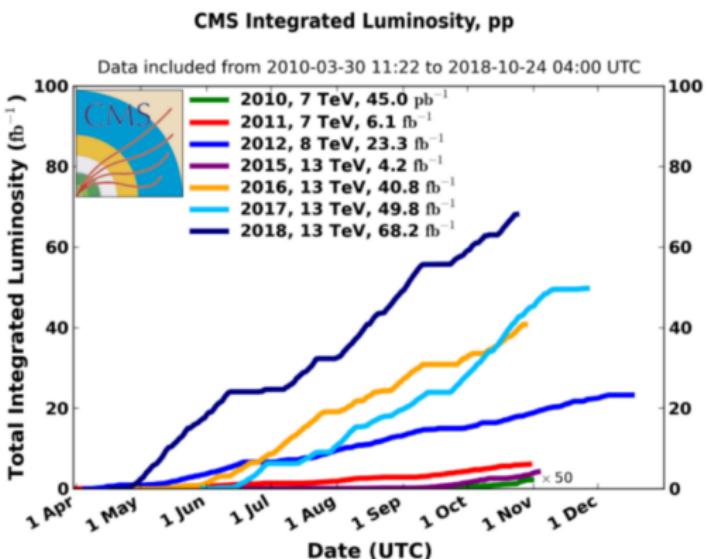
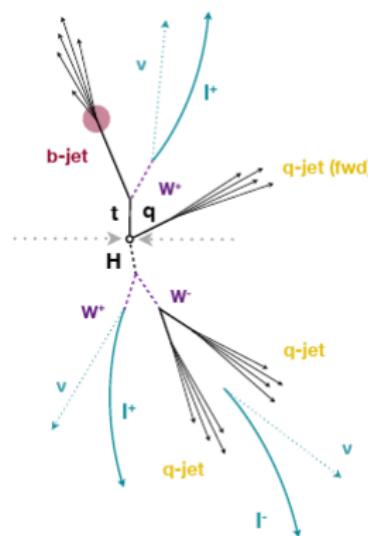


Figure: Integrated luminosity for CMS experiment

Topology of events tH

The characteristics of the signal tH:

- $t \rightarrow W^+ b \rightarrow l^+ \nu_\mu b$
- $H \rightarrow W^+ W^-$
 - $W \rightarrow l^\pm \nu$, where l can be μ^\pm, e^\pm, τ^\pm
 - $W \rightarrow q\bar{q}$
 - W bosons decay leptonically resulting in a signature of two same-sign leptons
- a light-flavor quark
- b quark jet [2]



Event selection

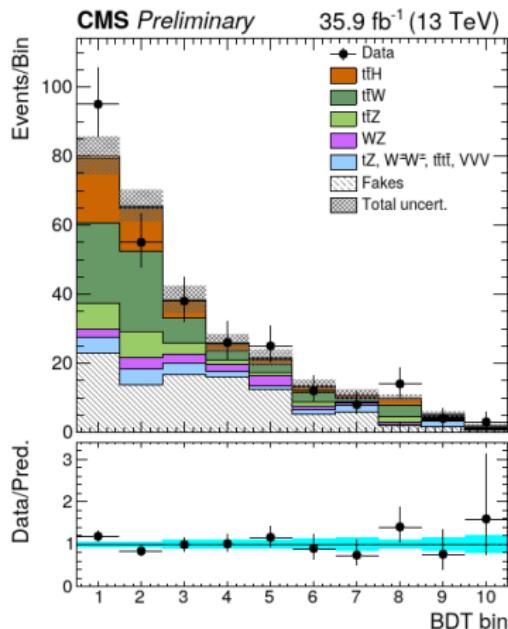
The main analysis strategy is to obtain a selection of events compatible with certain characteristics of the signal at pre-selection level.

It is required:

- The events are selected those that contain two leptons (I^+I^-) with the same sign.
 - Transverse moment $p_t > 25$ and 15 GeV, for the muons.
 - A foward jet with $p_t > 40$ GeV, $| \eta | > 2.4$
 - One or more b-jets with ($| \eta | < 2.4$)[1]

Previous results for $t\bar{t}H$ production

Process	$\mu\mu$
$t\bar{t}W^\pm$	68.03 ± 0.61
$t\bar{t}Z/t\bar{t}\gamma$	25.89 ± 1.12
WZ	15.07 ± 1.19
ZZ	1.16 ± 0.29
$W^\pm W^\pm qq$	3.96 ± 0.52
$W^\pm W^\pm(\text{DPS})$	2.48 ± 0.42
VVV	2.99 ± 0.34
$ttt\bar{t}$	2.32 ± 0.45
tZq	5.77 ± 2.24
tZW	2.13 ± 0.13
γ conversions	—
Non-prompt	80.94 ± 2.02
Charge flips	—
Total Background	210.74 ± 3.61
$t\bar{t}H$	24.18 ± 0.48
tHq (SM)	1.43 ± 0.04
tHW (SM)	0.71 ± 0.03
Total SM	237.06 ± 3.64
tHq ($\kappa_V = 1 = -\kappa_t$)	18.48 ± 0.22
tHW ($\kappa_V = 1 = -\kappa_t$)	7.72 ± 0.17
Data	280



Post-fit categorized BDT classifier outputs as used in the maximum likelihood fit for the $\mu\mu$ channel for 35.9 fb^{-1} . In the box below each distribution, the ratio of the observed and predicted event yields is shown.[2]

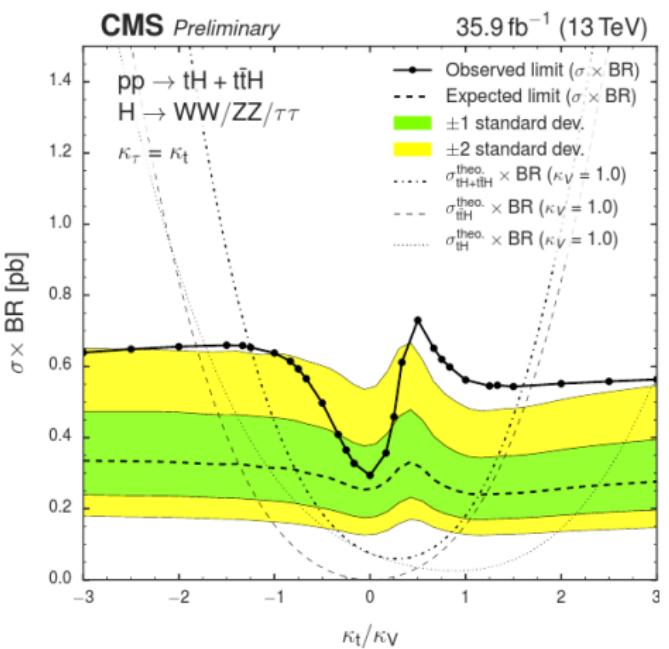


Figure: Observed and expected 95% C.L. upper limit on the $tH + t\bar{t}H$ cross section times $H \rightarrow WW + \tau\tau + ZZ$ branching fraction for different 20 values of the coupling ratio κ_t/κ_v . The expected limit is derived from a background-only MC dataset.

