

Q1) From Schrodinger's equation,

$$a) \frac{-\hbar^2}{2m} \nabla^2 \psi + V(x) \psi = i\hbar \frac{d}{dt} \psi$$

$$\Rightarrow \nabla^2 \psi = (V(x) \psi - i\hbar \frac{d}{dt} \psi) \frac{2m}{\hbar^2}$$

\therefore A complex wavefunction can be expressed

by $\psi = u + iV$, we know that $u_x = v_x$ and $u_t = -v_x$.

$$\Rightarrow \nabla^2 \psi^* = (V(x) \psi^* + i\hbar \frac{d}{dt} \psi^*) \frac{2m}{\hbar^2}$$

Now

So we proceed,

$$\begin{aligned} J &= \left(\frac{i\hbar}{2m} \right) \left(\frac{2m}{\hbar^2} \right) \left\{ \psi^* (V(x) \psi - i\hbar \frac{d}{dt} \psi) \right. \\ &\quad \left. - \psi (V(x) \psi^* + i\hbar \frac{d}{dt} \psi^*) \right\} \\ &= -\frac{i}{\hbar} \left\{ \psi^* V(x) \psi - \psi V(x) \psi^* \right. \\ &\quad \left. - \psi^* i\hbar \frac{d}{dt} \psi - \psi i\hbar \frac{d}{dt} \psi^* \right\} \\ &= -\frac{i}{\hbar} \left\{ -\psi^* i\hbar \frac{d}{dt} \psi - \psi i\hbar \frac{d}{dt} \psi^* \right\} \\ &= -\left\{ \psi^* \frac{\partial}{\partial t} \psi + \psi \frac{\partial}{\partial t} \psi^* \right\} \\ &= -\frac{\partial}{\partial t} (\psi^* \psi) = -\frac{\partial}{\partial t} |\psi|^2 = -\frac{\partial}{\partial t} p \end{aligned}$$

$$\begin{aligned} \rightarrow \nabla \cdot J &= \nabla \cdot \left\{ \frac{i\hbar}{2m} (\psi^* \nabla \psi - \psi \nabla \psi^*) \right\} \\ &= -\frac{i\hbar}{2m} \left(\nabla \psi^* \nabla \psi + \psi^* \nabla^2 \psi \right. \\ &\quad \left. - \nabla \psi \nabla \psi^* - \psi \nabla^2 \psi^* \right) \\ &= -\frac{i\hbar}{2m} (\psi^* \nabla^2 \psi - \psi \nabla^2 \psi^*) \end{aligned}$$

$$\Rightarrow -\frac{\partial}{\partial t} p = \nabla \cdot J \Rightarrow \frac{\partial}{\partial t} p + \nabla \cdot J = 0.$$

5) The Hamiltonian for a free particle is

$$H = \frac{\hat{p}^2}{2m},$$

the time evolution operator yields,

$$e^{-\frac{i\hat{p}^2 t}{2m}} \langle x | p \rangle = \frac{1}{\sqrt{2\pi\hbar}} e^{-\frac{i}{\hbar} \left(px - \frac{p^2 t}{2m} \right)}$$

So we calculate,

$$\begin{aligned} J_x &= -i\hbar \left\{ \frac{d}{dx} \left(\psi^* \frac{d\psi}{dx} - \psi \frac{d\psi^*}{dx} \right) \right\} \\ &= -i\hbar \left\{ \frac{1}{2m} \left(\frac{1}{\sqrt{2\pi\hbar}} \right)^2 e^{-\frac{i}{\hbar} \left(px - \frac{p^2 t}{2m} \right)} \left(\frac{ip}{\hbar} \right) e^{\frac{i}{\hbar} \left(px - \frac{p^2 t}{2m} \right)} \right. \\ &\quad \left. - \left(\frac{1}{\sqrt{2\pi\hbar}} \right)^2 e^{\frac{i}{\hbar} \left(px - \frac{p^2 t}{2m} \right)} \left(\frac{-ip}{\hbar} \right) e^{-\frac{i}{\hbar} \left(px - \frac{p^2 t}{2m} \right)} \right\} \\ &= \left(-\frac{i\hbar}{2m} \right) \left(\frac{1}{\pi\hbar} \right) \left(\frac{ip}{\hbar} + \frac{-ip}{\hbar} \right) \end{aligned}$$

$$J_x = \frac{p}{2\pi m\hbar}$$

$$\text{Moreover, since } p = \hbar|\vec{p}| = \psi^* \vec{p} = \frac{1}{\hbar} \nabla \psi^* \rightarrow \text{evidently,}$$

