# An Improved Numerical Approximation of the Horizontal Gradients in a Terrain-Following Coordinate System

## Y. MAHRER<sup>1</sup>

Cooperative Institute For Research in the Atmosphere (CIRA), Colorado State University, Fort Collins, CO 80523

(Manuscript received 6 July 1983, in final form 14 January 1984)

### **ABSTRACT**

The use of a fine vertical grid resolution near the surface may lead to a numerically inconsistent approximation of the horizontal gradient terms in a terrain-following coordinate system. This occurs when the distance between two vertical grid points is smaller than the elevation difference between two horizontally adjacent (in the terrain-following coordinate system) points. In this paper an improved numerical procedure is proposed which eliminates this inconsistency and significantly increases the accuracy of the numerical approximation. Results are compared with those obtained with the conventional forward and centered schemes.

### 1. Introduction

The use of the terrain-following coordinate system (TFCS), in meteorological numerical models, has been shown to be very effective when topographic features are considered. Generally, this coordinate system is defined by using the transformation

$$x^* = x$$
,  $y^* = y$ ,  $z^* = s \frac{z - Z_G}{s - Z_G}$ 

where x, y and z are the horizontal and vertical coordinates in a Cartesian coordinate system;  $x^*$ ,  $y^*$ , and  $z^*$  are the horizontal and vertical coordinates in the TFCS;  $Z_G$  is the terrain height and s is a reference height (usually the top of the model). This form of a TFCS has been used in recent years in large-scale and mesoscale numerical models (e.g., Kasahara, 1974; Mahrer and Pielke, 1975; Gal-Chen and Somerville, 1975; McNider and Pielke, 1981; Yamada, 1983; among others).

It will be shown in Section 2, that when using a numerical scheme for the horizontal gradients in the TFCS, care must be taken that the distance between two vertical grid points will not be less than the elevation difference between two horizontally (in the TFCS) adjacent points. This situation is mostly pertinent in mesoscale models that use a relatively fine vertical grid resolution in the atmospheric boundary layer.

It is also worth noting that models which use the sigma  $(\sigma)$  coordinate system, where the vertical co-

ordinate  $\sigma$  is defined as a function of pressure are subject to a similar minimum vertical grid size requirement. The  $\sigma$  coordinate system was first introduced by Phillips (1957) and has been adopted in mesoscale models by Anthes and Warner (1978), Alpert et al. (1982), among others.

In some of the aforementioned models, the investigators have used relatively fine grid resolution near the surface, assuming implicitly that the variables vary linearly with height within the elevation difference between two horizontally adjacent points. For example: in Mahrer and Pielke (1975) the smallest grid interval was 50 m while the maximum elevation difference between two horizontal grid points was 400 m; in Yamada (1983) the corresponding values were 2 m and 100 m, while in Gal-Chen and Somerville (1975) they were approximately 100 m and 120 m respectively.

Several investigators (e.g., Smagorinsky et al., 1967; Gary, 1973; Janjic, 1977) have recognized, in the  $\sigma$  coordinate system, an increased truncation error in the pressure gradient terms in the presence of very steep topography. Janjic (1977) specifically pointed out the inconsistency in the pressure gradient calculation in cases of very steep slopes of surfaces and thin layers. This inconsistency, however, could be tolerated in Janjic's experiments because his smallest vertical spacing was 100 mb ( $\sim$ 1000 m).

Although the errors associated with a TFCS are most significant in the pressure gradient terms, they may be also important in other horizontal gradient terms, mainly within the planetary boundary layer.

In the present study, it is demonstrated that by using an improved numerical expression for the horizontal gradient terms in the TFCS, errors in the horizontal gradients of the meteorological fields are reduced substantially when the vertical grid interval is very small.

<sup>&</sup>lt;sup>1</sup> Permanent affiliation: The Hebrew University of Jerusalem, Seagram Centre for Soil and Water Science, Faculty of Agriculture, Rehovot, 76100 Israel.

Results are compared with those obtained by the forward and centered conventional schemes.