Main backgrounds

In the leptonic channels, the main backgrounds are expected to arise from the production of top quarks

- In the dominant $t\bar{t}$ mode, where same-sign dilepton signatures can occur when a non-prompt lepton from heavy-flavor decay passes the signal selection, or in associated production with a W/Z or Higgs boson.
 - Processes with single top quarks also contribute, mostly in the associated production with a Z boson (tZ) or when produced with both a W and a Z boson (WZ)

Main backgrounds

- $t\bar{t}W^\pm$ and $t\bar{t}Z$ ($t\bar{t}V$): Backgrounds are estimated directly from simulated events
 - WZ: Diboson production with leptonic Z decays and additional jet radiation in the final state can lead to signatures very similar to that of the signal. Due to the larger cross section, the main contribution arises from WZ production.
 - $t\bar{t}H$: tHq production which decays to same sign dileptons.
 - Rare SM ($tZ, VVV, WW, tttt$): Due to small event yields, all these processes are grouped as one.
 - Non prompt leptons b quark decays and spurious lepton signatures from hadronic jets.

Main backgrounds

Background	$\sigma[\text{pb}]$	Decay process
$t\bar{t}H$	0.2151	$t\bar{t}H \rightarrow W^+b W^+b WW \rightarrow \mu^+\nu_\mu b\mu^-\bar{\nu}_\mu \bar{b} \mu^+\mu^- \mu^+\mu^-$
$t\bar{t}W$ and	0.2043	$t\bar{t}W \rightarrow W^+b W^+b\mu^+\mu^- \rightarrow \mu\nu_\mu b\mu\nu_\mu b\mu^+\mu^-$
$t\bar{t}Z$	0.2529	$t\bar{t}Z \rightarrow W^+b W^-b\mu^+\mu^- \rightarrow \mu^+\nu_\mu b \mu^-\bar{\nu}_\mu \bar{b}$
W^+W^-	1.64	$WW \rightarrow \mu^+\mu^-\mu^+\mu^-$
tZq	0.0758	$tZq \rightarrow W^+b\mu^+\mu^- q \rightarrow \mu^+\nu_\mu b \mu^+\mu^- q$
$tttt$	0.0091	$tttt \rightarrow W^+bW^+bW^+bW^+b \rightarrow \mu^+\nu_\mu b\mu^+\nu_\mu b q\bar{q}b q\bar{q}b$
WZZ	0.1651	$WZZ \rightarrow \mu^+\mu^-\mu^+\mu^- l^+l^-$
ZZZ	0.0139	$ZZZ \rightarrow \mu^+\mu^-\mu^+\mu^- l^+l^-$
WZZ	0.0556	$WZZ \rightarrow \mu^+\mu^-\mu^+\mu^- l^+l^-$
WZ	4.42965	$WZ \rightarrow \mu^+\mu^-\mu^+\mu^-$
ZZ	1.256	$ZZ \rightarrow \mu^+\mu^-\mu^+\mu^-$

- Rare SM: tttt, WWW, WWZ, WZZ, WW, tZq
 - WZ:WZ, ZZ

Boosted decision tree (BDT)

- A decision tree takes a set of input features and splits input data recursively based on those features. Boosting is a method of combining many weak learners (trees) into a strong classifier. The features can be a mix of categorical and continuous data.
 - The BDT training is performed using several event variables for signal and background.
 - BDT training uses MC samples for $t\bar{t}W^\pm$ and $t\bar{t}Z$ ($t\bar{t}V$), because $t\bar{t}V$ is one of the biggest backgrounds

BDT Variables

- Trailing lepton p_t
 - Total charge of tight leptons
 - $\min \Delta R$ (lepton pairs)
 - $\Delta\phi$ between highest p_t lepton pair
 - Number of jets with $|\eta| < 2.4$
 - Number of non b-tagged jets with $|\eta| > 1.0$
 - Maximum $|\eta|$ for jets
 - $\Delta\eta$ (most forward light jet, closest lepton)
 - $\Delta\eta$ (most forward light jet, hardest loosely b-tagged jet)
 - $\Delta\eta$ (most forward light jet, 2nd hardest loosely b-tagged jet)

Sources of uncertainty on event yields

- Luminosity measurement: 2.6%.
 - Data/MC scale factors for lepton selection (ID, iso) and trigger efficiencies 5% per lepton.
 - Choice of PDF set:
 - 3.7% for tHq
 - 4% for tHW, t \bar{t} W, t \bar{t} Z, t \bar{t} H
 - Scale uncertainties: 12%, 4 for t \bar{t} W, 10% for t \bar{t} Z, +5.8/-9.2% for t \bar{t} H.
 - Background: WZ, ZZ sample modelling and statistics: 50%.
 - Rare SM (tZ, tri-bosons, WWqq, tt \bar{t}) : 50%
 - Fake rate estimation: The predicted event yield has a normalization uncertainty of 30-50% [7]

Uncertainty values

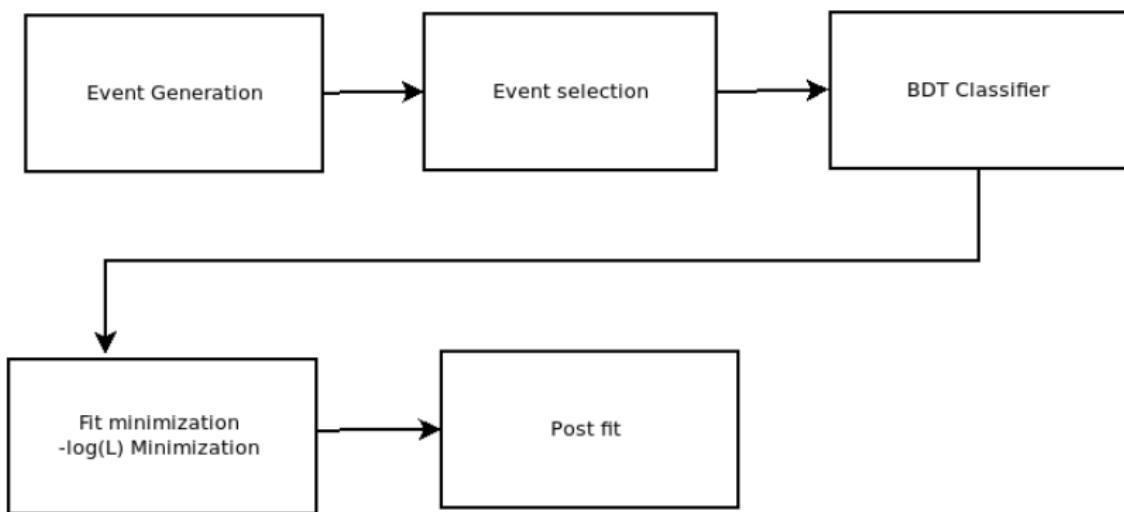
Common uncertainties for MC :

- Muon identification 6%
- Dimuon Trigger 1%

Background	Sources of systematic uncertainty	Total Systematic Uncertainty
$t\bar{t}H$	<ul style="list-style-type: none"> • PDF 4% • QCD scale 5.8% • Higgs BR 1% 	9.36%
$t\bar{t}Z$	<ul style="list-style-type: none"> • PDF 4% • QCD scale 10% 	12.36 %
$t\bar{t}W$	<ul style="list-style-type: none"> • PDF 4% • QCD scale 12% 	14.03 %

Background	Sources of systematic uncertainty	Total Systematic Uncertainty
Rares SM	<ul style="list-style-type: none"> ● Rare SM 50% 	50.36 %
WZ	<ul style="list-style-type: none"> ● Sample modelling and statistics 50% 	50.36%
Non prompt leptons	<ul style="list-style-type: none"> ● Fake rate estimation 50% ● Closure 7% 	50.48%