(Q6,3)

a) We use the property of the translation operator:
 $|\psi'\rangle = T(\delta x)|\psi\rangle \Rightarrow \psi'(x) = \psi(x - \delta x)$
 $\Rightarrow (\psi'(x))^* = \psi^*(x - \delta x)$

Then by definition

$$\langle x | \psi' \rangle = \langle \psi' | x | \psi \rangle, \text{ inserting } \int dx_1 \psi(x_1)$$

$$= \int dx \psi'^*(x) \times \psi'(x)$$

$$= \int dx \psi^*(x - \delta x) \times \psi(x - \delta x)$$

Let the substitution $x \mapsto x + \delta x \rightarrow dx = dx$

$$\therefore \langle x \rangle = \int dx \Psi^*(x) (x + \delta x) \Psi(x)$$

$$= \int dx \Psi^*(x) x \Psi(x) + \int dx \delta x \Psi^*(x) \Psi(x)$$

since $\int dx \Psi^* \Psi = 1 \Rightarrow \langle x \rangle = \langle x \rangle + \delta x$

5) By definition, $\hat{T}(\delta x) = e^{-\frac{i\hat{p}\delta x}{\hbar}}$. We then proceed,

$$\langle p_x \rangle = \langle \Psi' | p_x | \Psi' \rangle$$

$$= \int dx \Psi^*(x) e^{\frac{i\hat{p}\delta x}{\hbar}} p_x \Psi(x) e^{-\frac{i\hat{p}\delta x}{\hbar}}$$

$$= \int dx \Psi^*(x) e^{\frac{i\hat{p}\delta x}{\hbar}} (-i\hbar) \frac{\partial}{\partial x} (\Psi(x)) e^{-\frac{i\hat{p}\delta x}{\hbar}}$$

$$= \int dx \Psi^*(x) e^{\frac{i\hat{p}\delta x}{\hbar}} (-i\hbar) \left\{ \frac{\partial}{\partial x} \Psi(x) e^{-\frac{i\hat{p}\delta x}{\hbar}} + \Psi(x) \underbrace{\frac{\partial}{\partial x} \left(-\frac{i\hat{p}\delta x}{\hbar} \right)}_{\rightarrow 0} e^{-\frac{i\hat{p}\delta x}{\hbar}} \right\}$$

$$= \int dx \Psi^*(x) e^{\frac{i\hat{p}\delta x}{\hbar}} (-i\hbar) \frac{\partial}{\partial x} \Psi(x) e^{-\frac{i\hat{p}\delta x}{\hbar}}$$

$$= \int dx \Psi^*(x) p_x \Psi(x)$$

$$= \langle p_x \rangle$$

(Q6,5)

a) From normalization, $1 = \langle \Psi | \Psi \rangle = \int dP |\Psi|^2$

$$= |N|^2 \int_{-\rho/2}^{\rho/2} dP$$

$$\Rightarrow 1 = |N|^2 \rho \Rightarrow N = 1/\sqrt{\rho}$$

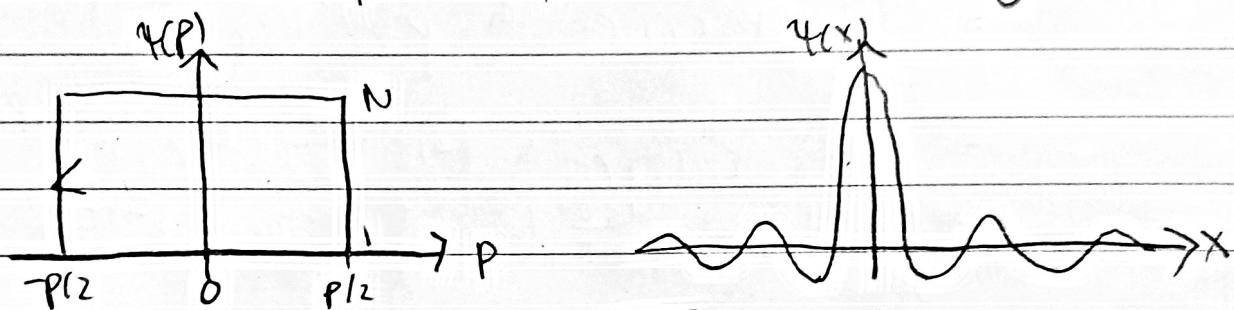
5) By definition, $\Psi(x) = \int dP \Psi(p) e^{\frac{ipx}{\hbar}} \frac{1}{\sqrt{2\pi\rho}}$