# 2. Evaluation of the numerical difference scheme for the horizontal gradients in TFCS

The transform equations of the partial derivatives with respect to x and z from the Cartesian coordinate system to the TFCS are:

$$\frac{\partial}{\partial x}\bigg|_{z} = \frac{\partial}{\partial x}\bigg|_{z^{*}} + \frac{\partial z^{*}}{\partial x}\bigg|_{z} \frac{\partial}{\partial z^{*}}$$

$$= \frac{\partial}{\partial x}\bigg|_{z^{*}} + \frac{z^{*} - s}{s - Z_{G}} \frac{\partial Z_{G}}{\partial x} \frac{\partial}{\partial z^{*}}, \qquad (1)$$

$$\frac{\partial}{\partial z} = \frac{\partial z^*}{\partial z} \frac{\partial}{\partial z^*}.$$
 (2)

A commonly used finite differencing scheme for (1) in TFCS is a forward-in-space approximation, namely:

$$\frac{\partial F}{\partial x}\Big|_{z} = \frac{F(i+1,j) - F(i,j)}{\Delta x} + \frac{\partial z^{*}}{\partial x}\Big|_{i,j} \frac{F(i+1,j) - F(i+1,j-1)}{\Delta z^{*}}, \quad (3)$$

where F represents the meteorological field, and

$$\left. \frac{\partial z^*}{\partial x} \right|_{i,j} = \frac{z^*(j) - s}{s - Z_G(i+1)} \frac{Z_G(i+1) - Z_G(i)}{\Delta x} \,.$$

As illustrated in Fig. 1, in which the vertical reso-

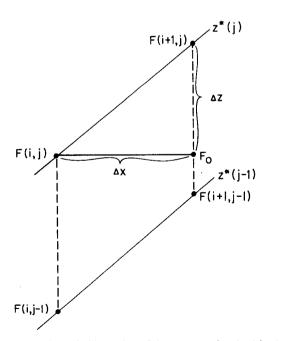


Fig. 1. A schematic illustration of the geometry involved in the calculation of the horizontal derivative in the terrain-following coordinate system.

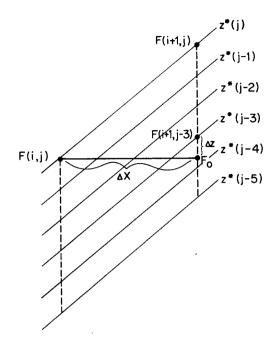


FIG. 2. As in Fig. 1, but with a five times finer vertical resolution.

lution is relatively coarse, the first-order approximation for  $F_0$  can be written as

$$F_0 = F(i+1,j) - \frac{\partial F}{\partial z} \Big|_{i+1,j} \Delta z. \tag{4}$$

Since the forward difference scheme for the horizontal gradient in the Cartesian coordinate system is given by

$$\frac{\partial F}{\partial x}\Big|_{z} = \frac{F_0 - F(i, j)}{\Delta x},$$
 (5)

Eq. (4) can be written as

$$\frac{\partial F}{\partial x}\Big|_{z} = \frac{F(i+1,j) - F(i,j)}{\Delta x} + \frac{\Delta z}{\Delta x} \frac{F(i+1,j) - F(i+1,j-1)}{z(j) - z(j-1)}.$$
(6)

After substituting  $z^*$  for z in Eq. (6) we obtain Eq. (3).

However, let us consider now a case in which a finer vertical resolution, as illustrated in Fig. 2, is adopted. Here  $\Delta x$  is the same as in Fig. 1 but, the vertical grid distances were reduced five times. In this situation the best first-order approximation for  $F_0$  will be

$$F_0 = F(i+1, j-3) - \frac{\partial F}{\partial z} \Big|_{i+1, j-3} \Delta z \tag{7}$$

and not the relation given in Eq. (4). Under the current refined vertical grid resolution, Eq. (4) provides a first-order approximation to  $F_0$  only when the field is linear

between the points j and j-4. Eq. (4), however, is indirectly commonly used in numerical models over complex terrain, since, in general the numerical scheme given by Eq. (3) is utilized.