Fitting



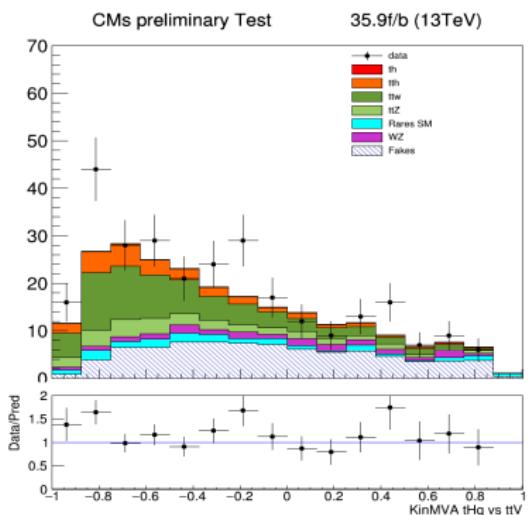
Likelihood model

The likelihood function is the product of Poisson probabilities for all bins

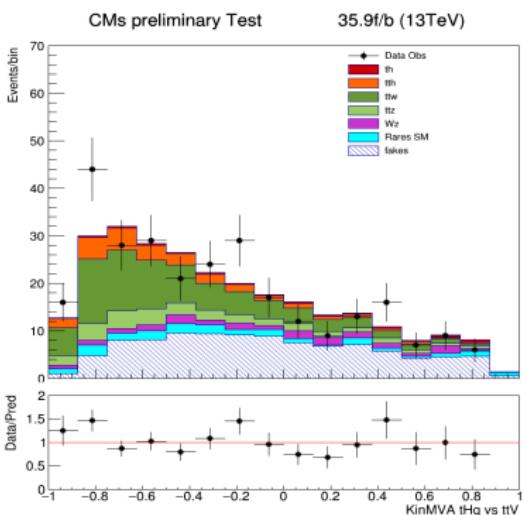
$$L(\mu, \theta) = \prod_{j=1}^N \frac{(\mu s_j + b_j)^{n_j}}{n_j!} e^{-(\mu s_j + b_j)} \quad (6)$$

- N=number of bins
- μ =parameter of signal
- s =signal
- b =background
- n =number of events

Results



Pre-fit signal and background yields for tH process. In the box below each distribution, the ratio of the observed and predicted event yields is shown



Post-fit signal and background yields for tH process. In the box below each distribution, the ratio of the observed and predicted event yields is shown

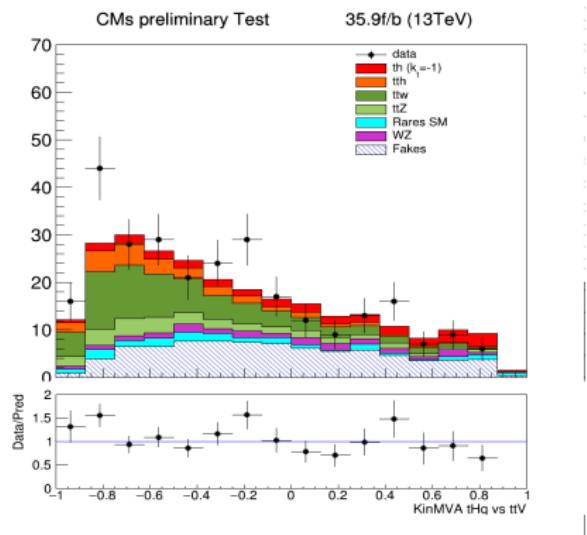
Results

Table: Prefit and postfit table for each yield.

Process	Number of events prefit	Number of events Postfit
tH	2.13 ± 0.21	5.36 ± 27.40
t <bar>t>H</bar>	24.18 ± 2.26	24.81 ± 2.24
t <bar>t>W</bar>	68.03 ± 9.54	77.01 ± 8.76
t <bar>t>Z</bar>	25.89 ± 3.20	26.74 ± 3.18
Rares SM	17.17 ± 8.64	18.77 ± 8.66
WZ	16.23 ± 8.17	17.41 ± 8.04
fakes	80.94 ± 40.86	100.15 ± 29.30

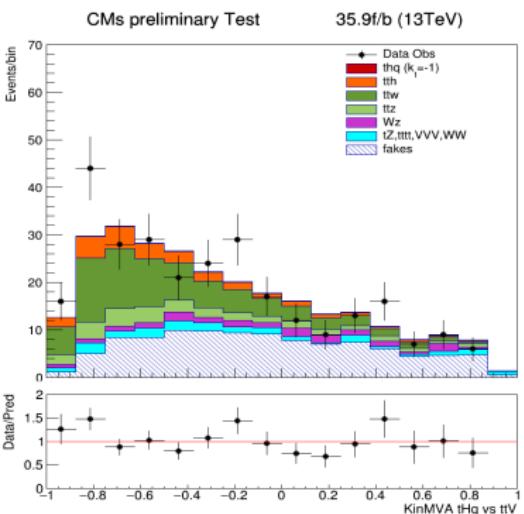
Prefit uncertainty is statistical only. Postfit uncertainty is statistical + systematic.

Results



Pre-fit signal and background yields for tH process for $k_t = -1$.

In the box below each distribution, the ratio of the observed and predicted event yields is shown



Post-fit signal and background yields for tH process for $k_t = -1$.

In the box below each distribution, the ratio of the observed and predicted event yields is shown

Results

Table: Prefit and postfit table for each yield. $k_t = -1$.

Process	Number of events prefit	Number of events Postfit
tH	26.2 ± 2.60	1.83 ± 26.63
t <bar>t>H</bar>	24.18 ± 2.26	24.82 ± 2.27
t <bar>t>W</bar>	68.03 ± 9.54	77.07 ± 8.99
t <bar>t>Z</bar>	25.89 ± 3.20	26.76 ± 3.18
Rares SM	17.17 ± 8.64	18.90 ± 8.37
WZ	16.23 ± 8.17	17.54 ± 8.15
fakes	80.94 ± 40.86	102.97 ± 29.51

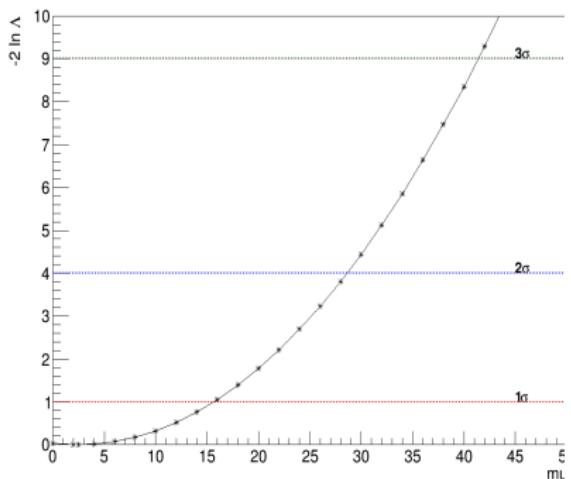
Prefit uncertainty is statistical only. Postfit uncertainty is statistical + systematic.

