

Homework 8 PHYS230

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1 Problem 8.40

As an external force is applied to the paper to displace it, friction is exerted on the bowling :

$$\sum F = ma = F_f$$

Using the fact that moment of inertia for a solid sphere is

$$I = \frac{2MR^2}{5}$$

The conservation of torque around the CM of the bowling ball gives:

$$\begin{aligned}\vec{\tau}_{tot} &= I\alpha \\ F_f R &= \frac{2MR^2}{5}\alpha\end{aligned}$$

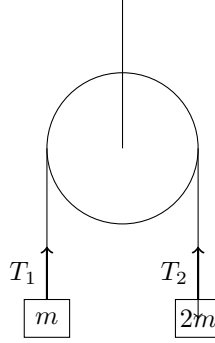
Since $\frac{F_f}{m} = a$, we can substitute and solve for a

$$\begin{aligned}\frac{5F_f}{2m} &= R\alpha \\ \rightarrow \frac{5a}{2} &= R\alpha\end{aligned}$$

Consider a point p at the contact point between the paper and the ball. Then by definition of no-slipping condition,

$$\begin{aligned}\vec{a}_p &= \vec{a}_{CM} + \vec{a}_{CM,p} \\ \Leftrightarrow a_0 &= a + R\omega \\ \Rightarrow a_0 &= a + \frac{5a}{2} = \frac{7a}{2} \\ \Rightarrow a &= \frac{2a_0}{7}\end{aligned}$$

2 Problem 8.44



Since $m_2 \geq m_1$, and acceleration is same for both objects (conservation of string), we'll assume the acceleration on the right side is positive and that on the left side is negative. Thus, $F = ma$ on both sides yields

$$\begin{aligned} T_1 - mg &= ma \\ T_2 - 2mg &= 2m(-a) \rightarrow 2mg - T_2 = 2ma \end{aligned}$$

Moment of inertia for a pulley of radius R rotating along its center is defined as $I = \frac{mR^2}{2}$. Therefore using conservation of torque we get

$$\begin{aligned} -T_1 r + T_2 r - I\alpha &= 0 \\ \rightarrow (T_2 - T_1)r &= \frac{mr^2\alpha}{2} \\ \rightarrow (T_2 - T_1) &= \frac{m(r\alpha)}{2} \end{aligned}$$

Since we have a no slipping condition, $v = R\omega \implies a = R\alpha$, so we can substitute this result

$$\begin{aligned} (T_2 - T_1) &= \frac{ma}{2} \\ \text{Therefore, } T_1 &= T_2 - \frac{ma}{2} \end{aligned}$$

Substituting the previous equation in the initial $F = ma$ equation gives

$$\begin{aligned} T_1 = ma + mg &= T_2 - \frac{ma}{2} \\ \rightarrow T_2 &= ma + mg + \frac{ma}{2} \\ \rightarrow 2mg - 2ma - mg &= \frac{3ma}{2} \\ \rightarrow mg = \frac{7ma}{2} &\implies \boxed{a = \frac{2g}{7}} \end{aligned}$$

3 Problem 11.29

Since Δx and t' are known, we can directly find v . Relativity equations are symmetric between light years and light seconds therefore

$$\Delta x = c \quad \text{and} \quad t' = 1$$

$$t = \gamma t' \\ \rightarrow v = \frac{\Delta x}{\gamma t'} = \frac{\Delta x \sqrt{1 - \frac{v^2}{c^2}}}{t'}$$

$$v^2(t')^2 = (\Delta x)^2 \left[1 - \frac{v^2}{c^2} \right] \\ = (\Delta x)^2 - \frac{(\Delta x v)^2}{c^2}$$

Using $\Delta x = c$ and $t' = 1$,

$$v^2 + \frac{c^2 v^2}{c^2} = c^2 \implies \boxed{v = \frac{c}{\sqrt{2}}}$$

4 Problem 11.32

4.1 a)

The length of the train in the ground's frame is L so the relativistic L_0 is

$$L_0 = \frac{L_0}{\gamma} = L \sqrt{1 - \frac{v^2}{c^2}} = L \sqrt{1 - \frac{(3c/5)^2}{c^2}} \\ \rightarrow \frac{4L}{5} = L_0$$

The train travels a distance of L but also it's whole length $\frac{4L}{5}$, so in total a distance of $\frac{9L}{5}$, thus the time it takes in the ground frame is

$$\frac{\Delta x}{\Delta t} = v \implies \frac{\Delta x}{v} = \Delta t \\ \rightarrow \frac{\left(\frac{9L}{5}\right)}{\left(\frac{3c}{5}\right)} = t = \boxed{\frac{3L}{c}}$$

4.2 b)

The person travels a total distance of L at a speed $\frac{3L}{c}$ with respect to the ground frame, thus

$$\frac{\Delta x}{\Delta t} = v \implies \frac{L}{\left(\frac{3L}{c}\right)} = \boxed{\frac{c}{3}}$$

4.3 c)

Since the person is travelling at a speed of $\frac{c}{3}$ with respect to the ground over a distance L

$$\Delta t = \frac{\Delta x}{v} = \frac{L}{\left(\frac{c}{3}\right)} = \frac{3L}{c}$$

Moreover,

$$\begin{aligned} t' &= \frac{t}{\gamma_{1/3}} = \frac{3L}{c} \sqrt{1 - \frac{(c/3)^2}{c^2}} \\ &= \frac{3L}{c} \frac{\sqrt{8}}{3} = \boxed{\frac{\sqrt{8}L}{c}} \end{aligned}$$