By substituting  $F_0$  according to Eq. (7) into Eq. (5) and replacing z with  $z^*$  we obtain

$$\frac{\partial F}{\partial x}\Big|_{z} = \frac{F(i+1, j-3) - F(i, j)}{\Delta x} + \left[\frac{\partial z^{*}}{\partial x}\Big|_{i,j} - \frac{z^{*}(j-3) - z^{*}(j)}{\Delta x}\right] \times \frac{F(i+1, j-3) - F(i+1, j-4)}{z^{*}(j-3) - z^{*}(j-4)}, \quad (8)$$

and in the general case

$$\frac{\partial F}{\partial x}\Big|_{z} = \frac{F(i+1,M) - F(i,j)}{\Delta x} + \left[\frac{\partial z^{*}}{\partial x}\Big|_{i,j} - \frac{z^{*}(M) - z^{*}(j)}{\Delta x}\right] \times \frac{F(i+1,M) - F(i+1,M-1)}{z^{*}(M) - z^{*}(M-1)}. \quad (9)$$

Notice that the elevation at grid point (i, j) satisfies the condition

$$z(i+1, M-1) \le z(i, j) \le z(i+1, M)$$

and that M may be larger, smaller or equal to j, depending on the slope of the topography. Eq. (9) will be referred to later as the modified scheme.

Similarly, the expression for the horizontal gradient can be easily modified when using a centered or a backward differencing scheme. For the centered scheme, two points M and N (to the left and to the right of the point in question) have to be determined.

In the aforementioned numerical approximations, special care must be taken when M=1 (Fig. 3), since then,  $F_0$  does not exist (it is an imaginary point inside the topography). Under these conditions the distance of a point to the boundary must be calculated so that

$$\left. \frac{\partial F}{\partial x} \right|_z = \frac{F(i,j) - F_e}{\text{DIS}}$$
.

Here  $F_e$  is the boundary value of F, and DIS is the horizontal distance (in the Cartesian coordinate system) of point (i, j) to the boundary.

It should be pointed out, that when using a numerical scheme (like Eq. (3), for example) for the horizontal derivative at point (i, 3), as illustrated in Fig. 3, a value for  $F_0$  is extrapolated and the actual boundary condition is not imposed.

Conclusion of the evaluation in this section suggests that in order that the numerical approximation given in Eq. (3) will be appropriate, the smallest vertical grid spacing,  $(\Delta z)$ , should satisfy the condition:

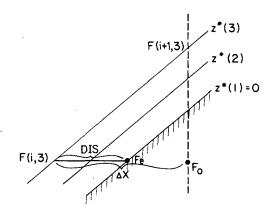


FIG. 3. A schematic illustration of the geometry involved in the calculation of the horizontal derivative in the terrain-following coordinate system near the terrain boundary with a fine vertical grid resolution.

$$\Delta z \geqslant \frac{s-z^*}{s} \Delta Z_G$$
 or  $\Delta z^* \geqslant \frac{s-z^*}{s-Z_G} \Delta Z_G$ ,

where  $\Delta Z_G$  is the height change of the terrain per one horizontal grid interval. Thus, at the surface, where  $z^* = 0$ , the smallest vertical grid interval must fulfill the requirement  $\Delta z \ge \Delta Z_G$ .

With the *sigma* coordinate system where  $\sigma$  is defined as

$$\sigma = \frac{P - P_t}{P_s - P_t},$$

where P is the pressure,  $P_s$  and  $P_t$  are the surface and top pressures respectively. Here, the condition of the minimum vertical grid interval will be

$$\Delta \sigma \geqslant \frac{\sigma}{P_s - P_t} \Delta P_s$$

where  $\Delta P_s$  is the change in the surface pressure per one horizontal grid interval.