Test statistic q_μ for upper limits

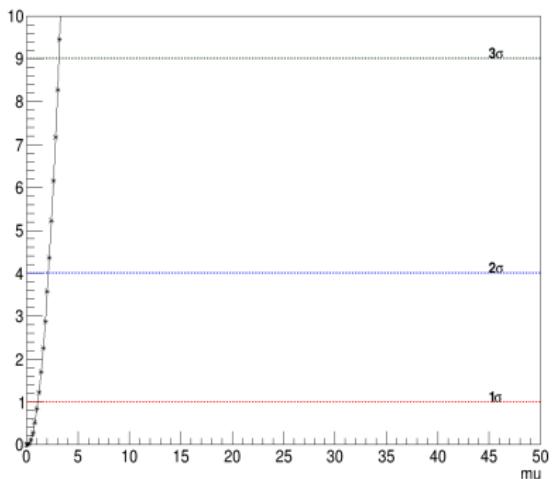
For purposes of establishing an upper limit on the strength parameter μ , we consider two closely related test statistics. First, we may define

$$q_\mu = \begin{cases} -2 \ln \lambda(\mu) & \hat{\mu} \leq \mu \\ 0 & \hat{\mu} < 0 \end{cases}$$

The reason for setting $q_\mu = 0$ for $\hat{\mu} > \mu$ is that when setting an upper limit, one would not regard data with $\hat{\mu} > \mu$ as representing less compatibility with μ than the data obtained, and therefore this is not taken as part of the rejection region of the test. From the definition of the test statistic one sees that higher values of q_μ represent greater incompatibility between the data and the hypothesized value of μ .

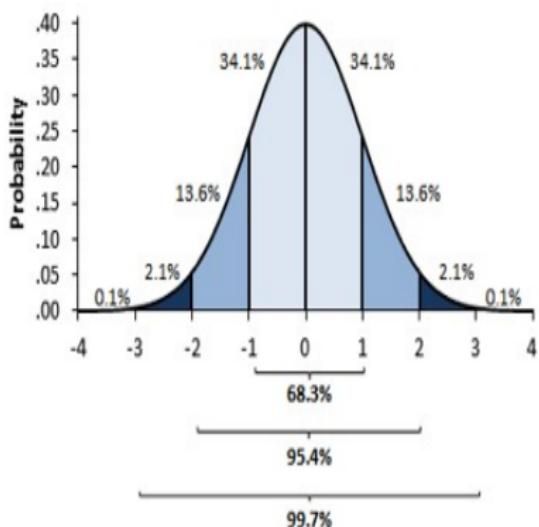
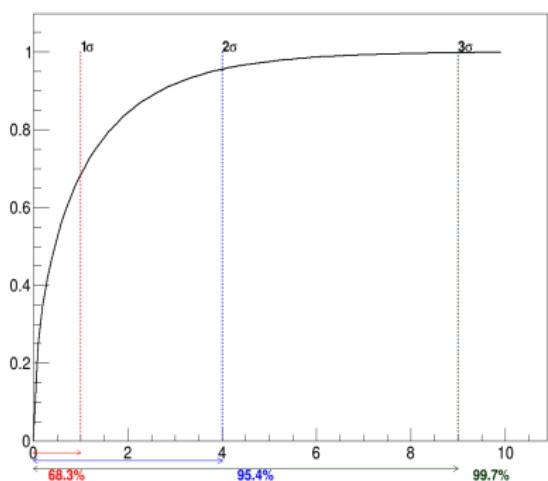


Likelihood scan for $k_t=1$ (SM)



Likelihood scan for $k_t = -1$

Test statistic q_μ for upper limits



Test statistic q_μ for upper limits

Table for estimation of μ for fitting and upper limit with luminosity of 35.9 fb $^{-1}$

Model	μ	μ upper limit
SM ($k_t = 1$)	$2.4975\text{e}+00 \pm 1.29\text{e}+01$	28.047
$k_t = -1$	$6.9361\text{e}-02 \pm 1.02\text{e}+00$	1.97

Luminosity

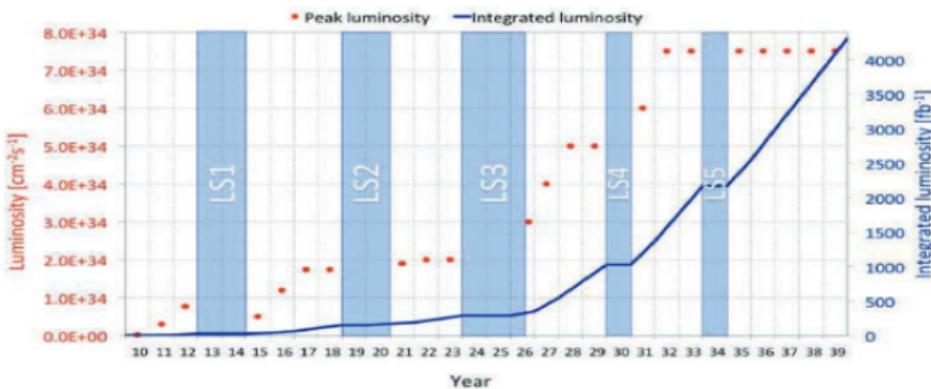
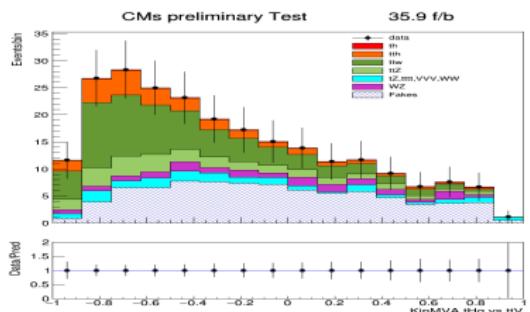
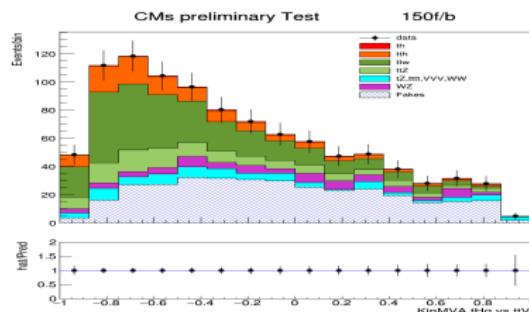
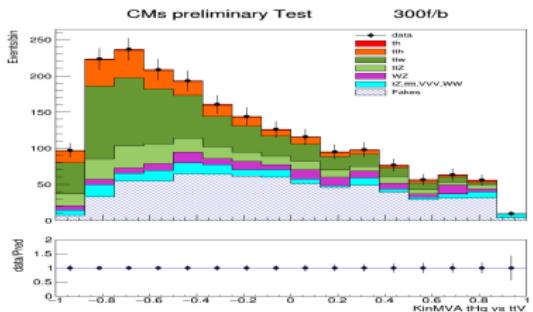


Table of luminosity of CERN measured and expected through years[8]