$$= \int_{-\rho/2}^{\rho/2} dP \frac{1}{\sqrt{\rho}} e^{\frac{ipx}{\hbar}} \frac{1}{\sqrt{2\pi\rho}}$$

here we treat $1/\sqrt{\rho}$ as a constant since it's the normalization. So then,

$$\begin{aligned}
 &= \frac{1}{\sqrt{p\sqrt{2\pi\hbar}}} \int_{-p/2}^{p/2} e^{\frac{i p x}{\hbar}} dp \\
 &= \frac{1}{\sqrt{p\sqrt{2\pi\hbar}}} \frac{1}{ix} \left(e^{\frac{ipx}{2\hbar}} - e^{-\frac{ipx}{2\hbar}} \right) \\
 \Rightarrow \psi(x) &= \frac{2\hbar}{x} \sin\left(\frac{px}{2\hbar}\right) \frac{1}{\sqrt{p\sqrt{2\pi\hbar}}} = \frac{1}{\sqrt{2\pi\hbar}} \frac{\sin(px/2\hbar)}{px/2\hbar}
 \end{aligned}$$

c) $\psi(p)$ is just a square well function.
 $\psi(x)$ has the form of a Dirichlet integral so



$$\Rightarrow \Delta p_x = \frac{1}{\sqrt{p}}, \quad \Delta x = \sqrt{\frac{p}{2\pi\hbar}}$$

$$\Rightarrow \Delta x \Delta p_x = 1/\sqrt{2\pi\hbar}, \text{ which is independent of } p.$$

Q6(b)

a) if $\psi \in \text{Real} \Rightarrow \psi(x) = \psi^*(x)$

$$\langle p_x \rangle = \langle \psi | p_x | \psi \rangle, \text{ here we insert } \int dx |x\rangle \langle x|$$

$$= \int dx \langle \psi | p_x | x \rangle \langle x | \psi \rangle$$

$$= \int dx \psi^* \underbrace{p_x}_{\in \mathbb{R}} \psi(x)$$

$$= \int dx \psi^* \frac{d}{dx} \psi(x)$$

$$= (-i\hbar) \underbrace{\int dx \psi \frac{d}{dx} \psi}_{\in \mathbb{R}}$$

$\in \mathbb{C}$

$$\Rightarrow \langle p_x \rangle \in \mathbb{C} \Rightarrow \operatorname{Re} \{ \langle p_x \rangle \} = 0.$$

b) Given $\psi'(x) \mapsto e^{\frac{i p_0 x}{\hbar}} \psi(x)$,

$$\begin{aligned}\langle x' \rangle' &= \langle \psi' | x | \psi' \rangle \\ &= \langle e^{-ip_0 x'/\hbar} \psi | x | e^{ip_0 x/\hbar} \psi \rangle \\ &= \langle \psi | e^{-ip_0 x'/\hbar} x e^{ip_0 x/\hbar} | \psi \rangle\end{aligned}$$

since $[e^{i\alpha x}, x] = 0 \quad \forall \alpha$, then

$$\begin{aligned}&= \langle \psi | x | \psi \rangle \\ &= \langle x \rangle\end{aligned}$$

$$\begin{aligned}\langle p_x \rangle &= \langle \psi' | p_x | \psi' \rangle \text{ we insert } \int dx |x\rangle \langle x| \\ &= \int dx \psi^*(x) e^{-ip_0 x/\hbar} p_x \psi(x) e^{ip_0 x/\hbar} \\ &= \int dx \psi^*(x) e^{-ip_0 x/\hbar} (-ih) \frac{d}{dx} (\psi(x) e^{ip_0 x/\hbar}) \\ &= \int dx \psi^*(x) e^{-ip_0 x/\hbar} (-ih) \left\{ \frac{d}{dx} \psi(x) e^{ip_0 x/\hbar} + \left(\frac{ip_0}{\hbar} \right) \psi(x) e^{ip_0 x/\hbar} \right\} \\ &= \int dx \psi^*(x) e^{-ip_0 x/\hbar} (-ih) \frac{d}{dx} \psi(x) e^{ip_0 x/\hbar} \\ &\quad + \int dx \psi^*(x) e^{-ip_0 x/\hbar} p_0 \psi(x) e^{ip_0 x/\hbar} \\ &= \int dx \psi^*(x) (-ih) \frac{d}{dx} \psi(x) + p_0 \int dx \psi^*(x) \psi(x) \\ &= \int dx \psi^*(x) p_x \psi(x) + p_0 \\ &= \langle p \rangle + \langle p_0 \rangle.\end{aligned}$$