### 3. Numerical test evaluations

In order to compare the accuracy of the "old" (Eq. 3) and the "modified" (Eq. 9) schemes, the horizontal gradients of hydrostatic pressure and temperature fields were calculated in the presence of topography. The calculations were performed in a two-dimensional (x, z) domain, consisting of  $30 \times 20$  grid points. A bell-shaped mountain of the form

$$h(X) = \frac{H}{1 + (X - X_0)^2},$$

with H=1 km and centered at  $X=X_0$ , was considered. The vertical  $z^*$  grid points were set at the heights of 0, 5, 15, 50, 100, 200, 300, 500, 700, 1000, 1500, 2000, 2500, . . . , 4500, and 5000 m. An initial horizontally uniform (in the Cartesian coordinate system) temperature field with a lapse rate of  $0.01^{\circ}$ C m<sup>-1</sup> in the first

TABLE 1. Horizontal pressure differences (mb) along one grid interval at selected heights near the mountain crest as obtained with: (a) forward differencing scheme (Eq. 3), (b) centered differencing scheme, (c) modified forward scheme (Eq. 4), (d) modified centered scheme.

Scheme	Height (m)										
	5	15	50	100	200	300	500	700	1000	2000	
a	1.065	0.962	0.967	0.882	0.938	0.722	0.831	0.601	0.507	0.313	
b	0.298	0.266	0.278	0.219	0.283	0.129	0.252	0.103	0.176	0.139	
c	0.019	0.024	0.039	0.059	0.082	0.086	0.044	`0.026	0.023	0.039	
d	0.017	0.019	0.025	0.030	0.045	0.041	0.022	0.015	0.006	0.003	

1000 m and of 0.0065°C m<sup>-1</sup> above 1000 m was prescribed. With this thermal stratification the horizontal gradients of the pressure and temperature fields should be identically zero.

The numerical evaluations for pressure and temperature differences along one horizontal grid interval with the forward and centered schemes are summarized in Tables 1 and 2 at grid points (14, j) where the largest errors occurred. It can be seen that with the "old" forward scheme horizontal pressure differences were as large as 1.065 mb per one horizontal grid interval, while with the "modified" forward scheme they were at least 10 times smaller. For temperatures, the differences with the modified forward scheme were practically zero, while with the "old" forward scheme they were larger than 1°C at the lower layers. In the upper layers, where the prescribed temperature distribution is continuously linear, both schemes evaluations were correct.

With the centered scheme, the errors were somewhat reduced (Tables 1 and 2, b) but were still substantially larger than the errors with the "modified" centered scheme (Tables 1 and 2, d).

In order to test the effect of increasing the vertical grid interval on the accuracy of the schemes the horizontal pressure gradients were evaluated with a coarser vertical grid spacing. The lowest  $z^*$  points were set at the heights of 0, 5, 100, 300, and 700 m. From Table 3 it can be seen that with the centered scheme the horizontal differences were somewhat reduced at equivalent heights. With the modified scheme, however, errors have increased, but were still less than with the centered scheme. These results show that the mod-

ified scheme is consistent with convergence of numerical approximations to the true value as the grid interval is reduced. The "old" scheme, however, suffers from the problem that the errors do not necessarily decrease when the grid interval is reduced (Mesinger, 1982).

### 4. Conclusions

It has been shown that the proposed numerical schemes for the approximation of the horizontal gradients in a terrain-following coordinate system significantly improves the accuracy when the vertical grid is highly refined. This method works well with very steep slopes and when the vertical grid interval is smaller than the elevation difference between two horizontally adjacent terrain points. Further improvement can be made, when the vertical resolution near the surface is very small, by using the actual horizontal distance between a grid point and the terrain boundary, instead of the original horizontal grid interval. In this paper the sources of the numerical deficiencies associated with the application of conventional schemes in TFCS are pointed out. The proposed procedure can be easily applied when using the forward, backward or centered differencing schemes. It is important to notice that reducing the horizontal grid interval can be a way to eliminate the errors. However, when steep slopes are considered this will not be practical because of the significant increase in computational expenses. Finally, it is worth noting that models that use the sigma coordinate system are also subject to a similar error.