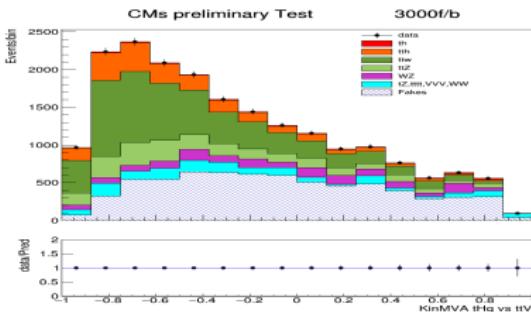
Extrapolation of luminosity for SM



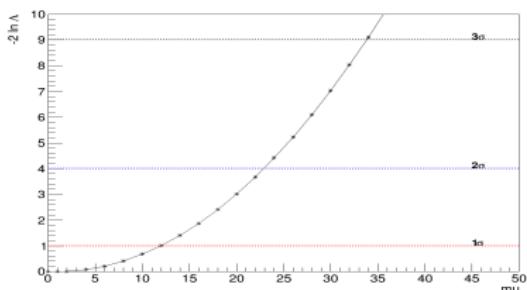
35.9 fb⁻¹



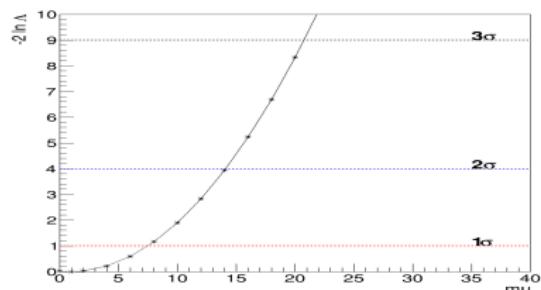
150 fb⁻¹



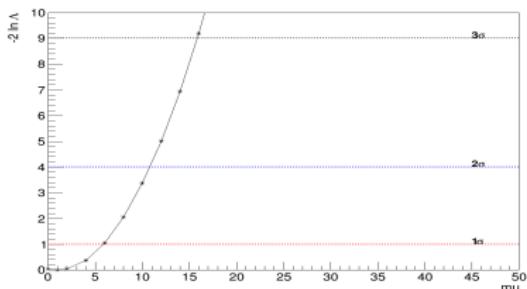
Likelihood scan for SM



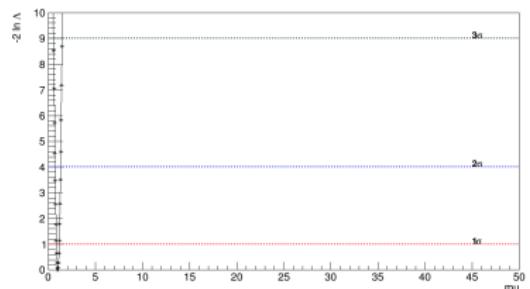
35.9 fb⁻¹



150 fb⁻¹

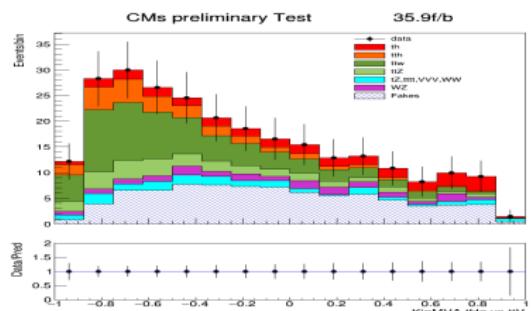


300 fb⁻¹

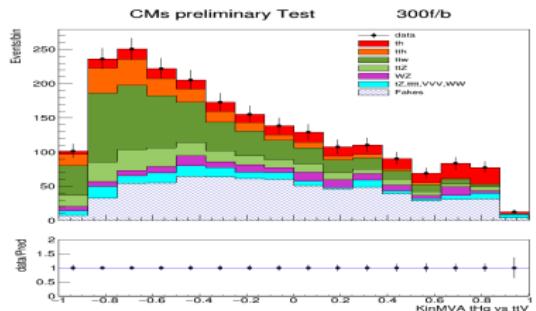


3000fb⁻¹

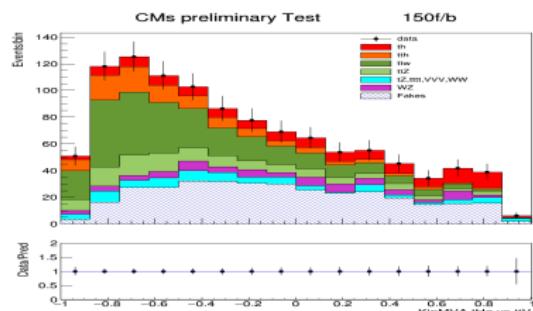
Extrapolation of luminosity for $k_t = -1$



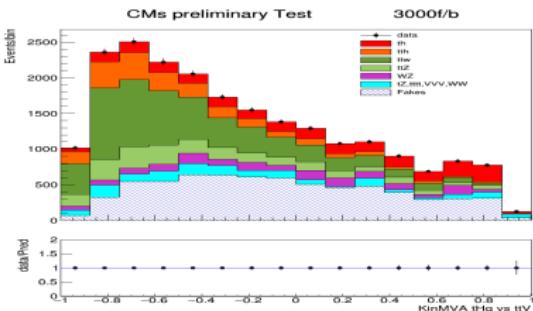
35.9 fb⁻¹



300 fb⁻¹

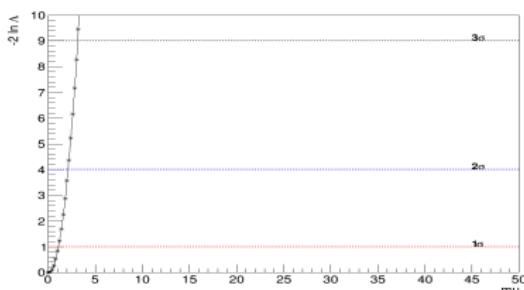


150 fb⁻¹

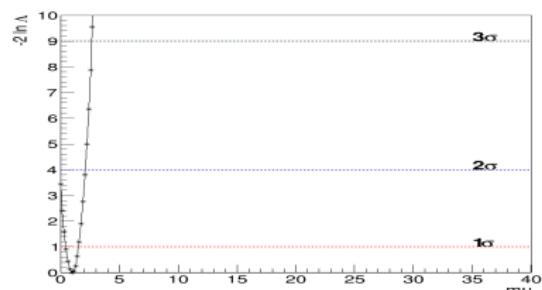


$\sim 3000 \text{ fb}^{-1}$

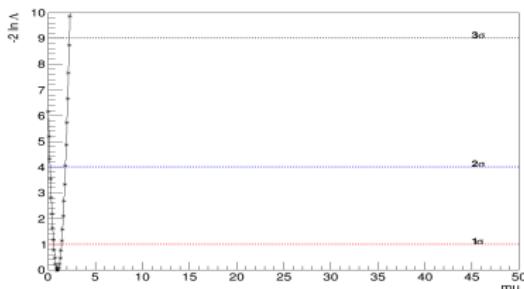
Likelihood scan for $k_t = -1$



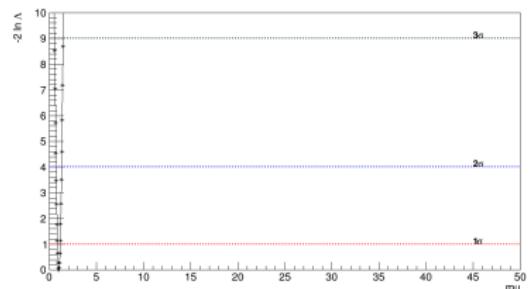
35.9 fb⁻¹



150 fb⁻¹



300 fb⁻¹



$\sim 3000 \text{ fb}^{-1}$

Results

Estimation of μ and upper limits for extrapolations

Luminosity (fb $^{-1}$)	μ	μ for $k_t = -1$	μ upper limit	μ upper limit for $k_t = 1$
35.9	1.00 ± 8.32	1.00 ± 0.249	22.328	2.777
150	1.00 ± 6.44	1.00 ± 0.544	12.619	0.915
300	1.00 ± 4.83	1.00 ± 0.407	9.427	0.649
3000	1.00 ± 1.54	1.00 ± 0.151	3.442	0.25

Conclusions

- We analyzed the tH process produced from PP collisions for the production of $\mu\mu$ final states
- We discussed about events selections for the 2016 and by using BDT process, separate signal and background data.
- The creation of a model with systematic uncertainties to make a minimization (fit) and obtain a fit which is compatible with Standard Model.
- two models are included for comparison
- Generation of likelihood scan for the exclusion of data and obtain the probability of detection of a Higgs boson.
- Generating simulations for predict results with higher luminosities, according to the future experiments.