TABLE 2. Horizontal temperature differences (°C) along one grid interval at selected heights near the mountain crest as obtained with: (a) forward differencing scheme (Eq. 3), (b) centered differencing scheme, (c) modified forward scheme (Eq. 4), (d) modified centered scheme.

Scheme	Height (m)										
	5	15	50	100	200	300	500	700	1000	2000	
a	1.734	1.702	1.590	1.429	1.108	0.787	0.146	0.000	0.000	0.000	
b	0.867	0.851	0.794	0.715	0.554	0.394	0.073	0.018	0.000	0.000	
c	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.008	0.007	0.000	
d	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.000	0.002	0.000	

TABLE 3. Horizontal pressure differences (mb) along one grid interval at selected heights near the mountain crest as obtained for the coarse vertical grid spacing with: (a) centered differencing scheme, (b) modified centered scheme.

	Height (m)								
5	100	300	700	1000	2000				
0.209	0.174	0.113	0.096	0.150	0.139 0.003				
	5 0.209 0.024	0.209 0.174	5 100 300 0.209 0.174 0.113	5     100     300     700       0.209     0.174     0.113     0.096	5 100 300 700 1000 0.209 0.174 0.113 0.096 0.150				

Acknowledgments. The author thanks Mordecay Segal for his useful comments and suggestions which led to the successful completion of this study. Bob Kessler is thanked for reading and commenting on the manuscript. This work was supported by the Israel Academy of Sciences and Humanities, Basic Research Foundation, by the Israel Ministry of Energy and Infrastructure, Department of Research and Development and by NSF Grant ATM-8242931. Computations in this study were performed using the National Center for Atmospheric Research computer.

#### REFERENCES

Alpert, P., A. Cohen, J. Neumann and E. Doron, 1982: A model simulation of the summer circulation from the eastern Medi-

- terranean past Lake Kinnret in the Jordan Valley. Mon. Wea. Rev., 110, 994-1006.
- Anthes, R. A., and T. T. Warner, 1978: Development of hydrodynamic models suitable for air pollution and other mesometeorological studies. Mon. Wea. Rev., 106, 1045-1078.
- Gal-Chen, T., and R. C. J. Somerville, 1975: Numerical solution of the Navier-Stokes equations with topography. J. Comput. Phys., 17, 276-310.
- Gary, J. M., 1973: Estimate of truncation error in transformed coordinate, primitive equation atmospheric models. J. Atmos. Sci., 30, 223-233.
- Janjic, Z. I., 1977: Pressure gradient force and advection scheme used for forecasting with steep and small scale topography. *Beitr. Phys. Atmos.*, 50, 186-199.
- Kasahara, A., 1974: Various vertical coordinate systems used for numerical weather prediction. Mon. Wea. Rev., 102, 509-522.
- Mahrer, Y., and R. A. Pielke, 1975: A numerical study of the air flow over mountains using the two-dimensional version of the University of Virginia mesoscale model. J. Atmos. Sci., 32, 2144-2155.
- McNider, R. T., and R. A. Pielke, 1981: Diurnal boundary-layer development over sloping terrain. *J. Atmos. Sci.*, **38**, 2198–2212.
- Mesinger, F., 1982: On the convergence and error problems of the calculation of the pressure gradient force in sigma coordinate models. Geophys. Astrophys. Fluid Dyn., 19, 105-117.
- Phillips, N. A., 1957: A coordinate system having some special advantages for numerical forecasting. *J. Meteor.*, **14**, 184–185.
- Smagorinsky, J. R. F. Strickler, W. E. Sangster, S. Manabe, J. L. Holloway, Jr. and G. D. Hembree, 1967: Prediction experiments with a general circulation model. *Proc. Int. Symp. on Dynamics of Large Scale Atmospheric Processes*. Moscow, 70–134.
- Yamada, T., 1983. Simulations of nocturnal drainage flows by a q<sup>2</sup>l turbulence closure model. J. Atmos. Sci., 40, 91-106.