References

- Gross F. *Relativistic quantum mechanics and field theory*, 1994 ,WILEY-VCH Verlag GmbH & Co. KGaA
 - Griffiths, D. *Introduction to Elementary Particles*, 2^o edition, 2008,WILEY-VCH Verlag GmbH & Co. KGaA
 - The CMS collaboration, *Search for production of a Higgs boson and a single top quark in multilepton final states in proton collisions at 13 TeV*, CMS-PAS-HIG-17-005 .arXiv:1811.09696
 - Verkerke W *Dealing with systematic uncertainties* 2014. From https://indico.cern.ch/event/287744/contributions/1641261/attachments/535763/738679/Verkerke_Statistics_3.pdf

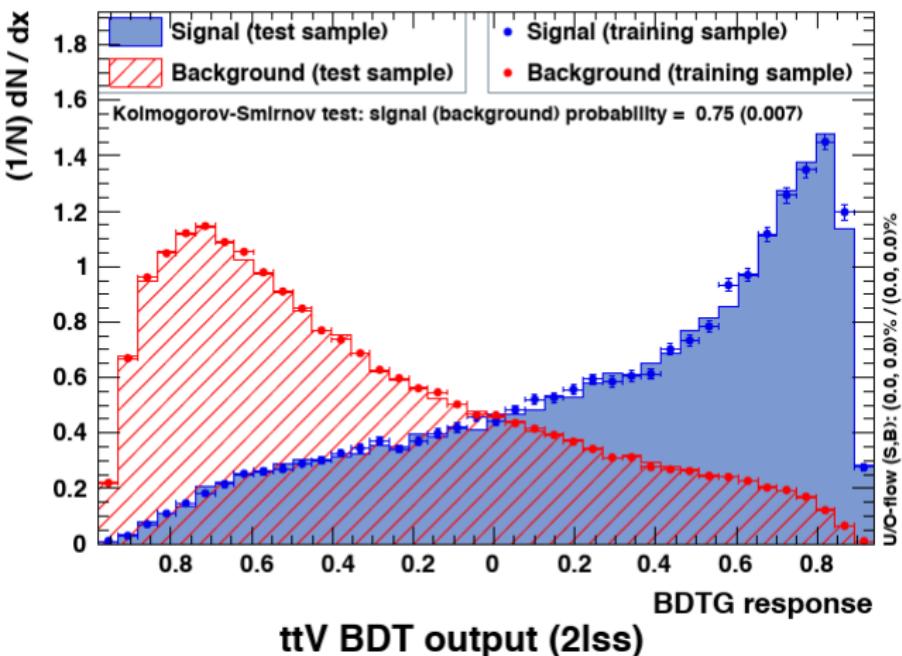
-  ATLAS and CMS Collaborations, *Combined measurement of the Higgs boson mass in pp collisions at $\sqrt{s} = 7$ and 8 TeV with the ATLAS and CMS experiments*, Phys. Rev. Lett. 114 (2015) doi:10.1103/PhysRevLett.114.191803, arXiv:1503.07589
-  Cowan G. , Cranmer K., Gross E. , Vitells O. *Asymptotic formulae for likelihood-based tests of new physics* 2013 , arXiv:1007.1727
-  CMS collaboration *Search for tHq production in multilepton final states at 13 TeV* ,2017,CMS AN-16-378
-  CMS collaboration *High-Luminosity Large Hadron Collider (HL-LHC)*, 2017,CERN-2017-007-M

Back up

BDT parameters

- Gradient boosted (BDTG)
 - No. of trees = 800
 - No. of cuts = 50
 - Maximum depth = 3
 - Found to be most discriminating, minimal overtraining.

TMVA overtraining check for classifier: BDTG



List of histograms used in the analysis

Data taken from the file plots-thq-2lss-kinMVA.root 2016
CERN

```
thqMVA_ttv_2lss_40_tZq
thqMVA_ttv_2lss_40_ttZ
thqMVA_ttv_2lss_40_VVV
thqMVA_ttv_2lss_40_ttW
thqMVA_ttv_2lss_40_data_fakes
thqMVA_ttv_2lss_40_tth
thqMVA_ttv_2lss_40_thW_hww
thqMVA_ttv_2lss_40_WWss
thqMVA_ttv_2lss_40_tttt
thqMVA_ttv_2lss_40_WZ
thqMVA_ttv_2lss_40_thQ_hww
```

Backgrounds and signal histograms

tHq: Signal (tH)

- thqMVA_ttv_2lss_40_tHq_hww
 - thqMVA_ttv_2lss_40_tHW_hww

Backgrounds

t[̄]W

- thqMVA ttv 2lss 40 ttW

t[̄]tZ

- thqMVA_ttv_2lss_40_ttZ

wz

- thqMVA_ttv_2lss_40_WZ

Backgrounds and signal histograms

Backgrounds

tZ, VVV, tttt, WW;

- thqMVA_ttv_2lss_40_tZq
 - thqMVA_ttv_2lss_40_WWss
 - thqMVA_ttv_2lss_40_VVV
 - thqMVA_ttv_2lss_40_tttt

t[̄]H

- thqMVA_ttv_2lss_40_ttH

Non prompt leptons (fakes)

- thqMVA_ttv_2lss_40_data_fakes

Closure: hay jets en los eventos de tt (gluon-gluon \rightarrow tt + gluon) que se pasan como muones. jet \rightarrow muon = fake

Fakes: proceso de QCD que genera muchos jets (por ejemplo gluon-gluon \rightarrow gluons, quarks) : jet \rightarrow muons fakes are estimated data

1

¹Due to existence of many uncertainties, it is necessary to sum all the uncertainties. When there is no Correlation, the uncertainties must be summed as the square root of the squares of each uncertainty

Sources of systematic uncertainty

Detector-simulation related uncertainty

- Calibrations (electron, jet energy scale)
 - Efficiencies (particle ID, reconstruction)
 - Resolutions (jet energy, muon momentum)

Theoretical uncertainties

- Factorization/Normalization scale of MC generators
 - Choice of MC generator (ME and/or PS, e.g. Herwig vs Pythia)

Monte Carlo Statistical uncertainties

- Statistical uncertainty of simulated samples[2]

Observed events

- After applying the event pre-selection on the dataset, 280 events are observed in the same-sign $\mu\mu$ channel
 - The events are then sorted into ten categories depending on the output of the two BDT classifiers according to an optimized binning strategy, resulting in a one-dimensional histogram with ten bins.
 - In each point, the tH and $t\bar{t}H$ production cross sections and the Higgs decay branching ratios are modified with the Higgs-top (κ_t) and Higgs-vector boson (κ_v) coupling strength.
 - The Higgs-tau coupling strength modifier (κ_τ) is assumed to be equal to κ_t .
 - All other parameters are assumed to be at the values predicted by the standard model[1].

α values for post fit

Floating Parameter	FinalValue	+/- Error
Lumi	1.0082e+00	+/- 2.42e-02
alpha_sample_tth_sys	1.9100e-01	+/- 9.94e-01
alpha_sample_ttw_sys	8.7537e-01	+/- 9.18e-01
alpha_sample_tt_z_sys	2.0012e-01	+/- 9.95e-01
alpha_sample_tz_sys	1.6749e-01	+/- 1.00e+00
alpha_sample_wz_sys	1.2732e-01	+/- 9.84e-01
alpha_sample_fakes_sys	4.5030e-01	+/- 7.17e-01
mu	2.4975e+00	+/- 1.29e+01

α values for post fit thq $k_t = -1$

Floating Parameter	FinalValue	+/- Error
Lumi	1.0085e+00	+/- 2.40e-02
alpha_sample_tth_sys	1.9290e-01	+/- 1.00e+00
alpha_sample_ttw_sys	8.7911e-01	+/- 9.42e-01
alpha_sample_ttz_sys	2.0370e-01	+/- 9.94e-01
alpha_sample_tz_sys	1.8269e-01	+/- 9.69e-01
alpha_sample_wz_sys	1.4244e-01	+/- 9.99e-01
alpha_sample_fakes_sys	1.4243e-01	+/- 9.98e-01
mu	6.9361e-02	+/- 1.02e+00

Results

Likelihood scan

- Likelihood function (often simply the likelihood) is a function of the parameters of a statistical model, given specific observed data.
 - Likelihood functions play a key role in frequentist inference, especially methods of estimating a parameter from a set of statistics.
 - In informal contexts, "likelihood" is often used as a synonym for probability.

Results

Likelihood scan

To test a hypothesized value of μ we consider the profile likelihood ratio

$$\lambda(\mu) = \frac{L(\mu, \hat{\theta})}{L(\hat{\mu}, \hat{\theta})} \quad (7)$$

- Here $\hat{\theta}$ in the numerator denotes the value of θ that maximizes L for the specified μ , it is the conditional maximum-likelihood (ML) estimator of $\hat{\theta}$ (and thus is a function of μ). The denominator is the maximized (unconditional) likelihood function, i.e., $\hat{\mu}$ and $\hat{\theta}$ are their ML estimators[4].
 - The presence of the nuisance parameters broadens the profile likelihood as a function of μ relative to what one would have if their values were fixed. This reflects the loss of information about μ due to the systematic uncertainties[4][6].

Statistical test

It is convenient to use the statistic

$$t_\mu = -2 \ln \lambda(\mu) \quad (8)$$

as the basis of a statistical test. Higher values of t_μ thus correspond to increasing incompatibility between the data and μ . We may define a test of a hypothesized value of μ by using the statistic t_μ directly as measure of discrepancy between the data and the hypothesis, with higher values of t_μ correspond to increasing disagreement[6]

Statistical test

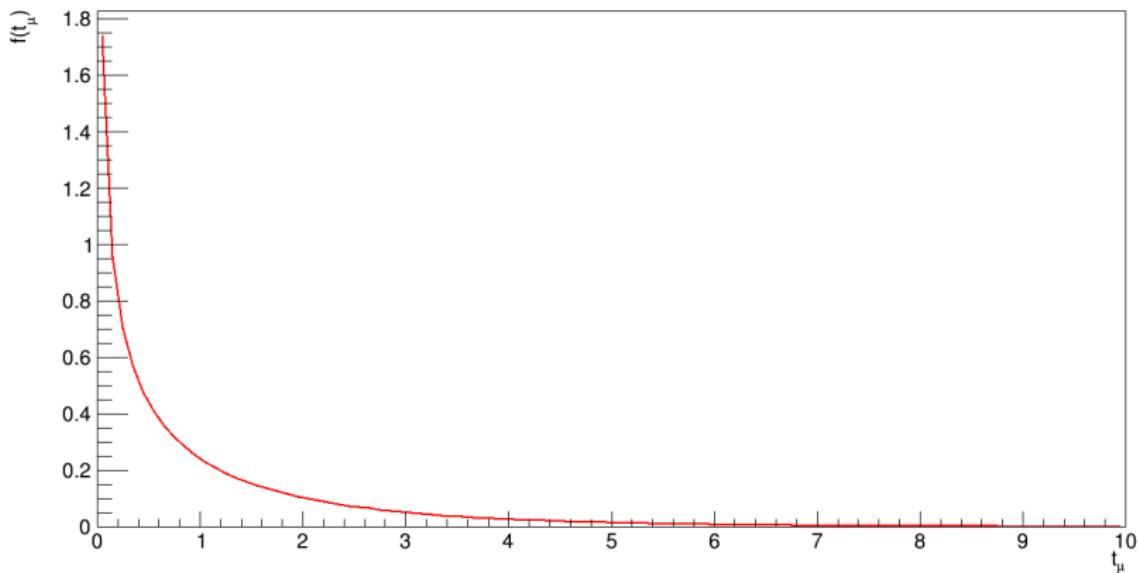


Figure: Statistic test plot $f(t_\mu)$ vs t_μ with $t_\mu = -2 \ln \lambda(\mu)$

P value

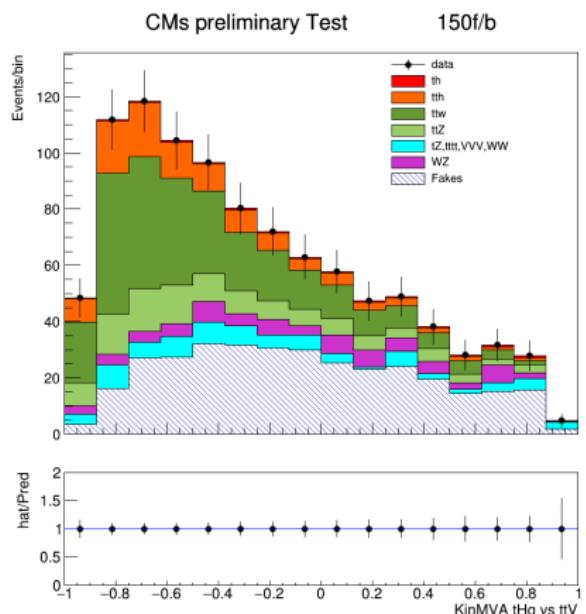
To quantify the level of disagreement we compute the P-value

$$P_\mu = \int_{t_\mu}^{\infty} f(t\mu|\mu) dt_\mu \quad (9)$$

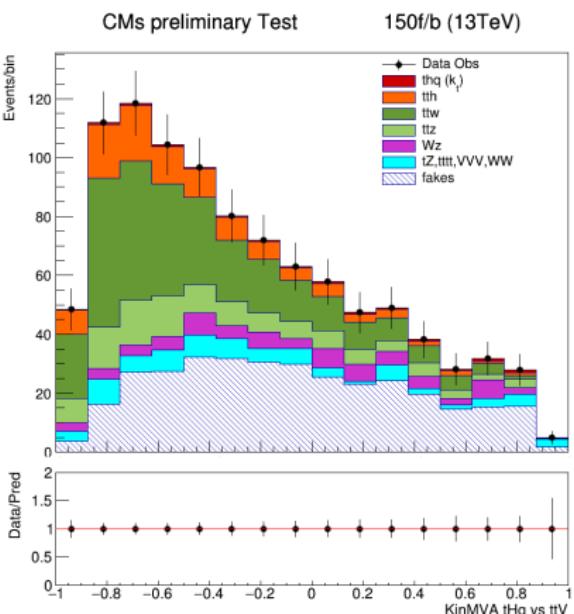
where $t \mu$ is the value of the statistic $t\mu$ observed from the data and $f(t\mu|\mu)$ denotes the PDF (Probability density function) of $t \mu$ under the assumption of the signal strength μ [6]

Luminosity scale

150 fb



Prefit



Post fit

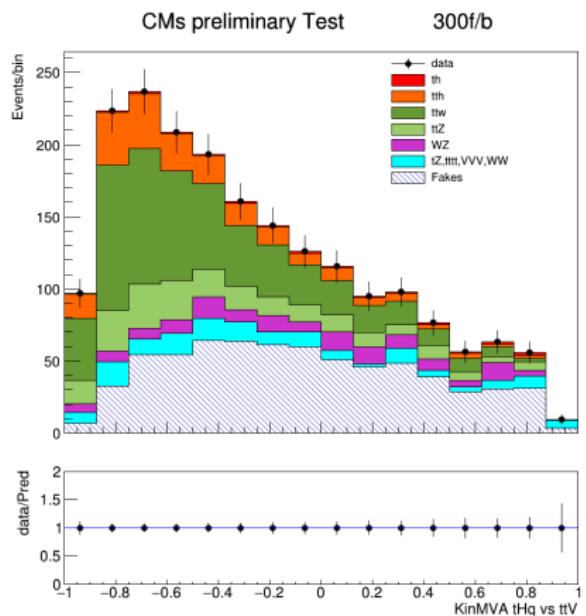
Luminosity scale

150 fb

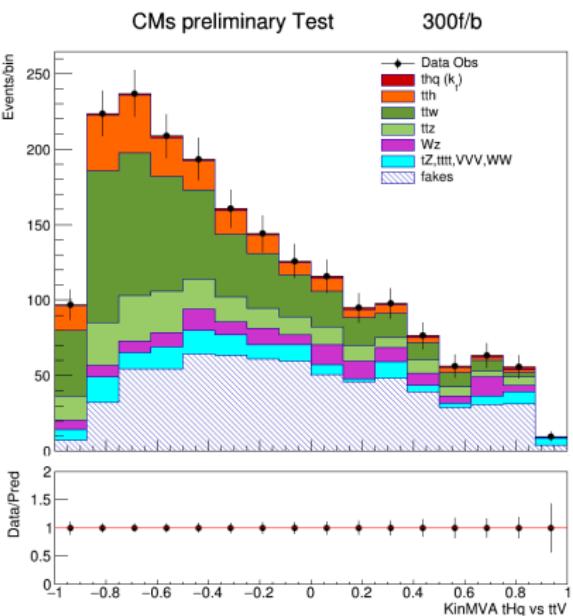
event	N prefit	N Postfit
tH	8.89972 ± 0.885511	8.89972 ± 57.2784
ttH	101.031 ± 9.45	101.031 ± 9.27757
ttW	284.248 ± 39.8961	284.248 ± 31.3604
ttZ	108.175 ± 13.3806	108.175 ± 13.1161
tZ	71.7409 ± 34.1567	71.7409 ± 31.8423
WZ	$67.8134 \pm$	67.8134 ± 31.8861
fakes	338.189 ± 170.744	338.189 ± 84.302

Luminosity scale

300 fb



Prefit



post fit

Luminosity scale

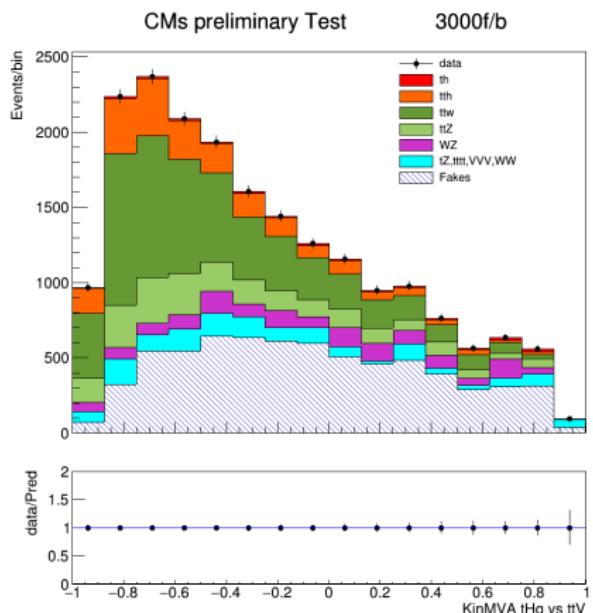
300 fb

Table of yields for background and signal

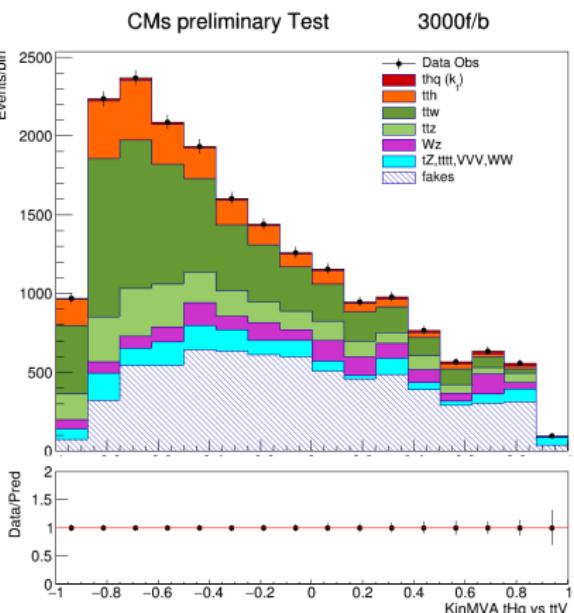
tH	$17.7994 \pm$	17.7994 ± 85.9204
ttH	202.061 ± 18.9162	202.061 ± 18.4728
ttW	568.496 ± 79.7922	568.496 ± 56.3133
ttZ	216.351 ± 26.7611	216.351 ± 26.099
tZ	143.482 ± 72.2699	143.482 ± 58.5044
WZ	135.627 ± 68.3133	135.627 ± 60.7892
fakes	676.379 ± 341.488	676.379 ± 133.725

Luminosity scale

3000 fb



Prefit



Post fit

Luminosity scale

3000 fb

Table of yields for background and signal

event	N prefit	N Postfit
tH	5076.72 \pm 505.128	5076.72 \pm 643.706
ttH	2020.61 \pm 189.162	2020.61 \pm 183.997
ttW	5684.96 \pm 797.922	5684.96 \pm 561.487
ttZ	2163.51 \pm 267.611	2163.51 \pm 256.289
tZ	1434.82 \pm 722.699	1434.82 \pm 303.361
WZ	1356.27 \pm 683.133	1356.27 \pm 383.271
fakes	6763.79 \pm 3414.88	6763.79 \pm 591.597